

# Neutrino Physics: Lecture 13

## Interactions and masses in the Standard Model

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- 1 The Standard Model Lagrangian: a reminder
- 2 Mass terms in the SM

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# The electroweak Lagrangian in compact form

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4}\vec{\mathbf{W}}_{\mu\nu} \cdot \vec{\mathbf{W}}^{\mu\nu} - \frac{1}{4}B_{\mu\nu}B^{\mu\nu} \\ & + \bar{L}\gamma^\mu \left( i\partial_\mu - \frac{g}{2}\vec{\mathbf{W}}_\mu \cdot \vec{\tau} - \frac{g'}{2}B_\mu Y \right) L \\ & + \bar{R}\gamma^\mu \left( i\partial_\mu - \frac{g'}{2}B_\mu Y \right) R \\ & + \left| \left( i\partial_\mu - \frac{g}{2}\vec{\mathbf{W}}_\mu \cdot \vec{\tau} - \frac{g'}{2}B_\mu Y \right) \Phi \right|^2 - V(\Phi) \\ & - (\lambda_d \bar{L}\Phi R + \lambda_u \bar{L}\Phi^C R + H.c.)\end{aligned}$$

# Only particles, no interactions

$$\begin{aligned}\mathcal{L} = & \\ & + \bar{L} \gamma^\mu \left( i \partial_\mu \right) L \\ & + \bar{R} \gamma^\mu \left( i \partial_\mu \right) R \\ & + \left| \left( i \partial_\mu \right) \Phi \right|^2 - V(\Phi)\end{aligned}$$

# Adding the weak hypercharge $U(1)_Y$

## $U(1)_Y$ gauge symmetry

$$\psi(x) \rightarrow e^{i \frac{g'}{2} \alpha(x) Y} \psi(x)$$

$$B_\mu \rightarrow B_\mu - \frac{g'}{2} (\partial_\mu \alpha)$$

$$\partial_\mu \rightarrow \partial_\mu + i \frac{g'}{2} B_\mu Y$$

## Y eigenvalues

$$\begin{pmatrix} \nu_{eL} \\ e_L \end{pmatrix} : -1, \quad e_R : -2, \quad \begin{pmatrix} u_L \\ d_L \end{pmatrix} : \frac{1}{3}, \quad u_R : \frac{4}{3}, \quad d_R : -\frac{2}{3}, \quad \Phi : 1$$

## The new gauge boson $B_\mu$

- No mass term
- Kinetic term  $(-1/4) B_{\mu\nu} B^{\mu\nu}$  ( $B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu$ )

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# SM Lagrangian with $U(1)_Y$

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} \\ & +\bar{L}\gamma^\mu\left(i\partial_\mu - \frac{g'}{2}B_\mu Y\right)L \\ & +\bar{R}\gamma^\mu\left(i\partial_\mu - \frac{g'}{2}B_\mu Y\right)R \\ & +\left|\left(i\partial_\mu - \frac{g'}{2}B_\mu Y\right)\Phi\right|^2 - V(\Phi)\end{aligned}$$

# Adding the weak isospin $SU(2)_L$

## $SU(2)_L$ gauge symmetry

$$\psi(x) \rightarrow e^{i \frac{g}{2} \vec{\alpha}(x) \cdot \vec{\tau}} \psi(x)$$

$$\vec{W}_\mu \rightarrow \vec{W}_\mu - \frac{g}{2} (\partial_\mu \vec{\alpha})$$

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## $SU(2)_L$ eigenvalues ( $T_3$ )

$$+\frac{1}{2} : \nu_{eL}, u_L, \Phi^+, \quad -\frac{1}{2} : e_L, d_L, \Phi^0, \quad 0 : e_R, u_R, d_R$$

## The new vector gauge boson $\vec{W}_\mu$

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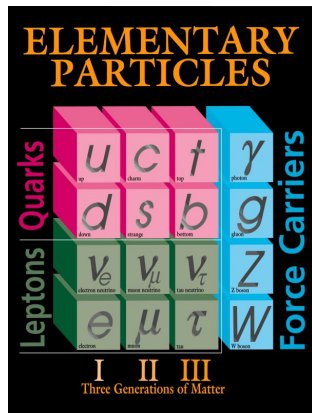
$$(\vec{W}_{\mu\nu} = \partial_\mu \vec{W}_\nu - \partial_\nu \vec{W}_\mu - g \vec{W}_\mu \times \vec{W}_\nu)$$

# SM Lagrangian with $SU(2)_L \times U(1)_Y$

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4}\vec{W}_{\mu\nu} \cdot \vec{W}^{\mu\nu} - \frac{1}{4}B_{\mu\nu}B^{\mu\nu} \\ & + \bar{L}\gamma^\mu \left( i\partial_\mu - \frac{g}{2}\vec{W}_\mu \cdot \vec{\tau} - \frac{g'}{2}B_\mu Y \right) L \\ & + \bar{R}\gamma^\mu \left( i\partial_\mu - \frac{g'}{2}B_\mu Y \right) R \\ & + \left| \left( i\partial_\mu - \frac{g}{2}\vec{W}_\mu \cdot \vec{\tau} - \frac{g'}{2}B_\mu Y \right) \Phi \right|^2 - V(\Phi)\end{aligned}$$

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# Only possible interactions: Yukawa



- $q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix} : (2, \frac{1}{3})$
- $u_R : (1, \frac{4}{3}), d_R : (1, -\frac{2}{3})$
- $l_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix} : (2, -1)$
- $e_R : (1, -2)$
- $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} : (2, 1)$
- $\Phi^C = i\tau^2 \Phi^* = \begin{pmatrix} \phi_0^* \\ -\phi_- \end{pmatrix} : (2, -1)$

- Only possible interaction terms:  
 $\bar{q}_L \Phi^c u_R, \quad \bar{q}_L \Phi d_R, \quad \bar{l}_L \Phi e_R \quad (+ \text{h.c.})$

expand...

- **No masses yet...**

# Mass terms for bosons and fermions

## Equation of motion from lagrangian

$$\partial_\alpha \left( \frac{\partial \mathcal{L}}{\partial (\partial_\alpha \Psi)} \right) = \frac{\partial \mathcal{L}}{\partial \Psi}$$

## Complex scalar field

- Lagrangian:  $\mathcal{L} = (\partial_\mu \phi)^\dagger (\partial^\mu \phi) - m^2 \phi^2$
- Equation of motion:  $\partial_\alpha \partial^\alpha \phi - m^2 \phi = 0$  Klein-Gordon

## Gauge bosons

- Lagrangian:  $\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m^2 A_\mu A^\mu$
- Equation of motion:  $\partial_\alpha \partial^\alpha A_\beta - m^2 A_\beta = 0$  Klein-Gordon

## Fermions

- Lagrangian:  $\mathcal{L} = \bar{\psi} \gamma^\mu i \partial_\mu \psi - m \bar{\psi} \psi$
- Equation of motion:  $i \partial_\alpha \gamma^\alpha \psi - m \psi = 0$  Dirac



# Gauge boson and Fermion mass terms forbidden

## Gauge bosons

- $\frac{1}{2} m_B^2 B_\mu B^\mu$  violates  $U(1)_Y$
- $\frac{1}{2} m_W^2 \vec{W}_\mu \cdot \vec{W}^\mu$  violates  $SU(2)_L$

## Fermion masses

- $m_u \bar{u}_L u_R, m_d \bar{d}_L d_R, m_e \bar{e}_L e_R$
- Violate both,  $SU(2)_L$  as well as  $U(1)_Y$

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# Electroweak symmetry breaking

$$V(\Phi) = \mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1 + i\phi_2 \\ \phi_3 + i\phi_4 \end{pmatrix}$$

$$\mu^2 < 0 \Rightarrow V(\Phi)_{min} \text{ at } \mu^2 = -2\lambda \Phi^\dagger \Phi$$

Minima on the surface of a 3-sphere:

$$\Phi^\dagger \Phi = \frac{1}{2} (\phi_1^2 + \phi_2^2 + \phi_3^2 + \phi_4^2)$$

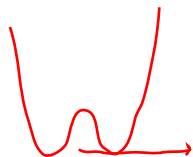
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# Symmetry broken: $SU(2)_L \times U(1)_Y \rightarrow U(1)_Q$

$$\Phi_{\text{vac}} = \begin{pmatrix} 0 \\ v \end{pmatrix} \frac{1}{\sqrt{2}}$$

$$\phi = \begin{pmatrix} 0 \\ \underline{v} + h(x) \end{pmatrix} \frac{1}{\sqrt{2}}$$

$v \neq 0$  in vacuum  $\Rightarrow$

- $T_3$  not conserved [ $T_3(\phi^0) = -1/2$ ]
- $Y$  not conserved [ $Y(\phi^0) = 1$ ]
- However, conserved quantity in new vacuum:

$$T_3(\phi^0) + \frac{Y(\phi^0)}{2} = 0$$

- Electric charge  $Q = T_3 + \frac{Y}{2}$  conserved by vacuum

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# Terms relevant for EWSB

$$\mathcal{L} = -\frac{1}{4}\vec{\mathbf{W}}_{\mu\nu} \cdot \vec{\mathbf{W}}^{\mu\nu} - \frac{1}{4}B_{\mu\nu} \cdot B^{\mu\nu}$$

$$+ \left| \left( -\frac{g}{2}\vec{\mathbf{W}}_{\mu} \cdot \vec{\tau} - \frac{g'}{2}B_{\mu} Y \right) \Phi \right|^2 - V(\Phi)_{min}$$

$$\downarrow$$
$$\Phi = \begin{pmatrix} 0 \\ v \end{pmatrix} \frac{1}{\sqrt{2}}$$

# Mass for $W^+$ and $W^-$

$$\begin{aligned} & \left| \left( -\frac{g}{2} \vec{W}_\mu \cdot \vec{\tau} - \frac{g'}{2} B_\mu Y \right) \Phi \right|^2 \\ &= \frac{1}{8} \left| \begin{pmatrix} gW_3 + g'B & g(W_1 + iW_2) \\ g(W_1 - iW_2) & -gW_3 + g'B \end{pmatrix} \begin{pmatrix} 0 \\ v \end{pmatrix} \right|^2 \\ &= \frac{1}{2} \left( \frac{1}{2}gv \right)^2 |W_1|^2 + \frac{1}{2} \left( \frac{1}{2}gv \right)^2 |W_2|^2 \\ & \quad + \frac{1}{8}v^2 \begin{pmatrix} W_3 & B \end{pmatrix} \begin{pmatrix} g^2 & -gg' \\ -gg' & g'^2 \end{pmatrix} \begin{pmatrix} W_3 \\ B \end{pmatrix} \end{aligned}$$

Diagonalize  $\Rightarrow$

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Diagonalize  $\Rightarrow$

# Mass for $Z$ , no mass for $A$

- Eigenvalues:  $0, (g^2 + g'^2)$

$$\Rightarrow \frac{1}{8}v^2 \begin{pmatrix} A & Z \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & (g^2 + g'^2) \end{pmatrix} \begin{pmatrix} A \\ Z \end{pmatrix}$$

- Eigenvectors:

$$\begin{pmatrix} A \\ Z \end{pmatrix} = \begin{pmatrix} \cos \theta_W & \sin \theta_W \\ -\sin \theta_W & \cos \theta_W \end{pmatrix} \begin{pmatrix} B \\ W_3 \end{pmatrix}$$

- Mixing angle:  $\tan 2\theta_W = 2gg'/(g^2 - g'^2) \Rightarrow \tan \theta_W = g'/g$

$$M_W = \frac{1}{2}gv, \quad M_Z = \frac{1}{2}v\sqrt{g^2 + g'^2}, \quad \rho = \frac{M_W^2}{M_Z^2 \cos^2 \theta_W} = 1$$

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*↓  
Weinberg*

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# Fermion masses in SM: Dirac masses

Interaction terms:

$$- \lambda_u \bar{q}_L \Phi^c u_R, \quad - \lambda_d \bar{q}_L \Phi d_R, \quad - \lambda_e \bar{l}_L \Phi e_R$$

After electroweak symmetry breaking:

$$(\lambda_u v) \bar{u}_L u_R, \quad (\lambda_d v) \bar{d}_L d_R, \quad (\lambda_e v) \bar{e}_L e_R$$

Dirac masses

$$m_u = \lambda_u v, \quad m_d = \lambda_d v, \quad m_e = \lambda_e v$$



# Fermion masses in SM: Dirac masses

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After electroweak symmetry breaking:

$$- \frac{(\lambda_u v)}{\sqrt{2}} \bar{u}_L u_R, \quad - \frac{(\lambda_d v)}{\sqrt{2}} \bar{d}_L d_R, \quad - \frac{(\lambda_e v)}{\sqrt{2}} \bar{e}_L e_R$$

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$$\frac{(\lambda_u V)}{\sqrt{2}} \bar{u}_L u_R, \quad \frac{(\lambda_d V)}{\sqrt{2}} \bar{d}_L d_R, \quad \frac{(\lambda_e V)}{\sqrt{2}} \bar{e}_L e_R$$

Dirac masses

$$m_u = \frac{\lambda_u V}{\sqrt{2}}, \quad m_d = \frac{\lambda_d V}{\sqrt{2}}, \quad m_e = \frac{\lambda_e V}{\sqrt{2}}$$

# No mass for neutrinos yet !