# Low Missing Energy and New Physics: To Miss or Not To Miss

Biswarup Mukhopadhyaya Regional Centre for Accelerator-based Particle Physics Harish-Chandra Research Institute Allahabad, India

May 1, 2012

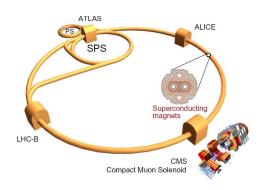
 New physics search at the LHC is often associated with large missing energy (MET): A big motivation: dark matter

- New physics search at the LHC is often associated with large missing energy (MET): A big motivation: dark matter
- But it is important to look for signals with low MET as well, and perhaps without compromise on the dark matter issue

- New physics search at the LHC is often associated with large missing energy (MET): A big motivation: dark matter
- But it is important to look for signals with low MET as well, and perhaps without compromise on the dark matter issue
- Distinguishing among various scenarios: a challenge

- New physics search at the LHC is often associated with large missing energy (MET): A big motivation: dark matter
- But it is important to look for signals with low MET as well, and perhaps without compromise on the dark matter issue
- Distinguishing among various scenarios: a challenge
- Useful signals to study at the LHC: same-sign tri-and four-leptons

## The LHC...



$$p \Rightarrow \Leftarrow p$$
 $7/8/14 \text{ TeV}$ 

#### Goals of the LHC....

- To discover the Higgs boson and complete the Standard Model of electroweak interactions
- To know more about top and bottom quark properties
- To understand strong interaction better
- To look for quark-gluon plasma
- Physics beyond the standard electroweak theory

#### The current state of affairs for electroweak physics....

• About 5  $fb^{-1}$  of data with 7 TeV already taken, and partly analysed

## nt state of affairs for electroweak physics....

- About 5  $fb^{-1}$  of data with 7 TeV already taken, and partly analysed
- Higgs mass range narrowed down, with some tantalising but inconclusive hints....

## nt state of affairs for electroweak physics....

- About 5  $fb^{-1}$  of data with 7 TeV already taken, and partly analysed
- Higgs mass range narrowed down, with some tantalising but inconclusive hints....
- The theory of weak + electromagnetic interactions is almost certainly the proposed  $SU(2) \times U(1)$  gauge theory, but the origin of masses not yet certain

## nt state of affairs for electroweak physics....

- About 5  $fb^{-1}$  of data with 7 TeV already taken, and partly analysed
- Higgs mass range narrowed down, with some tantalising but inconclusive hints....
- The theory of weak + electromagnetic interactions is almost certainly the proposed  $SU(2) \times U(1)$  gauge theory, but the origin of masses not yet certain
- The search for new physics is on, mostly accepting the 'Higgs lore', but no positive evidence yet

BUT....

• Why do we think there should be new physics ?

#### BUT....

- Why do we think there should be new physics ?
- Why should new laws be manifest at the LHC energy?



## Phenomenological dissatisfactions (unexplained features):

- Large number of unrelated free parameters
- Replication of fermion families
- The pattern of fermion masses
- Maximal P but small CP violation

#### Theoretical/Philosophical questions:

- No way to unify with strong interaction
- No clue on a quantum theory of gravity
- Divergent higher-order contributions to the Higgs boson mass

#### Sporadic/seasonal/volatile issues:

- The muon anomalous magnetic moment (3 3.5  $\sigma$  inconsistency)
- $\bullet$  PAMELA (excess positrons  $\sim$  10 80 GeV from galactic halo)
- $\bullet$  ATIC (excess galactic cosmic-ray electrons  $\sim$  300 -800 GeV)
- Tevatron multimuon events (excess multimuons, inexplicable from b-decays)
- Top quark forward-backward asymmetry at Tevatron
- Wjj events in the CDF experiment



#### **Concrete and persistent problems:**

- Neutrino masses and mixing
- Cold dark matter (no particle physics explanation)
- Matter-antimatter asymmetry in the universe
- A positive cosmological constant (!)

#### Physics beyond the standard model...

Effective energy scale to be probed at the LHC:  $\simeq 1$  -2 TeV

Out of the many motivations listed, which ones definitely suggest 'something new' at this energy?

## Why new physics at the TeV scale?

- The issue of Grand Unification (very indirect!)
- To understand why the Higgs should be within a TeV (relatively pressing!)
- Finding a cold(warm?) dark matter candidate (Somewhat imperative)

Thus the Dark matter issue is rather central to new physics search at the LHC

#### Searches for new physics at the LHC...

Events with large missing- $E_T$  (MET) or resonances are the first to be looked for

Popular belief: dark matter candidates ⇒ MET

Theories often proposed with a  $Z_2$  symmetry to accommodate a stable particle (dark matter candidate)

A candidate theory: supersymmetry (SUSY) with R-parity ( $R=(-)^{(3B+L+2S)}$ ) A lot of progress has taken place in SUSY search—mSUGRA-based cMSSM ruled out upto  $\simeq$  850 GeV (with  $4.7fb^{-1}$ )

There can be other  $Z_2$ -endowed scenarios: (Universal extra dimensions with KK parity, Little Higgs with T-parity)

## Distinguishing among models with $Z_2$ ....

'Inverse problem' within SUSY— mapping from signature space to parameter space
Arkani-Hamed et al. (2006)
Larger number of observables studied

⇒ degeneracy in the parameter space better lifted

SUSY vs. other scenarios with large MET General distinction strategies (LHC, dark matter search....):

A.K. Datta, G. Kane, M. Toharia (2005)

D. Hooper, G. Zaharijas (2006)

A. K Datta, P. Dey, S. Gupta, BM, A. Nyffeler (2007)

J. Hubisz et al. (2008)

M. Burns et al., 2008

W. Ehrenfeld et al. (2009)

B. Bhattacharjee it et al. (2009)



#### But is is also important to remember that....

So long as large MET signals elude us,

One needs to think of scenarios where such  $Z_2$  symmetry is broken

Some of the resulting theories may still accommodate dark matter candidates

Also, some scenarios with  $Z_2$  still yield low MET

We need criteria to point towards them and to discriminate among various low-MET scenarios in general

## Examples....

- (a) SUSY with R-parity violation (SUSY-RPV)
- (b) Little Higgs theories with broken T-parity (LHT-TPV)
- (c) Universal extra dimensions with conserved Kaluza-Klein parity (UED-KKC)
- (d) Universal extra dimensions with Kaluza-Klein parity violated (UED-KKV)
- (e) SUSY with a compressed spectrum

The present discussion includes cases (a) - (d)

#### SUSY-RPV....

The MSSM superpotential: (source of interaction of all chiral  $(f, \tilde{f})$  supermultiplets)  $W_{MSSM} = Y_{ij}^l L_i H_1 E_j^c + Y_{ij}^d Q_i H_1 D_j^c + Y_{ij}^u Q_i H_1 U_j^c$  (Assuming that baryon and lepton number are conserved)

When  $R = (-)^{L+3B+2S}$  violated via L or B, one can write

$$W=W_{MSSM}+W_{RPV}$$
, with  $W_{RPV}=\lambda_{ijk}L_iL_jE_k^c+\lambda_{ijk}^{'}L_iQ_jD_k^c+\epsilon_iL_iH_2+\lambda_{ijk}^{''}U_i^cD_j^cD_k^c$ 

We consider here L-violating  $W_{RPV}$ :  $\lambda_{ijk}^{"}=0$ 

Offers mechanisms for neutrino mass generation



#### SUSY-RPV....

Most phenomenological studies: one type of R-parity violating coupling at a time

Result: the MSSM-LSP (say, the lightest neutralino) has two/three body decays with at least one lepton in the final state

The gravitino or the axino may be the dark matter candidate

#### LHT-TPV....

The Higgs is the pesudo-goldstone boson of a broken approximate global symmetry

The breaking scale stabilises the electroweak scale

In a minimal (littlest) form,  $SU(5) \rightarrow SO(5)$ Underlying electroweak gauge group :  $[SU(2) \times U(1)]^2$ , with an exchange symmetry (T-parity) – a  $Z_2$  symmetry (stabilises, for example, the  $m_W/m_Z$  ratio)

Result: a division into T-even (SM) and T-odd (new) particles  $[SU(2) \times U(1)]^2 \longrightarrow SU(2) \times U(1)$  at scale f New particles include heavy T-odd fermions  $(Q_H, L_H)$ , heavy gauge bosons  $(W_H, Z_H, A_h)$ , a Higgs triplet...... The lightest T-odd particle (LTP) is stable: (Usually the  $A_H$ )

#### LHT-TPV....

 $Z_2$  symmetry  $\Rightarrow$  LTP is the dark matter candidate (A neutral, weakly interacting particle) The spectrum and the interactions are controlled by  $f = \mathcal{O}(\text{TeV}), f \kappa_{ii} = \text{matrix deciding heavy fermion masses}$ But T-parity can be broken... by Wess-Zumino-Witten anomaly terms Results: terms  $\sim \frac{N}{48\pi^2 f^2} \epsilon_{\mu\nu\alpha\beta} V^{\mu}_{H} V^{\nu} \partial^{\alpha} V^{\beta}$ The LTP becomes unstable: For example,  $A_H \longrightarrow WW^{(*)}$ ,  $ZZ^{(*)}$ leading to tree-level or loop-induced decays such as  $\ell\nu\ell\nu$ ,  $\ell\ell$ ,...

#### UED-KKC....

At least one spacelike compact extra dimension, of radius R, where all fields can propagate

New particles are Kaluza-Klein towers, with same spin as in the zero-mode SM states

The extra dimension 'orbifolded' about the axis from  $\phi = 0 - \pi$ : a  $Z_2$  symmetry (for ensuring proper fermion chiralities)  $\Rightarrow$  A conserved 'Kaluza-Klein parity'

The lightest KK-odd particle (LKP) is stable due to the  $Z_2$  symmetry: dark matter candidate (A neutral, weakly interacting particle): Usually the first excitation  $(A_1)$  of the photon

The spectrum is decided by  $R^{-1}$ , and the cut-off scale A highly compressed spectrum in general



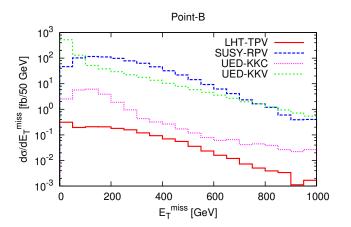
#### UED-KKV...

KK-parity can be broken by additional asymmetric operators at the orbifold fixed points ⇒ The LKP is again unstable

Claim: all of these four scenarios, including UED-KKC, can lead to similar MET signals at the LHC How to distinguish?

K. Ghosh, S. Mukhopadhyay, BM (2010)

#### MET distributions for all four scenarios...



 $M_s = Strongly interacting particle mass \simeq 1 TeV SUSY-RPV: <math>\lambda$ -type with one coupling



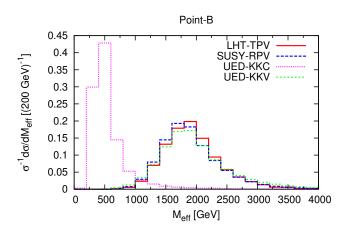
#### Distinguishing among the various scenarios...

Prescription: use multileptons  $\ell=e,\mu$ Dileptons: limited discriminating power Trileptons: sometimes effective, but.... four-and five-lepton final states can be useful discriminators

Isolated and central leptons, with appropriate transverse momentum requirements

Sufficient statistics required for discrimination: 14 TeV run necessary For  $M_s=600$  GeV, 5  $fb^{-1}$  at 14 TeV is enough For  $M_s=1$  TeV, 30  $fb^{-1}$  is required for  $5\sigma$  significance for all scenarios

#### Distinguishing among the various scenarios...



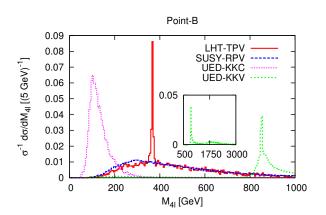
Effective mass distribution with  $M_s=1$  TeV: UED-KKC stands out  $(M_{eff}=\Sigma_i p_{\scriptscriptstyle T}^i + MET)$ 



## Four-lepton invariant mass...

$$M_{4l} = (p_{l_1} + p_{l_2} + p_{l_3} + p_{l_4})^2$$
  
If the  $4\ell$  all come from one particle,  
then  $M_{4l} = mass$  of the parent particle

#### Four-lepton invariant mass...



#### $4\ell$ invariant mass distributions: explanation

UED-KKC:

$$M_{4\ell} \leq \sqrt{2 \frac{(M_{Z_1}^2 - M_{L_1}^2)(M_{L_1}^2 - M_{\gamma_1}^2)}{M_{L_1}^2} + 4(E_1 E_3 + E_1 E_4 + E_2 E_3 + E_2 E_4)}$$
compressed spectrum  $\Rightarrow M_{4\ell}$  within a narrow band

#### $4\ell$ invariant mass distributions: explanation

• UED-KKC:

$$\begin{array}{l} \textit{M}_{4\ell} \leq \\ \sqrt{2\frac{(\textit{M}_{Z_1}^2 - \textit{M}_{L_1}^2)(\textit{M}_{L_1}^2 - \textit{M}_{\gamma_1}^2)}{\textit{M}_{L_1}^2}} + 4(\textit{E}_1\textit{E}_3 + \textit{E}_1\textit{E}_4 + \textit{E}_2\textit{E}_3 + \textit{E}_2\textit{E}_4)} \\ \textit{compressed spectrum} \Rightarrow \textit{M}_{4\ell} \textit{ within a narrow band} \end{array}$$

• LHT-TPV: most events from  $A_H A_H \longrightarrow 4W^{(*)}$ On the top is a narrow peak from  $A_H \longrightarrow 2Z^{(*)}$ 

#### $4\ell$ invariant mass distributions: explanation

• UED-KKC:

$$M_{4\ell} \leq \sqrt{2 \frac{(M_{Z_1}^2 - M_{L_1}^2)(M_{L_1}^2 - M_{\gamma_1}^2)}{M_{L_1}^2} + 4(E_1 E_3 + E_1 E_4 + E_2 E_3 + E_2 E_4)}$$
compressed spectrum  $\Rightarrow M_{4\ell}$  within a narrow band

- LHT-TPV: most events from  $A_H A_H \longrightarrow 4W^{(*)}$ On the top is a narrow peak from  $A_H \longrightarrow 2Z^{(*)}$
- UED-KKV: Peak from one side:  $Z_1 \longrightarrow \ell L_1 \longrightarrow \ell \ell \gamma_1 \longrightarrow 4\ell$

#### $4\ell$ invariant mass distributions: explanation

UED-KKC:

$$\begin{array}{l} \textit{M}_{4\ell} \leq \\ \sqrt{2\frac{(\textit{M}_{Z_1}^2 - \textit{M}_{L_1}^2)(\textit{M}_{L_1}^2 - \textit{M}_{\gamma_1}^2)}{\textit{M}_{L_1}^2}} + 4(\textit{E}_1\textit{E}_3 + \textit{E}_1\textit{E}_4 + \textit{E}_2\textit{E}_3 + \textit{E}_2\textit{E}_4)} \\ \textit{compressed spectrum} \Rightarrow \textit{M}_{4\ell} \textit{ within a narrow band} \end{array}$$

- LHT-TPV: most events from  $A_H A_H \longrightarrow 4W^{(*)}$ On the top is a narrow peak from  $A_H \longrightarrow 2Z^{(*)}$
- UED-KKV: Peak from one side:  $Z_1 \longrightarrow \ell L_1 \longrightarrow \ell \ell \gamma_1 \longrightarrow 4\ell$
- SUSY-RPV: no peak at all

#### Some more discrimination criteria....

• Pairwise opposite-sign lepton invariant masses and their correlations in 4ℓ events

#### Some more discrimination criteria....

- Pairwise opposite-sign lepton invariant masses and their correlations in 4ℓ events
- Angular correlation of each lepton with the nearest jet (for SUSY-RPV, there is often peaking in the forward direction)

#### Some more discrimination criteria....

- Pairwise opposite-sign lepton invariant masses and their correlations in 4ℓ events
- Angular correlation of each lepton with the nearest jet (for SUSY-RPV, there is often peaking in the forward direction)
- $N(5\ell)/N(4\ell)$ : SUSY-RPV and LHT-TPV show higher ratios than the two other cases
  Reason: compressed spectrum tends to soften leptons from cascade

# Same-sign trileptons(SS3 $\ell$ ): unexplored potential...

Lepton sign: seriously used in the search for same-sign dilepton (SSD) events

Majorana fermions enhance SSD rates  $p_{\pm}$  + isolation

Majorana fermions enhance SSD rates,  $p_T$  + isolation cuts reduce backgrounds (mostly from  $t\bar{t}$ )

Leptons of higher multiplicity and same sign: SM backgrounds extremely small

Theories with L-violation + self-conjugate fields: unsuppressed signals

A very discriminating check on scenarios with low-MET

SUSY-RPV stands out by contributing to SS3 $\ell$  (and also SS4 $\ell$ ) (Even in the early run)

Also, the dynamics of R-parity violation can be probed thereby S. Mukhopadhyay, BM (2010, 2011)

# Same-sign trileptons(SS3 $\ell$ ): unexplored potential...

Standard model contribution to  $\sigma(\text{SS}3\ell)$ : with appropriate cuts,  $\simeq 2.5 \times 10^{-3}$  fb (  $\simeq 7.0 \times 10^{-4}$  fb ) at 14 (7) TeV Even smaller backgrounds for SS4 $\ell$ 

If high-MET new physics signals continue to elude us, Low-MET ones must be looked for SS3 $\ell \Rightarrow$  a discriminating signature of specific scenario(s) In SUSY-RPV, LSP-pair decays (with no branching ratio suppression) can yield two same-sign leptons, and one more comes from the cascade

• SUSY with R-parity: rates are always very low due to lack of L-violation and branching fraction suppression

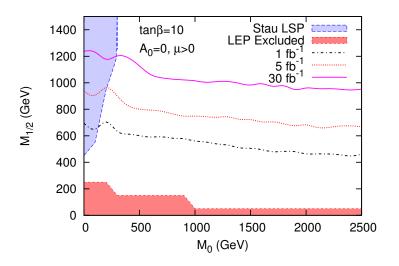
- SUSY with R-parity: rates are always very low due to lack of L-violation and branching fraction suppression
- SUSY with compressed spectrum: rates even lower

- SUSY with R-parity: rates are always very low due to lack of L-violation and branching fraction suppression
- SUSY with compressed spectrum: rates even lower
- LHT-TPV: Decay of the LTP to one charged lepton has low branching ratio

- SUSY with R-parity: rates are always very low due to lack of L-violation and branching fraction suppression
- SUSY with compressed spectrum: rates even lower
- LHT-TPV: Decay of the LTP to one charged lepton has low branching ratio
- UED: SS3ℓ occurs even more rarely

- SUSY with R-parity: rates are always very low due to lack of L-violation and branching fraction suppression
- SUSY with compressed spectrum: rates even lower
- LHT-TPV: Decay of the LTP to one charged lepton has low branching ratio
- UED: SS3ℓ occurs even more rarely
- Thus SUSY-RPV stands out

# Coverage in mSUGRA parameter space at 14 TeV



5-event contours with background << 1: One  $\lambda$ -type coupling included

#### Rates at 7 TeV

Case	$\sigma_{SS3\ell}$ (fb)
$\lambda$ -type: $m_{\tilde{g}} \simeq 660  GeV$ , neutralinoLSP	19.82
$\lambda$ -type: $m_{\tilde{g}} \simeq 000  \text{GeV}$ , neutralinoLSP	4.29
$\lambda$ -type: $m_{\tilde{g}} \simeq 770  GeV$ , stauLSP	30.74
$\lambda$ -type: $m_{ ilde{g}} \simeq 1  TeV$ , stauLSP	3.35
$\lambda'$ -type: $m_{\tilde{g}} \simeq 660  GeV$ , neutralinoLSP	2.07

#### In a purely phenomenological R-parity violating SUSY

#### Some general conclusions about SS3 $\ell$ :

- $M_{2(1)} \ge 2M_{1(2)} \Rightarrow Signal \ rate \ enhanced$
- $M_1 \simeq M_2 \Rightarrow Signal \ suppressed \ unless \ M_3 >> M_{1,2}$
- For  $m_{\tilde{g}}$ ,  $m_{\tilde{q}} \simeq 1$  TeV, 5 background-free events possible with 0.6 5.5 fb<sup>-1</sup> at 7 TeV, and 0.1 3.0 fb<sup>-1</sup> at 14 TeV
- The nature of SUSY-RPV can be extracted for one type of RPV coupling present at a time

#### Some numbers in CMSSM with $m_{ ilde{p}} \simeq m_{ ilde{p}} \simeq 1$ TeV: ss3l

$$\tan \beta = 5$$

Cut	SM	S	Sig(S)
Lepton selection $+ MET > 30 GeV$	$7.01 \times 10^{-4}$	2.41	5.9
MET > 50 GeV		2.23	5.6
MET > 100 GeV		1.65	4.7
$m_{ ext{eff}}^{\ell} > 100~ ext{GeV}$		2.39	5.8
$m_{ ext{eff}}^{\ell} > 200~ ext{GeV}$		1.57	4.6
$m_{ m eff} > 150~{\it GeV}$		2.40	5.9
$m_{eff} > 250 \; GeV$		2.10	5.4

Table: Same-sign trilepton rates (in fb) and the expected signal significance, for  $7\,\text{TeV}$  LHC and  $1~\text{fb}^{-1}$  of integrated luminosity.  $(\lambda_{123}=10^{-3})$ 

# After lepton acceptance + MET cut, 3 events with





# Some numbers in CMSSM with $m_{\tilde{p}} \simeq m_{\tilde{p}} \simeq 1$ TeV: SSD

$$\tan \beta = 5$$

Cut	SM	S	Sig(S)
Lepton selection $+ MET > 30 GeV$	10.7	11.61	3.1
MET > 50 GEV		10.95	
MET > 100  GeV		8.66	
$m_{eff}^{\ell} > 100~GeV$	6.4	10.59	3.5
$m_{ ext{eff}}^{ ilde{\ell}^{\prime}} > 200 \; GeV$	1.0	5.50	3.7
$m_{\it eff} > 150~\it GeV$	7.4	11.41	3.5
$m_{eff} > 250 \; GeV$	1.8	9.46	4.7

Table: Same-sign dilepton rates (in fb) and the expected signal significance, for  $7\,TeV$  LHC and  $1\,fb^{-1}$  of integrated luminosity. ( $\lambda_{123}=10^{-3}$ )

After lepton acceptance + MET cut,  $5\sigma$  discovery in the SSD channel with  $2.6fb^{-1}$ 



Define 
$$x = \sigma_{SS3\ell}/(\sigma_{SS3\ell} + \sigma_{MS3\ell})$$
, and  $y = \sigma_{SS4\ell}/(\sigma_{SS4\ell} + \sigma_{MS4\ell})$ 

The dynamics is reflected in x and y

x and y depend only on how charginos and neutralinos decay to leptons of either sign, and are independent of the spectrum Assumption:

- (a) neutralino LSP
- (b)only one L-violating coupling at a time



• For  $\lambda$ -type coupling,

 $x \simeq 0.12$ 

In actual simulations,  $x = 0.11 \pm 0.02$ 

• For  $\lambda$ -type coupling,

$$x \simeq 0.12$$

In actual simulations,  $x = 0.11 \pm 0.02$ 

• With  $\lambda'$ -type coupling,

Let 
$$B(\chi_1^0 \longrightarrow l^{\pm}q\bar{q}') = \alpha$$
 ( $\alpha \simeq 0.5$ ) Then

$$x = \alpha^2/4 + 4y(1/\alpha - 1)$$

For the bilinear terms  $\epsilon_i L_i H_2$  side by side with  $\mu H_1 H_2$ ,

The  $\epsilon_i$  can be rotated away from the superpotential,

RPV is then driven by sneutrino vev in the scalar potential

Then correct neutrino masses  $\Rightarrow \langle \tilde{\nu} \rangle \simeq 100 \text{keV}$  in that basis

$$\chi_1^0 \longrightarrow \ell W, \nu Z$$

(BR's fixed unless sneutrinos are closely degenerate with the Higgs)

#### Then

$$x = 3.53y + 0.06$$

(Including backgrounds, the relations are satisfied upto 10 - 20 %)

• Several well-motivated scenarios can have signals where MET is not the main thing

- Several well-motivated scenarios can have signals where MET is not the main thing
- Some of these may even contain dark matter candidates

- Several well-motivated scenarios can have signals where MET is not the main thing
- Some of these may even contain dark matter candidates
- Multilepton final states can help in distinguishing among different scenarios

- Several well-motivated scenarios can have signals where MET is not the main thing
- Some of these may even contain dark matter candidates
- Multilepton final states can help in distinguishing among different scenarios
- Same-sign trileptons: a very clear signal of SUSY with L-violation

- Several well-motivated scenarios can have signals where MET is not the main thing
- Some of these may even contain dark matter candidates
- Multilepton final states can help in distinguishing among different scenarios
- Same-sign trileptons: a very clear signal of SUSY with L-violation
- The early run has interesting prospects

- Several well-motivated scenarios can have signals where MET is not the main thing
- Some of these may even contain dark matter candidates
- Multilepton final states can help in distinguishing among different scenarios
- Same-sign trileptons: a very clear signal of SUSY with L-violation
- The early run has interesting prospects
- SS3\ell and SS4\ell can differentiate among various R-parity breaking terms

"It is the mark of an educated mind to be able to be able to entertain a thought without accepting it"

—Aristotle