

Lattice QCD

Lattice QCD : Some Topics

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Basic Lattice Gauge Theory

Phase Diagram

Quark Number Susceptibility

Screening Lengths

Summary

Basic Lattice Gauge Theory

- Discrete space-time : Lattice spacing a UV Cut-off.
- Matter fields $\psi(x)$, $\bar{\psi}(x)$ on lattice sites.
- Gauge transformation : $\psi'(x) = V_x \psi(x)$, $V_x \in SU(3)$.
- Gauge Fields on links $U'_\mu(x) = V_x U_\mu(x) V_{x+\hat{\mu}}^{-1}$.
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 - Staggered Fermions (partial chiral and flavour symmetry),
 - Wilson fermions (only flavour symmetry),

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Typically, we need to evaluate

$$\langle \Theta(m_v) \rangle = \frac{\int DU \exp(-S_G) \Theta(m_v) \text{Det } M(m_s)}{\int DU \exp(-S_G) \text{Det } M(m_s)}, \quad (1)$$

where M is the Dirac matrix in x , color, spin, flavour space for fermions of mass m_s , S_G is the gluonic action, and the observable Θ may contain fermion propagators of mass m_v .

Since $\langle \Theta \rangle$ is computed by averaging over a set of configurations $\{U_\mu(x)\}$ which occur with probability $\propto \exp(-S_G) \cdot \text{Det } M$, the complexity of evaluation of $\text{Det } M \implies$ approximations : Quenched ($m_s = \infty$ limit), Partially Quenched (low $m_s = m_u = m_d$), and Full (including a heavier s quark).

Q $\rightarrow$ PQ $\rightarrow$ Full $\rightsquigarrow$ Computer time $\uparrow$ and Precision $\downarrow$.

Phase Diagram

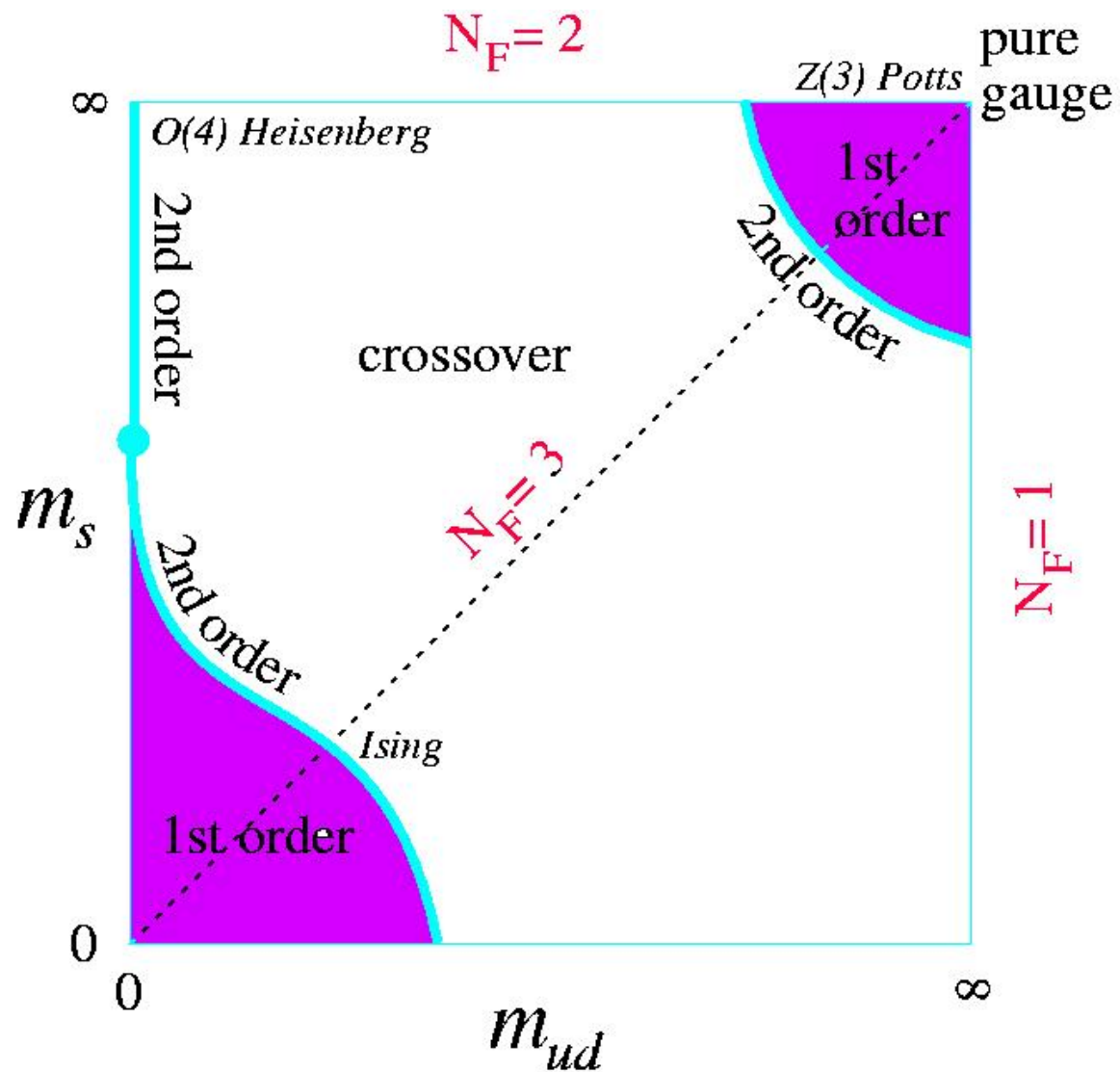
- Lattice details :
 - $N_s^3 \times N_t$ Lattice, $N_s \gg N_t$ for $T \neq 0$,
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- Theoretical expectations based on effective models :



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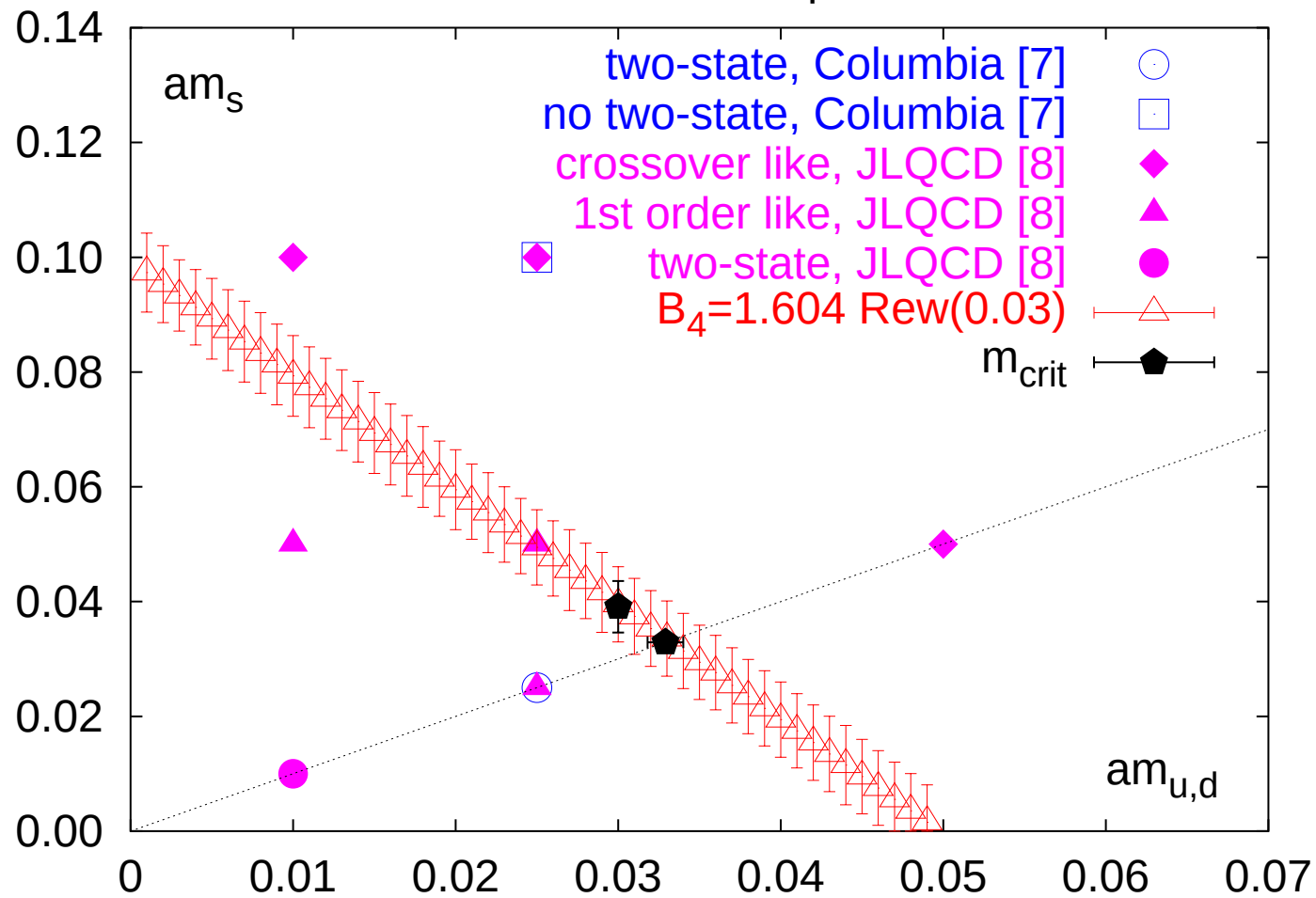
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- Theoretical expectations on the Phase diagram work out too.

Lattice'02 : Schmidt et al., hep-lat/0209009



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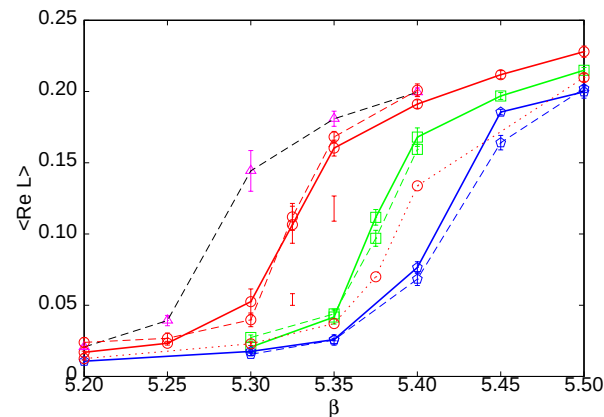
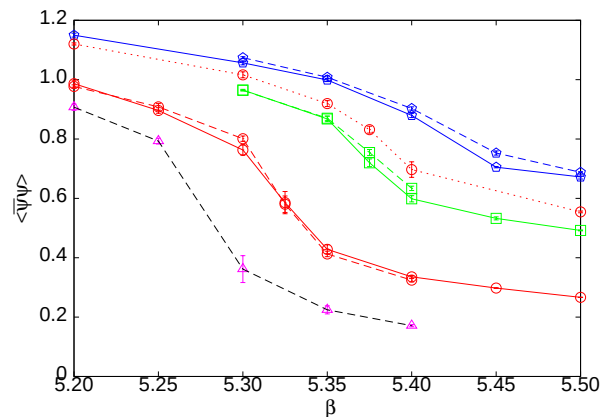
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$8^3 \times 4$ Lattice; $m_d/m_u = 1$ (dashed), 2(full) and 10 (dotted).



$$\mu_B \neq 0$$

- QCD at nonzero baryon density may $\longrightarrow$ Color Superconductivity, (Strange) Quark Stars, etc.
- Phase Problem : $\text{Det } M(\mu)$ is complex for $\mu \neq 0$.
- Exciting results in recent past for small μ .
 - Re-weighting Method (Fodor & Katz, JHEP '02)
 - Imaginary μ (de Forcrand & Philipsen, May '02, D'Elia & Lombardo, May '02)
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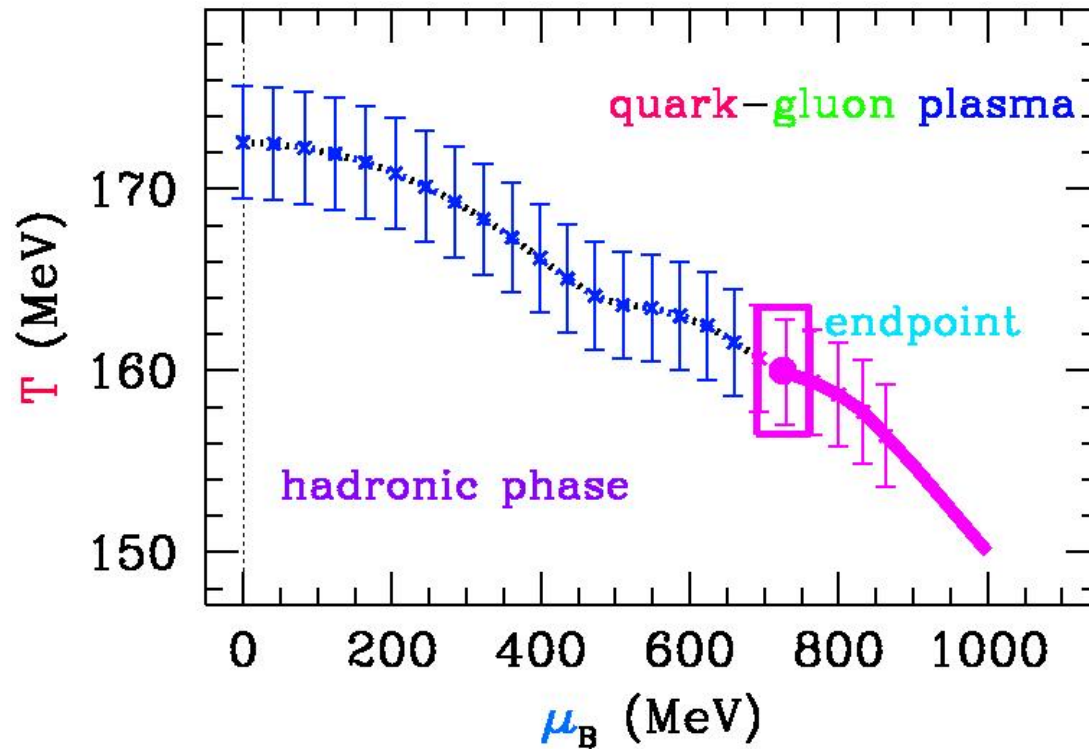
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Fodor-Katz Results



$N_s^3 \times 4$ Lattices,
 $N_s = 4, 6, 8$;
 Bit heavy u,d quarks.
 Critical Endpoint :
 $T = 160(4)$ MeV,
 $\mu = 725(35)$ MeV

As $m_{ud} \downarrow$, does $\mu_E \downarrow$? Larger N_t ??

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$$n_i = \frac{T}{V} \frac{\partial \ln \mathcal{Z}}{\partial \mu_i}, \quad \chi_{ij} = \frac{T}{V} \frac{\partial^2 \ln \mathcal{Z}}{\partial \mu_i \partial \mu_j}$$

Setting $\mu_i = 0$, $n_i = 0$ but χ_{ij} are nontrivial. Diagonal χ 's are

$$\chi_0 = \frac{1}{2}[\mathcal{O}_1(m_u) + \frac{1}{2}\mathcal{O}_2(m_u)] \quad (3)$$

$$\chi_3 = \frac{1}{2}\mathcal{O}_1(m_u) \quad (4)$$

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Here $\mathcal{O}_i$ trace of products of M^{-1} , M' and M'' and are estimated by a stochastic method:

$\text{Tr } A = \sum_{i=1}^{N_v} R_i^\dagger A R_i / 2N_v$, and $(\text{Tr } A)^2 = 2 \sum_{i>j=1}^L (\text{Tr } A)_i (\text{Tr } A)_j / L(L-1)$, where R_i is a complex vector from a set of N_v , subdivided in L independent sets.

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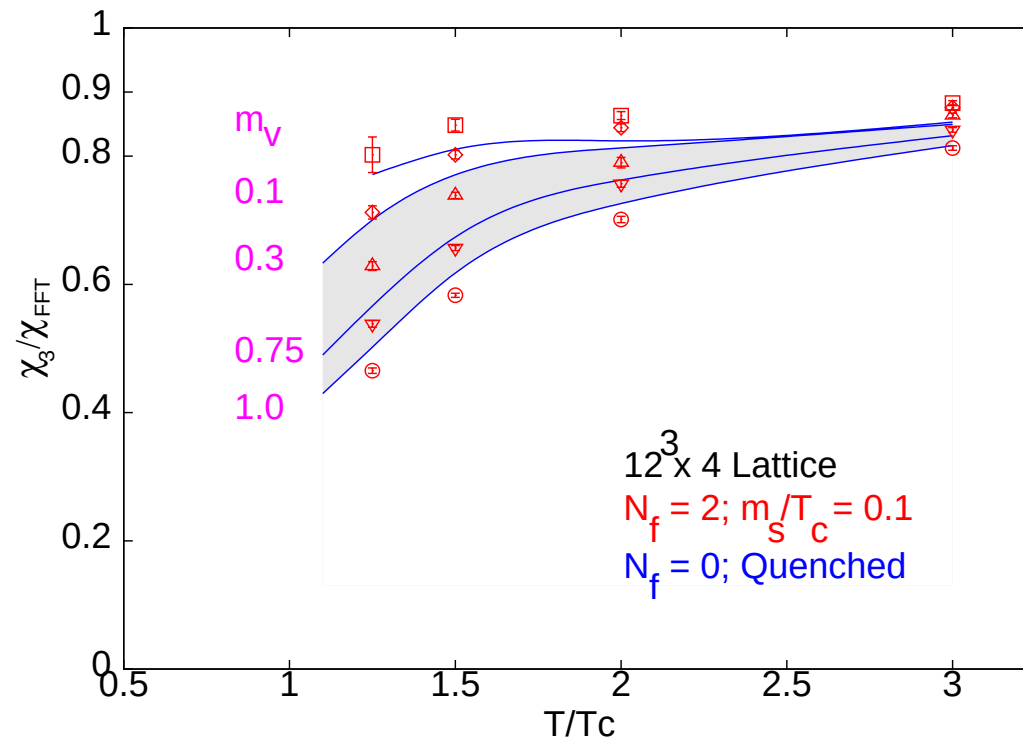
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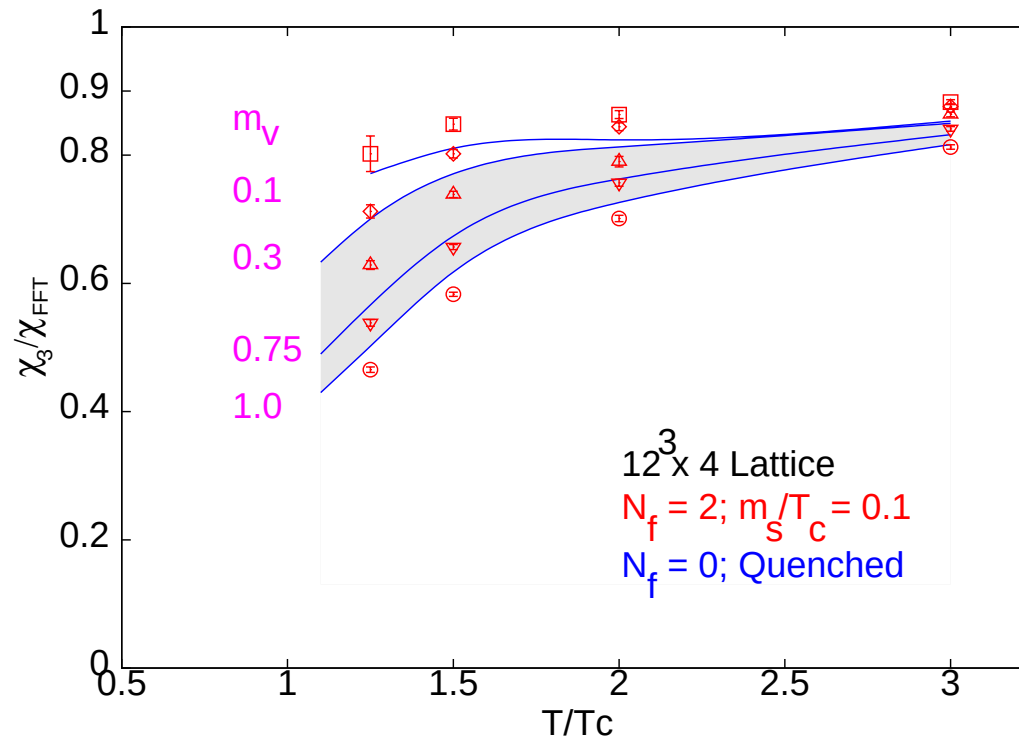
♠ Earlier results : Only close to T_c & for fixed ma .

Gavai & Gupta PR D '01; Gavai, Gupta & Majumdar, PR D 2002

χ_{FFT} — Ideal gas results for same Lattice.



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Note that PDG values for strange quark mass $\implies$

$$m_v^{strange}/T_c \simeq 0.3-0.7 \ (N_f=0); \ 0.45-1.0(N_f=2).$$

Perturbation Theory

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Weak coupling expansion gives:

$$\frac{\chi}{\chi_{FFT}} = 1 - 2\left(\frac{\alpha_s}{\pi}\right) + 8\sqrt{1 + 0.167N_f}\left(\frac{\alpha_s}{\pi}\right)^{\frac{3}{2}}$$

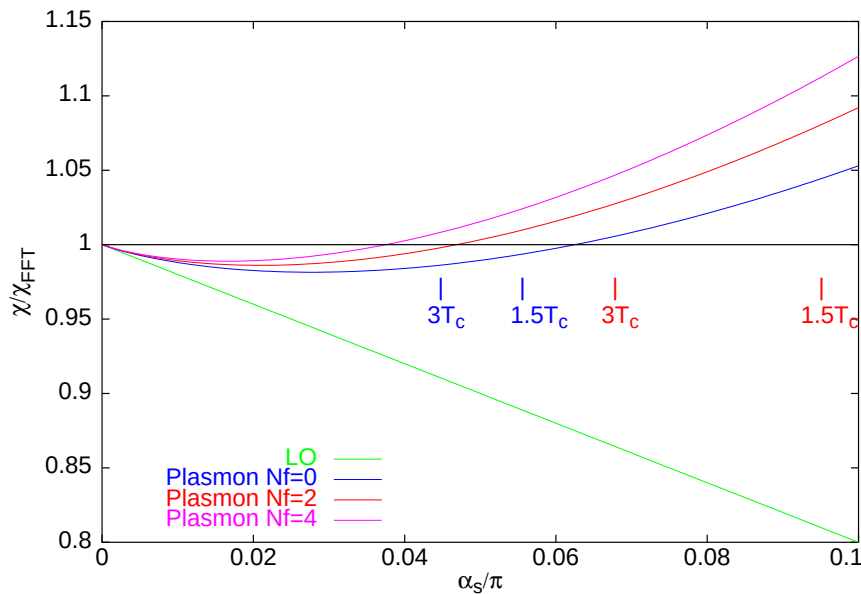
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- ♣ Minm 0.981 (0.986) at 0.03 (0.02) for $N_f = 0$ (2).
- ♣ For $1.5 \leq T/T_c \leq 3$ pert. theory $\longrightarrow$ 0.99-0.98 (1.08=1.03) for $N_f = 0$ (2).

Resummed Perturbation Theory

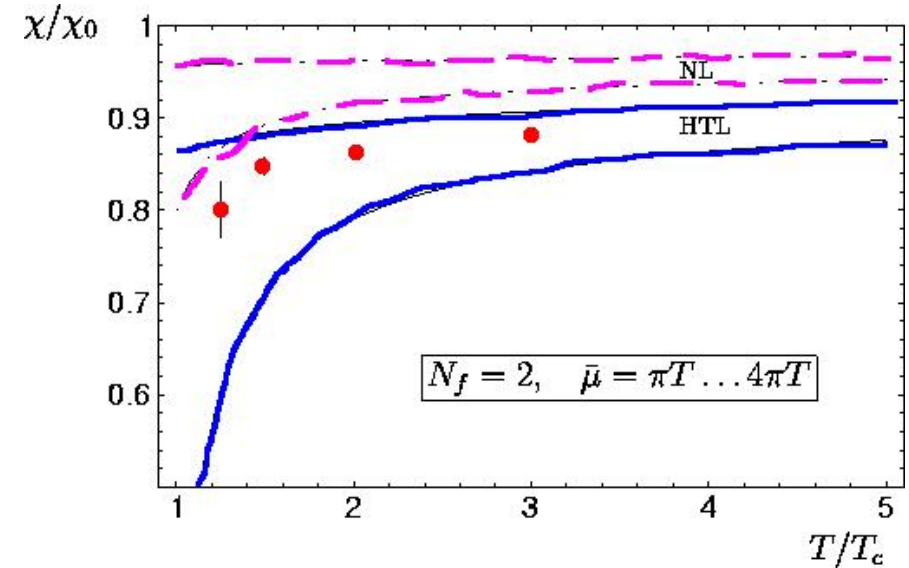
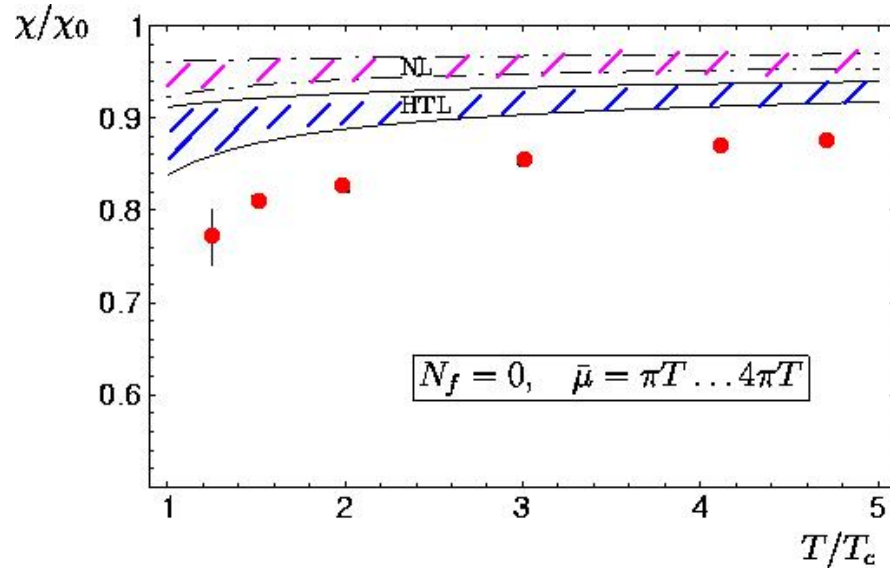
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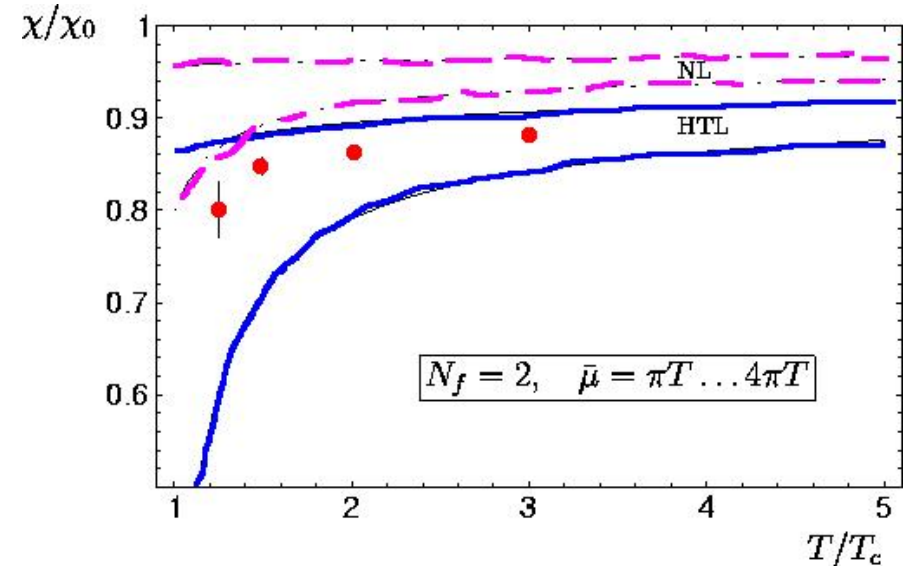
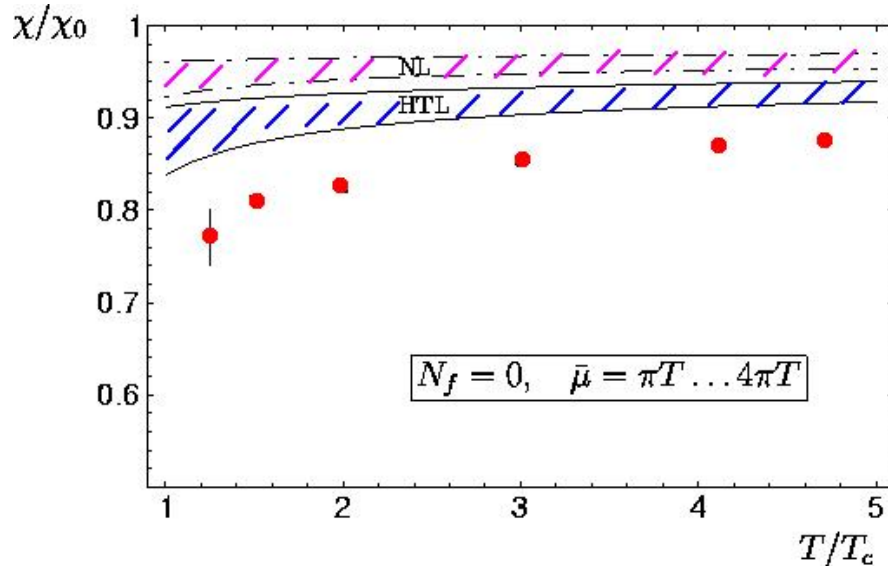
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Our results for $N_t = 4 \rightsquigarrow$ Lattice artifacts ?
Check for larger N_t and improved actions.

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♡ Identically zero for Ideal gas but $O(\alpha_s^3)$ in P.T.

Using the same scale and α_s as for $\chi_3 \longrightarrow \chi_{ud} \sim O(10^{-4})$!!

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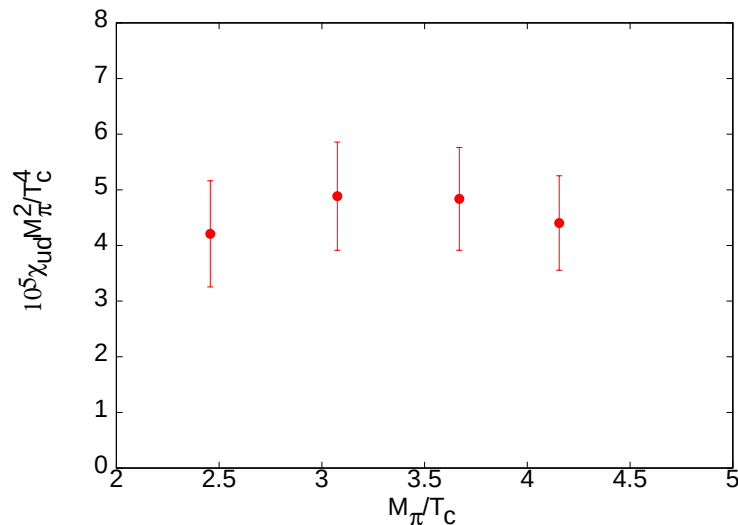
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♣ $12^3 \times 4$ Lattice; Quenched.
 ♣ $T = 0.75 T_c$
 ♣ Gavai, Gupta & Majumdar,
 PR D 2002

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♠ Naik action : Improved by $O(a)$ compared to Staggered.
Introduction of μ nontrivial but straightforward.

(Naik, NP B 1989; Gavai, hep-lat/0209008)

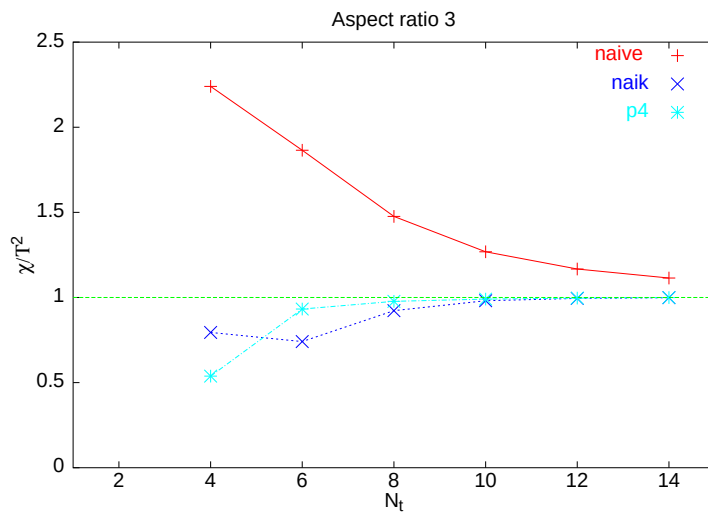
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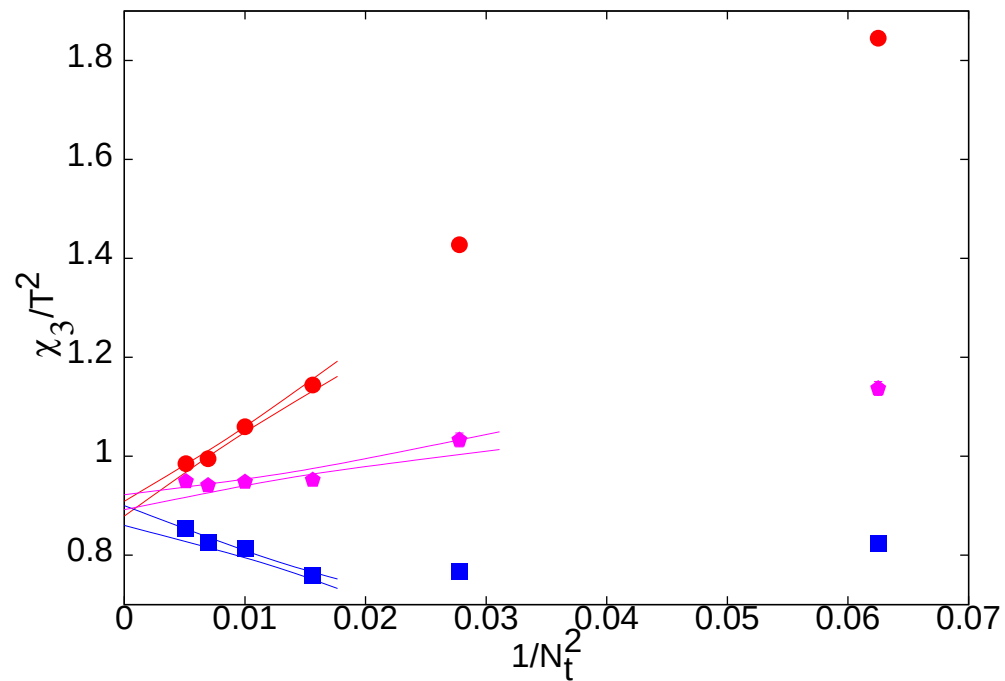
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Introduction of μ nontrivial but straightforward.

(Naik, NP B 1989; Gavai, hep-lat/0209008)

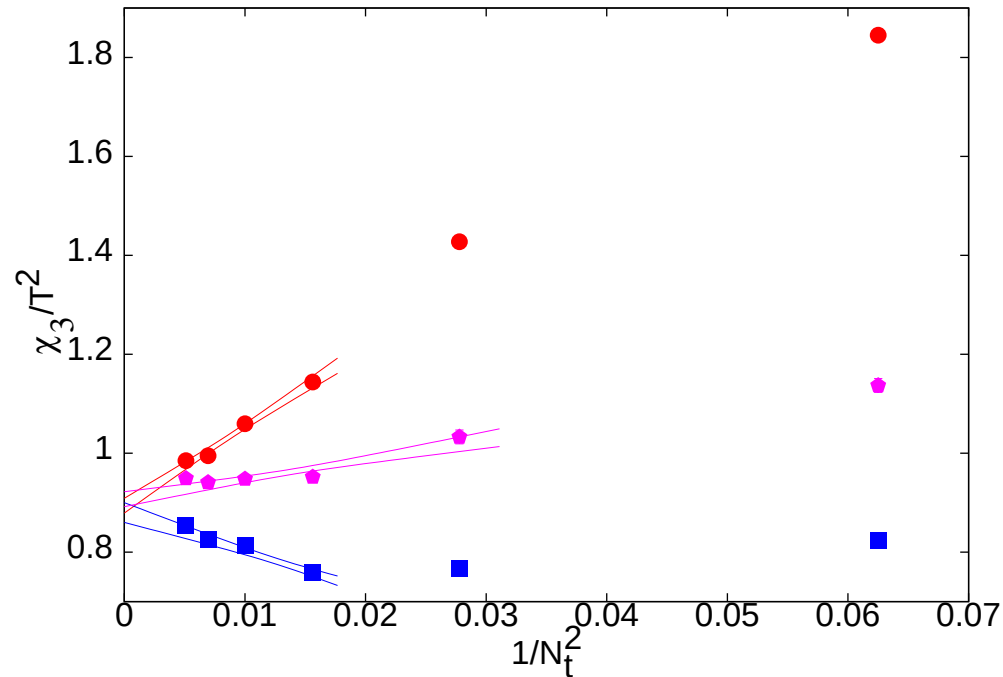


♠ Does improve the N_t -dependence of the free fermions.

Results at $2T_c$:

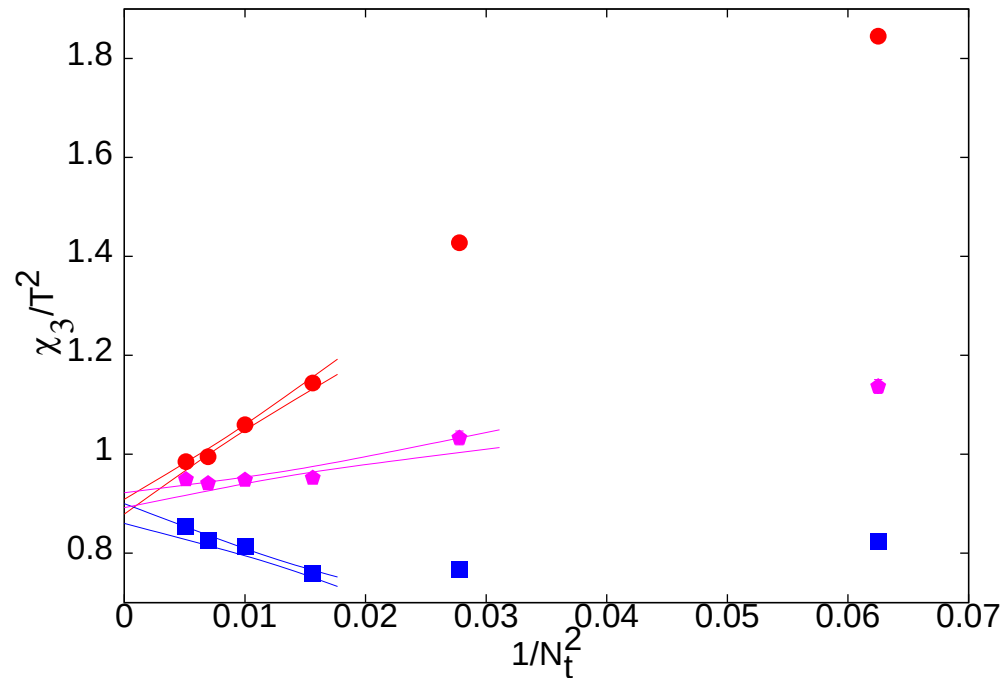


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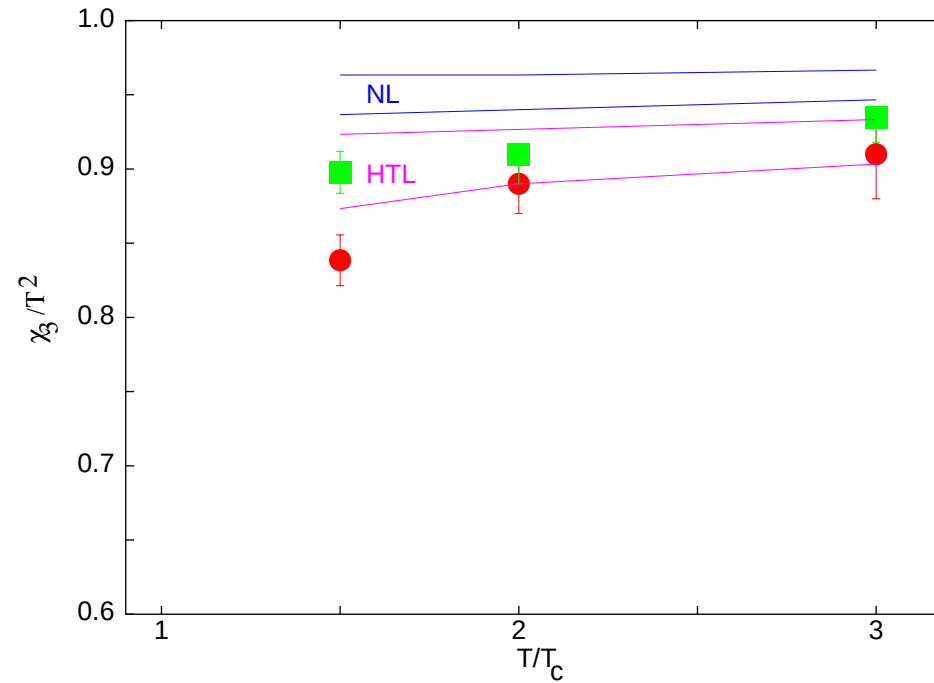


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◇ Milder $N_t^{-2} \sim a^2$ -dependence for Naik fermions.

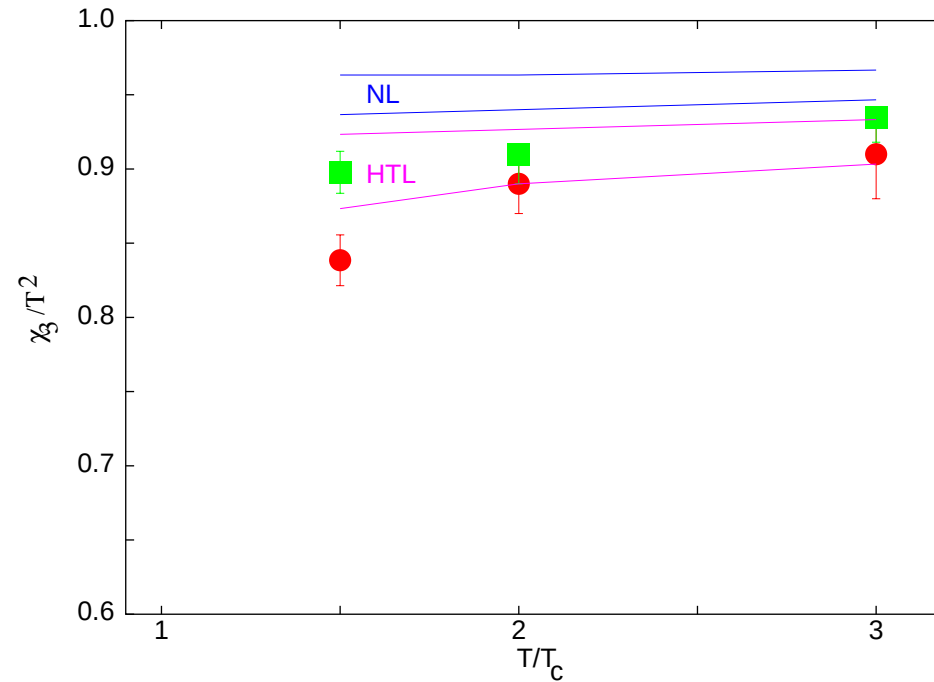
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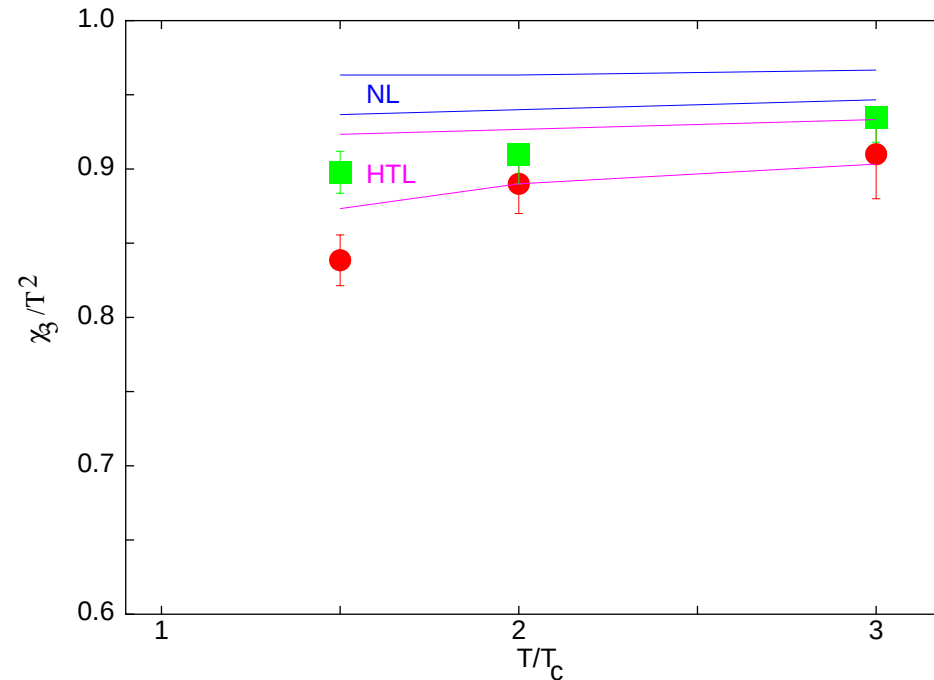
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♡ χ_{ud} behaves the same way for ALL N_t and both fermions, leading to the same $O(10^{-6})$ values in continuum too.

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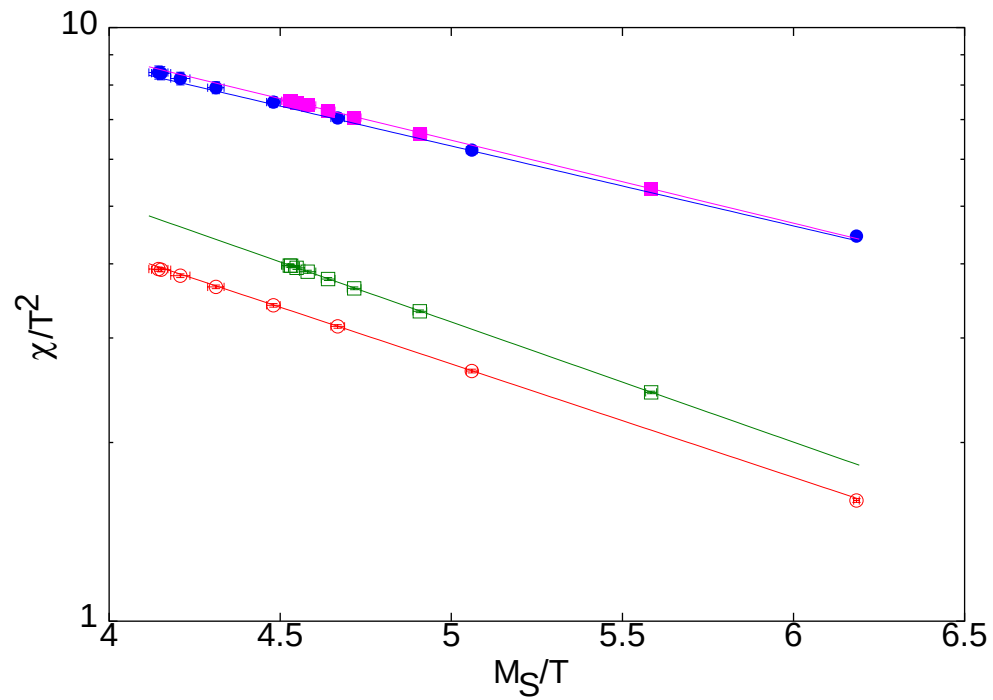
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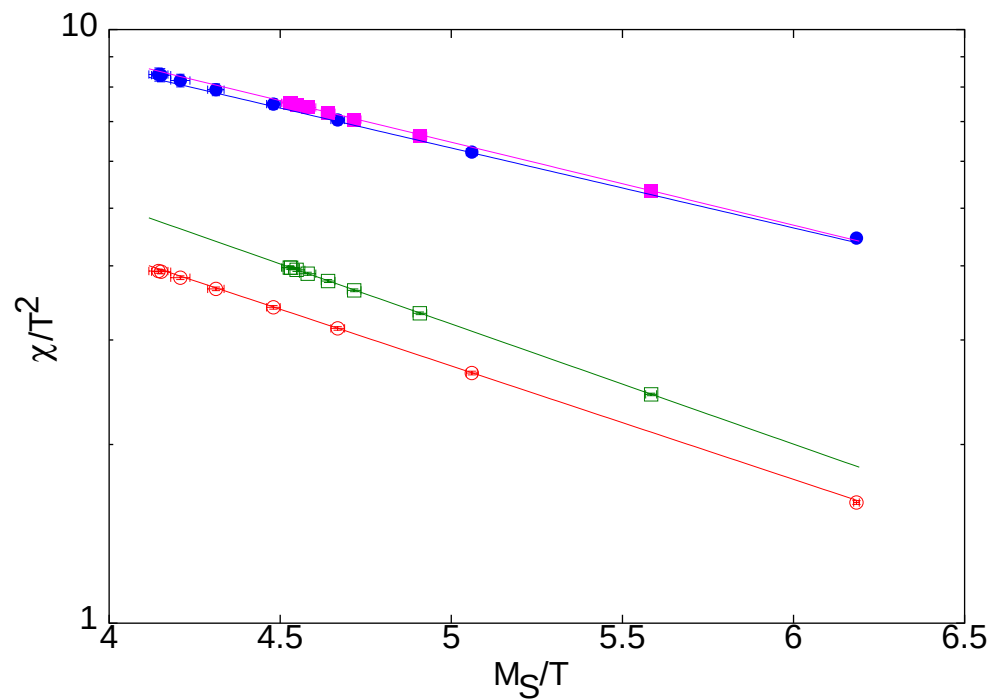
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- Summing up the C_{Γ} for pion $\rightarrow$ Pion susceptibility.

$N_t = 4$ Lattices with $N_z = 16$.
 $4\chi_3/T^2$ (open symbols) and $\chi_\pi/10T^2$ (filled)
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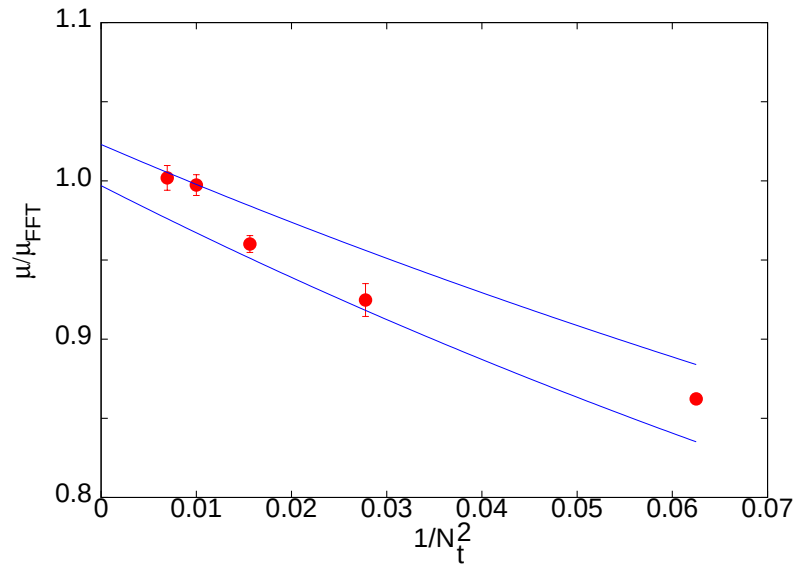


Why ? $\chi_3 \sim \sum$ propagator of nonlocal vector meson.

Again Taking Continuum Limit

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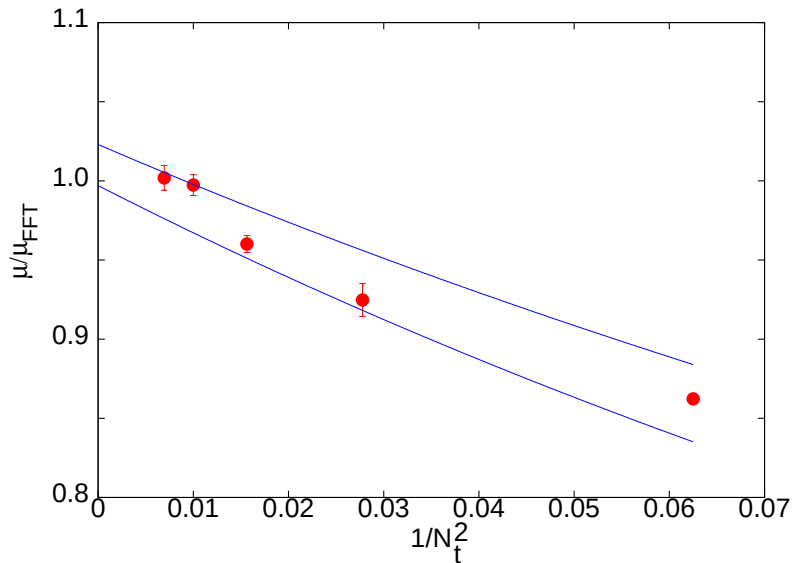
On finer lattices, $a = 1/8T - 1/12T$, Pion screening lengths become degenerate with those of ρ , i.e, also close to FFT!! (Gavai & Gupta, hep-lat/0211015)



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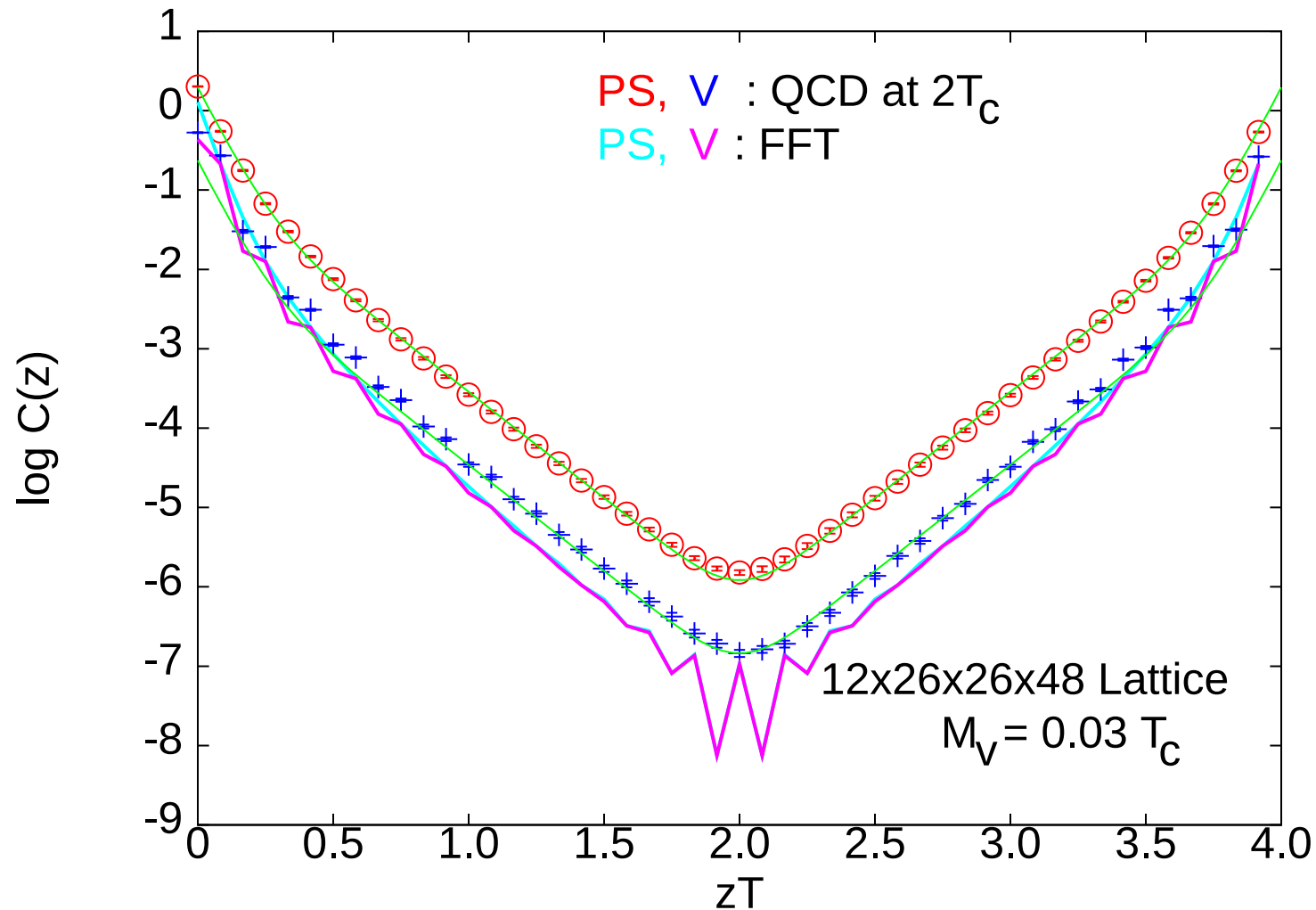
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However, chiral condensate, $\langle \bar{\psi}\psi \rangle$ differs from FFT by 2, as do the detailed shapes of the correlators.

Note that both PS and V have SAME fit (green line) with changed normalization.



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- Many questions still for full 2+1 QCD : Order, Large N_t , $\dots$.