



# Strong Interaction In The 21st Century



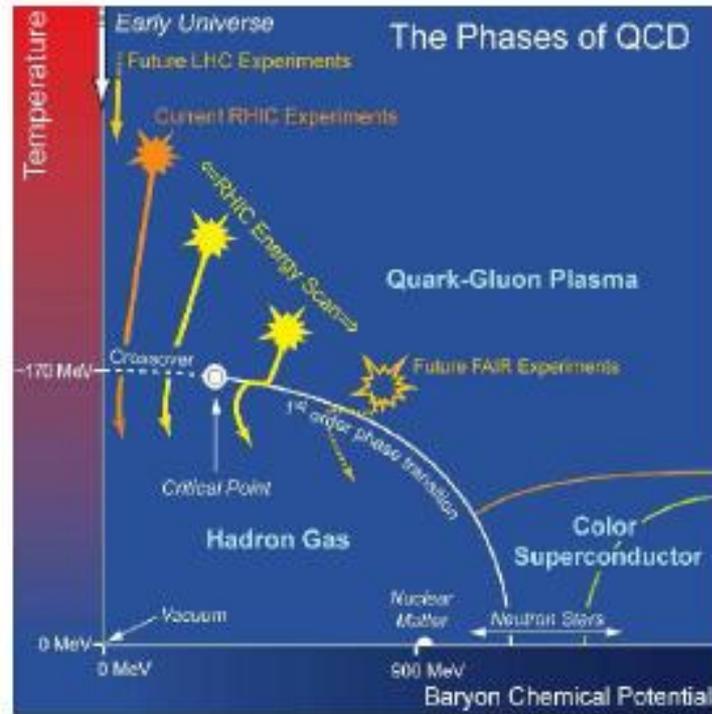
## Strong Interactions and Finite Baryon Density: The Lattice Approach



*Maria Paola Lombardo*  
*Laboratori Nazionali di Frascati*

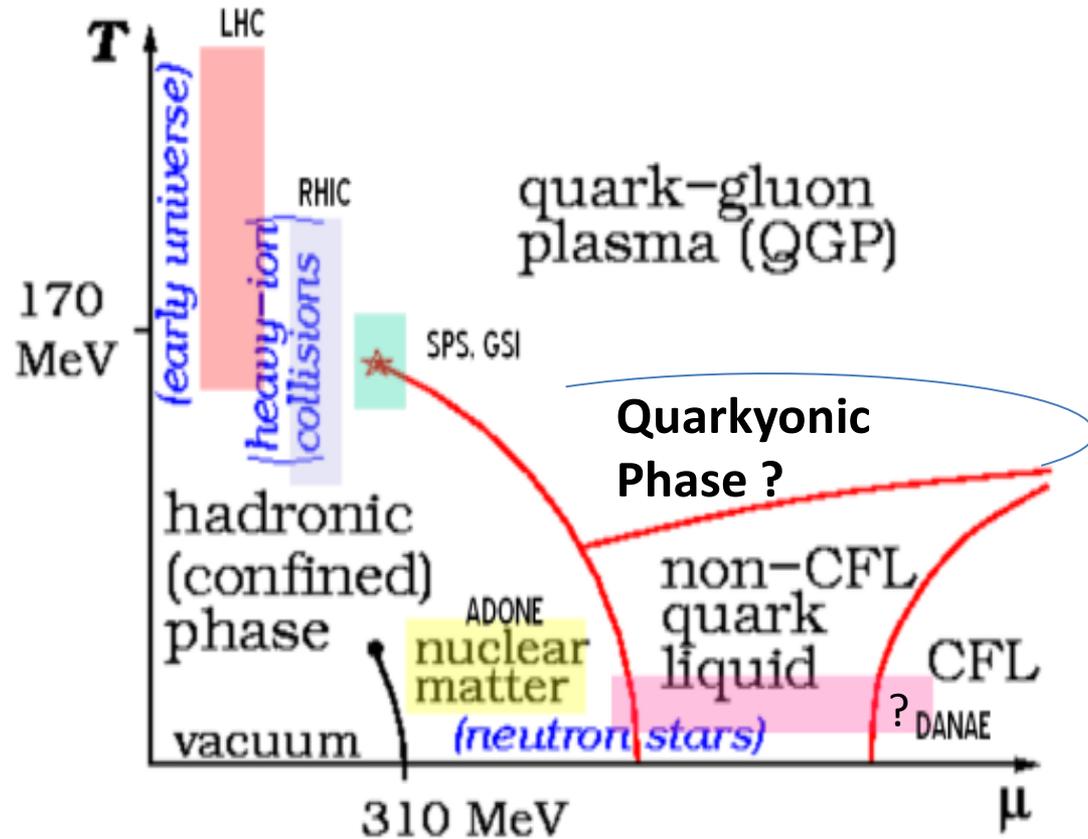


# The QCD Phase Diagram

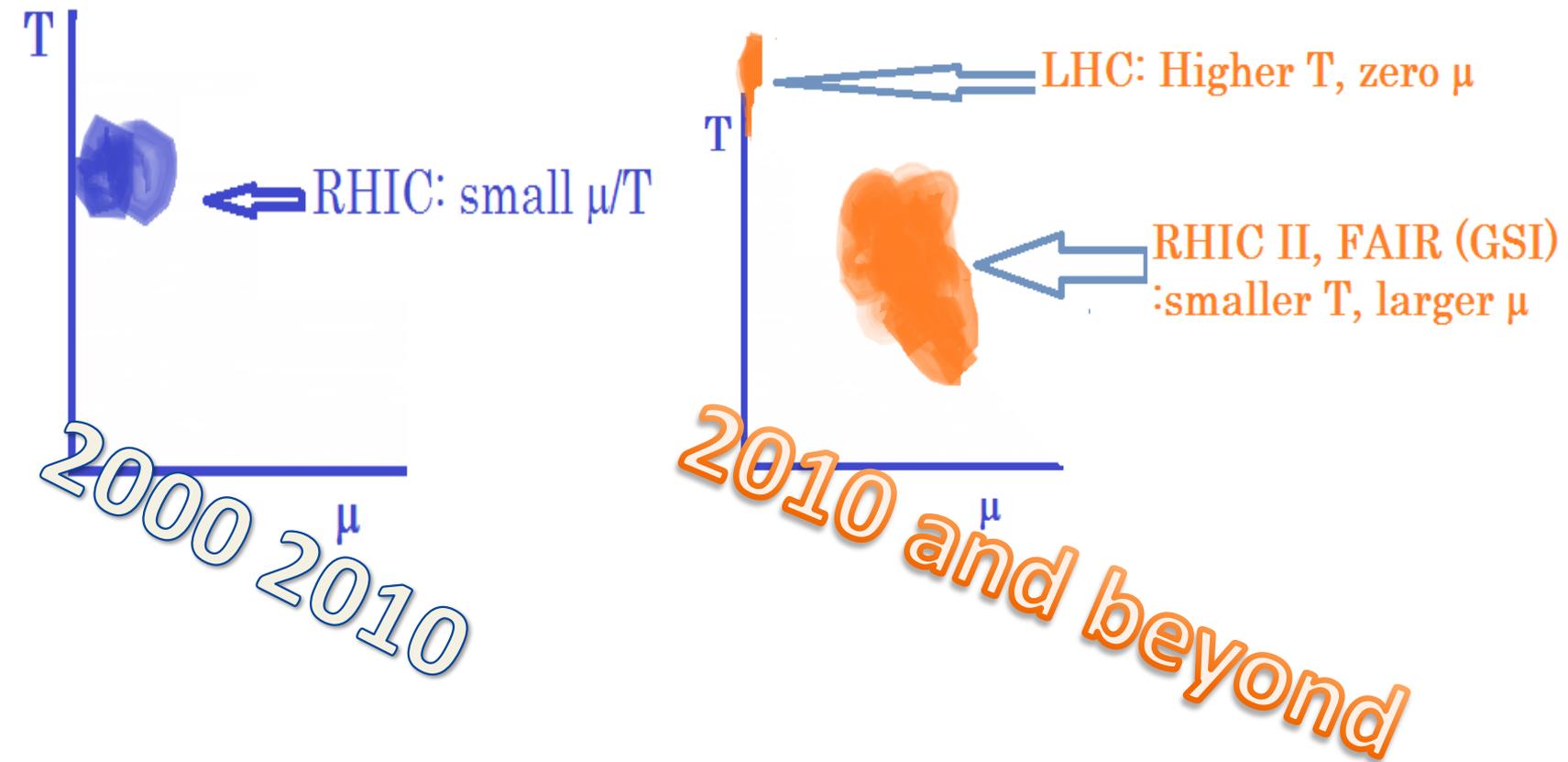


# Nonzero baryon density tells us about general properties of strong interactions

- Asymptotic freedom
  - Confinement, Chiral Symmetry
  - Quarkyonic Phase
- Mc Lerran



# Strong Interactions and Finite Baryon density in the XXI Century: shift of the focus at the turn of the decade



# OUTLINE

INTRODUCTION

RESULTS ON THE PHASE DIAGRAM:

CRITICAL ENDPOINT

PHYSICS OF THE FREEZOUT LINE

THE QUARKYONIC PHASE

# THE THEORETICAL APPARATUS: QCD, THE FIELD THEORY OF STRONG INTERACTIONS

$$\mathcal{L} = \mathcal{L}_{YM} + \bar{\psi}(i\gamma_{\mu}D_{\mu} + m + \mu\gamma_0)\psi$$

## LATTICE QCD ALLOWS FIRST PRINCIPLES CALCULATIONS FROM THE QCD LAGRANGIAN

$$\mathcal{L} = \mathcal{L}_{YM} + \bar{\psi}(i\gamma_{\mu}D_{\mu} + m + \mu\gamma_0)\psi$$

We can tune physical parameter, as in real experiments: baryon chemical potential, temperature, isospin chemical potential, strangeness,...

We can also play with number of color and number of flavor.

We can address phenomenological issues as well as theoretical questions.

## COMPUTATIONAL SCHEMES

$$\mathcal{Z} = \int d\phi d\bar{\psi} dU e^{-S(\phi, \bar{\psi}, U)}; S(\phi, \bar{\psi}, U) = \int_0^{1/T} dt \int d^d x \mathcal{L}(\phi, \bar{\psi}, U)$$

$$\mathcal{L}_{QCD} = \mathcal{L}_{YM} + \bar{\psi}(i\gamma_\mu D_\mu + m)\psi + \mu\bar{\psi}\gamma_0\psi$$

Two options:

1. Integrate out gluons first:

$$\mathcal{Z}(T, \mu, \bar{\psi}, \psi, U) \simeq \mathcal{Z}(T, \mu, \bar{\psi}, \psi) \rightarrow$$

effective **approximate** fermion models

2. Integrate out fermions **exactly** as S is bilinear in  $\psi, \bar{\psi}$

$$S = S_{YM}(U) + \bar{\psi}M(U)\psi$$

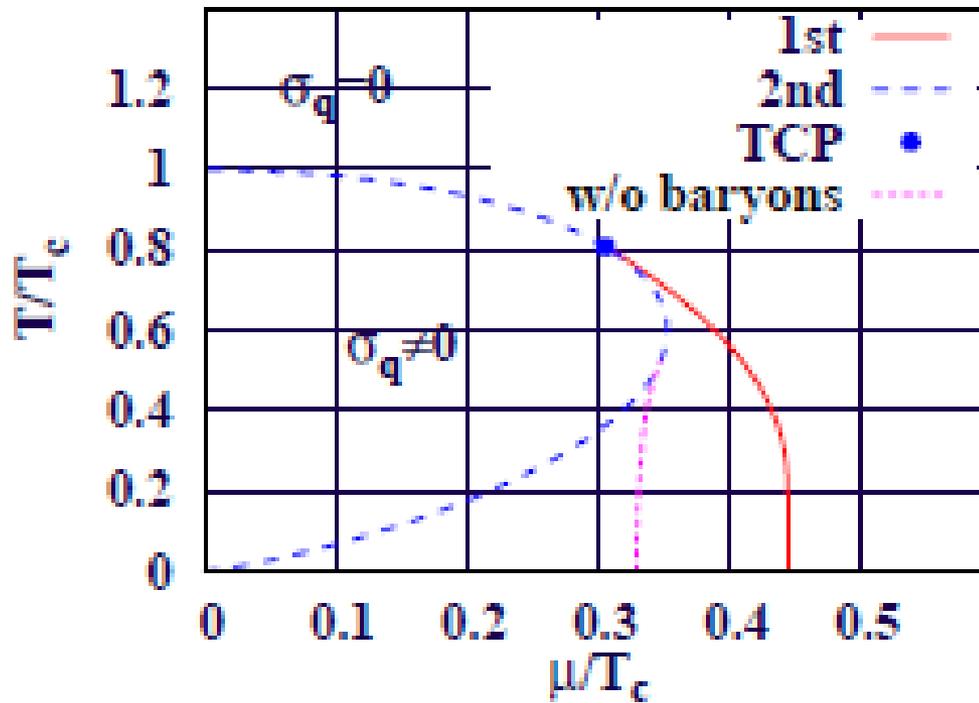
$$\mathcal{Z}(T, \mu, U) = \int dU e^{-(S_{YM}(U) - \log(\det M))} \rightarrow$$

starting point for numerical calculations

---

# Option1: The strong coupling expansion

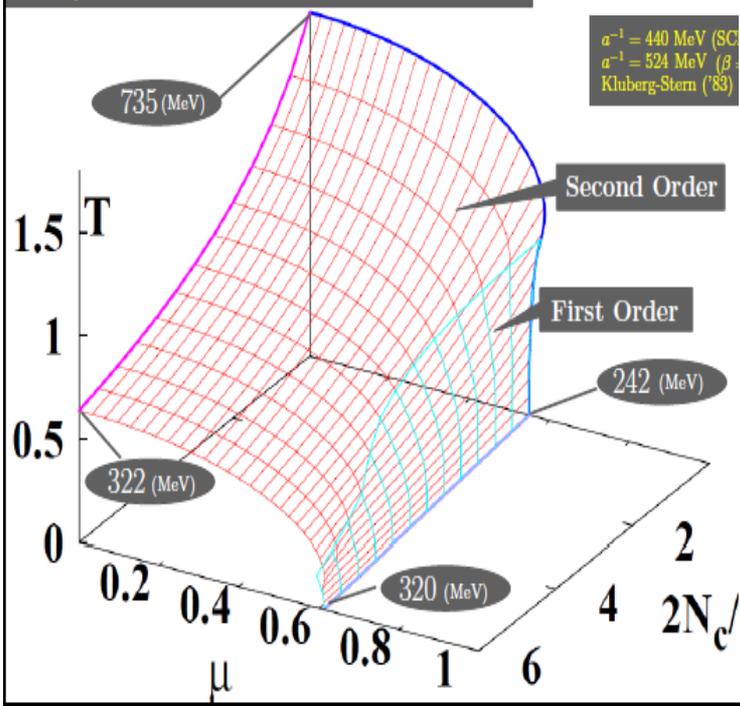
A long history..



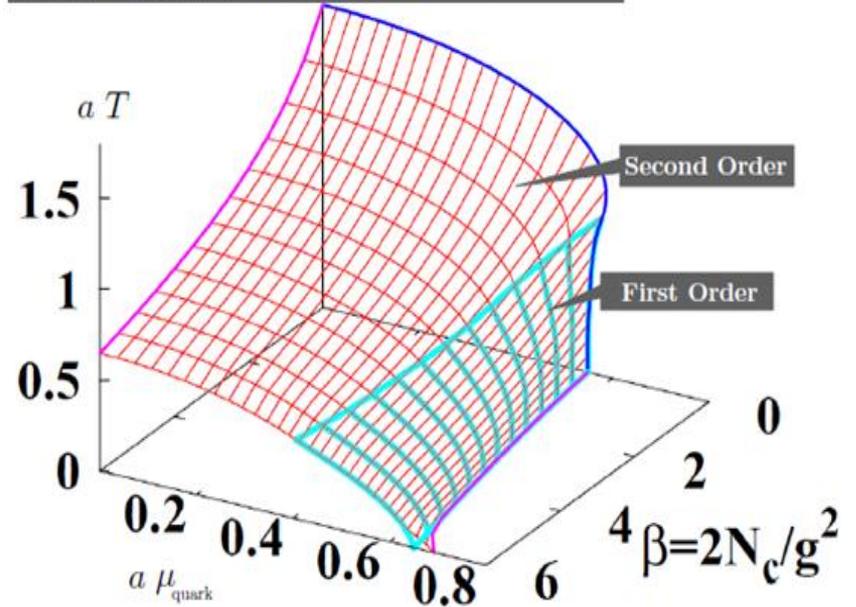
Kawamoto,  
Miura, Onishi  
2007

# The Strong Coupling Expansion approaching the continuum limit

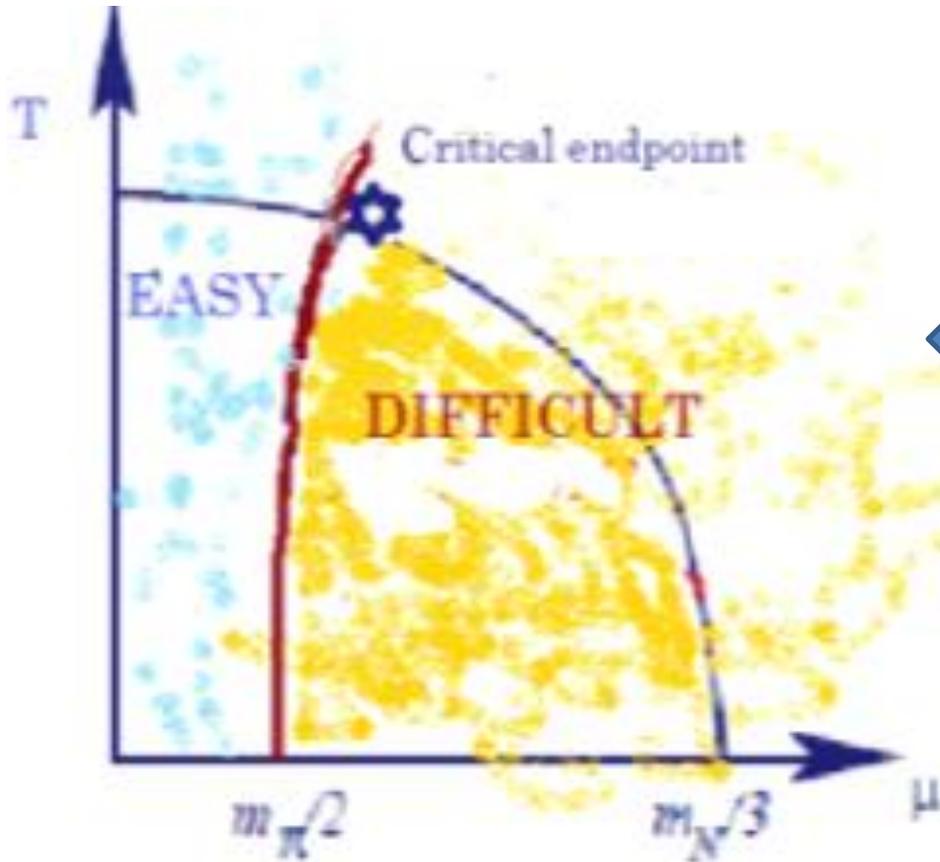
Phase Diagram Evolution with Increasing  $\beta = 2N_c/g^2$   
Next-to-Leading Order of  $1/g^2$  Expansion  
K. Miura, T.-Z. Nakano, A. Ohnishi and N. Kawamoto  
Phys. Rev. D 80 (2009), 074034.  
Courtesy Kohtaroh Miura



Phase Diagram Evolution with Increasing  $\beta$   
Next-to-Next-to-Leading Order of  $1/g^2$  Expansion  
arXiv:0911.3453, T.-Z. Nakano, K. Miura and A. Ohnishi  
Courtesy Kohtaroh Miura



Option 2 : Integrate over fermions and ..  
The  $m_\pi/2$  barrier



Summary  
Of our  
efforts!!

# THE CHALLENGE

## IMPORTANCE SAMPLING AND THE POSITIVITY ISSUE

$$\mathcal{Z}(T, \mu, U) = \int dU e^{-(S_{YM}(U) - \log(\det M))}$$

## $\det M > 0 \rightarrow$ IMPORTANCE SAMPLING MONTECARLO SIMULATIONS

To assess sign problem consider  $M^\dagger(\mu_B) = -M(-\mu_B)$

- $\mu = 0 \rightarrow \det M$  is **real**  
Particles-antiparticles symmetry : MC Simulations OK
- Imaginary  $\mu \neq 0 \rightarrow \det M$  is **real**  
(Real) Particles-antiparticles symmetry : MC Simulations OK
- Real  $\mu \neq 0$  Particles-antiparticles asymmetry  
 $\rightarrow \det M$  is **complex** in QCD

*QCD with a real baryon chemical potential:  
use information from the accessible region*

$$\text{Real} \mu = 0, \text{Im} \mu \neq 0$$

---

# TOWARDS THE REAL SOLUTION

COMPLEX LANGEVIN?

→Gert Aarts

DENSITY OF STATE METHODS?

→Christian Schmidt

NEW ALGORITHMS?

→Shailesh Chandrasekharan

LEARNING FROM SIMPLER SYSTEMS?

→David Kaplan

*Many discussions  
at this meeting  
..not covered in  
this talk*

# THE 'EASY' SOLUTION : EXTRAPOLATE

## QCD AT NONZERO BARYON DENSITY: METHODS

Multiparameter Reweighting ( $\mu = 0$ ):

*Fodor, Katz, Csikor, Egri, Szabo, Toth*

Derivatives ( $\mu = 0$ ):

*Gupta, Gvai and collaborators; MILC; QCD-Taro*

Expanded Reweighting ( $\mu = 0$ )

*Bielefeld-Swansea*

Analytic continuation from Imaginary  $\mu$

Strong Coupling QCD MpL

Dim. Reduced QCD *Laine, Hart, Philipsen*

QCD de Forcrand, Philipsen, Kratochvila

*D'Elia, MpL, Di Renzo*

*Azcoiti, Di Carlo, Galante, Laliena, Staggered*

*Luo et al. Wilson*

Models *Giudice, Papa; de Forcrand, Kim....*

The results on the phase diagram presented in the following were obtained by these 'easy' approaches

**CRITICAL ENDPOINT**

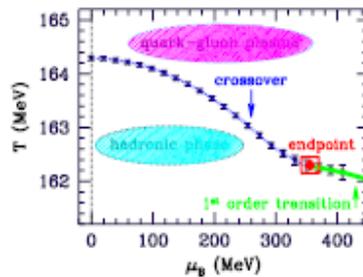
## THE CRITICAL ENDPOINT



OR

BOTH SCENARIO ARE COMPATIBLE  
WITH MODEL CALCULATIONS AND UNIVERSALITY

## STRATEGY 0 : FODOR KATZ , REWEIGHTING FROM $\mu = 0$



### CRITICISM:

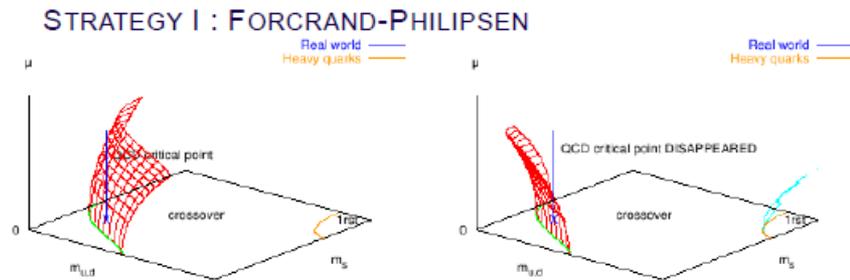
Critical point is close to the phase quenched threshold where reweighting fails at  $T=0$

### HOWEVER :

Important contribution from the phase does not necessarily hamper reweighting : overlap might still be large or correlation with the phase might be small.

Splitterff, Verbaarschot, MpL , in progress

# CHALLENGING THE ENDPOINT



Scenario I or Scenario II ? To decide, measure slope  $K$  in

$$\frac{m_c(\mu)}{m_c(0)} = 1 + K \left(\frac{\mu}{T}\right)^2 + \dots$$

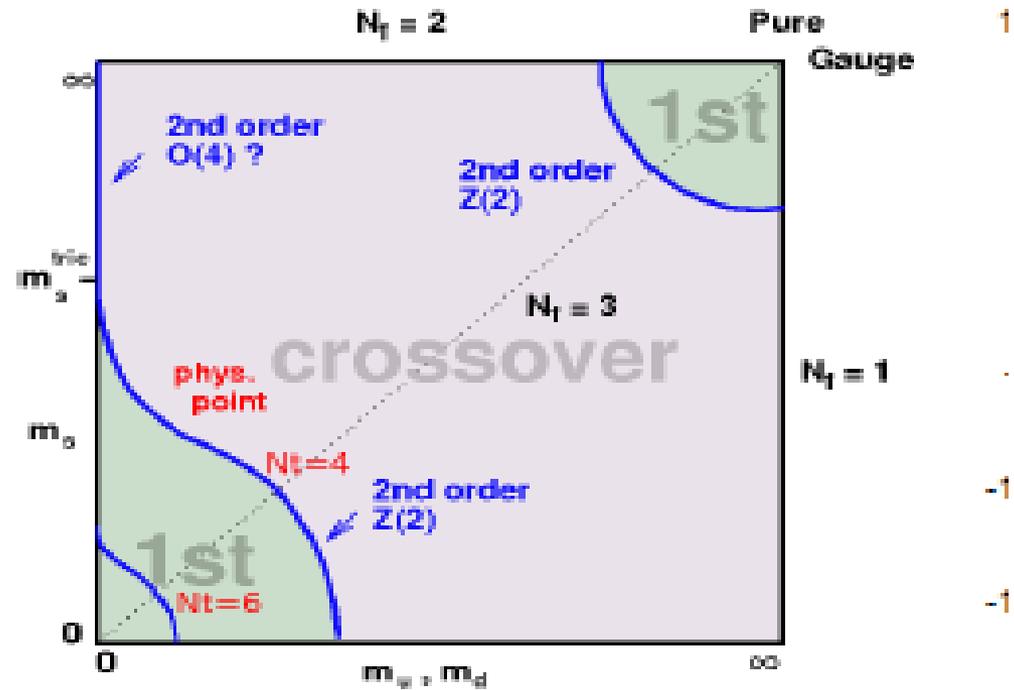
$K > 0$  : Scenario I , critical endpoint at small  $\mu_B$

$K < 0$  : Scenario II, NO critical endpoint at small  $\mu_B$

CURRENT RESULTS SUGGEST NO CRITICAL ENDPOINT FOR  $\mu_B < 600 MeV$

NB: assume that endpoint is part of the critical surface at  $m=0$

# Towards the continuum



# RESCUING THE ENDPOINT

## STRATEGY II : GAVAI AND GUPTA, BIELEFELD-RBC

Series expansion for the pressure:

$$P(T, \mu_B) = P(T) + \frac{1}{2}\chi_B^{(2)}(T)\mu_B^2 + \frac{1}{4!}\chi_B^{(4)}(T)\mu_B^4 + \frac{1}{6!}\chi_B^{(6)}(T)\mu_B^6 + \frac{1}{8!}\chi_B^{(8)}(T)\mu_B^8 + \dots,$$

The quark number susceptibility has the expansion

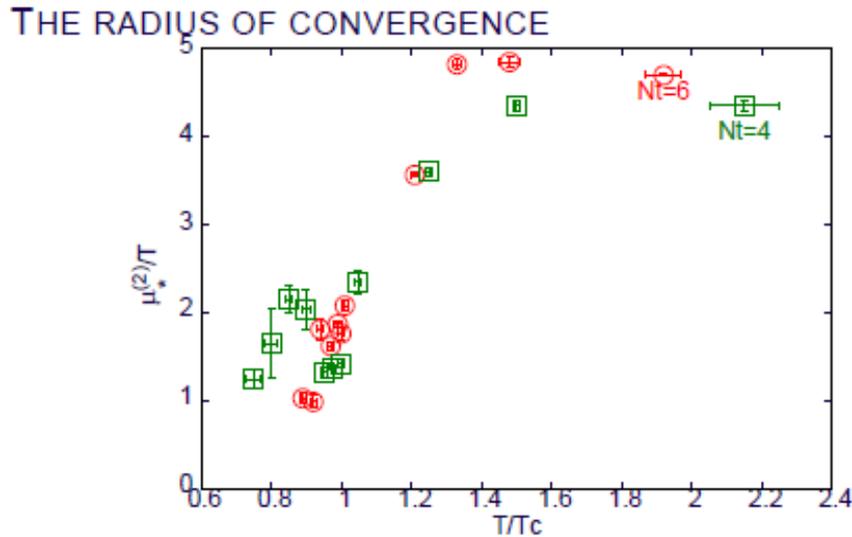
$$\chi_B(T, \mu_B) = \chi_B^{(2)}(T) + \frac{1}{2}\chi_B^{(4)}(T)\mu_B^2 + \frac{1}{4!}\chi_B^{(6)}(T)\mu_B^4 + \frac{1}{6!}\chi_B^{(8)}(T)\mu_B^6 + \dots.$$

THIS SERIES IS EXPECTED TO DIVERGE AT THE QCD CRITICAL END POINT. RADIUS OF CONVERGENCE IS

$$\lim_{n \rightarrow \infty} \mu_*^{(n)} = \sqrt{\frac{1}{n(n-1)} \frac{\chi_B^{(n+2)}}{\chi_B^{(n)}}}.$$

*The endpoint is the first singularity in the complex  $\mu$  plane occurring at real  $\mu$ . Coefficients should be all positive at large  $n$ .*

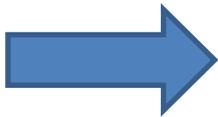
# Finite radius of convergence : ENDPOINT



$$\frac{T^E}{T_c} = 0.94 \pm 0.01, \text{ and } \frac{\mu_B^E}{T^E} = 1.8 \pm 0.1.$$

Extrapolation of this result to the thermodynamic limit,  $L \rightarrow \infty$  on the coarse lattice

:

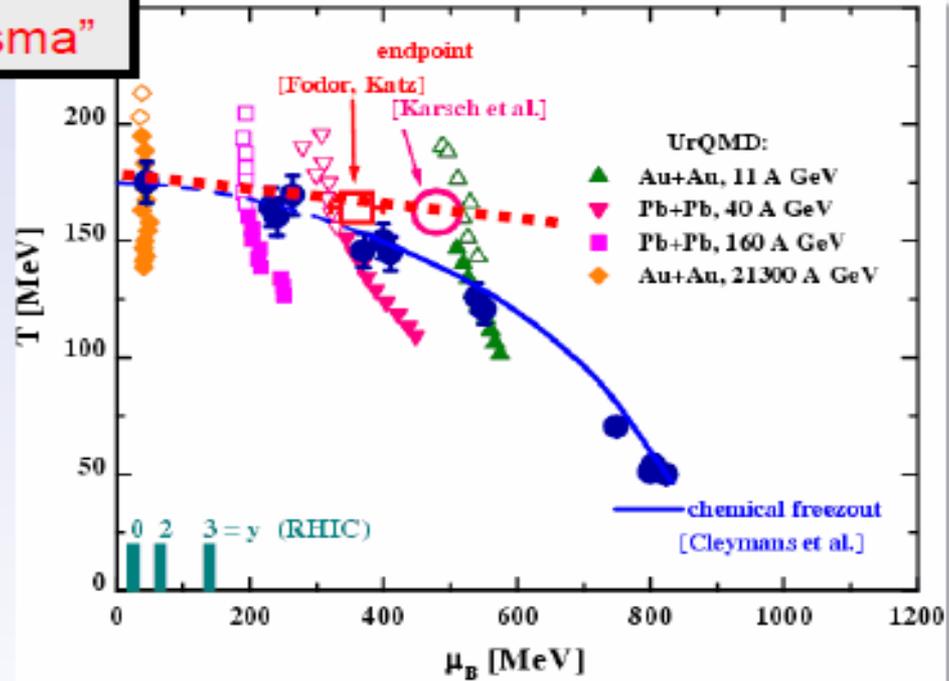


$$\frac{T^E}{T_c} = 0.94 \pm 0.01, \text{ and } \frac{\mu_B^E}{T^E} = 1.1 \pm 0.1.$$

Gvai Gupta 2007--2009

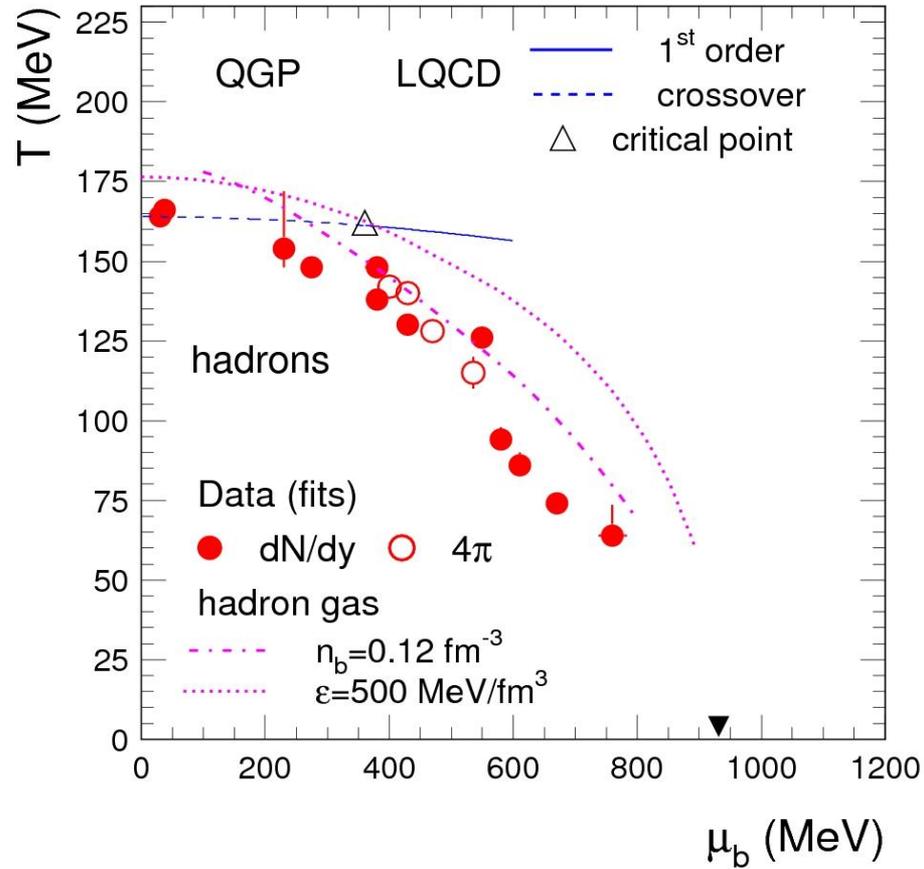
**FREEZOUT**

plasma"



# Freezout

Andronic, Braun-Munzinger, Stachel 2009 –  
Courtesy of the Authors



## FREEZOUT

Values of  $\mu_q^F/T$  at freezeout for the temperatures used in the lattice simulations.

**Table 1: Freezeout parameters**

$T/T_c$	$\mu_B^F$ (GeV)	$\mu_q^F/T$
0.81	0.48	1.16
0.87	0.38	0.85
0.90	0.3	0.65
0.96	0.15	0.30

Previous analysis have shown that for this range of temperatures the Hadron Gas parametrization is satisfied by the first coefficients.

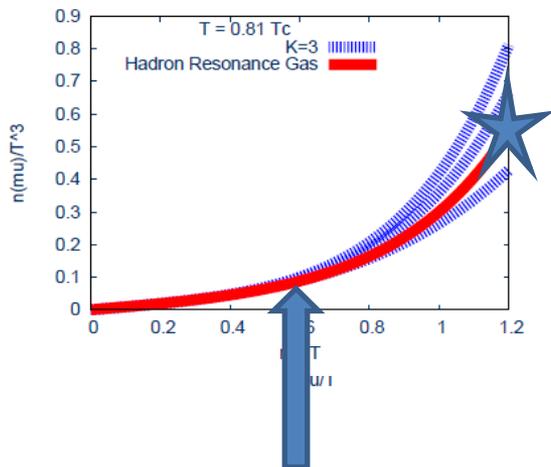
Then, to assess the extent of the convergence, we can directly contrast  $n_q^3(T, \mu_{qI})/T^3$  and  $n_q^{HG}(T, \mu_{qI})/T^3$ , with  $F(T) = \frac{2}{3}c_2$ .

# Lower bound on the radius of convergence And freezout point

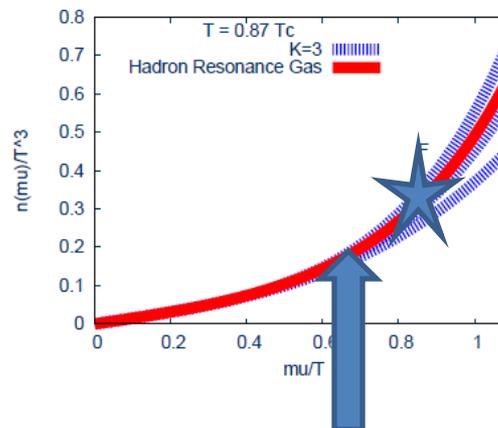
Data from RBC Collaboration  
Courtesy E. Laermann and C.  
Schmidt. C. Ratti and MpL QM09

★ Freezout point

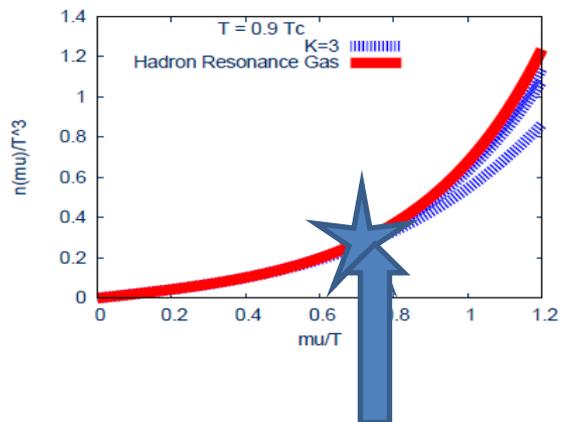
$T = 0.81 T_c$



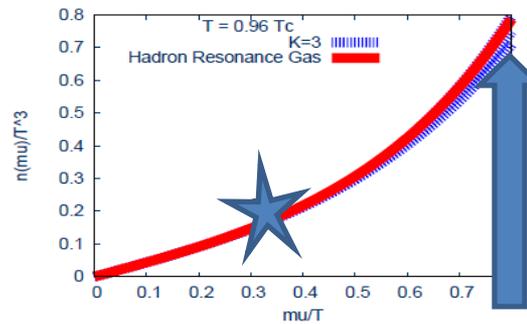
$T = 0.87 T_c$



$T = 0.90 T_c$

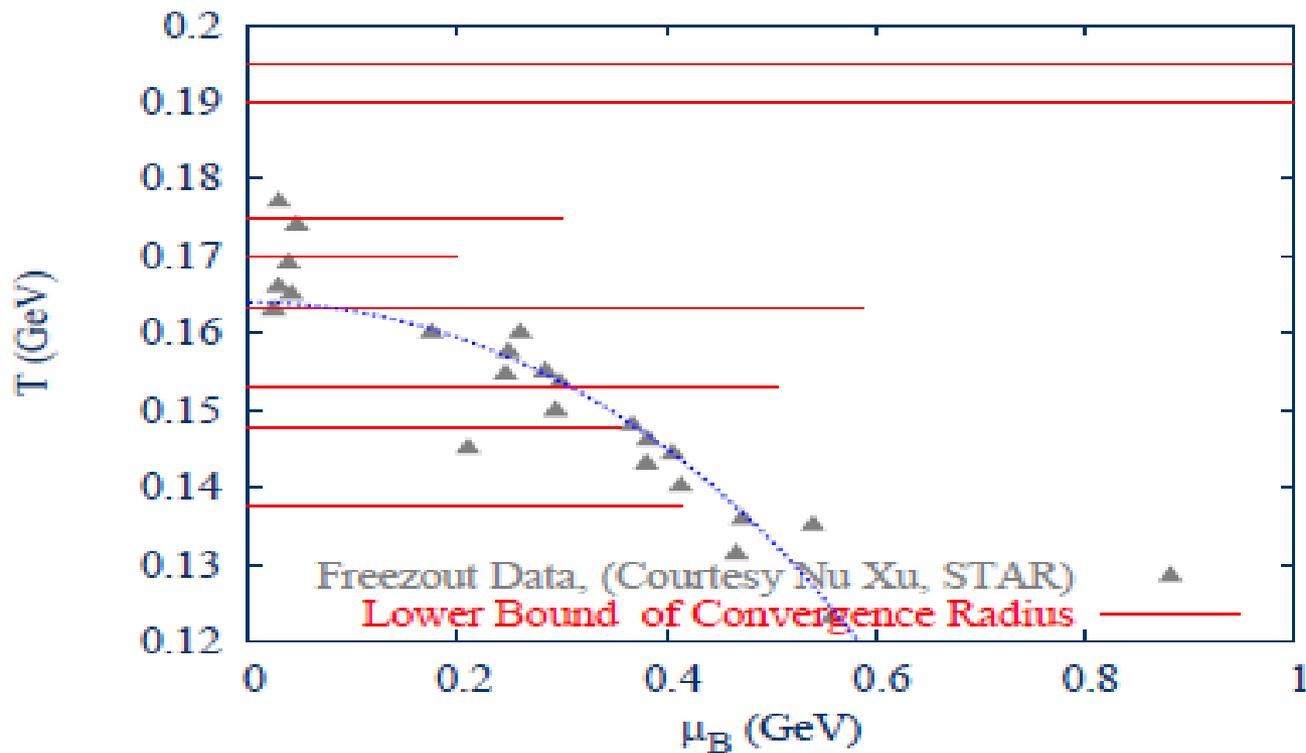


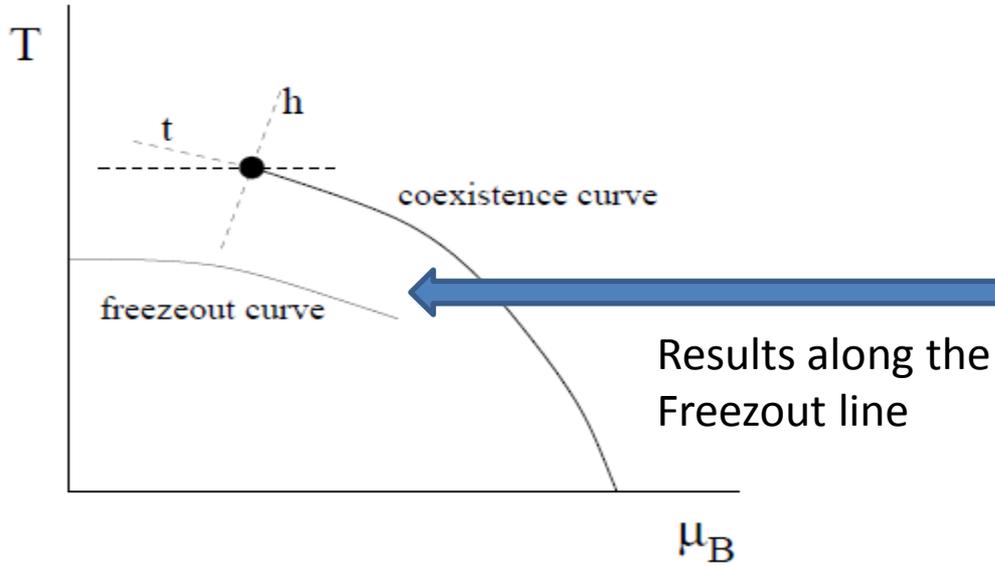
$T = 0.96 T_c$



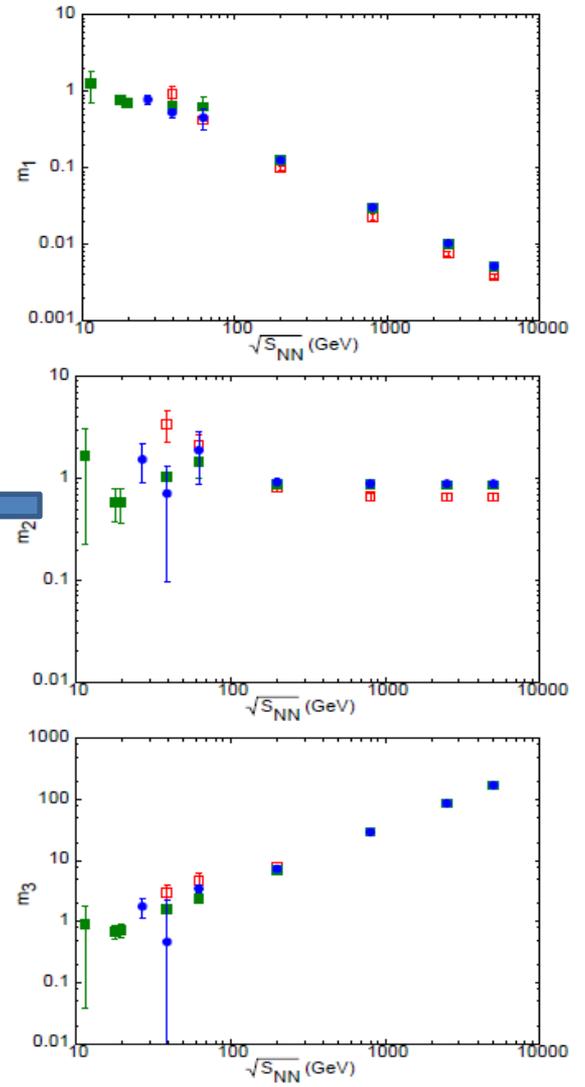
# Freezout line might well be amenable to a lattice study

Lattice data from RBC-Bielefeld  
Collaboration – C. Ratti MpL QM09





Gavai Gupta 2010



# THE QUARKYONIC PHASE

# The quarkyonic phase

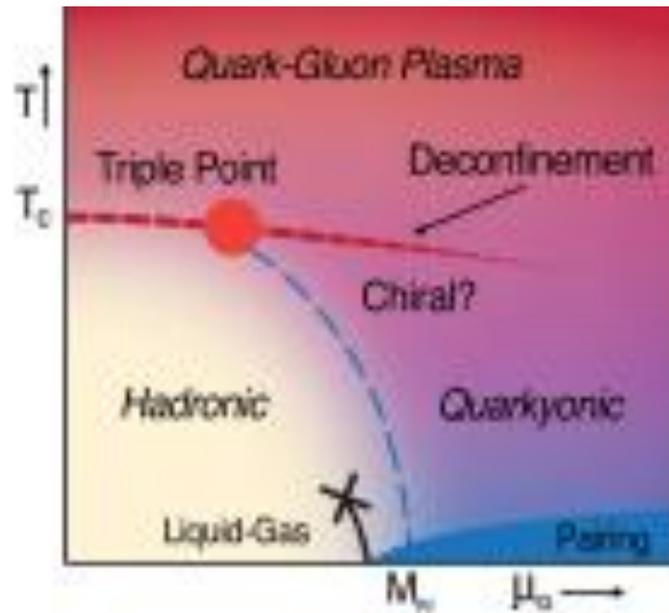
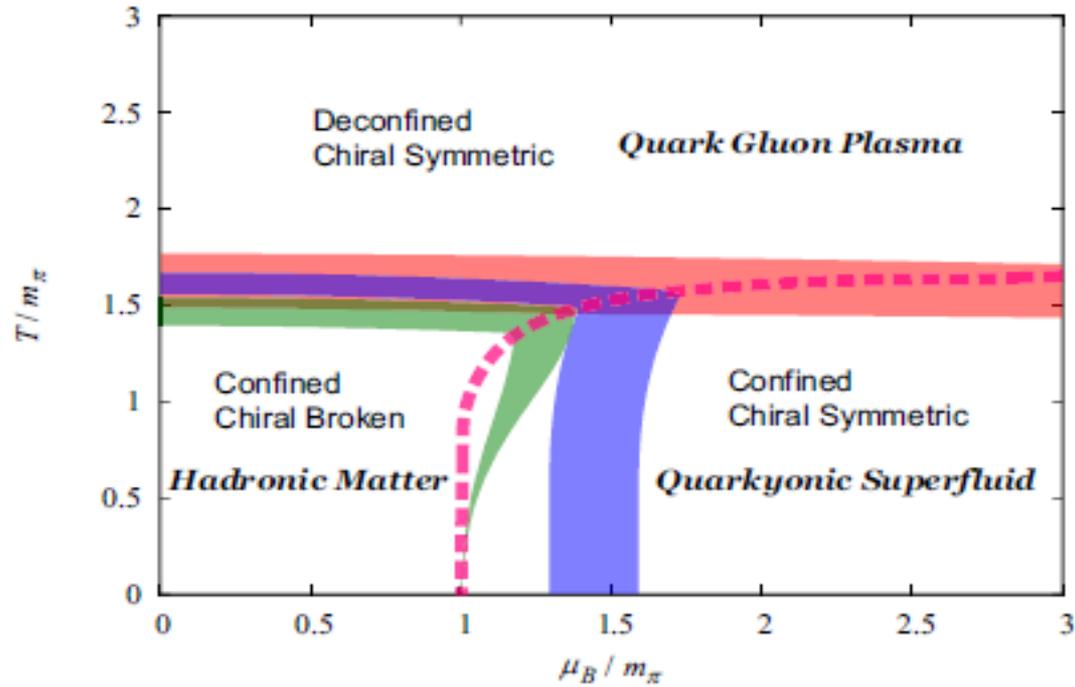


Fig. 5. The phase diagram of strongly interacting matter.

**Hadron Production in Ultra-relativistic Nuclear Collisions: Quarkyonic Matter and a Triple Point in the Phase Diagram of QCD**

A. Andronic<sup>a</sup>, D. Blaschke<sup>b,c</sup>, P. Braun-Munzinger<sup>a,d,e,f</sup>,  
J. Cleymans<sup>g</sup>, K. Fukushima<sup>h</sup>, L.D. McLerran<sup>ij</sup>, H. Oeschler<sup>g</sup>,  
R.D. Pisarski<sup>l</sup>, K. Redlich<sup>a,b,k</sup>, C. Sasaki<sup>l,f</sup>, H. Satz<sup>k</sup>, and  
J. Stachel<sup>m</sup>

# Quarkyonic phase – Two color



Brauner, Fukushima, Hidaka 2009

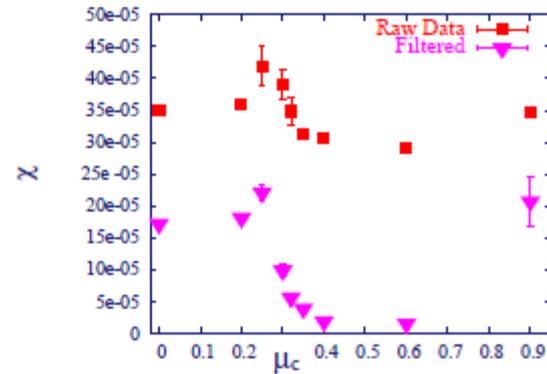
# Superfluid phase still confined

## GLUONIC OBSERVABLES IN THE BEC PHASE of QC<sub>2</sub>D

$0^{++}$  Glueball : lighter in the BEC phase

Susceptibility:  $\chi = \langle P^2 \rangle - \langle P \rangle^2$  peaks at  $\mu_c$

Normal Phase	
$m_\pi/m_\rho$	$m_0^{++}/m_\rho$
0.40	1.07
0.42	1.26
BEC	
0.64	0.80
0.80	0.23



# A Quarkyonic Phase in Dense Two Color Matter

Simon Hands

*Department of Physics, Swansea University, Singleton Park, Swansea SA2 8PP, U.K.*

Seyong Kim

*Department of Physics, Sejong University, Seoul 143-747, Korea.*

Jon-Ivar Skullerud

*Department of Mathematical Physics, National University of Ireland Maynooth, Maynooth, County Kildare, Ireland.*

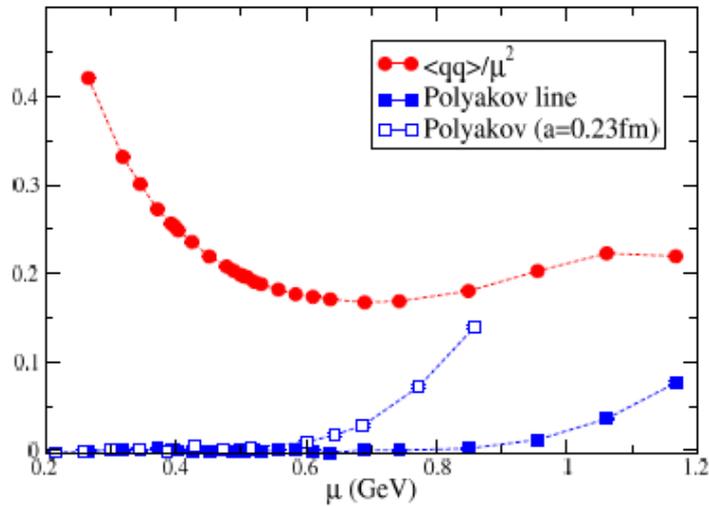


FIG. 4: (Color online) Superfluid order parameter  $\langle qq \rangle / \mu^2$  and Polyakov line versus  $\mu$ .

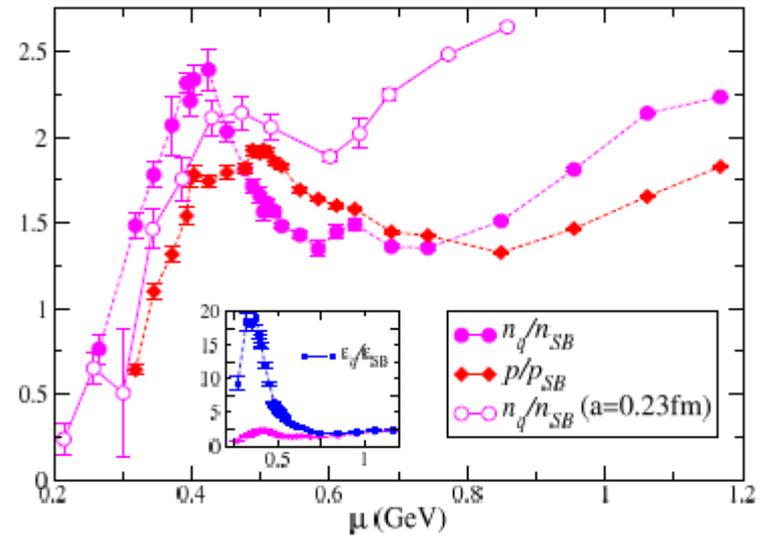


FIG. 1: (Color online)  $n_q/n_{SB}$  and  $p/p_{SB}$  vs.  $\mu$  for QC<sub>2</sub>D. Inset shows  $\epsilon_q/\epsilon_{SB}$  for comparison.

# Strong Interactions and Finite Baryon density in the (second decade of) XXI Century:

