

High energy neutrinos from astrophysical sources: a self-consistent approach

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in collaboration with M. Kachelrieß

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Outline

1 Motivation

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- 2 Astrophysical sources
 - Optically thick sources as neutrino sources

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 - Neutrino Yield
 - Unnormalized neutrino fluxes
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Motivation

Interest of high energy neutrinos (HE ν)

- penetrating power + no deflection → **neutrino astronomy**
 - **interior of stars** [G. G. Raffelt, 1996, R. C. Schirato and G. M. Fuller, 2002]
 - **distant cosmological epochs** → relic neutrino absorption spectroscopy
[B. Eberle, A. Ringwald, L. Song and T. Weiler, 2004]
 - **origin of UHECRs** → characterize astrophysical sources
[L. A. Anchordoqui, H. Goldberg, F. Halzen and T. J. Weiler, 2005, T. Kashti and E. Waxman, 2005,
M. Ahlers, L. A. Anchordoqui, H. Goldberg, F. Halzen, A. Ringwald, and Th. J. Weiler, 2005]

Motivation

Interest of high energy neutrinos (HE ν)

- penetrating power + no deflection $\rightarrow$ **neutrino astronomy**
- ν conversion from the source to the detector $\rightarrow$ **ν properties**
 - **mixing angles** [H. Athar, M. Jezabek and O. Yasuda, 2000, P. Bhattacharjee and N. Gupta, 2005, P. D. Serpico and M. Kachelrieß, 2005, M. C. González-García, F. Halzen and M. Maltoni, 2005, P. Serpico, 2005]
 - **CP phase** [P. D. Serpico and M. Kachelrieß, 2005, W. Rodejohann, 2007]
 - **ν decay** [J. F. Beacom et al, 2002, G. Barenboim and C. Quigg, 2003, D. Meloni and T. Ohlsson, 2006]
 - + terrestrial experiments $\rightarrow$ **mass hierarchy** [W. Winter, 2006]

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Main problem: low statistics

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Requirements

- **big ν detection volume** : ICECUBE, Auger, Anita, ...
 - **flavor determination**

[D. Fargion, 1997, S. I. Dutta, M. H. Reno and I. Sarcevic, 2000, J. F. Beacom, N. F. Bell, D. Hooper, S. Pakvasa and T. J. Weiler, 2003]

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- **large HE ν fluxes**
 - top-down models:
 - topological defects: **disfavoured** by EGRET and spectral shape of UHECRs
 - SHDM : **disfavoured** by the absence of γ at high energies and galactic anisotropy

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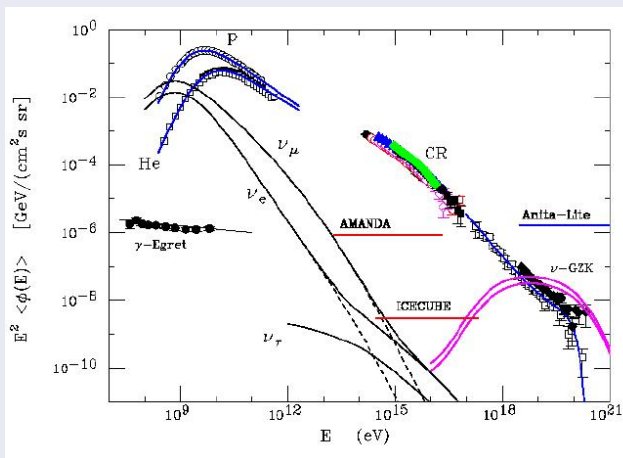
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 - top-down models:
 - topological defects: **disfavoured**
 - SHDM : **disfavoured**
 - bottom-up models $\rightarrow$ **astrophysical sources** $\rightarrow$ **optically thick**

Astrophysical sources

Limits



[P. Lipari, 2006]

Astrophysical sources

Limits

Waxman-Bahcall upper limit on high energy ν fluxes [E. Waxman and J. N. Bahcall, 1999]

$$\rightarrow E^2 \phi_\nu(E) \lesssim 10^{-8} \text{ GeV/cm}^2\text{s sr}$$

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 - thick ($\tau \gtrsim 1$) $\rightarrow$ multiple scattering suppresses CR flux
 - not high density $\rightarrow$ not to suppress π and K decays
- strongly energy dependent flavor ratio $\rightarrow$ a bump shows up

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Optically thick sources

- **Transparent sources:** accelerated protons **interact once** with the surrounding material $\rightarrow$ relationship $\nu \leftrightarrow \text{UHECR} \leftrightarrow \gamma$'s
optical depth $\tau \equiv L/l_{\text{int}} < 1$

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 - **cocooned black hole** [V. S. Berezinsky and V. L. Ginzburg, 1987]
 - **cores of AGNs** [F. Stecker, C. Done, M. Salamon and P. Sommers, 1991]
 - **collapsing galactic nucleus** [V. S. Berezinsky and V. I. Dokuchaev, 2001]

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Optically thick sources as neutrino sources

Our Monte Carlo simulation

source: phenomenological characterization

- injected **protons** $dN/dE \propto E^{-\alpha} \longrightarrow (\alpha = 2.2, E_{\max} = 10^{24} \text{ eV})$
- slab with $\left\{ \begin{array}{l} \text{--protons} \\ \text{--thermal } \textcolor{red}{\text{photons}} \end{array} \right. \longrightarrow \left\{ \begin{array}{l} \text{-- size } \textcolor{violet}{L} \\ \text{-- temperature } \textcolor{violet}{T} \\ \text{-- magnetic field negligible} \end{array} \right.$

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interactions: SOPHIA [A. Mücke *et al.*, 2000] and HERWIG [G. Corcella *et al.*, 2001]

$$p + \gamma \rightarrow \text{secondary mesons} \left\{ \begin{array}{l} \pi \rightarrow \mu + \nu_{\mu} \rightarrow e + \nu_e + \nu_{\mu} + \nu_{\mu} \\ K^{\pm} \rightarrow \mu + \nu_{\mu}, \pi^{\pm} + \pi^0, \dots \\ K_L^0 \rightarrow \pi + e + \nu_e, \pi + \mu + \nu_{\mu}, \dots \\ \text{charm mesons decay} \end{array} \right.$$



High Energy Neutrinos

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Neutrino Yield: $Y_\nu = \frac{J_\nu(E)}{\tau J_p(E)}$ [V. S. Berezinsky and V. V. Volynsky, 1979, V. S. Berezinsky and A. Z. Gazizov, 1993]

transparent source $\tau < 1$

$$Y_\nu = Y_\nu(x), \quad \text{dimensionless energy } x \equiv \frac{E\omega}{m_p^2}, \quad \omega = 1.6T$$

- **contribution:** mainly from **π decays**: $\pi \rightarrow \mu + \nu_\mu \rightarrow e + \nu_e + 2\nu_\mu$

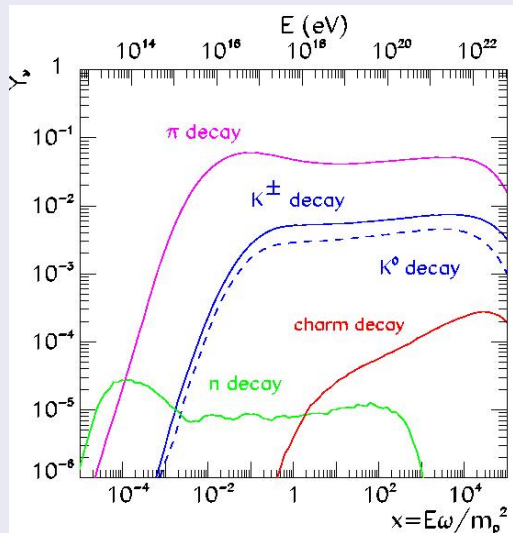
contributions

mainly

- π decay
- neutron decay

(at low energies)

$$T = 10^4 \text{ K}, \tau = 0.4$$

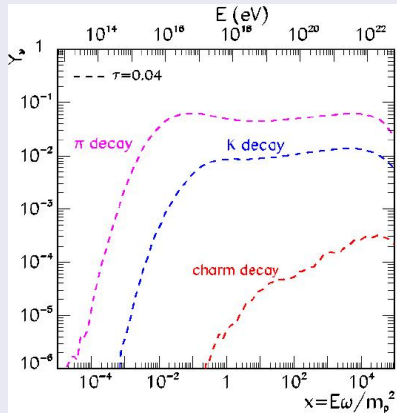
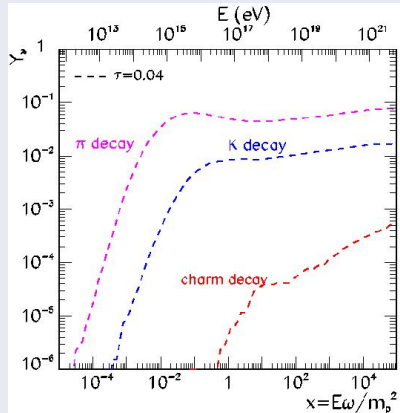


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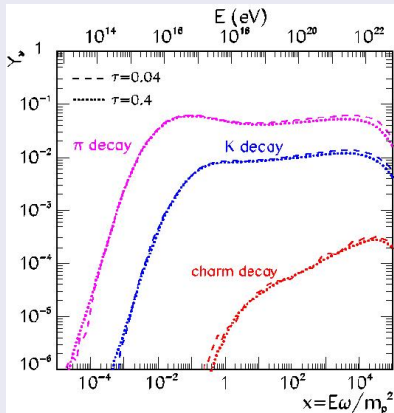
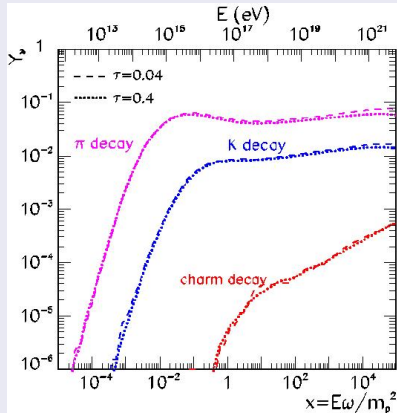
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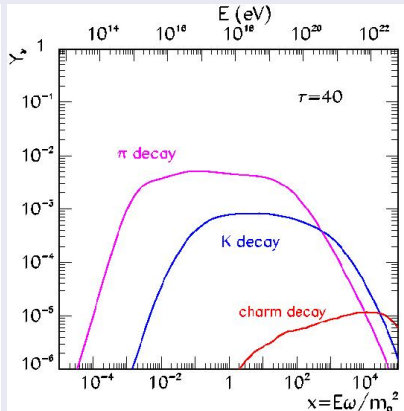
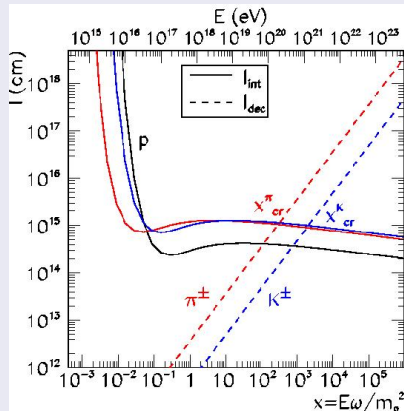
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optically thick sources $\tau \gtrsim 1$

multiple scattering important $\rightarrow$ **critical energy** $x_{\text{cr}}^{\pi,K}(T) : l_{\text{decay}} = l_{\text{int}}$

- **contribution** depends on the energy
 - $x < x_{\text{cr}}^\pi \rightarrow$ pion decay
 - $x_{\text{cr}}^\pi \lesssim x \lesssim x_{\text{cr}}^K \rightarrow K$ decay
 - $x_{\text{cr}}^K < x \rightarrow$ charm meson decay

$$T = 10^4 \text{ K}, \tau = 40$$



- $x < x_{cr}^{\pi} \rightarrow$ pion decay dominates
- $x_{cr}^{\pi} \lesssim x \lesssim x_{cr}^K \rightarrow \pi$'s scatter before decaying $\Rightarrow$ K decay important
- $x_{cr}^K < x \rightarrow$ K's scatter before decaying $\Rightarrow$ charm meson decay dominates

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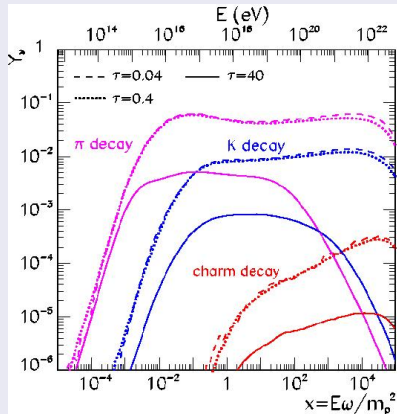
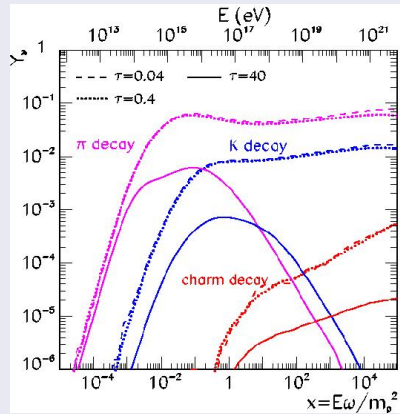
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- **Y_ν depends on T** : $x_{\text{cr}}^\pi \approx 2 \times 10^{10} / (T/\text{K})^2$

$T = 10^4 \text{ K}$

 $T = 10^5 \text{ K}$


Large HE ν flux if

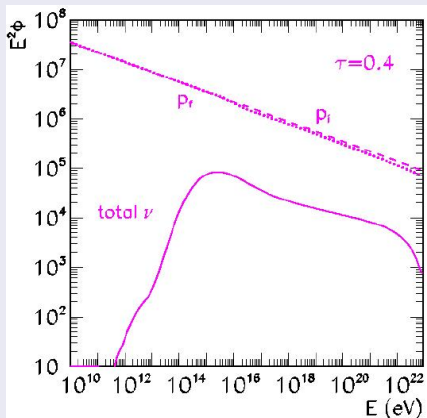
- **opaque sources** ($\tau \gtrsim 1 \Rightarrow L \gtrsim l_{\text{int}}^p$) to avoid the WB upper limit
- **not very high temperature** ($E_{\text{cr}} \propto 1/T^3$)

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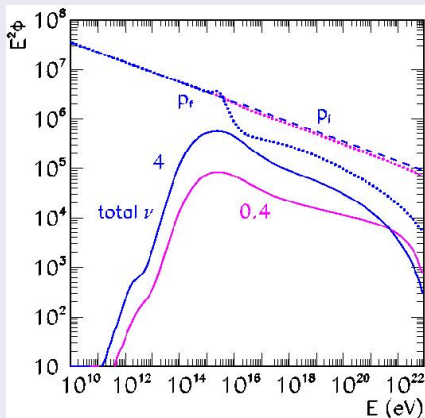
Unnormalized Neutrino Spectra: thickness

$$T = 10^5 \text{ K}$$



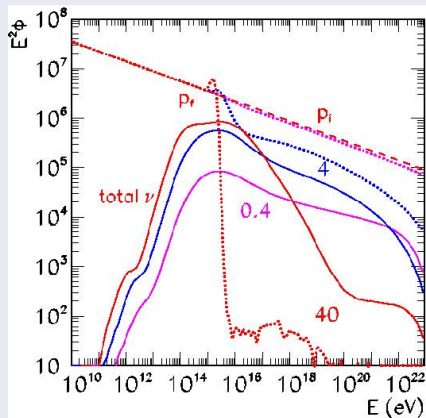
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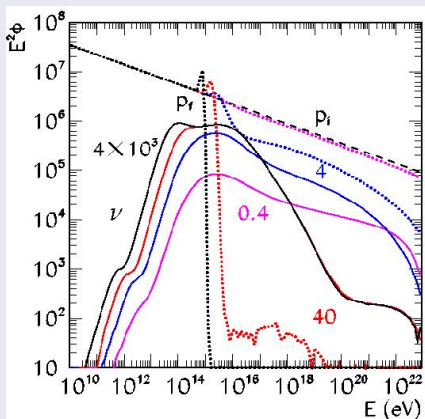
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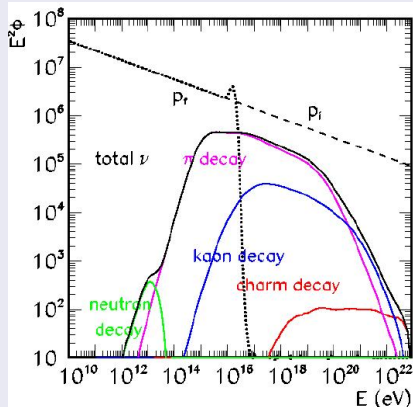
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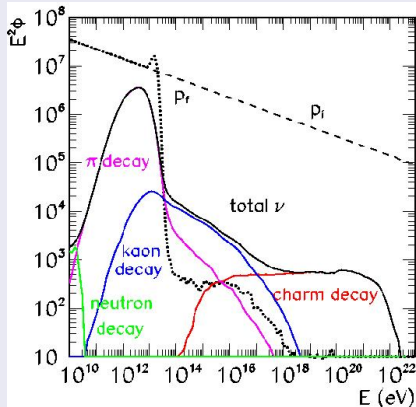
$\tau \gtrsim 1 \implies$ cosmic ray suppression

Unnormalized Neutrino Spectra: temperature

$T = 10^4$ K



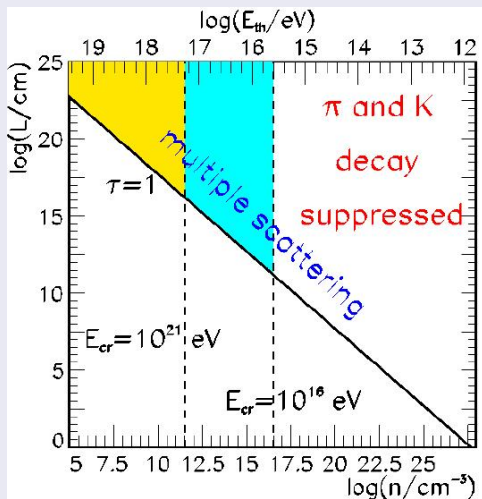
$T = 10^7$ K



high temperature $\Rightarrow$ high-energy ν flux suppression

Conditions for large HE_ν fluxes

opaque source ($\tau \gtrsim 1$) and not too dense



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Flavor content ($\phi_{\nu_e}, \phi_{\nu_\mu}, \phi_{\nu_\tau}$)

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- *measurable* in ν telescopes and extensive air shower experiments

[J. F. Beacom, N. F. Bell, D. Hooper, S. Pakvasa and T. J. Weiler, 2003, D. Fargion, 1997, S. I. Dutta, M. H. Reno and I. Sarcevic, 2000]

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- *interesting applications*

- *astrophysical diagnostics*: characterize sources

[L. A. Anchordoqui, H. Goldberg, F. Halzen and T. J. Weiler, 2005, T. Kashti and E. Waxman, 2005]

- *learn about ν properties* [H. Athar, M. Jezabek and O. Yasuda, 2000, P. Bhattacharjee and N. Gupta, 2005,

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Example

- pion decay $\rightarrow (1,2,0)$
- neutron beam source $\rightarrow (1,0,0)$ [L. A. Anchordoqui *et al.*, 2003, P. D. Serpico and M. Kachelrieß, 2005]
- muon-damped ν_μ sources from π decay $\rightarrow (0,1,0)$

[J. P. Rachen and P. Meszaros, 1998, T. Kashti and E. Waxman, 2005]

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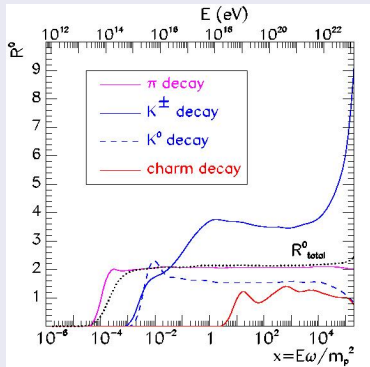
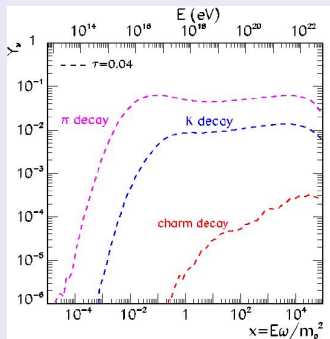
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- **optically thick sources** $\rightarrow (\phi_{\nu_e}, \phi_{\nu_\mu}, \phi_{\nu_\tau})$?

Flavor Content at the source

Flavor ratio: $R^0 = \phi_{\nu_\mu} / (\phi_{\nu_e} + \phi_{\nu_\tau})$

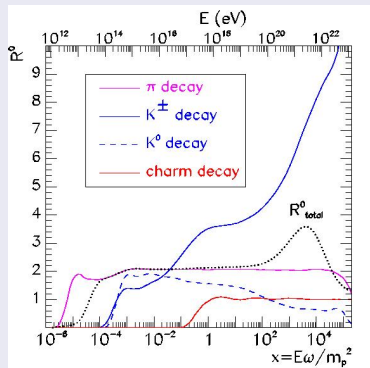
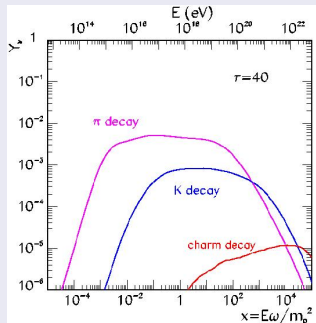
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Flavor Content at the source

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- optically **thick sources**: strong energy dependence
 - $x < x_{\text{cr}}^\pi \rightarrow R^0 = R_\pi \approx 2$
 - $x_{\text{cr}}^\pi \lesssim x \lesssim x_{\text{cr}}^K \rightarrow R^0 = R_K > 2 \Rightarrow$ bump
 - $x_{\text{cr}}^K < x \rightarrow R^0 = R_{\text{ch}} \approx 1$



Flavor Content at the source

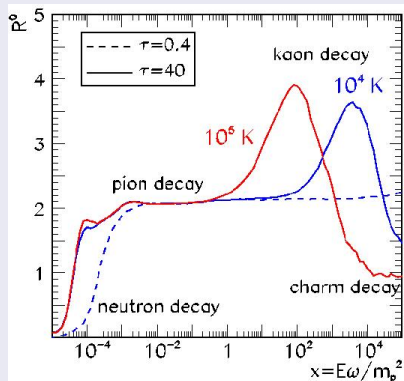
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- depends on the reaction $\rightarrow$ optically thin source: $R^0 = R_\pi \approx 2$
- optically **thick sources**: strong energy dependence

- $x < x_{\text{cr}}^\pi \rightarrow R^0 = R_\pi \approx 2$
- $x_{\text{cr}}^\pi \lesssim x \lesssim x_{\text{cr}}^K \rightarrow R^0 = R_K > 2$
- $x_{\text{cr}}^K < x \rightarrow R^0 = R_{\text{ch}} \approx 1$



bump at $x \approx x_{\text{cr}}^K$

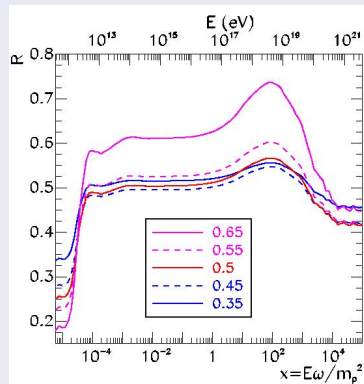


Flavor Content at the Earth

Flavor ratio: $R^0 + \nu$ oscillation $\rightarrow R = \phi_{\nu_\mu}^D / (\phi_{\nu_e}^D + \phi_{\nu_\tau}^D)$

$\phi_{\nu_\alpha}^D = \sum_\beta P_{\alpha\beta} \phi_{\nu_\beta} \Rightarrow R$ depends on the ν mixing parameters (θ_{ij}, δ_{CP})

mainly on θ_{23}



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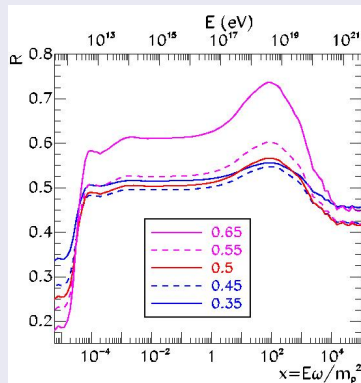
$\phi_{\nu_\alpha}^D = \sum_\beta P_{\alpha\beta} \phi_{\nu_\beta} \Rightarrow R$ depends on the ν mixing parameters (θ_{ij}, δ_{CP})

mainly on θ_{23}



presence of a **bump** $\rightarrow$ information on

- **source**: cross-over $\pi - K$ decay
- ν 's: sensitivity to the θ_{23} octant



Summary

High Energy Neutrinos

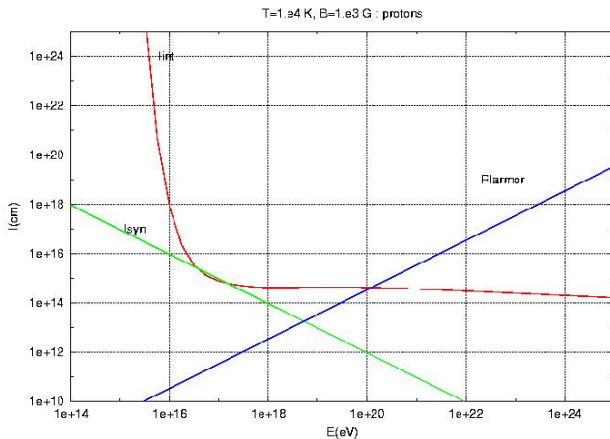
- penetrating power + no deflection → **ideal astronomical messengers**: UHECRs sources, ...
- complementary source of ν 's → **neutrino properties**: mixing, CP , ...



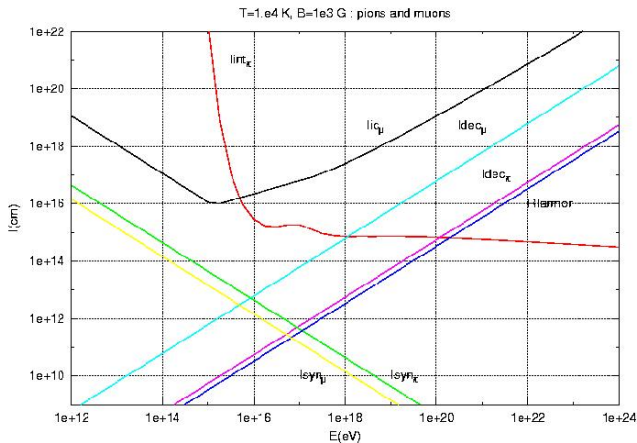
Optically Thick Sources

- Neutrino yield
 - multiple scattering $\Rightarrow$ **WB limit avoided**
 - if density is too high $\Rightarrow$ **HE ν 's from π and K decay suppressed**
- Flavor ratio depends strongly on the energy
 - **bump** shows up at the cross-over between π and K

Presence of magnetic fields (work in progress ...)



Presence of magnetic fields (work in progress ...)



Neutrino Spectra

Assumptions:

- **injected proton spectrum:**

$$dN_p/dE \propto E^{-\alpha} \begin{cases} \alpha = 2 & , E < 10^{-18} \text{ eV} \\ \alpha = 2.7 & , E > 10^{-18} \text{ eV} \end{cases}$$

[V. Berezhinsky, A. Gazizov and S. Grigorieva, 2004]

- **source:** homogeneous slab filled with thermal photons

$$\text{characterized by } \begin{cases} T = 10^4 \text{ K}, \dots, 10^7 \text{ K} \\ L = 10^{10} \text{ cm}, \dots, 10^{16} \text{ cm} \end{cases}$$

Diffuse neutrino spectra from uniform sources

$$F_\nu(E) = \frac{1}{4\pi} \int_0^\infty \frac{dz}{H(z)} \mathcal{L}_\nu(z, E_g)$$

- $\mathcal{L}_\nu(z, E_g)$: number of ν 's emitted per comoving volume, time and energy with generation energy $E_g = E(1+z)$

$$\mathcal{L}_\nu(z, E_g) = \eta(z) J_\nu(Eg)$$

- $J_\nu(Eg)$: neutrino spectra from a single source
- $\eta(z)$: number of sources per co-moving volume and per unit of time
- $H(z)$: Hubble parameter, $H(z)^2 = H_0^2[\Omega_m(1+z)^3 + \Omega_\Lambda]$

Constraints

- **UHECR spectrum** (AGASA) $\Rightarrow$ limits on the proton spectra emitted
 $\Rightarrow J_\nu(E) \rightarrow$ optically thin sources

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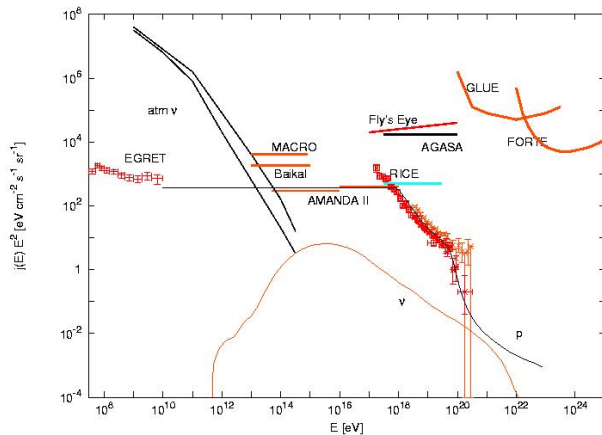
- **UHECR spectrum** (AGASA) $\Rightarrow$ limits on the proton spectra emitted
 $\Rightarrow J_\nu(E) \rightarrow$ optically thin sources
- **current limits on ν fluxes** (AMANDAII) $\rightarrow$ sources with $\tau \gtrsim 1$
- **EGRET observations** $\Rightarrow$ limits on the electromagnetic cascade produced:

$$\omega_{\text{cas}} = f_{\text{em}} E_{\text{tot}} \int_0^\infty \frac{dz}{H(z)} \eta(z) \leq \omega_{\text{obs}} = 2 \times 10^{-6} \text{ eV/cm}^3$$

[R. Aloisio, V. Berezhinsky and M. Kachelrieß, 2004]

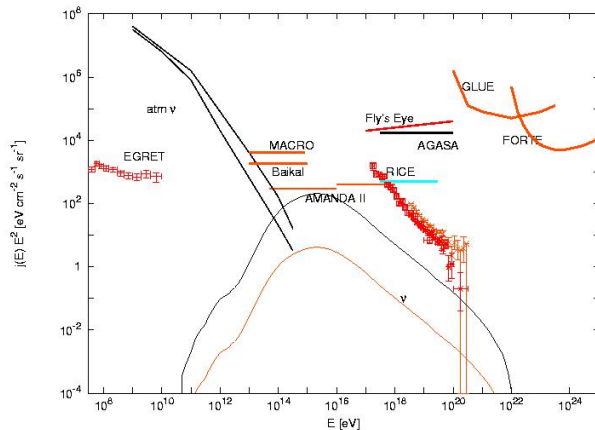
Constraints

$$T = 10^5 \text{ K and } \tau \ll 1$$



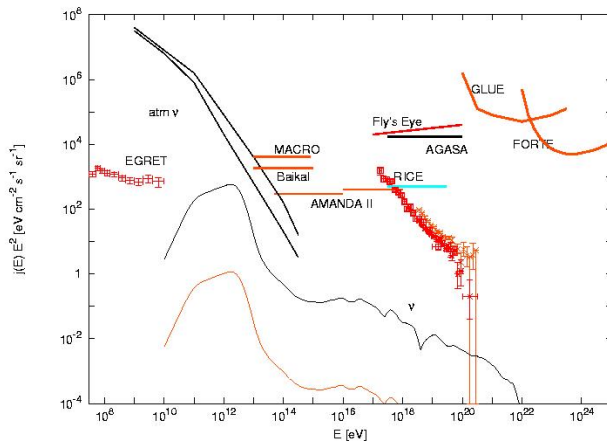
Constraints

$$T = 10^5 \text{ K and } \tau \gtrsim 1$$



Constraints

$$T = 10^7 \text{ K and } \tau \gg 1$$



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