Flavour physics: Lecture 2 Flavour data and New Physics

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Flavor physics: a wall of the SM edifice

Building up the Standard Model

- $\tau \theta$ puzzle \Rightarrow Parity violation
- Cabibbo angle ⇒
 weak coupling universality ⊕ quark mixing
- GIM mechanism ⇒ no FCNC at tree level, charm
- CKM paradigm ⇒ (at least) three quark families
- Large $B \overline{B}$ mixing \Rightarrow heavy top quark
- Rate of radiative B decay ⇒ top quark mass

The standard model looks complete now, with no confirmed signals of BSM physics at the LHC yet

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Is there just a desert beyond?



In the light of the high-energy and high-intensity colliders.

We may find some pugmarks

(Need not be of the camel we are seeking)

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Flavour physics: a window of the SM edifice

Precision tests of the Standard Model

- CKM elements: do they explain all CP violation?
- Rare decays: new particles contribute through loops?
- Asymmetries: are the predicted SM relations obeyed?
- What are the features of New Physics (if any)?

How to look at the multidimentional flavour data?

- Identify patterns of correlations
- Correlations within a meson system give information on BSM operator structure (V-A, tensor, LR, etc.)
- Correlations among different meson systems tell about underlying flavour symmetries (MFV, universality, etc.)



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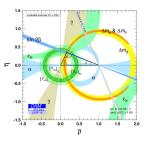
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 - CKM matrix elements
 - Mixing and decay in neutral mesons
 - Rare FCNC processes
- Quantifying constraints on NP
 - Model-independent constraints
 - Constraints on specific new physics models
- Concluding remarks

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Global fits to CKM elements

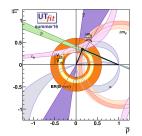
CKMfitter:



Constraints in the $\bar{\rho}$ - $\bar{\eta}$ plane:

- the ratio $|V_{ub}/V_{cb}|$
- ϵ_K from $K \to \pi\pi$
- Mass differences ΔM_d and ΔM_s
- Angles α, β, γ (or ϕ_2, ϕ_1, ϕ_3) of the unitarity triangle

UTfit: 2016-fits

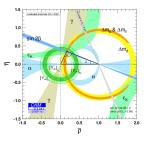


KM paradigm mostly vindicated!



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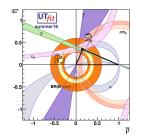
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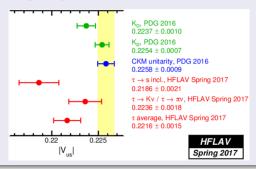
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Measurements for determination of CKM elements

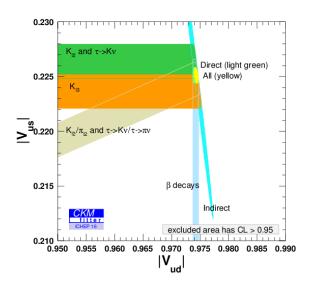
Precision measurements of $|V_{us}|$

Unitarity vs. semileptonic K decays vs. hadronic τ decays

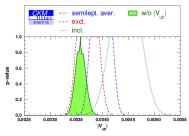
- Semileptonic K decays $\Rightarrow |V_{us}| = 0.2258 \pm 0.0009$
- Strange vs. non-strange hadronic τ decays $\Rightarrow |V_{US}| = 0.2216 \pm 0.0015$
- $> 2\sigma$ discrepancy!

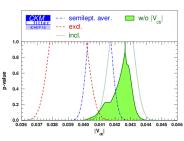


$K_{\ell 2}$ vs. $K_{\ell 3}$ tension in $|V_{us}|$



Inclusive vs exclusive $|V_{cb}|$ and $|V_{ub}|$





 $|V_{ub}| \times 10^3$:

• Excl: (3.72 ± 0.19)

• Incl: $(4.49 \pm 0.16 \pm 0.18)$

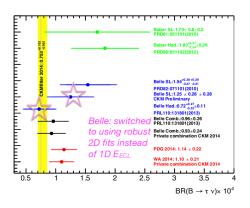
 $|V_{cb}| \times 10^3$:

• Excl: (39.2 ± 0.7)

• Incl: (42.2 ± 0.8)

CKMFitter

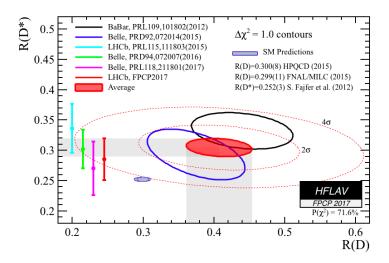
Effective $|V_{ub}|$: through $B(B \to \tau \nu)$



P. Urquio, CKM14

• Apart from the latest Belle measurement, all others indicated excess $B(B \to \tau \nu)$.

Effective $|V_{cb}|$: semileptonic $B \to D \tau \nu$ and $B \to D^* \tau \nu$



• Affect $b \to c \tau \nu$, indicate lepton-universality violation ?



Consolidating observations before going ahead

- Semileptonic decays $s \to u\ell\nu$ as well as $b \to u\ell\nu$ systematically give lower values of $|V_{uq}|$ as compared to those given by the leptonic decays of K and B.
- For semileptonic B decays $b \to u\ell\nu$ and $b \to c\ell\nu$, inclusive decay rates are systematically larger than the exclusive ones.
- Lepton non-universality at play in $b \to c \ell \nu$?
- Lepton non-universality is severely constrained in the first two generations, not so much for the third one. Models with H⁺/Z' are natural candidates.
- A single hint may not be sufficient, but overall trends may point the way..

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Mass and width differences: theory and experiment

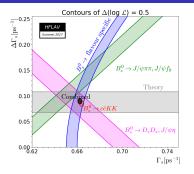
△M Measurements

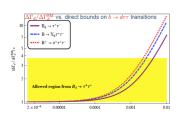
- $\Delta M_d/\Gamma_d = 0.769 \pm 0.003 \Rightarrow |V_{td}|$
- $\Delta M_s/\Gamma_s = 26.80 \pm 0.08 \Rightarrow |V_{ts}|$

$\Delta\Gamma_d$ and $\Delta\Gamma_s$: predictions and measurements

- In SM, $\Delta \Gamma_d / \Gamma_d = (42 \pm 8) \times 10^{-4}$
- Measurement: $\Delta\Gamma_d/\Gamma_d=(20\pm 100)\times 10^{-4}$ (HFAG)
- In SM, $\Delta\Gamma_s/\Gamma_s=0.137\pm0.027$
- Measurement: $\Delta\Gamma_s/\Gamma_s = 0.135 \pm 0.008$ (mainly from $B_s \rightarrow J/\psi \phi$ at LHCb)

Lifetime difference in B_s and B_d decays

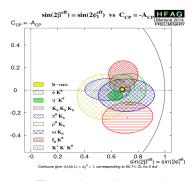


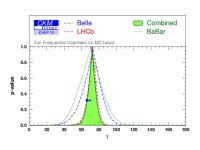


- $\Delta\Gamma$ measured through B_s to flavor-specific modes, $B_s \to K^+K^-, J/\psi f_0, B_s \to J/\psi \phi.$
- Some enhancement possible, has to be through $b \rightarrow s\tau\tau$ transitions.
- $\Delta\Gamma_d$ should also be controlled by new physics $b \rightarrow d\tau\tau$ transitions

Lenz, CKM2014

The angles of the unitarity triangle





- The value of β consistent across multiple modes
- ullet The value of γ consistent across experiments
- $\alpha = \pi \beta \gamma$ by definition. True test of unitarity:

$$\sin \beta_s = \left| \frac{V_{us}}{V_{ud}} \right|^2 \frac{\sin \beta \sin(\gamma + \beta_s)}{\sin(\beta + \gamma)} \left[1 + \mathcal{O}(\lambda^4) \right]$$
Aleksan et al, 1994

• Measurements of β_s needed



The tale of two betas in B_s - \bar{B}_s mixing

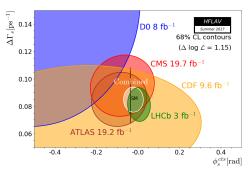
$eta_{\mathcal{S}}^{J/\psi\phi}$ from $\overline{B_{\mathcal{S}} o J/\psi\phi}$

- $\beta_s^{J/\psi\phi} \approx \frac{1}{2} \text{Arg} \left(-\frac{(V_{cb}V_{cs}^*)^2}{M_{12s}} \right)$
- $m{\bullet} \ eta_{s}^{J/\psi}(SM) = 0.019 \pm 0.001$

β_s^{sl} from a_{sl}

- $a_{sl} = (\Delta \Gamma_s / \Delta M_s) \tan \phi_s^{sl}$
- $Arg(\Gamma_{12}) \neq Arg(V_{cb}V_{cs}^*)^2$ since the (c-u) and (u-u) intermediate states contribute to Γ_{12} .
- $\phi_s^{sl}(SM) = 0.0041 \pm 0.0007$
- \bullet $\beta_s^{sl}(SM) = -0.0020 \pm 0.0003$

$\phi_{\mathcal{S}}^{J/\psi\phi}$: Angular analysis of $\mathcal{B}_{\mathcal{S}} o J/\overline{\psi\phi}$

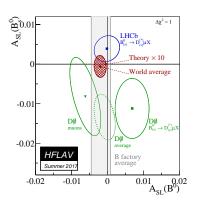


$$\phi_{\mathcal{S}} = -2eta_{\mathcal{S}}^{J/\psi\phi}$$

- Results for $\Delta\Gamma_s$ very close to SM now: enhancement by only a few tens of per cent possible.
- $\phi_s = -0.021 \pm 0.031$ rad Large (relative) enhancement in $\beta_s^{J/\psi\phi}$ is possible, also detectable since the SM value is small and precisely known.



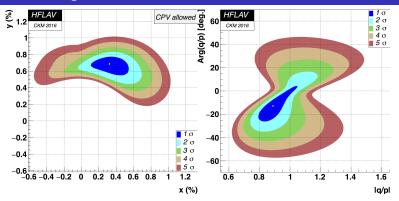
β_s^{sl} : Like-sign dimuon asymmetry



- SM \Rightarrow $A_{sl}^b = (-0.023^{+0.005}_{-0.006})\%$
- $A_{sl}^b = (-0.787 \pm 0.172 \pm 0.093)\%$ $\Rightarrow 3.9\sigma$ deviation (original D0 result)

- ullet $A_{\mathrm{CP}}(B_s^0)\equiv a_{sl}^s=rac{\Delta\Gamma_s}{\Delta M_s}\, an\phi_s^{sl}$
- Large $\Delta\Gamma_s$ and/or large ϕ_s !

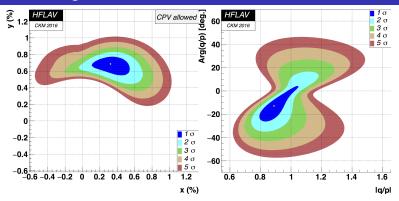
$D-\bar{D}$ mixing



- Both $\Delta m/\Gamma$ and $\Delta \Gamma/\Gamma$ significantly nonzero
- $D-\bar{D}$ mixing detected.
- Cannot be sure that it is all SM, long-distance contributions hard to calculate
- It is possible that |q/p| is significantly non-unity.
- CP violation through mixing possible, not detected yet



$D-\bar{D}$ mixing



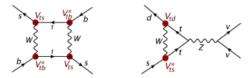
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Flavour changing neutral current processes

 Suppressed in the SM due to the loop factor, CKM hierarchy, chiral structure and GIM mechanism.



CKM hierarchy predicts specific pattern of effects in the SM

$$\underbrace{V_{ts}^* V_{td}}_{K \text{ system}} \sim 5 \cdot 10^{-4} \ll \underbrace{V_{tb}^* V_{td}}_{B_d \text{ system}} \sim 10^{-2} < \underbrace{V_{tb}^* V_{ts}}_{B_s \text{ system}} \sim 4 \cdot 10^{-2}$$

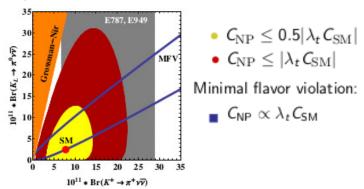
 $\succ K$ decays in general most sensitive to BSM physics

(Slide from M. Blanke)



The rare but clean decay $K \to \pi \nu \bar{\nu}$ ***

Not yet observed, but...

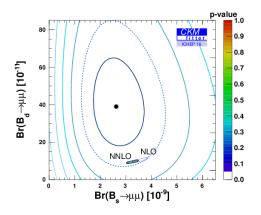


J. Brod, CKM14

• Models can change the relative BRs of $K_L \to \pi^0 \nu \bar{\nu}$ and $K^+ \to \pi^+ \nu \bar{\nu}$ to a large extent



Branching ratios of $B_s \to \mu^+\mu^-$ and $B_d \to \mu^+\mu^-$



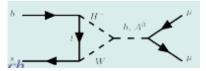
•
$$B(B_s \to \mu\mu) = (3.145^{+0.150}_{-0.069}) \times 10^{-9}$$

•
$$B(B_d \to \mu\mu) = (9.55^{+0.25}_{-0.44}) \times 10^{-11}$$



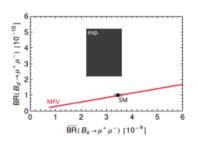
What NP can affect $B(B_{s/d} \rightarrow \mu\mu)$?

Sensitive to minimal SUSY model parameters:



$$B(B_q o \mu \mu) \propto |V_{tb}^* V_{tq}| rac{m_b^2 m_\ell^2 an^6 eta}{m_A^4}$$

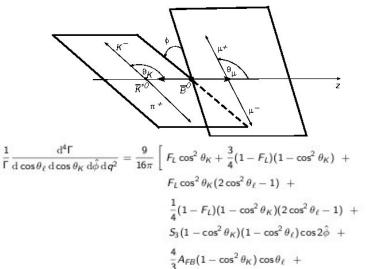
• Severely restricts large $\tan \beta$



- MFV models cannot account for the observed values
 Buras 2014
- Role for Z'?

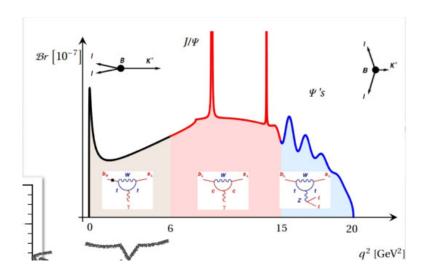


Angular distribution in $B \to K^* \mu^+ \mu^-$

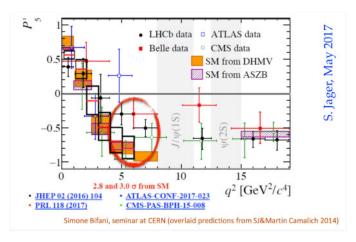


 $A_{lm}(1-\cos^2\theta_K)(1-\cos^2\theta_\ell)\sin 2\hat{\phi}$

q^2 distribution of $B o K^* \mu^+ \mu^-$



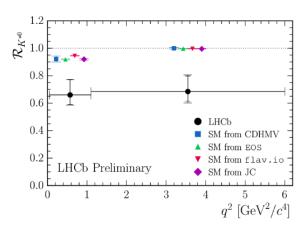
The P_5' anomaly



- $P_5' = \frac{S_5}{\sqrt{F_L(1-F_L)}}$, largely free from formfactor uncertainties
- Local discrepancy of 3.7σ in P'_5 .



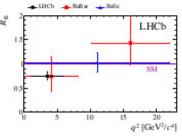
Lepton non-universality in $B \to K^* \ell \ell$



- ullet The ratio $R_K(q^2) \equiv B(B^+ o K^* \mu^+ \mu^-)/B(B^+ o K^* e^+ e^-)$
- $2.2 2.5 \sigma$ deviation in each bin

Lepton non-universality in $B \to K\ell\ell$

• The ratio $R_K \equiv B(B^+ \to K^+ \mu^+ \mu^-)/B(B^+ \to K^+ e^+ e^-)$ expected to be 1.00 in SM



- Measured value for 1 < q^2 < 9 GeV²: $R_K = 0.745^{+0.090}_{-0.074} \pm 0.035 \Rightarrow \sim 2.6\sigma$ deviation from SM
- Non-universality from Higgs not enough since Higgs contribution is m_{ℓ} -suppressed.
- Models with Z' that does not couple to electrons can explain this anomaly.



Possible enhancement of rare top FCNC decays



MSSM



RS flavor



	SM	2HDM	MSSM	RS
$t \rightarrow cZ$	$\sim 10^{-14}$	$\lesssim 10^{-6}$	$\lesssim 10^{-7}$	$\lesssim 10^{-5}$
$t o c\gamma$	$\sim 10^{-13}$	$\lesssim 10^{-7}$	$\lesssim 10^{-8}$	$\lesssim 10^{-9}$
t ightarrow cg	$\sim 10^{-11}$	$\lesssim 10^{-5}$	$\lesssim 10^{-7}$	$\lesssim 10^{-10}$
$t \rightarrow ch$	$\sim 10^{-15}$	$\lesssim 10^{-2}$	$\lesssim 10^{-5}$	$\lesssim 10^{-4}$

Atwood-Reina-Soni 1996, Cao et al 2009, Agashe-Contino 2009 Azatov et al 2009, Casagrande et al 2010

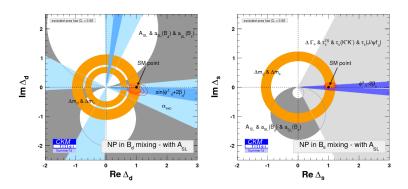
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NP contributions to $B-\overline{B}$ mixing



- $\Gamma_{12d/s}(NP) = 0$ assumed (not true in general)

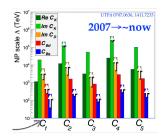


Effective operators for meson mixing



In general:

$$\mathcal{H}_{ ext{eff}} = \sum_{i=1}^{5} c_i(\mu) \mathcal{O}_i(\mu)$$



SM:

$$\begin{split} \mathcal{O}_1 &= (\bar{b}^\alpha \gamma_\mu L q^\alpha) \; (\bar{b}^\beta \gamma_\mu L q^\beta) \\ \mathcal{O}_2 &= (\bar{b}^\alpha L q^\alpha) \; (\bar{b}^\beta L q^\beta) \\ \mathcal{O}_3 &= (\bar{b}^\alpha L q^\beta) \; (\bar{b}^\beta L q^\alpha) \\ \end{split}$$

$$\mathcal{O}_5 &= (\bar{b}^\alpha L q^\beta) \; (\bar{b}^\beta R q^\alpha) \\ \mathcal{O}_7 &= (\bar{b}^\alpha L q^\beta) \; (\bar{b}^\beta R q^\alpha) \\ \mathcal{O}_8 &= (\bar{b}^\alpha L q^\beta) \; (\bar{b}^\beta R q^\alpha) \\ \end{split}$$

BSM:

$$\mathcal{O}_4 = (\bar{b}^{\alpha} L q^{\alpha}) (\bar{b}^{\beta} R q^{\beta})$$
$$\mathcal{O}_5 = (\bar{b}^{\alpha} L q^{\beta}) (\bar{b}^{\beta} R q^{\alpha})$$

 Meson mixing measurements give limits on $c_i^{\text{NP}} \equiv 1/\Lambda^2$

Lorentz structure of NP operators in $b \to s \mu \mu$

$$\mathcal{H}_{ ext{eff}}(b o s\mu^+\mu^-) = \mathcal{H}_{ ext{eff}}^{ extit{SM}} + \mathcal{H}_{ ext{eff}}^{ extit{VA}} + \mathcal{H}_{ ext{eff}}^{ extit{SP}} + \mathcal{H}_{ ext{eff}}^{ extit{T}} \; ,$$

$$\mathcal{H}_{\mathrm{eff}}^{SM} = -\frac{4G_{F}}{\sqrt{2}} V_{ls}^{*} V_{lb} \left\{ \sum_{i=1}^{5} C_{i}(\mu) \mathcal{O}_{i}(\mu) + C_{7} \frac{e}{16\pi^{2}} \left(\bar{s} \sigma_{\mu\nu} (m_{s} P_{L} + m_{b} P_{R}) b \right) F^{\mu\nu} \right.$$

$$+ C_{9} \frac{\alpha_{em}}{4\pi} \left(\bar{s} \gamma_{\mu} P_{L} b \right) \bar{\mu} \gamma_{\mu} \mu + C_{10} \frac{\alpha_{em}}{4\pi} \left(\bar{s} \gamma_{\mu} P_{L} b \right) \bar{\mu} \gamma_{\mu} \gamma_{5} \mu \right\}$$

$$\mathcal{H}_{eff}^{VA} = \frac{\alpha G_{F}}{\sqrt{2}\pi} V_{lb}^{*} V_{ls} \left\{ R_{V} \bar{s} \gamma_{\mu} P_{L} b \bar{\mu} \gamma_{\mu} \mu + R_{A} \bar{s} \gamma_{\mu} P_{L} b \bar{\mu} \gamma_{\mu} \gamma_{5} \mu \right.$$

$$+ R_{V}^{'} \bar{s} \gamma_{\mu} P_{R} b \bar{\mu} \gamma_{\mu} \mu + R_{A}^{'} \bar{s} \gamma_{\mu} P_{R} b \bar{\mu} \gamma_{\mu} \gamma_{5} \mu \right\} ,$$

$$\mathcal{H}_{eff}^{SP} = \frac{\alpha G_{F}}{\sqrt{2}\pi} V_{lb}^{*} V_{ls} \left\{ R_{S} \bar{s} P_{R} b \bar{\mu} \mu + R_{P} \bar{s} P_{R} b \bar{\mu} \gamma_{5} \mu \right.$$

$$+ R_{S}^{'} \bar{s} P_{L} b \bar{\mu} \mu + R_{P}^{'} \bar{s} P_{L} b \bar{\mu} \gamma_{5} \mu \right\} ,$$

$$\mathcal{H}_{eff}^{T} = \frac{\alpha G_{F}}{\sqrt{2}\pi} V_{lb}^{*} V_{ls} \left\{ C_{T} \bar{s} \sigma_{\mu\nu} b \bar{\mu} \sigma^{\mu\nu} \mu + i C_{TE} \bar{s} \sigma_{\mu\nu} b \bar{\mu} \sigma_{\alpha\beta} \mu \epsilon^{\mu\nu\alpha\beta} \right\} .$$

Characteristics of NP Lorentz structures

New SP operators: affect $B(B_s \to \mu^+ \mu^-)$

$$\begin{split} B(\bar{B}_s \to \mu^+ \, \mu^-) &= \frac{G_F^2 \alpha_{em}^2 m_{B_s}^5 f_{B_s}^2 \tau_{B_s}}{64 \pi^3} |V_{tb} V_{ts}^*|^2 \sqrt{1 - \frac{4 m_\mu^2}{m_{B_s}^2}} \times \\ &\left\{ \left(1 - \frac{4 m_\mu^2}{m_{B_s}^2} \right) \left| \frac{R_S - R_S'}{m_b + m_s} \right|^2 + \left| \frac{R_P - R_P'}{m_b + m_s} + \frac{2 m_\mu}{m_{B_s}^2} (C_{10} + R_A - R_A') \right|^2 \right\}. \end{split}$$

⇒ Strong bounds on Scalar and pseudoscalar operators

A typical B-decay rate calculation ($b \rightarrow s\mu\mu$)

The effective Hamiltonian: Operator Product Expansion

$$\mathcal{H}_{\text{eff}}^{SM} = -\frac{4G_{F}}{\sqrt{2}} V_{ts}^{*} V_{tb} \left\{ \sum_{i=1}^{6} C_{i}(\mu) \mathcal{O}_{i}(\mu) + C_{7} \frac{e}{16\pi^{2}} (\bar{s} \sigma_{\mu\nu} (m_{s} P_{L} + m_{b} P_{R}) b) F^{\mu\nu} + C_{9} \frac{\alpha_{em}}{4\pi} (\bar{s} \gamma_{\mu} P_{L} b) \bar{\mu} \gamma_{\mu} \mu + C_{10} \frac{\alpha_{em}}{4\pi} (\bar{s} \gamma_{\mu} P_{L} b) \bar{\mu} \gamma_{\mu} \gamma_{5} \mu \right\}$$

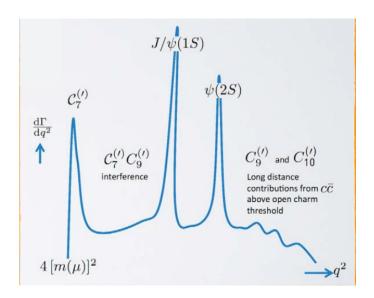
Decay rate:

$$\Gamma(B \to f) = [phase space)]|\langle f|H_{eff}^{SM}|B\rangle|^2$$

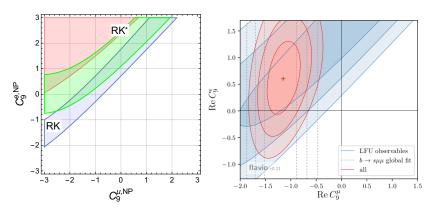
Quantities involved:

- Masses, Decay constants, Bag parameters,
- Wilson coefficients, Hadronic matrix elements (form factors),
- CKM elements

Wilson coefficients in $B o K^* \mu \mu$



Constraints on $C_9^{{ m NP},\mu}$ and $C_9^{{ m NP},e}$



Bhatia, AD, et al.

Altmannshofer et al.

Rare top decays

$$\mathcal{L}_{\mathsf{eff}} = \mathcal{L}_{\mathsf{SM}} + \sum_{i} rac{\mathcal{C}_{i}}{\Lambda^{2}} \mathcal{O}_{i}$$

Operators that are only weakly constrained by indirect probes

	t ightarrow c	$t \rightarrow u$
$SU(2)$ dipole \mathcal{O}_{LR}^W	$\Lambda \gtrsim 0.75 \text{ TeV}$	$\Lambda \gtrsim 0.75 \text{ TeV}$
SU(3) dipole \mathcal{O}_{RL}^{G}	$\Lambda \gtrsim 4.0 \text{ TeV}$	$\Lambda \gtrsim 5.8 \text{ TeV}$
Higgs penguin \mathcal{O}_{RL}^h	$\Lambda \gtrsim 0.73 \text{ TeV}$	$\Lambda \gtrsim 0.73 \text{ TeV}$

J. Brod

TeV scales are already being probed here.

Hints and constraints

- The hints: from precision measurements
 - CKM matrix elements
 - Mixing and decay in neutral mesons
 - Rare FCNC processes
- Quantifying constraints on NP
 - Model-independent constraints
 - Constraints on specific new physics models
- 3 Concluding remarks

Two-Higgs Doublet Model (Type-III) for R_D and R_{D^*}

$$\mathcal{L}_{\mathrm{Y}} = \overline{Q}_{fL}^{a} [Y_{fi}^{d} \epsilon_{ab} H_{d}^{b*} - \epsilon_{fi}^{d} H_{u}^{a}] d_{iR} - \overline{Q}_{fL}^{a} [Y_{fi}^{u} \epsilon_{ab} H_{u}^{b*} + \epsilon_{fi}^{u} H_{u}^{a}] u_{iR} + \mathrm{h.c.}$$

$$\begin{matrix} 1.0 \\ 0.5 \\ \hline E \\ 0.0 \\ 0.5 \end{matrix}$$

$$\begin{matrix} R(D) \\ 0.5 \\ \hline E \\ 0.0 \\ 0.5 \end{matrix}$$

$$\begin{matrix} R(D) \\ 0.5 \\ \hline E \\ 0.0 \end{matrix}$$

$$\begin{matrix} R(D) \\ 0.5 \\ \hline E \\ 0.0 \end{matrix}$$

$$\begin{matrix} R(D) \\ 0.5 \\ \hline E \\ 0.0 \end{matrix}$$

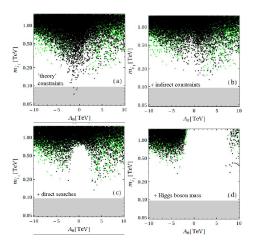
$$\begin{matrix} C_{S_{1}} \simeq \frac{1}{2\sqrt{2}G_{F}} \frac{m_{\tau}}{v} \epsilon_{33}^{d} \frac{\sin \beta \tan^{2} \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } BABAR! \end{matrix}$$

$$\begin{matrix} C_{S_{2}} \simeq \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{m_{\tau}}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{m_{\tau}}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{m_{\tau}}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{m_{\tau}}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{m_{\tau}}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{1}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{m_{\tau}}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{1}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{1}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{1}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}} \\ -\mathrm{disfavoured by } \frac{1}{2\sqrt{2}G_{F}V_{cb}} \frac{1}{v} \epsilon_{32}^{u*} \frac{\sin \beta \tan \beta}{M_{H^{\pm}}^{2}}$$

- R(D) and $R(D^*)$ may be accounted with a single parameter ϵ_{32}^u .
- Not possible for Type-I or Type-II models



Constraints on CMSSM with light stop $(m_{\tilde{i}} < 1.5 \text{ TeV})$

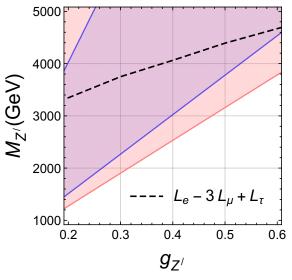


AD, Ghosh, Patel, Raychaudhuri 2013

- ullet Higgs mass constraints are dominant, at least at low tan eta
- At high $\tan \beta$, flavour constraints come into play again.



Constraints on specific Z' models



Bhatia, Chakraborty, AD, 2017

Hints and constraints

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- Concluding remarks

Concluding remarks

- B physics: a magnifying glass for testing SM Thanks to the lattice colleagues for the precision.
- Rare decays and precision measurements constrain specific NP models as well as indicate what classes of NP may be present
- While rare K decays have the largest precision reach, the large number of rare B decays offer probing power into the flavour structure.
- SUSY, multiple Higgses, Z', extra qyarks, leptoquarks, many possibilities
- Only data will tell, one has to look closely, though...

From beauty...

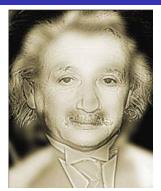


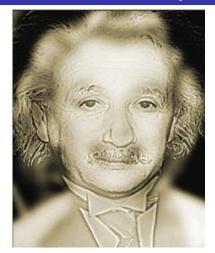
From beauty...

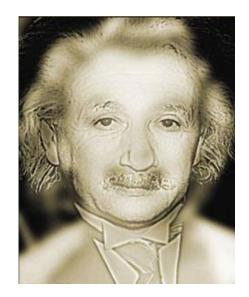


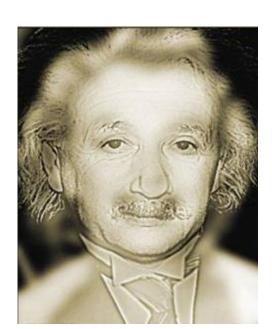












... to relativity!

So simply....



Follow the pugmarks.....