

Neutrino Physics: Lecture 9

Three flavour mixing

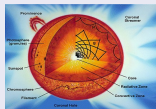
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Mar 22, 2010

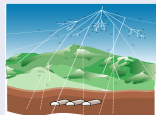
Two disconnected two-flavour mixings

Solar neutrino parameters



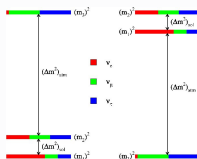
$$\Delta m_{\odot}^2 \approx \underline{8 \times 10^{-5} \text{ eV}^2}, \theta_{\odot} \approx 32^\circ \checkmark$$

Atmospheric neutrino parameters



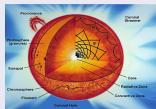
$$\Delta m_{\text{atm}}^2 \approx \underline{2.5 \times 10^{-3} \text{ eV}^2}, \theta_{\text{atm}} \approx 45^\circ \checkmark$$

Normal vs. inverted mass ordering / mass hierarchy



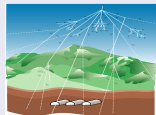
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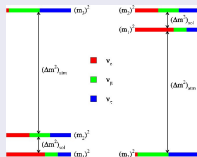
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$$\Delta m_{\text{atm}}^2 \approx 2.5 \times 10^{-3} \text{ eV}^2, \theta_{\text{atm}} \approx 45^\circ$$

Normal vs. inverted mass ordering / mass hierarchy



- 1 General formalism for three flavours
- 2 Matching with atmospheric and solar solution

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Flavour eigenstates and mass eigenstates

$$\left. \begin{array}{l} \nu_e \\ \nu_\mu \\ \nu_\tau \end{array} \right\} \leftrightarrow \left\{ \begin{array}{l} \nu_1 \\ \nu_2 \\ \nu_3 \end{array} \right.$$

Definitions of flavour eigenstates

- ν_e : Interacts to give e^-
- ν_μ : Interacts to give μ^-
- ν_τ : Interacts to give τ^-

Definitions of mass eigenstates

- ν_3 : The state that does not participate in solar oscillations
(the state “further away”)
- ν_2 : The heavier state involved in solar oscillations
- ν_1 : The lighter state involved in solar oscillations

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Assigning Δm^2 values from available data

Solar data

$$\Delta m_{\odot}^2 = \underline{m_2^2 - m_1^2} \equiv \Delta m_{21}^2 \quad \checkmark$$

- $\Delta m_{21}^2 > 0$ by definition

Atmospheric data

$$\Delta m_{atm}^2 \equiv m_3^2 - (m_2^2 + m_1^2)/2$$

- $\Delta m_{31}^2 \equiv m_3^2 - m_1^2$
- $\Delta m_{32}^2 \equiv m_3^2 - m_2^2$
- $\Delta m_{31}^2 \approx \Delta m_{32}^2 \approx \Delta m_{atm}^2$
- Sign of Δm_{atm}^2 unknown:
Measured quantity $P_{\mu\mu} = 1 - \sin^2 2\theta \sin^2(\Delta m_{atm}^2 L / (4E))$

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The mixing matrix

$$\nu_\alpha = \sum_i U_{\alpha i} \nu_i$$

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

PMNS matrix: elements in general complex

$$\nu_e = U_{e1} \nu_1 + U_{e2} \nu_2 + U_{e3} \nu_3$$

Vacuum oscillations with three neutrinos

Starting with a flavour eigenstate ν_α :

$$|\nu_\alpha(0)\rangle = U_{\alpha 1}|\nu_1\rangle + U_{\alpha 2}|\nu_2\rangle + U_{\alpha 3}|\nu_3\rangle$$

After propagation:

$$|\nu_\alpha(t)\rangle = U_{\alpha 1}|\nu_1\rangle e^{-i\frac{m_1^2 L}{2E}} + U_{\alpha 2}|\nu_2\rangle e^{-i\frac{m_2^2 L}{2E}} + U_{\alpha 3}|\nu_3\rangle e^{-i\frac{m_3^2 L}{2E}}$$

Projection on flavour eigenstate $|\nu_\beta\rangle$:

$$\langle\nu_\beta|\nu_\alpha(t)\rangle = U_{\beta 1}^* U_{\alpha 1} e^{-i\frac{m_1^2 L}{2E}} + U_{\beta 2}^* U_{\alpha 2} e^{-i\frac{m_2^2 L}{2E}} + U_{\beta 3}^* U_{\alpha 3} e^{-i\frac{m_3^2 L}{2E}}$$

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Conversion / survival probability

$$\langle \nu_\beta | \nu_\alpha(t) \rangle = U_{\beta 1}^* U_{\alpha 1} e^{-i \frac{m_1^2 L}{2E}} + U_{\beta 2}^* U_{\alpha 2} e^{-i \frac{m_2^2 L}{2E}} + U_{\beta 3}^* U_{\alpha 3} e^{-i \frac{m_3^2 L}{2E}}$$

$P_{\nu_\alpha \rightarrow \nu_\beta} \rightarrow P_{\alpha\beta} \equiv |\langle \nu_\beta | \nu_\alpha(t) \rangle|^2$

$$= |U_{\beta 1}^* U_{\alpha 1}|^2 + |U_{\beta 2}^* U_{\alpha 2}|^2 + |U_{\beta 3}^* U_{\alpha 3}|^2$$

$\square_{\alpha\beta 12}$ \rightarrow $+ U_{\beta 1}^* U_{\alpha 2} U_{\beta 2} U_{\alpha 1} e^{i \frac{m_2^2 - m_1^2}{2E} L} + \text{c.c.}$ $\rightarrow 2 \Delta_{21}$

$$+ U_{\beta 1}^* U_{\alpha 3} U_{\beta 3} U_{\alpha 1} e^{i \frac{m_3^2 - m_1^2}{2E} L} + \text{c.c.}$$
$$+ U_{\beta 2}^* U_{\alpha 3} U_{\beta 3} U_{\alpha 2} e^{i \frac{m_3^2 - m_2^2}{2E} L} + \text{c.c.}$$

$P_{\alpha\beta}$: further simplification

Plaquette \longrightarrow $\square_{\alpha\beta ij} \equiv U_{\alpha i} U_{\beta j} U_{\alpha j}^* U_{\beta i}^*$

$$P_{\alpha\beta} = |U_{\beta 1}^* U_{\alpha 1}|^2 + |U_{\beta 2}^* U_{\alpha 2}|^2 + |U_{\beta 3}^* U_{\alpha 3}|^2 \\ + 2\text{Re}(\square_{\alpha\beta 12} e^{2i\Delta_{21}}) + 2\text{Re}(\square_{\alpha\beta 13} e^{2i\Delta_{31}}) + 2\text{Re}(\square_{\alpha\beta 23} e^{2i\Delta_{32}})$$

$$P_{\alpha\beta} = |U_{\beta 1}^* U_{\alpha 1}|^2 + |U_{\beta 2}^* U_{\alpha 2}|^2 + |U_{\beta 3}^* U_{\alpha 3}|^2 \\ + \sum_{i < j} 2\text{Re}(\square_{\alpha\beta ij}) \cos(2\Delta_{ji}) - \sum_{i < j} 2\text{Im}(\square_{\alpha\beta ij}) \sin(2\Delta_{ji})$$

$$P_{\alpha\beta} = \delta_{\alpha\beta} - \sum_{i < j} 4\text{Re}(\square_{\alpha\beta ij}) \sin^2(\Delta_{ji}) - 2 \sum_{i < j} \text{Im}(\square_{\alpha\beta ij}) \sin(2\Delta_{ji})$$

$P_{\alpha\beta}$: further simplification

$$\square_{\alpha\beta ij} \equiv U_{\alpha i} U_{\beta j} U_{\alpha j}^* U_{\beta i}^*$$

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Features of $P_{\alpha\beta}$

$$P_{\alpha\beta} = \delta_{\alpha\beta} - \sum_{i<j} 4\text{Re}(\square_{\alpha\beta ij}) \sin^2(\Delta_{ji}) - 2 \sum_{i<j} \text{Im}(\square_{\alpha\beta ij}) \sin(2\Delta_{ji})$$

$$\square_{\alpha\beta ij} \equiv U_{\alpha i} U_{\beta j} U_{\alpha j}^* U_{\beta i}^*$$

$$\Delta_{ji} \equiv \frac{(m_j^2 - m_i^2)}{4E}$$

- No-oscillation term
- CP-conserving oscillation term
- CP-violating oscillation term

Parameterizing U_{PMNS} with three mixing angles

Neglect CP violation \Rightarrow Real U_{PMNS} matrix

$$\begin{aligned} U &= R_{23}(\theta_{23})R_{13}(\theta_{13})R_{12}(\theta_{12}) \\ &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13} \\ 0 & 1 & 0 \\ -s_{13} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13} & c_{12}c_{23} - s_{12}s_{23}s_{13} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13} & -c_{12}s_{23} - s_{12}c_{23}s_{13} & c_{23}c_{13} \end{pmatrix} \end{aligned}$$

$$c_j \equiv \cos \theta_j, s_j \equiv \sin \theta_j$$

- 1 General formalism for three flavours
- 2 Matching with atmospheric and solar solution

Reactor neutrino experiments

$\bar{\nu}_e \rightarrow \bar{\nu}_e$ with $\Delta_{32} \sim 1, \Delta_{21} \ll 1$

$$\Delta m_{atm}^2 / 4E$$

$$\square_{ee12} = U_{e1}U_{e2}U_{e3}^*U_{e1}^* \text{ etc}$$

$$P_{ee} = 1 - \cancel{4|U_{e1}|^2|U_{e2}|^2 \sin^2(\Delta_{21})} - 4|U_{e2}|^2|U_{e3}|^2 \sin^2(\Delta_{32}) - 4|U_{e1}|^2|U_{e3}|^2 \sin^2(\Delta_{31})$$

$$\approx \neq 1 - 4|U_{e3}|^2 \sin^2(\Delta_{31}) (|U_{e2}|^2 + |U_{e1}|^2)$$

$$P_{ee} \approx 1 - 4|U_{e3}|^2(1 - |U_{e3}|^2) \sin^2(\Delta_{31})$$

$$P_{ee} \approx 1 - \sin^2 2\theta_{13} \sin^2(\Delta_{31})$$

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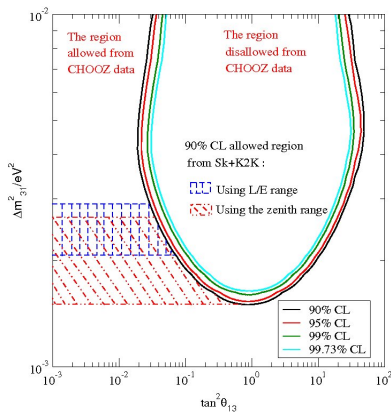
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$$P_{ee} = 1 - 4|U_{e1}|^2|U_{e2}|^2 \sin^2(\Delta_{21}) \\ - 4|U_{e2}|^2|U_{e3}|^2 \sin^2(\Delta_{32}) - 4|U_{e1}|^2|U_{e3}|^2 \sin^2(\Delta_{31})$$

$$P_{ee} \approx 1 - \underbrace{4|U_{e3}|^2(1 - |U_{e3}|^2)}_{4s_{13}^2 c_{13}^2} \sin^2(\Delta_{31})$$

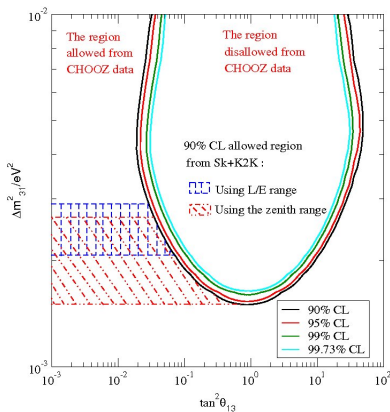
$$P_{ee} \approx 1 - \sin^2 2\theta_{13} \sin^2(\Delta_{31})$$

Results from CHOOZ



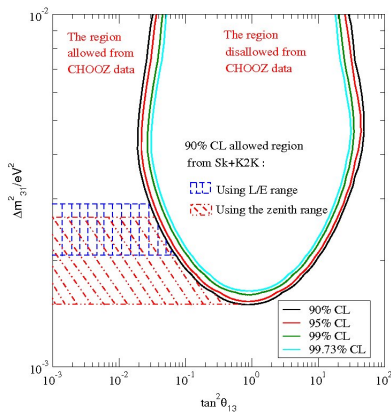
- $E_\nu \approx 3 \text{ MeV}, L \approx 1000 \text{ m} \Rightarrow \Delta_{32} \sim 1, \Delta_{21} \ll 1$
- Disappearance experiment: null result (no disappearance)
- $|U_{e3}|^2(1 - |U_{e3}|^2) < 0.05$, consistent with zero
- $|U_{e3}|^2 < 0.05$ (3σ) or $|U_{e3}|^2 > 0.95$ (3σ)
- $\sin^2 \theta_{13} < 0.05$ or $\sin^2 \theta_{13} > 0.95$

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Atmospheric neutrino conversion probability

$$U_{\mu\mu 12} = U_{\mu 1} U_{\mu 2} U_{\mu 2}^* U_{\mu 1}^* \text{ etc}$$

$$\underline{\underline{P_{\mu\mu}}} = 1 - \cancel{4|U_{\mu 1}|^2|U_{\mu 2}|^2 \sin^2(\Delta_{21})} - 4|U_{\mu 2}|^2|U_{\mu 3}|^2 \sin^2(\Delta_{32}) - 4|U_{\mu 1}|^2|U_{\mu 3}|^2 \sin^2(\Delta_{31})$$

$$\approx 1 - 4|U_{\mu 3}|^2 \sin^2(\Delta_{31}) (|U_{\mu 2}|^2 + |U_{\mu 1}|^2)$$

$$P_{\mu\mu} \approx 1 - 4|U_{\mu 3}|^2 (1 - |U_{\mu 3}|^2) \sin^2(\Delta_{31})$$

$$\approx 1 - 4c_{13}^2 s_{23}^2 (1 - c_{13}^2 s_{23}^2) \sin^2(\Delta_{31})$$

When $\sin \theta_{13} \approx 0$: $P_{\mu\mu} \approx 1 - \sin^2 2\theta_{23} \sin^2(\Delta m_{32}^2 L/4E)$

When $\sin \theta_{13} \approx 1$: $P_{\mu\mu} \approx 1 \Rightarrow$ No oscillations

$$\theta_{23} \approx \theta_{atm}, \quad \theta_{13} \approx 0$$

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$$\xrightarrow{\quad} 4s_{23}^2 c_{23}^2$$

When $\sin \theta_{13} \approx 0$: $P_{\mu\mu} \approx 1 - \sin^2 2\theta_{23} \sin^2(\Delta m_{32}^2 L/4E)$ ✓

When $\sin \theta_{13} \approx 1$: $P_{\mu\mu} \approx 1 \Rightarrow$ No oscillations ✗

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When $\sin \theta_{13} \approx 1$: $P_{\mu\mu} \approx 1 \Rightarrow$ No oscillations

$$\theta_{23} \approx \theta_{atm}, \quad \theta_{13} \approx 0$$

Solar neutrino problem with three flavours

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = R_{23}(\theta_{23}) R_{13}(0) R_{12}(\theta_{12}) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

New basis:

$$\begin{pmatrix} \nu_e \\ \nu_x \\ \nu_y \end{pmatrix} = R_{23}^\dagger(\theta_{23}) \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

$$\begin{pmatrix} \nu_e \\ \nu_x \\ \nu_y \end{pmatrix} = R_{23}^\dagger(\theta_{23}) R_{23}(\theta_{23}) R_{12}(\theta_{12}) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} = R_{12}(\theta_{12}) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

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New basis:

$$\begin{pmatrix} \nu_e \\ \nu_x \\ \nu_y \end{pmatrix} = R_{23}^\dagger(\theta_{23}) \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

$$\begin{pmatrix} \nu_e \\ \nu_x \\ \nu_y \end{pmatrix} = R_{23}^\dagger(\theta_{23}) R_{23}(\theta_{23}) R_{12}(\theta_{12}) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} = R_{12}(\theta_{12}) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Reducing solar neutrino problem to two flavours

$$\begin{pmatrix} \nu_e \\ \nu_x \\ \nu_y \end{pmatrix} = \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

- $\nu_3 = \nu_y$ decouples
- Two neutrino mixing:

$$\begin{pmatrix} \nu_e \\ \nu_x \end{pmatrix} = \begin{pmatrix} c_{12} & s_{12} \\ -s_{12} & c_{12} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

- Mixing angle θ_{12}

$$\theta_{12} = \theta_{\odot} \text{ (when } \theta_{13} = 0 \text{)}$$

Solar oscillations are $\nu_e \leftrightarrow \nu_x$

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- Mixing angle θ_{12}

✓ $\theta_{12} = \theta_{\odot}$ (when $\theta_{13} = 0$)

✓ Solar oscillations are $\nu_e \leftrightarrow \nu_x$

Net three-flavour picture (No CP violation)

Mass squared differences

- $\Delta m_{21}^2 = \Delta m_{\odot}^2 \approx 8 \times 10^{-5} \text{ eV}^2$ ✓
- $\Delta m_{31}^2 \approx \Delta m_{32}^2 \approx \Delta m_{atm}^2 \approx \pm 2.5 \times 10^{-3} \text{ eV}^2$ ✓

Mixing angles

- $\nu_{\alpha} = R_{23}(\theta_{23})R_{13}(\theta_{13})R_{12}(\theta_{12}) \nu_i$
- $\theta_{13} \approx 0$
- $\theta_{23} \approx \theta_{atm} \approx 45^{\circ}$
- $\theta_{12} \approx \theta_{\odot} \approx 32^{\circ}$

Neutrino mass-flavour spectrum

