



Neutrino Physics: an Introduction

Lecture 3: Neutrinos in astrophysics and cosmology

Amol Dighe

Department of Theoretical Physics
Tata Institute of Fundamental Research, Mumbai

SERC EHEP School 2017
NISER Bhubaneswar, Nov 20-22, 2017

Lecture 1: Neutrino detection and basic properties

- Unique properties
- Discovery of neutrino flavours
- Measuring mass, helicity, interactions

Lecture 2: Neutrino mixing and oscillations

- Solar and atmospheric puzzles and solutions
- Neutrino mixing, oscillations, flavour conversions
- The three-neutrino mixing picture

Lecture 3: Neutrinos in astrophysics and cosmology

- Low-energy (meV) cosmological neutrinos
- Medium-energy (MeV) supernova neutrinos
- High-energy ($> \text{TeV}$) astrophysical neutrinos

Neutrino Physics: an Introduction (Lecture 3)

- 1 Matter effects and future detectors
- 2 Low-energy (meV) cosmological neutrinos
- 3 Medium-energy (MeV) supernova neutrinos
- 4 High-energy ($> \text{TeV}$) astrophysical neutrinos

Effect of matter on mixing angle and Δm^2 , Resonance

Constant matter density: $\Delta \equiv \frac{\Delta m^2}{4E}$, $V_C \equiv \sqrt{2}G_F N_e$

Neutrinos

- $$\tan 2\theta_m = \frac{2\Delta \sin 2\theta}{2\Delta \cos 2\theta - V_C}$$

- $$\Delta_m = \sqrt{\left(\Delta \cos 2\theta - \frac{V_C}{2}\right)^2 + (\Delta \sin 2\theta)^2}$$

Antineutrinos

- $$\tan 2\theta_m = \frac{2\Delta \sin 2\theta}{2\Delta \cos 2\theta + V_C}$$

- $$\Delta_m = \sqrt{\left(\Delta \cos 2\theta + \frac{V_C}{2}\right)^2 + (\Delta \sin 2\theta)^2}$$

Conversion probability and CP violation

$$\begin{aligned}
 P_{e\mu} \simeq & \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2[(1 - \hat{A})\Delta]}{(1 - \hat{A})^2} \quad \alpha \equiv \frac{\Delta m_{21}^2}{\Delta m_{31}^2}, \Delta \equiv \frac{\Delta m_{31}^2 L}{4E}, \hat{A} \equiv \frac{2\sqrt{2}G_F n_e E}{\Delta m_{31}^2} \\
 & + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{\text{CP}} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 - \hat{A})\Delta]}{(1 - \hat{A})} \\
 & + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{\text{CP}} \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 - \hat{A})\Delta]}{(1 - \hat{A})} \\
 & + \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

(Cervera et al. 2000; Freund, Huber, Lindner, 2000; Huber, Winter, 2003; Akhmedov et al, 2004)

$$\Delta \equiv \Delta_{31} L$$

(Slide from W. Winter)

Conversion probability and CP violation

$$\begin{aligned}
 P_{e\mu} &\simeq \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})^2} \quad \alpha \equiv \frac{\Delta m_{21}^2}{\Delta m_{31}^2}, \Delta \equiv \frac{\Delta m_{31}^2 L}{4E}, \hat{A} \equiv \frac{2\sqrt{2}G_F n_e E}{\Delta m_{31}^2} \\
 &+ \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{\text{CP}} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\
 &+ \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{\text{CP}} \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\
 &+ \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

■ **Antineutrinos:** $P_{e\bar{\mu}} = P_{e\mu}(\delta_{\text{CP}} \rightarrow -\delta_{\text{CP}}, \hat{A} \rightarrow -\hat{A})$

(Cervera et al. 2000; Freund, Huber, Lindner, 2000; Huber, Winter, 2003; Akhmedov et al, 2004)

$$\Delta \equiv \Delta_{31} L \quad (\text{Slide from W. Winter})$$

Conversion probability and CP violation

$$P_{e\mu} \simeq \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})^2} \quad \alpha \equiv \frac{\Delta m_{21}^2}{\Delta m_{31}^2}, \Delta \equiv \frac{\Delta m_{31}^2 L}{4E}, \hat{A} \equiv \frac{2\sqrt{2}G_F n_e E}{\Delta m_{31}^2}$$

$$\begin{aligned} & \mp \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{\text{CP}} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\ & + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{\text{CP}} \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\ & + \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2} \end{aligned}$$

- **Antineutrinos:** $P_{e\bar{\mu}} = P_{e\mu}(\delta_{\text{CP}} \rightarrow -\delta_{\text{CP}}, \hat{A} \rightarrow -\hat{A})$
- **Magic baseline:** $\sin(\hat{A}\Delta) = 0 \Rightarrow \sqrt{2}G_F n_e(L)L = 2\pi$

(Cervera et al. 2000; Freund, Huber, Lindner, 2000; Huber, Winter, 2003; Akhmedov et al, 2004)

$$\Delta \equiv \Delta_{31} L$$

(Slide from W. Winter)

Conversion probability and CP violation

$$\begin{aligned}
 P_{e\mu} &\simeq \sin^2 2\theta_{13} \frac{\cos^2 \theta_{23}}{\sin^2 \theta_{23}} \frac{\sin^2[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})^2} \quad \alpha \equiv \frac{\Delta m_{21}^2}{\Delta m_{31}^2}, \Delta \equiv \frac{\Delta m_{31}^2 L}{4E}, \hat{A} \equiv \frac{2\sqrt{2}G_F n_e L}{\Delta m_{31}^2} \\
 &\pm \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{\text{CP}} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\
 &\pm \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{\text{CP}} \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\
 &+ \alpha^2 \frac{\sin^2 \theta_{23}}{\cos^2 \theta_{23}} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

- **Antineutrinos:** $P_{e\bar{\mu}} = P_{e\mu}(\delta_{\text{CP}} \rightarrow -\delta_{\text{CP}}, \hat{A} \rightarrow -\hat{A})$
- **Magic baseline:** $\sin(\hat{A}\Delta) = 0 \Rightarrow \sqrt{2}G_F n_e(L)L = 2\pi$
- **Silver:** $P_{e\tau} = P_{e\mu}(s_{23}^2 \leftrightarrow c_{23}^2, \sin 2\theta_{23} \rightarrow -\sin 2\theta_{23})$

(Cervera et al. 2000; Freund, Huber, Lindner, 2000; Huber, Winter, 2003; Akhmedov et al, 2004)

$$\Delta \equiv \Delta_{31} L$$

(Slide from W. Winter)

Conversion probability and CP violation

$$\begin{aligned}
 P_{e\mu} &\simeq \sin^2 2\theta_{13} \frac{\cos^2 \theta_{23}}{\sin^2 \theta_{23}} \frac{\sin^2[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})^2} \quad \alpha \equiv \frac{\Delta m_{21}^2}{\Delta m_{31}^2}, \Delta \equiv \frac{\Delta m_{31}^2 L}{4E}, \hat{A} \equiv \frac{2\sqrt{2}G_F n_e E}{\Delta m_{31}^2} \\
 &\quad + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{\text{CP}} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\
 &\quad + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{\text{CP}} \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\
 &\quad + \alpha^2 \frac{\sin^2 \theta_{23}}{\cos^2 \theta_{23}} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

- **Antineutrinos:** $P_{e\bar{\mu}} = P_{e\mu}(\delta_{\text{CP}} \rightarrow -\delta_{\text{CP}}, \hat{A} \rightarrow -\hat{A})$
- **Magic baseline:** $\sin(\hat{A}\Delta) = 0 \Rightarrow \sqrt{2}G_F n_e(L)L = 2\pi$
- **Silver:** $P_{e\tau} = P_{e\mu}(s_{23}^2 \leftrightarrow c_{23}^2, \sin 2\theta_{23} \rightarrow -\sin 2\theta_{23})$
- **Platinum, Superb.:** $P_{\mu e} = P_{e\mu}(\delta_{\text{CP}} \rightarrow -\delta_{\text{CP}})$

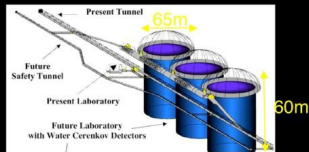
(Cervera et al. 2000; Freund, Huber, Lindner, 2000; Huber, Winter, 2003; Akhmedov et al, 2004)

$$\Delta \equiv \Delta_{31} L$$

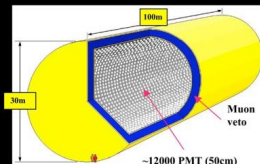
(Slide from W. Winter)

Future directions of multipurpose detectors

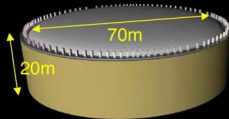
- Three types of large multi-purpose underground detectors with astrophysical program



Water Cherenkov ($\approx 0.5 \rightarrow 1$ Mton)
MEMPHYS



Liquid Scintillator ($\rightarrow 50$ kton)
LENA

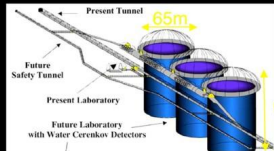


Liquid Argon ($\approx 10 \rightarrow 100$ kton)
GLACIER

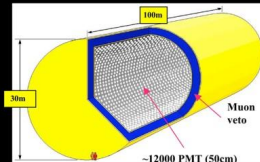
- Megaton water Cherenkov detectors
- 50 kiloton scintillator detectors
- 100 kiloton liquid Ar detectors

Future directions of multipurpose detectors

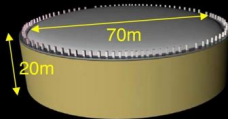
- Three types of large multi-purpose underground detectors with astrophysical program



Water Cherenkov ($\approx 0.5 \rightarrow 1$ Mton)
MEMPHYS



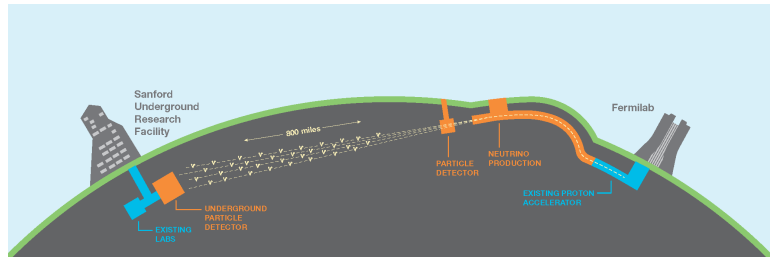
Liquid Scintillator ($\rightarrow 50$ kton)
LENA



Liquid Argon ($\approx 10 \rightarrow 100$ kton)
GLACIER

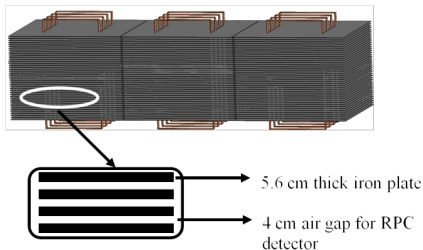
- Megaton water Cherenkov detectors
- 50 kiloton scintillator detectors
- 100 kiloton liquid Ar detectors
- Advantages ?

Long baseline experiments



- Deep Underground Neutrino Experiment (DUNE)
- Detector 1600 km away from source

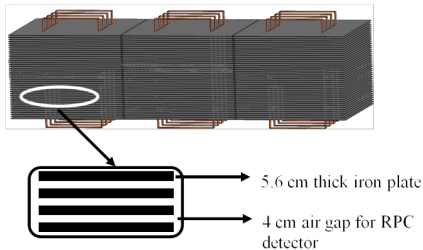
Coming soon inside a mountain near you: INO



India-based Neutrino Observatory

- 1 km rock coverage from all sides
- 50 kiloton of magnetized iron (50 000 000 kg)

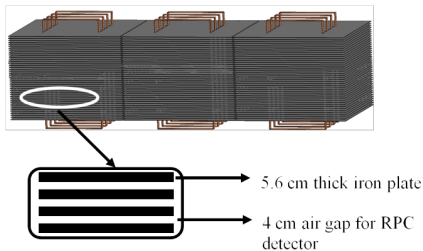
Coming soon inside a mountain near you: INO



India-based Neutrino Observatory

- 1 km rock coverage from all sides
- 50 kiloton of magnetized iron (50 000 000 kg)
- A large range of L and E available
- Can distinguish neutrinos from antineutrinos

Coming soon inside a mountain near you: INO

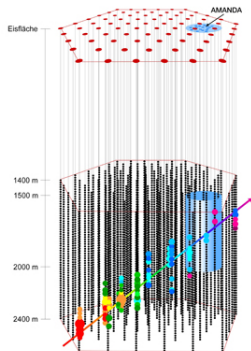
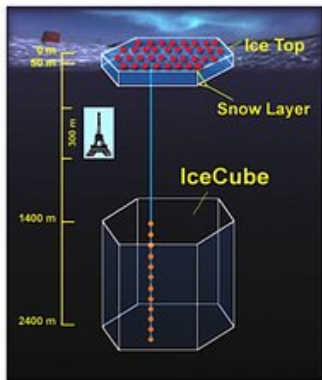


India-based Neutrino Observatory

- 1 km rock coverage from all sides
- 50 kiloton of magnetized iron (50 000 000 kg)
- A large range of L and E available
- Can distinguish neutrinos from antineutrinos
- Determining mass ordering from atmospheric neutrinos
- Can look for dark matter, magnetic monopoles, ...

Below the antarctic ice: Gigaton IceCube

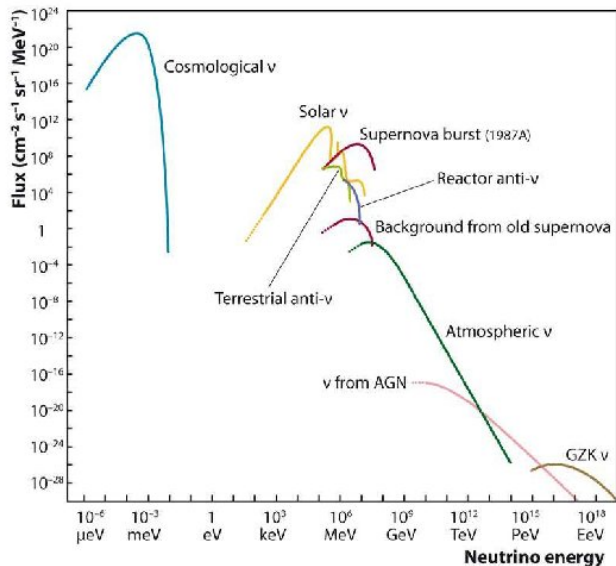
1 000 000 000 000 litres of ice



Neutrinos as messengers

- No bending in magnetic fields \Rightarrow point back to the source
- Minimal obstruction / scattering \Rightarrow can arrive directly from regions from where light cannot come.

Neutrino fluxes at different energies



Neutrino Physics: an Introduction (Lecture 3)

- 1 Matter effects and future detectors
- 2 Low-energy (meV) cosmological neutrinos
- 3 Medium-energy (MeV) supernova neutrinos
- 4 High-energy ($> \text{TeV}$) astrophysical neutrinos

Neutrino Physics: an Introduction (Lecture 3)

- 1 Matter effects and future detectors
- 2 Low-energy (meV) cosmological neutrinos
- 3 Medium-energy (MeV) supernova neutrinos
- 4 High-energy ($> \text{TeV}$) astrophysical neutrinos

Source: abundance and temperature

- Relic density: ~ 110 neutrinos /flavor /cm³
- Temperature: $T_\nu = (4/11)^{1/3} T_{\text{CMB}} \approx 1.95 \text{ K} = 0.16 \text{ meV}$
- The effective number of neutrino flavors:
 $N_{\text{eff}}(\text{SM}) = 3.074$. Planck $\Rightarrow N_{\text{eff}} = 3.30 \pm 0.27$.
- Contribution to dark matter density:

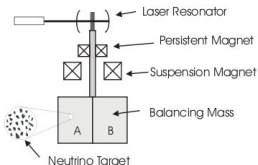
$$\Omega_\nu / \Omega_{\text{baryon}} = 0.5 \left(\sum m_\nu / \text{eV} \right)$$

- Looking really far back:

	Time	Temp	z
Relic neutrinos	0.18 s	$\sim 2 \text{ MeV}$	$\sim 10^{10}$
CMB photons	$\sim 4 \times 10^5 \text{ years}$	0.26 eV	1100

Lazauskas, Vogel, Volpe, 2008

Detection of relic neutrinos: the torsion balance idea



- De Brogli wavelength of relic neutrinos: $\lambda \approx h/p \approx 1.5\text{mm}$.
 - ν can interact coherently with a sphere of this size
 - Measure force on such “spheres” due to the relic neutrino wind
 - For iron spheres and 100 times local overdensity for ν , acceleration $a \lesssim 10^{-26} \text{ cm/s}^2$
- Shvartsman et al 1982
- $\gtrsim 10$ orders of magnitude smaller than the sensitivity of current torsion balance technology
 - If neutrinos are Majorana, a further suppression by $v/c \approx 10^3$ (polarized target), $(v/c)^2 \approx 10^{-6}$ (unpolarized)
- Hagmann, astro-ph/9901102
- The idea is essentially impractical.

The inverse beta reaction

- Need detection of low-energy neutrinos, so look for zero-threshold interactions
- Beta-capture on beta-decaying nuclei:



End-point region ($E > M_{N_1} - M_{N_2}$) background-free.
Energy resolution crucial.

Weinberg 1962, cocco, Mangano, Messina 2008, Lazauskas et al 2008, Hodak et al 2009

- Possible at ^3H experiments with 100 g of pure tritium but atomic tritium is needed to avoid molecular energy levels

Lazauskas, Vogel, Volpe 2009, Hodak et al 2011

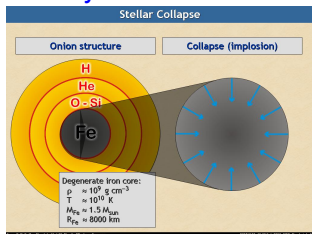
- PTOLEMY: Princeton Tritium Observatory for Light, Early-universe, Massive neutrino Yield

Neutrino Physics: an Introduction (Lecture 3)

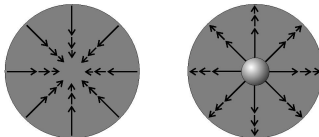
- 1 Matter effects and future detectors
- 2 Low-energy (meV) cosmological neutrinos
- 3 Medium-energy (MeV) supernova neutrinos**
- 4 High-energy ($> \text{TeV}$) astrophysical neutrinos

Supernova: the death of a star

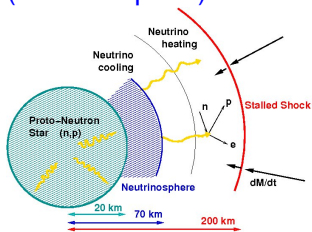
Gravity \Rightarrow



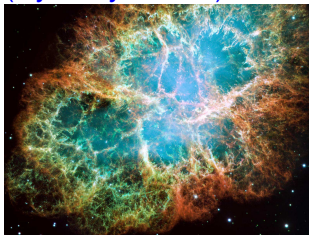
Strong nuclear force \Rightarrow



Weak nuclear force
(Neutrino push) \Rightarrow

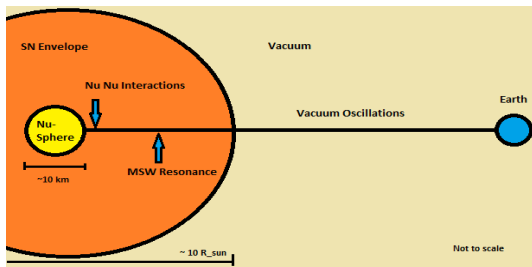


Electromagnetism
(Hydrodynamics) \Rightarrow



(Crab nebula, SN seen in 1054)

Neutrino oscillations in matter of varying density



Inside the SN: *flavour conversion*

Non-linear “collective” effects and resonant matter effects

Between the SN and Earth: *no flavour conversion*

Neutrino mass eigenstates travel independently

Inside the Earth: *flavour oscillations*

Resonant matter effects (*if detector is shadowed by the Earth*)

Can neutrino conversions affect SN explosions ?

- Simulations of light SN have started giving explosions with the inclusions of 2D/3D large scale convections and hydrodynamic instabilities
- More push to the shock wave is still desirable.

Can neutrino conversions affect SN explosions ?

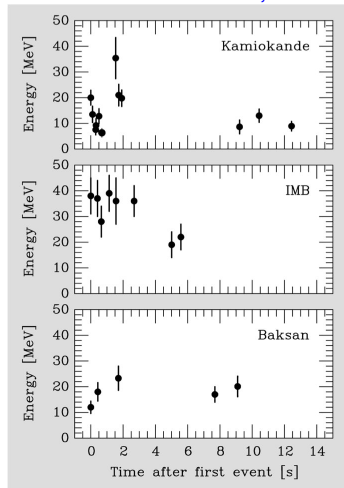
- Simulations of light SN have started giving explosions with the inclusions of 2D/3D large scale convections and hydrodynamic instabilities
- More push to the shock wave is still desirable.
- Non-electron neutrino primary spectra harder
 - ⊕ electron neutrino cross section higher
 - ⇒ After conversion, greater push to the shock wave
- Deeper the conversions, greater the neutrino push

Can neutrino conversions affect SN explosions ?

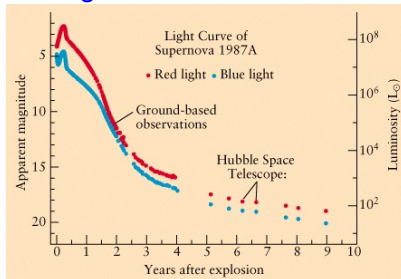
- Simulations of light SN have started giving explosions with the inclusions of 2D/3D large scale convections and hydrodynamic instabilities
- More push to the shock wave is still desirable.
- Non-electron neutrino primary spectra harder
 - ⊕ electron neutrino cross section higher
 - ⇒ After conversion, greater push to the shock wave
- Deeper the conversions, greater the neutrino push
- Neutrino flavour conversions in extremely dense media:
 - MSW resonances: 1000 km,
 - Neutrino-neutrino collective effects: 100 km
 - “Fast conversions”: 10 km [Angular anisotropies needed, but quite naturally possible]

SN1987A: neutrinos and light

Neutrinos: Feb 23, 1987

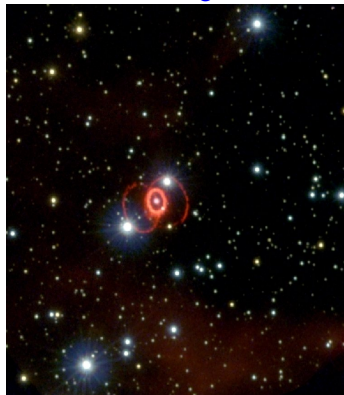


Light curve: 1987-1997



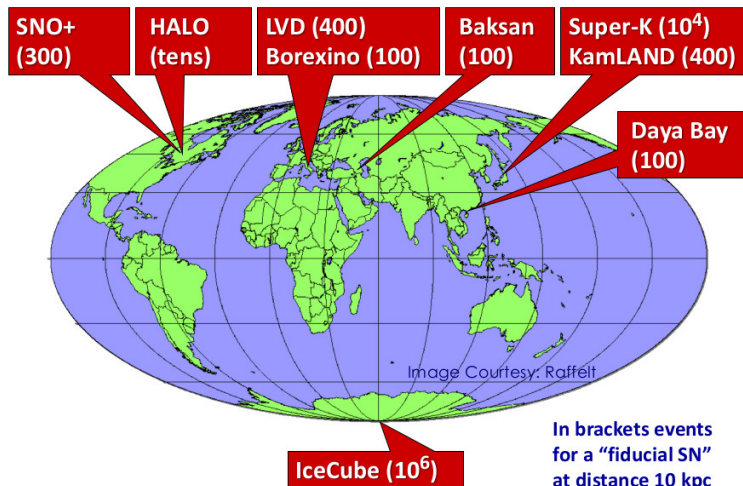
SN1987A: what did we learn ?

Hubble image: now



- Confirmed the SN cooling mechanism through neutrinos
- Number of events too small to say anything concrete about neutrino mixing
- Some constraints on SN parameters obtained
- Strong constraints on new physics models obtained (neutrino decay, Majorans, axions, extra dimensions, ...)

Supernova neutrino detectors



What supernova neutrinos can tell us

On neutrino masses and mixing

- Identify neutrino mass ordering: **normal** or **inverted**

What supernova neutrinos can tell us

On neutrino masses and mixing

- Identify neutrino mass ordering: **normal or inverted**

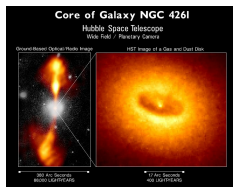
On supernova astrophysics

- Locate a supernova hours before the light arrives
- Track the shock wave through neutrinos while it is still inside the mantle (**Not possible with light**)
- How is a neutron star / black hole formed ? Is there a QCD phase transition ?
- How are heavy elements formed ?

Neutrino Physics: an Introduction (Lecture 3)

- 1 Matter effects and future detectors
- 2 Low-energy (meV) cosmological neutrinos
- 3 Medium-energy (MeV) supernova neutrinos
- 4 High-energy ($> \text{TeV}$) astrophysical neutrinos

Sources of high-energy neutrinos



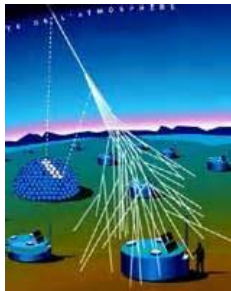
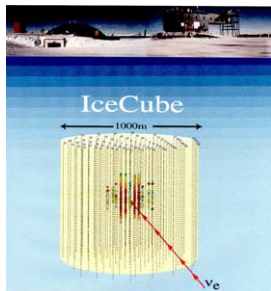
The origins

- Primary protons interacting within the source or with CMB photons $\Rightarrow \pi^{\pm} \Rightarrow$ Decay to ν
- Individual sources like AGNs and GRBs
- Diffused flux accumulated over the lifetime of universe

What we will learn

- Mechanisms of astrophysical phenomena
- Limits on neutrino decay, Lorentz violation, etc

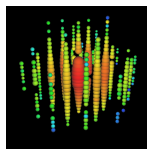
Detection of high energy neutrinos



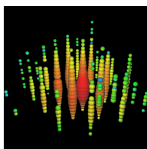
Detection techniques

- Water Cherenkov like IceCube: $10^{11} \text{ eV} \lesssim E \lesssim 10^{16} \text{ eV}$
- Cosmic ray arrays for $E \gtrsim 10^{17} \text{ eV}$
- Radio detection from balloon experiments (Askaryan)

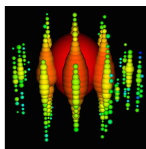
Highest energy neutrinos observed till now



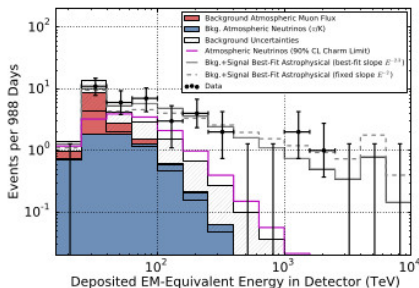
Bert



Ernie



Big Bird



- Three events at $\sim 1, 1.1, 2.2$ PeV energies found
- Cosmogenic ? X Glashow resonance? X atmospheric ?
Roulet et al 2013 ++ many
- IceCube analyzing 54 events from 30 TeV to 10 PeV

Flavor information from UHE neutrinos

Flavor ratios $\nu_e : \nu_\mu : \nu_\tau$ at sources

- Neutron source (nS): $1 : 0 : 0$
- Pion source (π S): $1 : 2 : 0$,
- Muon-absorbing sources (μ DS): $0 : 1 : 0$

Flavor information from UHE neutrinos

Flavor ratios $\nu_e : \nu_\mu : \nu_\tau$ at sources

- Neutron source (nS): $1 : 0 : 0$
- Pion source (π S): $1 : 2 : 0$,
- Muon-absorbing sources (μ DS): $0 : 1 : 0$

Flavor ratios at detectors

- Neutron source: $\approx 5 : 2 : 2$
- Pion source: $\approx 1 : 1 : 1$
- Muon-absorbing sources : $\approx 4 : 7 : 7$

Flavor information from UHE neutrinos

Flavor ratios $\nu_e : \nu_\mu : \nu_\tau$ at sources

- Neutron source (nS): $1 : 0 : 0$
- Pion source (π S): $1 : 2 : 0$,
- Muon-absorbing sources (μ DS): $0 : 1 : 0$

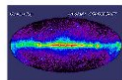
Flavor ratios at detectors

- Neutron source: $\approx 5 : 2 : 2$
- Pion source: $\approx 1 : 1 : 1$
- Muon-absorbing sources : $\approx 4 : 7 : 7$

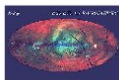
New physics effects

- Decaying neutrinos can skew the flavor ratio even further:
as extreme as $6 : 1 : 1$ or $0 : 1 : 1$
Ratio measurement \Rightarrow improved limits on neutrino lifetimes

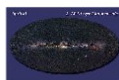
Dawn of multi-messenger astronomy



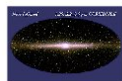
Gamma ray



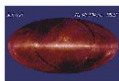
X-ray



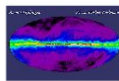
Visible



Near infrared



Infrared



Radio waves

