

CP violation and the third generation of quarks

Nobel Prize in Physics, 2008

Amol Dighe

Department of Theoretical Physics, TIFR



Makato Kobayashi



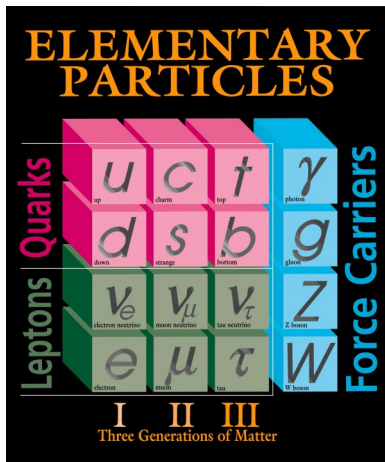
Toshihide Maskawa

*... for the discovery of the origin of the broken symmetry
which predicts the existence of at least three families of quarks in nature*

TIFR Colloquium, Nov 20, 2008

- 1 Before Kobayashi-Maskawa
 - The status of Particle Physics in 1972
 - The broken symmetry: charge-parity (CP)
- 2 The insight of Kobayashi and Maskawa
 - Prog. Theo. Phys. 49, 652 (1973)
 - In modern language: the CKM paradigm
- 3 Testing the Kobayashi-Maskawa predictions
 - The third generation
 - CP violation through the CKM mechanism
- 4 Concluding remarks

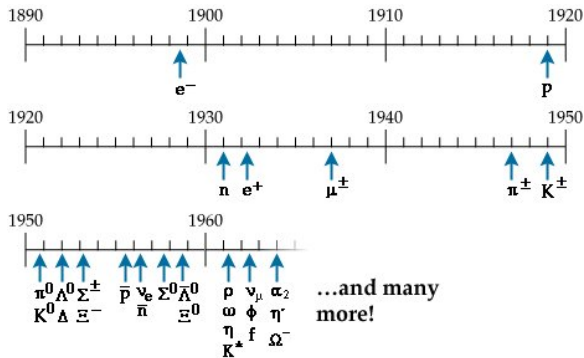
Standard model of particle physics: 2008



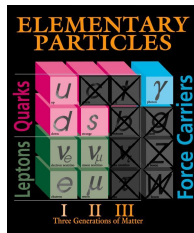
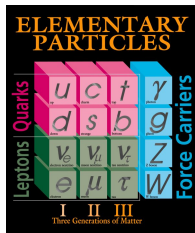
- Three families of quarks and leptons
- quarks:
 - up charmed top
 - down strange bottom
- Mixing between families

- 1 Before Kobayashi-Maskawa
 - The status of Particle Physics in 1972
 - The broken symmetry: charge-parity (CP)
- 2 The insight of Kobayashi and Maskawa
 - Prog. Theo. Phys. 49, 652 (1973)
 - In modern language: the CKM paradigm
- 3 Testing the Kobayashi-Maskawa predictions
 - The third generation
 - CP violation through the CKM mechanism
- 4 Concluding remarks

Known particles in 1972



- $p \equiv uud$
- $n \equiv udd$
- $\pi^+ \equiv u\bar{d}$
- $\pi^- \equiv d\bar{u}$
- $\pi^0 \equiv (u\bar{u} - d\bar{d})/\sqrt{2}$
- $K^+ \equiv u\bar{s}$
- $K^- \equiv s\bar{u}$
- $K^0 \equiv d\bar{s}$
- $\bar{K}^0 \equiv s\bar{d}$



Universality of weak interactions: Cabibbo angle



Interrelated coupling constants:

- (i) muon decay: $g_{e\mu}$
 $\mu^- \rightarrow \nu_\mu e^- \bar{\nu}_e$
- (ii) neutron decay : g_{ud}
 $n \rightarrow pe^- \bar{\nu}_e$ ($d \rightarrow ue^- \bar{\nu}_e$)
- (ii) kaon decay: g_{us}
 $K^- \rightarrow \pi^0 e^- \bar{\nu}_e$ ($s \rightarrow ue^- \bar{\nu}_e$)

$$|g_{e\mu}|^2 = |g_{ud}|^2 + |g_{us}|^2$$

Universality:

- There is only one coupling constant, $g = g_{e\mu}$
- u quark couples to only one combination of d and s :
 $d' \equiv \cos \theta_c \cdot d + \sin \theta_c \cdot s$
- Cabibbo angle θ_c : the first quark mixing angle

N. Cabibbo, "Unitary Symmetry and Leptonic Decays,"
Phys. Rev. Lett. **10**, 531 (1963)

Suppression of flavor-changing neutral currents

- Cabibbo angle unable to explain why

$$\Gamma(K_L \rightarrow \mu^+ \mu^-) \ll \Gamma(K^+ \rightarrow \mu^+ \nu_\mu)$$

- Possible explanation via another “c” quark:
charge $+2/3$, couples to

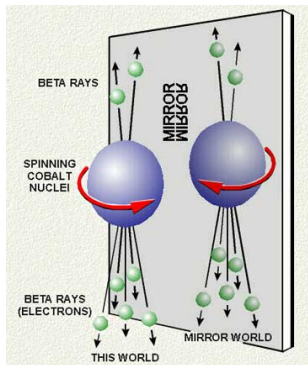
$$s' \equiv -\sin \theta_c \cdot d + \cos \theta_c \cdot c$$

- The $s \rightarrow u \rightarrow d$ and $s \rightarrow c \rightarrow d$ contribution cancel, leading to the suppression of FCNC $s \rightarrow d$
- GIM mechanism: existence of the “charmed” quark.

S. L. Glashow, J. Iliopoulos and L. Maiani,
“Weak Interactions with Lepton-Hadron Symmetry,”
Phys. Rev. D **2**, 1285 (1970)

- 1 Before Kobayashi-Maskawa
 - The status of Particle Physics in 1972
 - **The broken symmetry: charge-parity (CP)**
- 2 The insight of Kobayashi and Maskawa
 - Prog. Theo. Phys. 49, 652 (1973)
 - In modern language: the CKM paradigm
- 3 Testing the Kobayashi-Maskawa predictions
 - The third generation
 - CP violation through the CKM mechanism
- 4 Concluding remarks

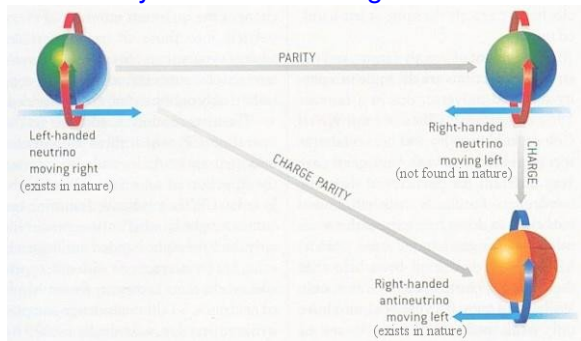
Discovery of parity violation: 1956-57



- Mirror world does not behave the same as the real world
- Theoretical possibility:
T.D.Lee and C.N.Yang,
Phys. Rev. **104**, 254 (1956)
- Experiments: 1957
 - Wu (^{60}Co)
 - Friedman-Telegdi
($\pi^+ \rightarrow \mu^+ \rightarrow e^+$)
- Nobel prize 1957: Lee–Yang

Can Charge \oplus Parity may be conserved ?

Parity: left handed \leftrightarrow right handed



- Neutrinos violate parity: they are only left-handed
- But antineutrinos are right-handed !
- Does that mean **C** and **P** violations cancel each other to give **CP** conservation ?

Prediction of CP violation in K decay

$$K^0 \equiv d\bar{s} \quad \bar{K}^0 \equiv s\bar{d}$$

- CP eigenstates:

$$K_1 \equiv (K^0 + \bar{K}^0)/\sqrt{2} \quad (\text{CP even})$$

$$K_2 \equiv (K^0 - \bar{K}^0)/\sqrt{2} \quad (\text{CP odd})$$

- CP even decay channel: $\pi\pi$

- CP odd decay channel: $\pi\pi\pi$

- CP conservation \Rightarrow

$$K_1 \rightarrow \pi\pi \text{ short-lived, } K_{\text{Short}}$$

$$K_2 \rightarrow \pi\pi\pi \text{ long-lived, } K_{\text{Long}}$$

- Original $K^0 = (K_{\text{Short}} + K_{\text{Long}})/\sqrt{2}$

Prediction of CP violation in K decay

$$K^0 \equiv d\bar{s} \quad \bar{K}^0 \equiv s\bar{d}$$

- CP eigenstates:

$$K_1 \equiv (K^0 + \bar{K})/\sqrt{2} \quad (\text{CP even})$$

$$K_2 \equiv (K^0 - \bar{K})/\sqrt{2} \quad (\text{CP odd})$$

- CP even decay channel: $\pi\pi$
- CP odd decay channel: $\pi\pi\pi$
- CP conservation \Rightarrow

$$K_1 \rightarrow \pi\pi \text{ short-lived, } K_{\text{Short}}$$

$$K_2 \rightarrow \pi\pi\pi \text{ long-lived, } K_{\text{Long}}$$

- Original $K^0 = (K_{\text{Short}} + K_{\text{Long}})/\sqrt{2}$

Prediction of CP violation in K decay

$$K^0 \equiv d\bar{s} \quad \bar{K}^0 \equiv s\bar{d}$$

- CP eigenstates:

$$K_1 \equiv (K^0 + \bar{K})/\sqrt{2} \quad (\text{CP even})$$

$$K_2 \equiv (K^0 - \bar{K})/\sqrt{2} \quad (\text{CP odd})$$

- CP even decay channel: $\pi\pi$
- CP odd decay channel: $\pi\pi\pi$
- CP conservation \Rightarrow

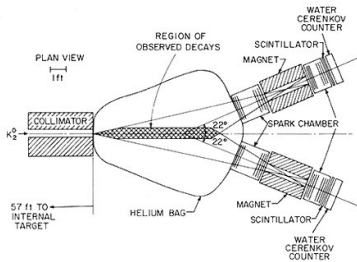
$$K_1 \rightarrow \pi\pi \text{ short-lived, } K_{\text{Short}}$$

$$K_2 \rightarrow \pi\pi\pi \text{ long-lived, } K_{\text{Long}}$$

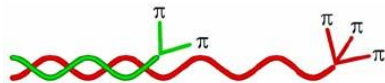
- Original $K^0 = (K_{\text{Short}} + K_{\text{Long}})/\sqrt{2}$

Discovery of CP violation: 1964

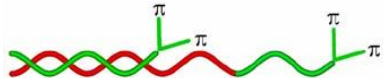
Cronin-Fitch experiment



Nobel prize 1980



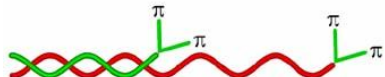
(a) Kaon Mixing



(b) Indirect CP Violation

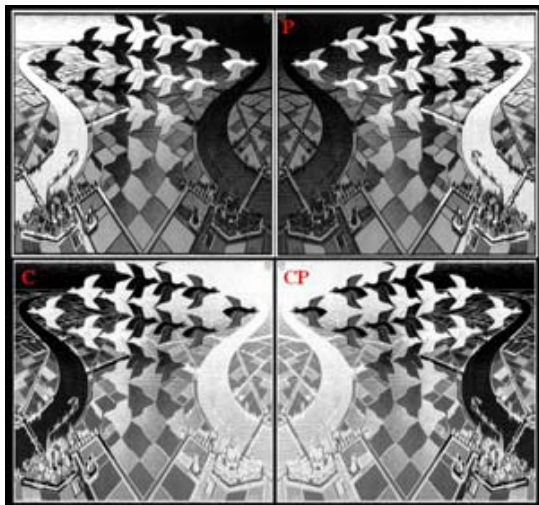


(c) Polarized Light Analogy



(d) Direct CP Violation

Charge-parity violated slightly



“Day and Night”, M.C.Escher

Questions raised by the discovery of CP violation

- Is it small or large ? Is CP an approximate symmetry ?
- Is the symmetry breaking spontaneous ?
- Where does it come from ? Are there extra interactions ?

- 1 Before Kobayashi-Maskawa
 - The status of Particle Physics in 1972
 - The broken symmetry: charge-parity (CP)
- 2 **The insight of Kobayashi and Maskawa**
 - **Prog. Theo. Phys. 49, 652 (1973)**
 - In modern language: the CKM paradigm
- 3 Testing the Kobayashi-Maskawa predictions
 - The third generation
 - CP violation through the CKM mechanism
- 4 Concluding remarks

652

Progress of Theoretical Physics, Vol. 49, No. 2, February 1973

***CP*-Violation in the Renormalizable Theory of Weak Interaction**

Makoto KOBAYASHI and Toshihide MASKAWA

Department of Physics, Kyoto University, Kyoto

(Received September 1, 1972)

In a framework of the renormalizable theory of weak interaction, problems of *CP*-violation are studied. It is concluded that no realistic models of *CP*-violation exist in the quartet scheme without introducing any other new fields. Some possible models of *CP*-violation are also discussed.



What the paper is about

- Consider various ways of putting (u_L, d_L, c_L, s_L) and (u_R, d_R, c_R, s_R) as doublets or singlets of $SU(2)_{\text{weak}}$
($4 = 2+2$, $4=2+1+1$ or $4 = 1+1+1+1$?)
- Experimental data \Rightarrow
 - (u_L, d_L) have to form a doublet: isospin symmetry
 - (c_L, s_L) must also form a doublet: FCNC suppression
- Now, how can one get CP violation ?

CP violation and complex coupling

- CP violation: $A(X \rightarrow Y) \neq A(\bar{X} \rightarrow \bar{Y})$
- If all amplitudes are real, $|A(X \rightarrow Y)|^2 = |A(\bar{X} \rightarrow \bar{Y})|^2$
- CP violation possible if complex numbers involved in

$$\begin{aligned} |A(X \rightarrow Y)|^2 &= |A(X \rightarrow w \rightarrow Y) + A(X \rightarrow z \rightarrow Y)|^2 \\ |A(\bar{X} \rightarrow \bar{Y})|^2 &= |A(\bar{X} \rightarrow \bar{w} \rightarrow \bar{Y}) + A(\bar{X} \rightarrow \bar{z} \rightarrow \bar{Y})|^2 \end{aligned}$$

- CP violation \Rightarrow Amplitudes complex \Rightarrow Couplings complex

Two generations of quarks are not enough

- The mixing matrix between up-type and down-type quarks has to be a 2×2 unitary matrix

$$\mathcal{L} \propto (\overline{u}_L, \overline{c}_L) \begin{pmatrix} \cos \theta e^{i\phi_1} & \sin \theta e^{i\phi_2} \\ -\sin \theta e^{i\phi_3} & \cos \theta e^{i(\phi_2 + \phi_3 - \phi_1)} \end{pmatrix} \begin{pmatrix} d_L \\ s_L \end{pmatrix}$$

- Can change three relative phases of quarks to get rid of all three complex phases ϕ_1, ϕ_2, ϕ_3

$$\mathcal{L} \propto (\overline{u}_L, \overline{c}_L e^{i(\phi_3 - \phi_1)}) \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} d_L e^{i\phi_1} \\ s_L e^{i\phi_2} \end{pmatrix}$$

- Mixing matrix real \Rightarrow no CP violation

Two generations of quarks are not enough

- The mixing matrix between up-type and down-type quarks has to be a 2×2 unitary matrix

$$\mathcal{L} \propto (\overline{u}_L, \overline{c}_L) \begin{pmatrix} \cos \theta e^{i\phi_1} & \sin \theta e^{i\phi_2} \\ -\sin \theta e^{i\phi_3} & \cos \theta e^{i(\phi_2 + \phi_3 - \phi_1)} \end{pmatrix} \begin{pmatrix} d_L \\ s_L \end{pmatrix}$$

- Can change three relative phases of quarks to get rid of all three complex phases ϕ_1, ϕ_2, ϕ_3

$$\mathcal{L} \propto (\overline{u}_L, \overline{c}_L e^{i(\phi_3 - \phi_1)}) \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} d_L e^{i\phi_1} \\ s_L e^{i\phi_2} \end{pmatrix}$$

- Mixing matrix real \Rightarrow no CP violation

Parameter counting for two generations

- 2×2 complex matrix \Rightarrow 4 real + 4 imaginary quantities
- Unitarity $U^\dagger U = I$: 3 real and 1 imaginary conditions
- 1 real and 3 imaginary parameters left
- Can choose the 3 relative phases between quarks to get rid of the 3 imaginary parameters
- The mixing matrix is completely real

Parameter counting for two generations

- 2×2 complex matrix \Rightarrow 4 real + 4 imaginary quantities
- Unitarity $U^\dagger U = I$: 3 real and 1 imaginary conditions
- 1 real and 3 imaginary parameters left
- Can choose the 3 relative phases between quarks to get rid of the 3 imaginary parameters
- The mixing matrix is completely real

Parameter counting for two generations

- 2×2 complex matrix \Rightarrow 4 real + 4 imaginary quantities
- Unitarity $U^\dagger U = I$: 3 real and 1 imaginary conditions
- 1 real and 3 imaginary parameters left
- Can choose the 3 relative phases between quarks to get rid of the 3 imaginary parameters
- The mixing matrix is completely real

Three generations work

- 3×3 matrix \Rightarrow 9 real + 9 imaginary quantities
- Unitarity $U^\dagger U = \mathbf{I}$: 6 real and 3 imaginary conditions
- 3 real and 6 imaginary parameters left
- Can choose the 5 relative phases between quarks to get rid of 5 imaginary parameters
- In addition to 3 real parameters (Euler angles of rotation), one imaginary quantity is unavoidable
- Mixing matrix complex \Rightarrow CP violation may be present

$$\begin{pmatrix} c_1 & -s_1 c_3 & -s_1 s_3 \\ s_1 c_2 & c_1 c_2 c_3 - s_2 s_3 e^{i\delta} & c_1 c_2 s_3 + s_2 c_3 e^{i\delta} \\ s_1 s_2 & c_1 s_2 c_3 + c_2 s_3 e^{i\delta} & c_1 s_2 s_3 - c_2 c_3 e^{i\delta} \end{pmatrix}$$

Three generations work

- 3×3 matrix \Rightarrow 9 real + 9 imaginary quantities
- Unitarity $U^\dagger U = I$: 6 real and 3 imaginary conditions
- 3 real and 6 imaginary parameters left
- Can choose the 5 relative phases between quarks to get rid of 5 imaginary parameters
- In addition to 3 real parameters (Euler angles of rotation), one imaginary quantity is unavoidable
- Mixing matrix complex \Rightarrow CP violation may be present

$$\begin{pmatrix} c_1 & -s_1 c_3 & -s_1 s_3 \\ s_1 c_2 & c_1 c_2 c_3 - s_2 s_3 e^{i\delta} & c_1 c_2 s_3 + s_2 c_3 e^{i\delta} \\ s_1 s_2 & c_1 s_2 c_3 + c_2 s_3 e^{i\delta} & c_1 s_2 s_3 - c_2 c_3 e^{i\delta} \end{pmatrix}$$

Three generations work

- 3×3 matrix \Rightarrow 9 real + 9 imaginary quantities
- Unitarity $U^\dagger U = I$: 6 real and 3 imaginary conditions
- 3 real and 6 imaginary parameters left
- Can choose the 5 relative phases between quarks to get rid of 5 imaginary parameters
- In addition to 3 real parameters (Euler angles of rotation), one imaginary quantity is unavoidable
- Mixing matrix complex \Rightarrow CP violation may be present

$$\begin{pmatrix} c_1 & -s_1 c_3 & -s_1 s_3 \\ s_1 c_2 & c_1 c_2 c_3 - s_2 s_3 e^{i\delta} & c_1 c_2 s_3 + s_2 c_3 e^{i\delta} \\ s_1 s_2 & c_1 s_2 c_3 + c_2 s_3 e^{i\delta} & c_1 s_2 s_3 - c_2 c_3 e^{i\delta} \end{pmatrix}$$

Three generations work

- 3×3 matrix \Rightarrow 9 real + 9 imaginary quantities
- Unitarity $U^\dagger U = I$: 6 real and 3 imaginary conditions
- 3 real and 6 imaginary parameters left
- Can choose the 5 relative phases between quarks to get rid of 5 imaginary parameters
- In addition to 3 real parameters (Euler angles of rotation), one imaginary quantity is unavoidable
- Mixing matrix complex \Rightarrow CP violation may be present

$$\begin{pmatrix} c_1 & -s_1 c_3 & -s_1 s_3 \\ s_1 c_2 & c_1 c_2 c_3 - s_2 s_3 e^{i\delta} & c_1 c_2 s_3 + s_2 c_3 e^{i\delta} \\ s_1 s_2 & c_1 s_2 c_3 + c_2 s_3 e^{i\delta} & c_1 s_2 s_3 - c_2 c_3 e^{i\delta} \end{pmatrix}$$

- 1 Before Kobayashi-Maskawa
 - The status of Particle Physics in 1972
 - The broken symmetry: charge-parity (CP)
- 2 The insight of Kobayashi and Maskawa
 - Prog. Theo. Phys. 49, 652 (1973)
 - In modern language: the CKM paradigm
- 3 Testing the Kobayashi-Maskawa predictions
 - The third generation
 - CP violation through the CKM mechanism
- 4 Concluding remarks

Flavor basis vs. mass basis

$$U' \equiv \begin{pmatrix} u \\ c \\ t \end{pmatrix}, \quad D' \equiv \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

- Charged current in the basis of flavor eigenstates:

$$\mathcal{L}_{CC} = \frac{g}{\sqrt{2}} \overline{U}'_L \gamma^\mu D'_L W_\mu^+ + h.c.$$

- Charged current in the basis of mass eigenstates:

$$\mathcal{L}_{CC} = \frac{g}{\sqrt{2}} \overline{U}_L \gamma^\mu (V_{UL}^\dagger V_{DL}) D_L W_\mu^+ + H.c.$$

V_{UL}, V_{DL} : unitary matrices that change the basis

- Coupling between U_L and D_L : $(g/\sqrt{2}) V_{CKM}$

$$V_{CKM} \equiv V_{UL}^\dagger V_{DL}$$

V_{CKM} is the Cabibbo-Kobayashi-Maskawa (CKM) matrix

Flavor basis vs. mass basis

$$U' \equiv \begin{pmatrix} u \\ c \\ t \end{pmatrix}, \quad D' \equiv \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

- Charged current in the basis of flavor eigenstates:

$$\mathcal{L}_{CC} = \frac{g}{\sqrt{2}} \overline{U}'_L \gamma^\mu D'_L W_\mu^+ + h.c.$$

- Charged current in the basis of mass eigenstates:

$$\mathcal{L}_{CC} = \frac{g}{\sqrt{2}} \overline{U}_L \gamma^\mu (V_{UL}^\dagger V_{DL}) D_L W_\mu^+ + H.c.$$

V_{UL}, V_{DL} : unitary matrices that change the basis

- Coupling between U_L and D_L : $(g/\sqrt{2}) V_{CKM}$

$$V_{CKM} \equiv V_{UL}^\dagger V_{DL}$$

V_{CKM} is the Cabibbo-Kobayashi-Maskawa (CKM) matrix

Flavor basis vs. mass basis

$$U' \equiv \begin{pmatrix} u \\ c \\ t \end{pmatrix}, \quad D' \equiv \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

- Charged current in the basis of flavor eigenstates:

$$\mathcal{L}_{CC} = \frac{g}{\sqrt{2}} \overline{U}'_L \gamma^\mu D'_L W_\mu^+ + h.c.$$

- Charged current in the basis of mass eigenstates:

$$\mathcal{L}_{CC} = \frac{g}{\sqrt{2}} \overline{U}_L \gamma^\mu (V_{UL}^\dagger V_{DL}) D_L W_\mu^+ + H.c.$$

V_{UL}, V_{DL} : unitary matrices that change the basis

- Coupling between U_L and D_L : $(g/\sqrt{2})V_{CKM}$

$$V_{CKM} \equiv V_{UL}^\dagger V_{DL}$$

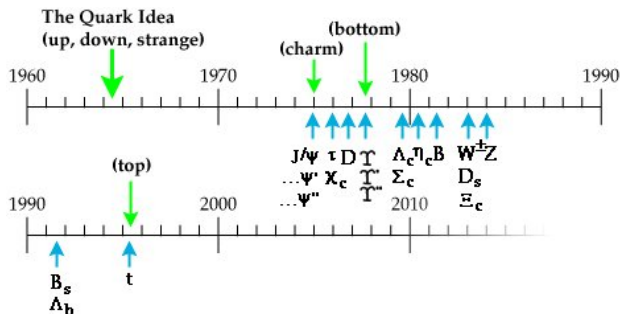
V_{CKM} is the Cabibbo-Kobayashi-Maskawa (CKM) matrix

Predictive power of the CKM mechanism

- A third family of quarks exists (never suspected before)
- All CP violation can be described in terms of a single complex number: Jarlskog invariant $J \equiv s_1 s_2 s_3 c_1^2 c_2 c_3 s_\delta$

- 1 Before Kobayashi-Maskawa
 - The status of Particle Physics in 1972
 - The broken symmetry: charge-parity (CP)
- 2 The insight of Kobayashi and Maskawa
 - Prog. Theo. Phys. 49, 652 (1973)
 - In modern language: the CKM paradigm
- 3 Testing the Kobayashi-Maskawa predictions**
 - The third generation**
 - CP violation through the CKM mechanism
- 4 Concluding remarks

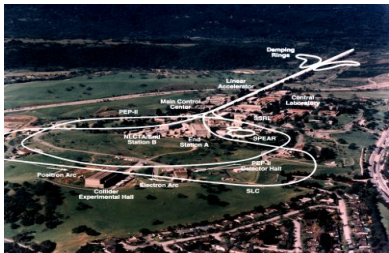
Experimental discovery of the third generation



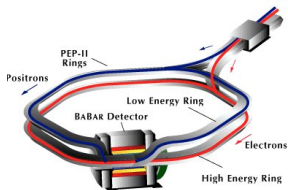
- Discovery of τ : 1976
- Υ , B , B_s , λ_b contain b quark
- Top quark: 1995
- The last element, ν_T , discovered in 2000.

- 1 Before Kobayashi-Maskawa
 - The status of Particle Physics in 1972
 - The broken symmetry: charge-parity (CP)
- 2 The insight of Kobayashi and Maskawa
 - Prog. Theo. Phys. 49, 652 (1973)
 - In modern language: the CKM paradigm
- 3 Testing the Kobayashi-Maskawa predictions**
 - The third generation
 - **CP violation through the CKM mechanism**
- 4 Concluding remarks

B factories: $B \equiv \bar{b}u, \bar{b}d, \bar{b}s, \bar{b}c, \quad \bar{B} \equiv \bar{u}b, \bar{d}b, \bar{s}b, \bar{c}b$



Babar (SLAC, USA)



BELLE (KEK, Japan)



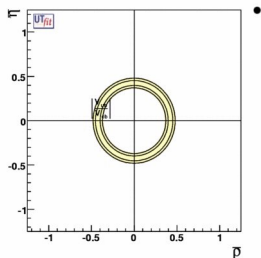
$$e^+ e^- \rightarrow B\bar{B} \rightarrow \text{decay products}$$

Wolfenstein parametrization of the CKM matrix

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \\ = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

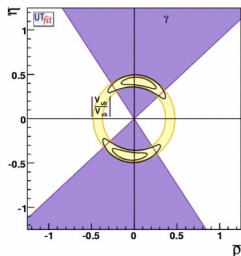
- λ : Cabibbo angle
- η : the imaginary component of V_{CKM}
- η/ρ large \Rightarrow CP violation is large, not approximate

More and more stringent tests of the CKM mechanism



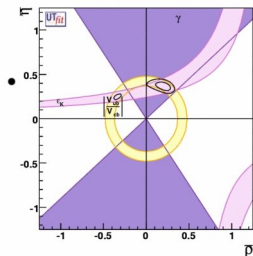
- Semileptonic decay $B \rightarrow D\ell\nu$

More and more stringent tests of the CKM mechanism



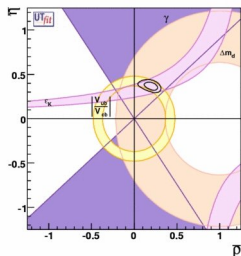
- Semileptonic decay $B \rightarrow D\ell\nu$
- “Charmed” decays $B \rightarrow DK$

More and more stringent tests of the CKM mechanism



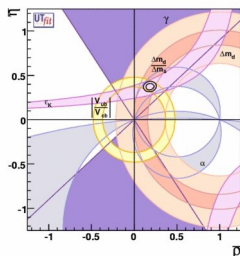
- Semileptonic decay $B \rightarrow D\ell\nu$
- “Charmed” decays $B \rightarrow DK$
- CP violation in K mesons

More and more stringent tests of the CKM mechanism



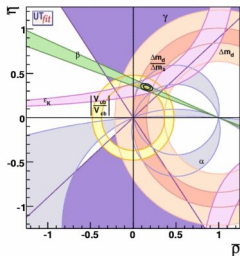
- Semileptonic decay $B \rightarrow D\ell\nu$
- “Charmed” decays $B \rightarrow DK$
- CP violation in K mesons
- ΔM in $B_d-\bar{B}_d$ system

More and more stringent tests of the CKM mechanism



- Semileptonic decay $B \rightarrow D\ell\nu$
- “Charmed” decays $B \rightarrow DK$
- CP violation in K mesons
- ΔM in $B_d-\bar{B}_d$ system
- ΔM in $B_s-\bar{B}_s$ system
- Decays to π and K

More and more stringent tests of the CKM mechanism



- Semileptonic decay $B \rightarrow D l \nu$
- “Charmed” decays $B \rightarrow DK$
- CP violation in K mesons
- ΔM in $B_d - \bar{B}_d$ system
- ΔM in $B_s - \bar{B}_s$ system
- Decays to π and K
- CP asymmetry in $B \rightarrow J/\psi K_S$

Concluding remarks

- No deviation from the CKM predictions has been observed
- Constrains many new physics models
- Future expts: LHC (ATLAS, CMS, LHC-b), super-B factory
- CP violation required for baryon asymmetry, but the CKM is not enough, so there will be life beyond CKM
- The successful prediction of Kobayashi and Maskawa still inspiring theoretical as well as experimental research