

Supernova Shock Waves and Sterile Neutrinos

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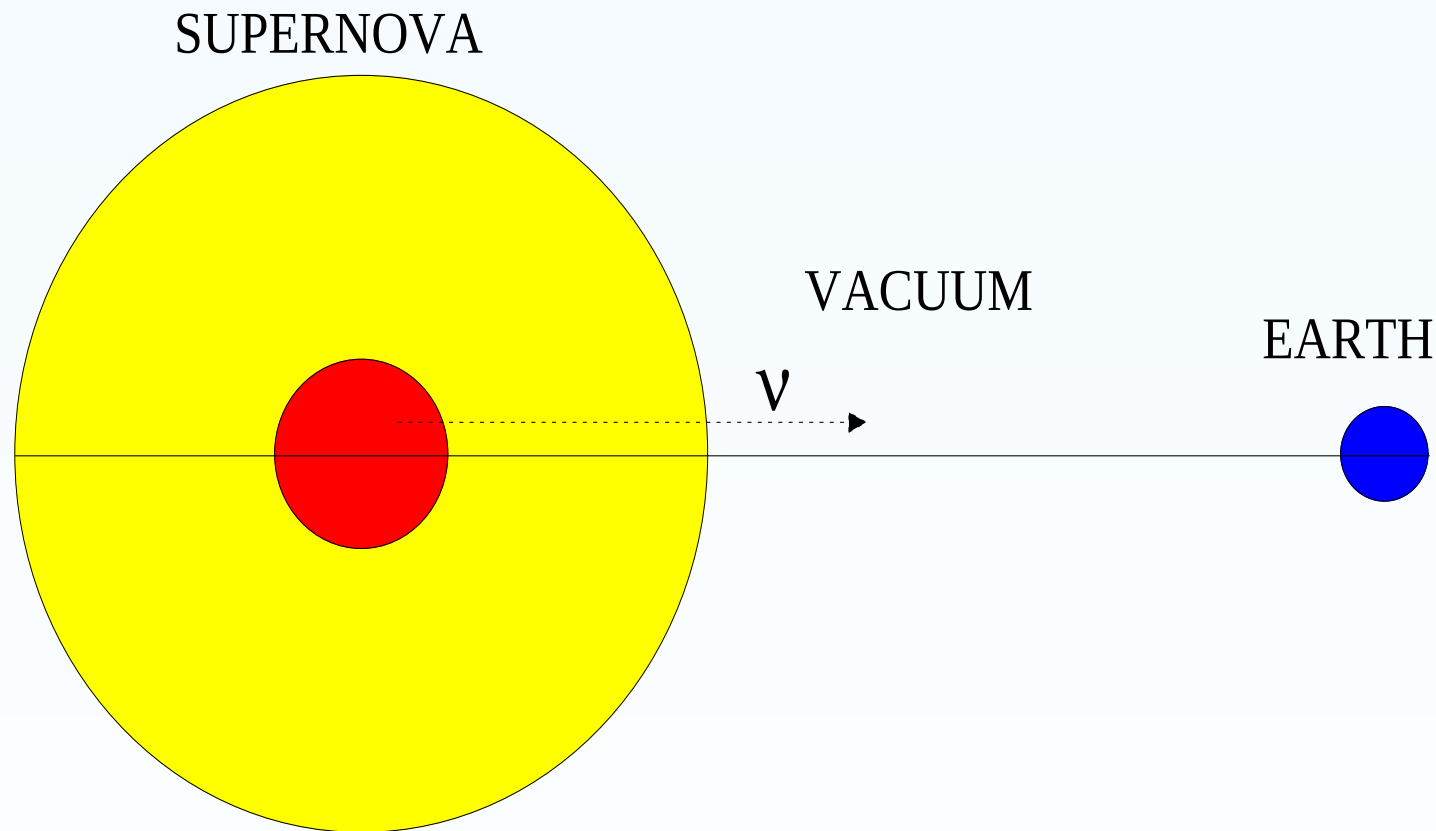
S.C., Nicholas Harries and Graham Ross

Phys. Rev. D74, 053010 (2006) and hep-ph/0702xxx

Flavor Oscillations of Supernova Neutrinos



Flavor Oscillations of Supernova Neutrinos



- ✓ Neutrino cross the mantle and envelope of the SN
- ✓ Neutrino **might** cross the earth
- ✓ Matter effects are **crucial**
- ✓ In this talk we will assume that they do not cross the earth

Neutrino Flavor Oscillations in Matter

✓ Evolution equation of neutrinos in matter is given by:

$$i \frac{d}{dt} \begin{pmatrix} \nu_\alpha \\ \nu_\beta \end{pmatrix} = \frac{1}{4E} \mathcal{M}_F^m \begin{pmatrix} \nu_\alpha \\ \nu_\beta \end{pmatrix}$$

✓ For oscillations between active neutrinos:

$$\mathcal{M}_F^m = \begin{pmatrix} -\Delta m^2 \cos 2\theta + 2\sqrt{2}G_F N_e E & \Delta m^2 \sin 2\theta \\ \Delta m^2 \sin 2\theta & \Delta m^2 \cos 2\theta - 2\sqrt{2}G_F N_e E \end{pmatrix}$$

✓ Mixing angle in matter changes to:

$$\tan 2\theta_m = \frac{\Delta m^2 \sin 2\theta}{\Delta m^2 \cos 2\theta - 2\sqrt{2}G_F N_e E}$$

Neutrino Flavor Oscillations in Matter

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$$\Delta m^2 \cos 2\theta = 2\sqrt{2}G_F N_e E \Rightarrow \text{Maximal Mixing}$$

→ MSW Matter Enhanced Resonance

L. Wolfenstein, Phys. Rev. D **17**, 1978

S.P.Mikhyev, A.Yu.Smirnov, SJNP **42** 1985

✓ Resonance condition possible *only* for $\Delta m^2 > 0$

Neutrino Flavor Oscillations in Matter

✓ For oscillations between active anti-neutrinos:

$$\mathcal{M}_F^m = \begin{pmatrix} -\Delta m^2 \cos 2\theta - 2\sqrt{2}G_F N_e E & \Delta m^2 \sin 2\theta \\ \Delta m^2 \sin 2\theta & \Delta m^2 \cos 2\theta + 2\sqrt{2}G_F N_e E \end{pmatrix}$$

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$$\Delta m^2 < 0$$

Neutrino Flavor Oscillations in Matter

✓ For oscillations between ν_μ and $\nu_{sterile}$:

$$\mathcal{M}_F^m = \begin{pmatrix} -\Delta m^2 \cos 2\theta - \sqrt{2}G_F N_n E & \Delta m^2 \sin 2\theta \\ \Delta m^2 \sin 2\theta & \Delta m^2 \cos 2\theta + \sqrt{2}G_F N_n E \end{pmatrix}$$

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Neutrino Flavor Oscillations in Matter

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Neutrino Flavor Oscillations in Matter

✓ For oscillations between $\bar{\nu}_e$ and $\bar{\nu}_{sterile}$

$$\mathcal{M}_F^m = \begin{pmatrix} -\Delta m^2 \cos 2\theta - 2\sqrt{2}G_F N_e E + \sqrt{2}G_F N_n E & \Delta m^2 \sin 2\theta \\ \Delta m^2 \sin 2\theta & \Delta m^2 \cos 2\theta + 2\sqrt{2}G_F N_e E - \sqrt{2}G_F N_n E \end{pmatrix}$$

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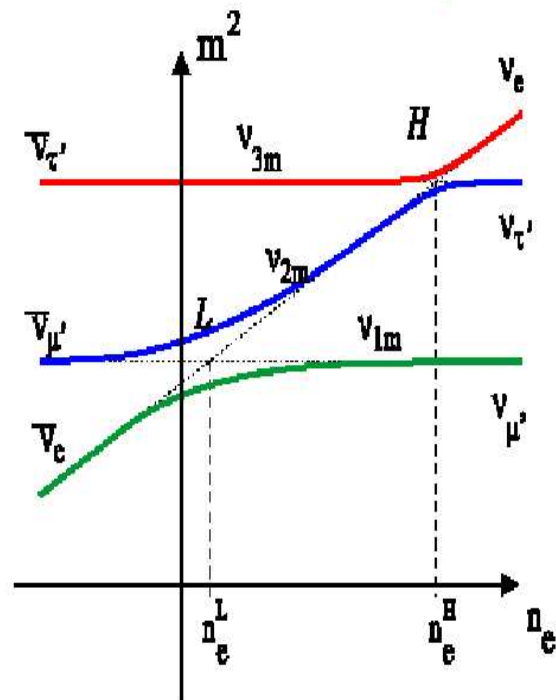
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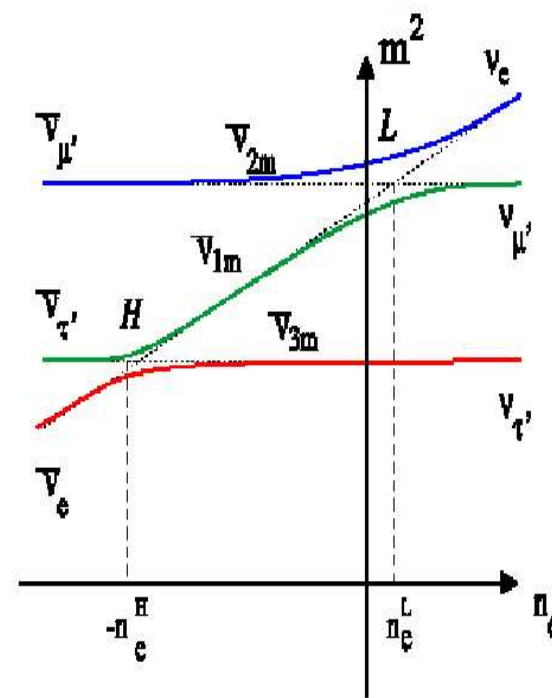
Flavor Oscillations of Supernova Neutrinos

✓ Three flavor scenario

Normal mass hierarchy



Inverted mass hierarchy



Dighe and Smirnov (2000)

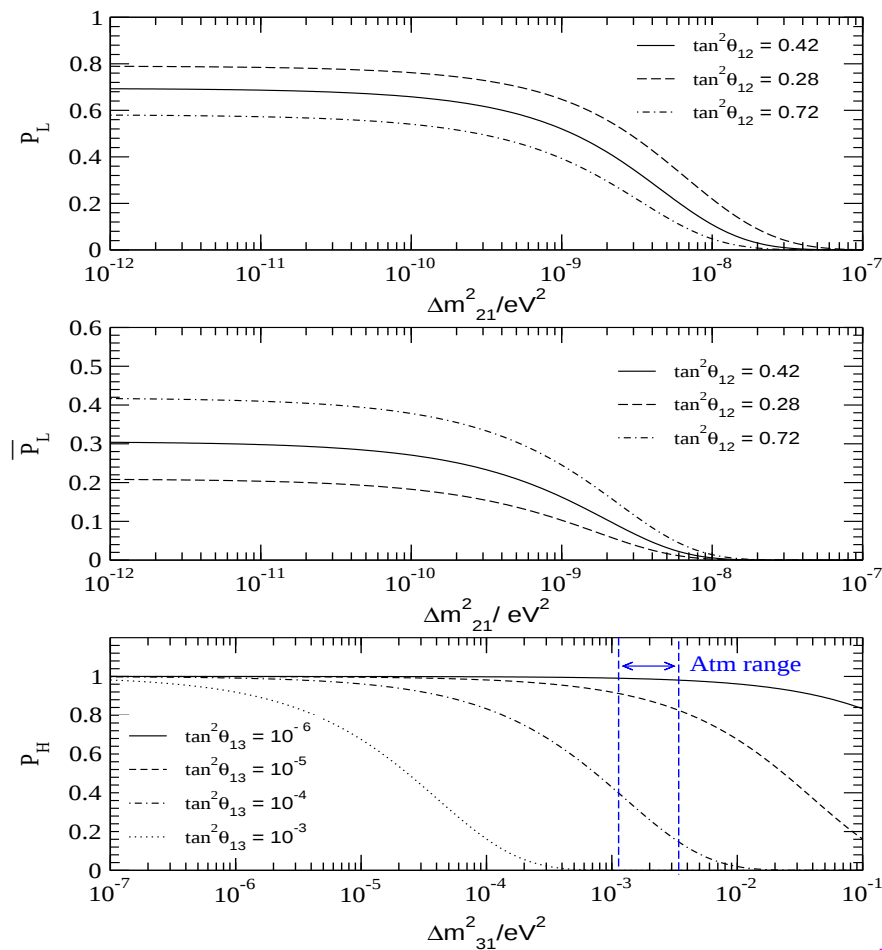
- ✓ L always lies in the **neutrino** channel
- ✓ H will be in the **(anti)neutrino** channel for $\Delta m_{31}^2 (<) > 0$
- ✓ Net flavor conversion depends on the “jump probability” P_J

Flavor Oscillations of Supernova Neutrinos

$$P_J \approx \exp(-\gamma \sin^2 \theta); \quad \gamma = \pi \frac{\Delta m^2}{E} \left| \frac{d \ln N_e}{dx} \right|_{x=x_{\text{res}}}^{-1}$$

Flavor Oscillations of Supernova Neutrinos

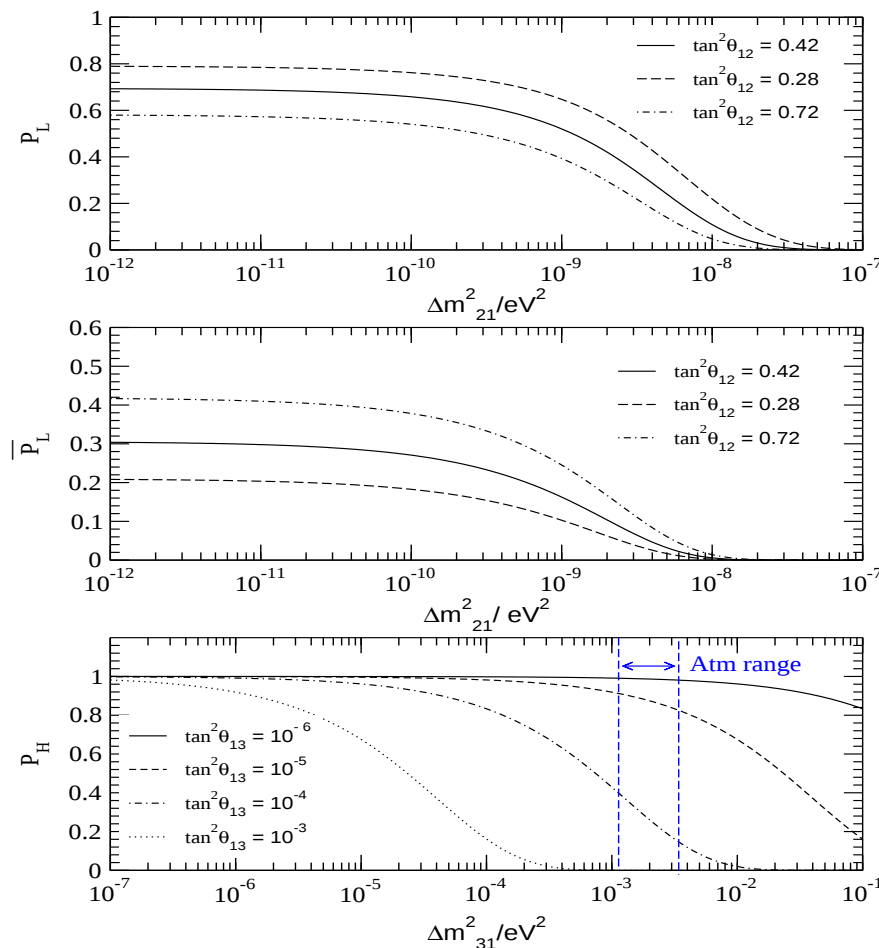
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Bandyopadhyay, S.C., Goswami, Kar, hep-ph/0312315

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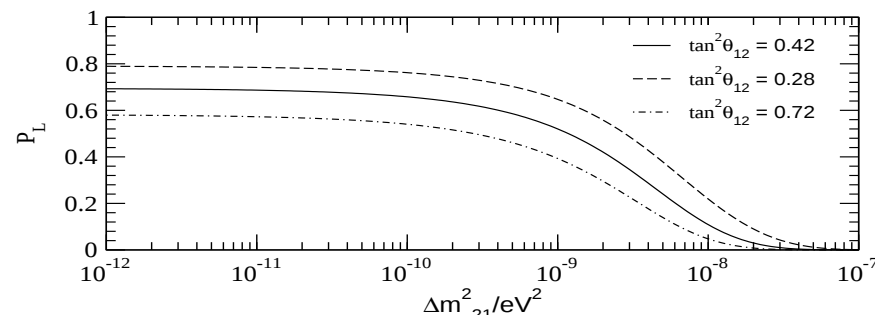


$\checkmark P_L = 0$ and $\bar{P}_L = 0$ **always**
 $\checkmark P_H$ depends on $\sin^2 \theta_{13}$

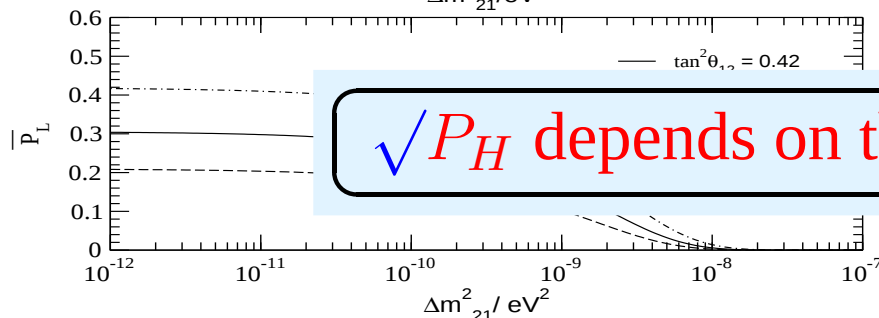
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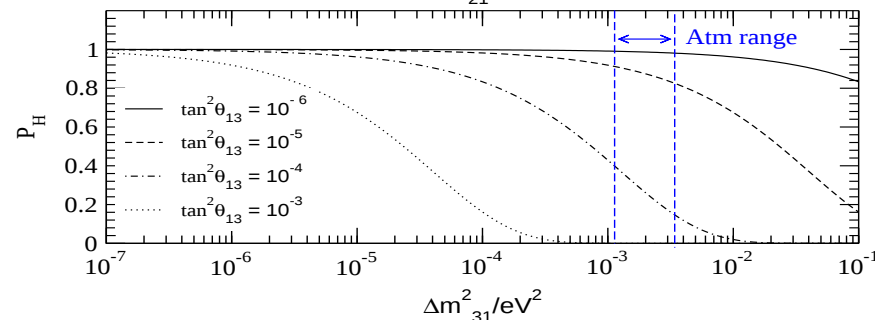
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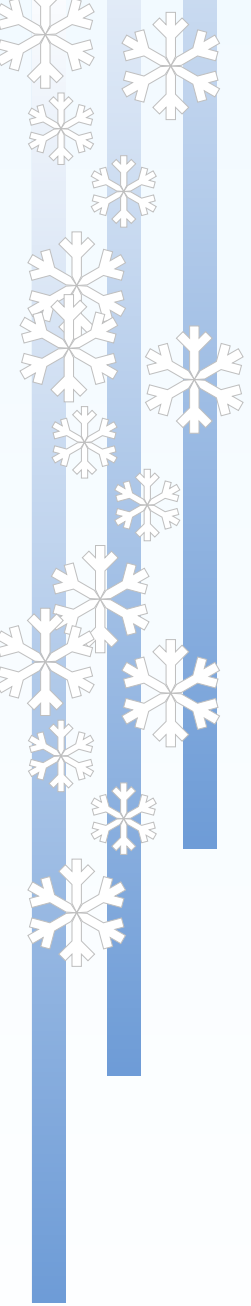


$\checkmark P_H$ depends on the SN density profile



Bandyopadhyay, S.C., Goswami, Kar, hep-ph/0312315

Explaining LSND

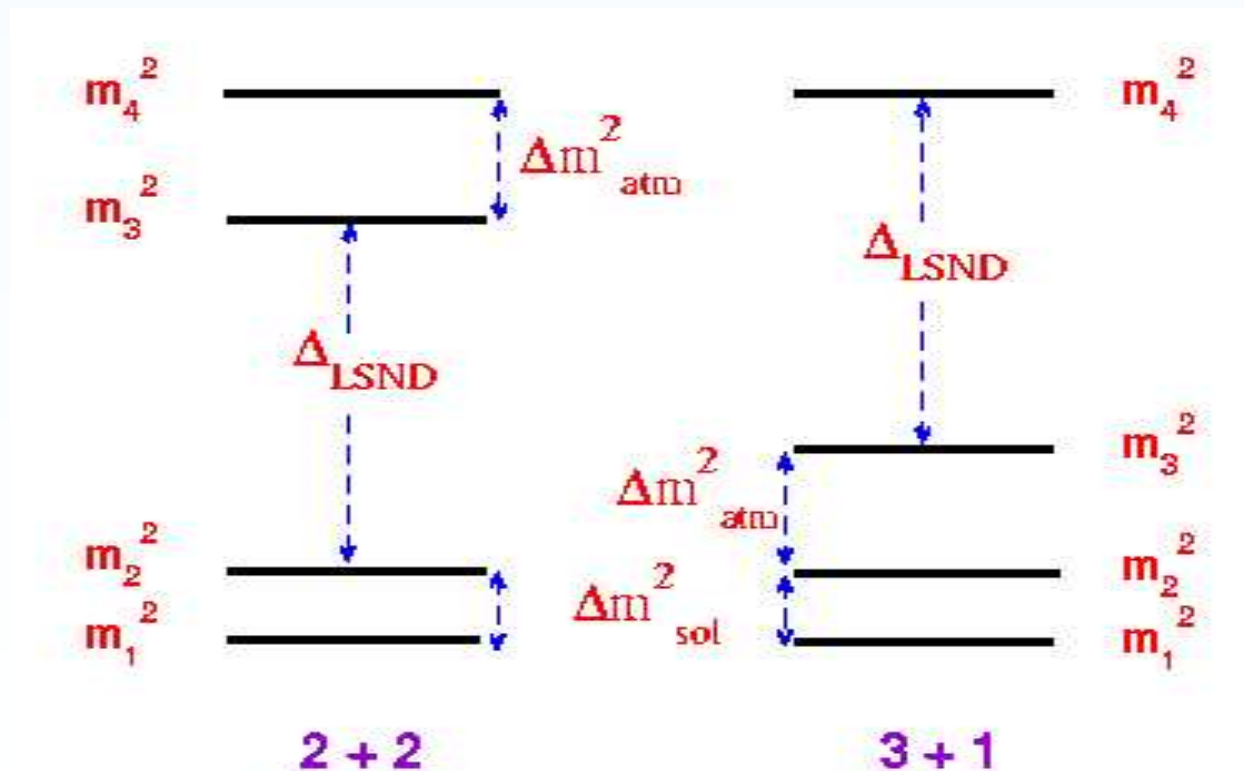


Explaining LSND

- ✓ Is there an evidence for oscillations on very short length scales? $\Rightarrow \Delta m^2 \sim \text{eV}^2$
- ✓ Can it be accommodated with the solar and atm data?

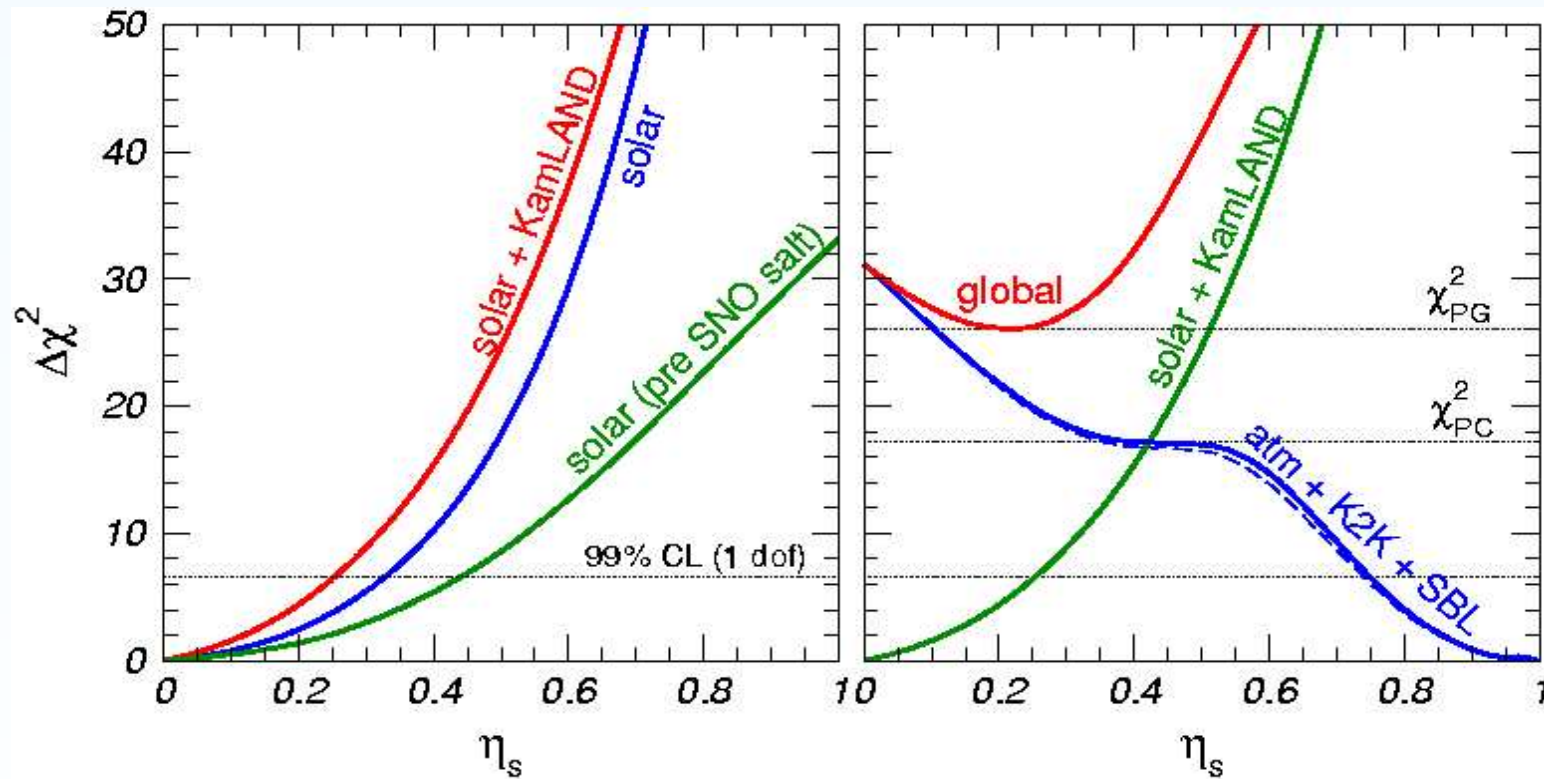
Explaining LSND

- ✓ Is there an evidence for oscillations on very short length scales? $\Rightarrow \Delta m^2 \sim \text{eV}^2$
- ✓ Can it be accommodated with the solar and atm data?
- ✓ Add one sterile neutrino – **2+2 and 3+1 mass schemes**



Explaining LSND

- ✓ The 2+2 scheme is ruled out by the solar and atm data

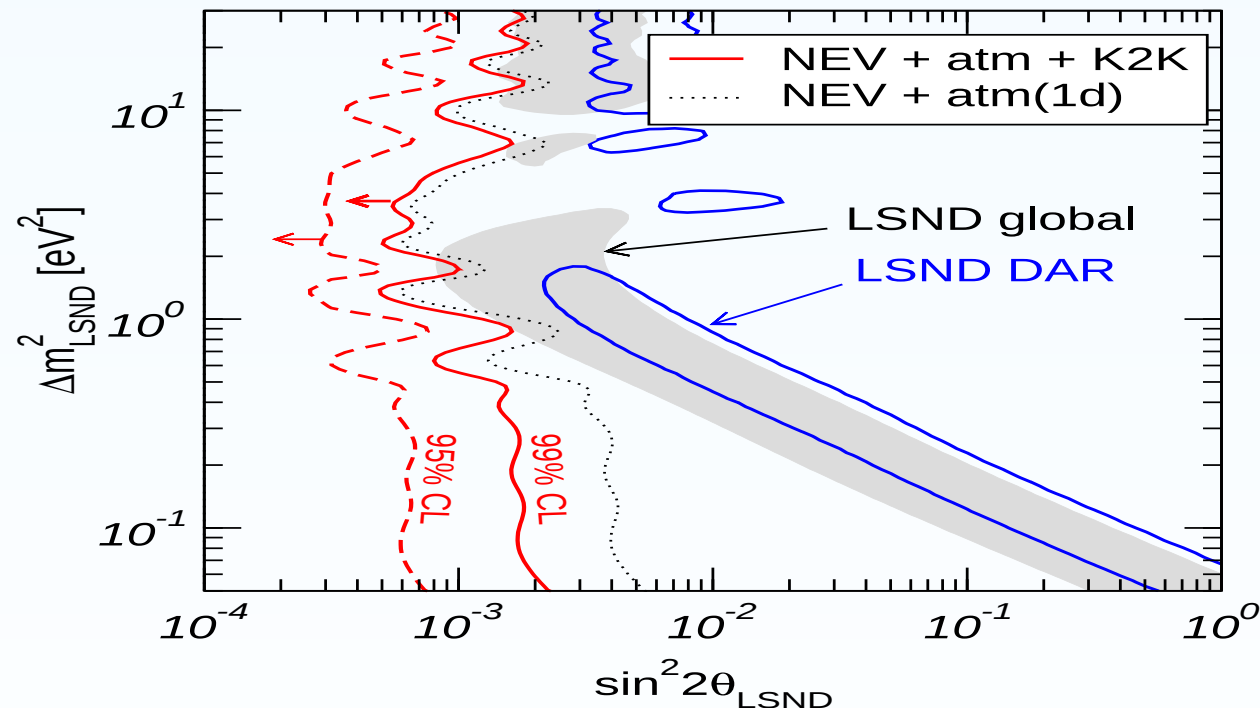


Maltoni et al, hep-ph/0405172

- ✓ Oscillations to pure sterile states is strongly disfavored by both the solar and atmospheric data

Explaining LSND

- ✓ The 3+1 scheme is strongly disfavored by the SBL data

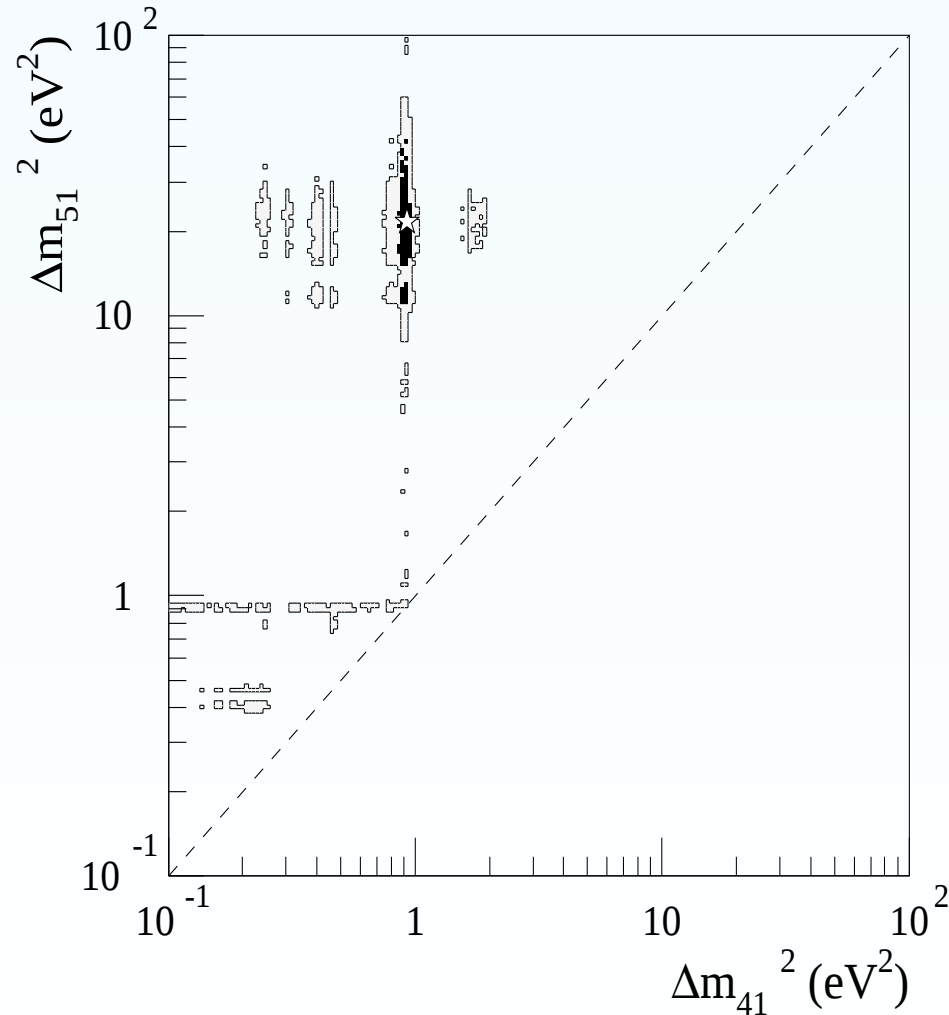


Maltoni et al, hep-ph/0405172

- ✓ Very small area allowed at the 99% C.L.

The (3+2) Scenario to Explain LSND

✓ Addition of a 2nd sterile ν state reduces the tension



✓ Best-Fit:

$$\Delta m_{41}^2 = 0.92 \text{ eV}^2,$$

$$\Delta m_{51}^2 = 22 \text{ eV}^2,$$

$$U_{e4} = 0.121,$$

$$U_{\mu 4} = 0.204,$$

$$U_{e5} = 0.036,$$

$$U_{\mu 5} = 0.224$$

✓ BF2:

$$0.46 \text{ eV}^2$$

$$0.89 \text{ eV}^2$$

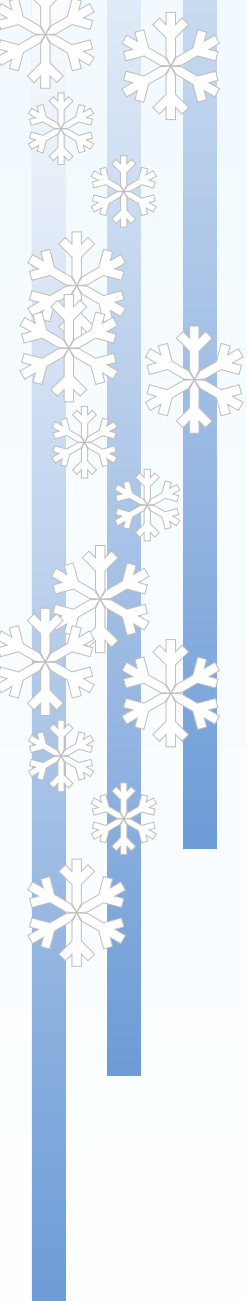
$$0.090$$

$$0.226$$

$$0.125$$

$$0.160$$

Sorel, Conrad, Shaevitz, hep-ph/0305255



Survival Probability of SN $\bar{\nu}_e$ in (3+2) Scenario

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✓ Mass Spectra Possible within (3+2)

$$\text{N2} + \text{N3} : \quad \Delta m_{31}^2 > 0 \quad , \Delta m_{41}^2 > 0 \text{ and } \Delta m_{51}^2 > 0 \quad ,$$

$$\text{N2} + \text{I3} : \quad \Delta m_{31}^2 < 0 \quad , \Delta m_{41}^2 > 0 \text{ and } \Delta m_{51}^2 > 0 \quad ,$$

$$\text{H2} + \text{N3} : \quad \Delta m_{31}^2 > 0 \quad , \Delta m_{41}^2 > 0 \text{ and } \Delta m_{51}^2 < 0 \quad ,$$

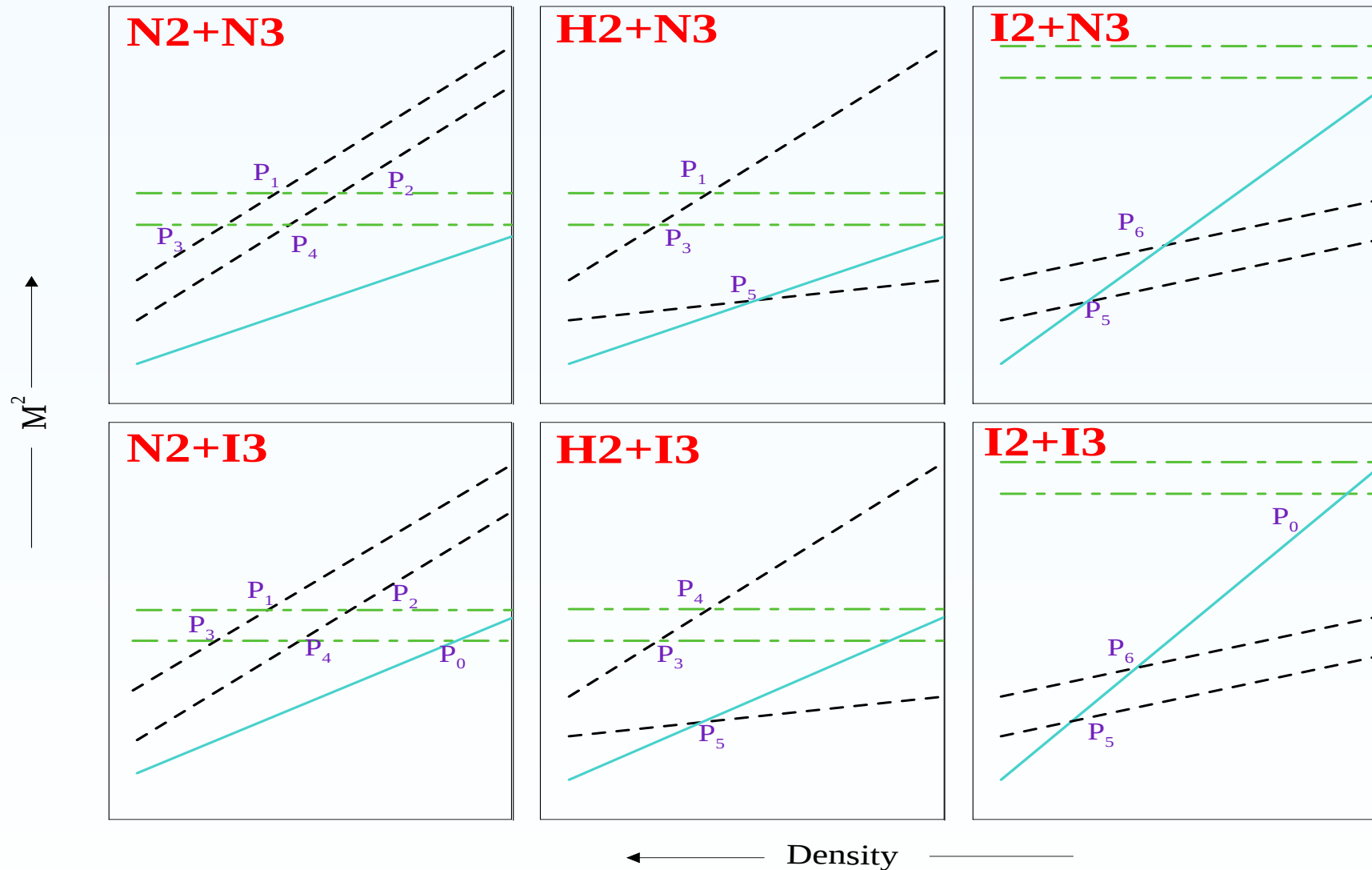
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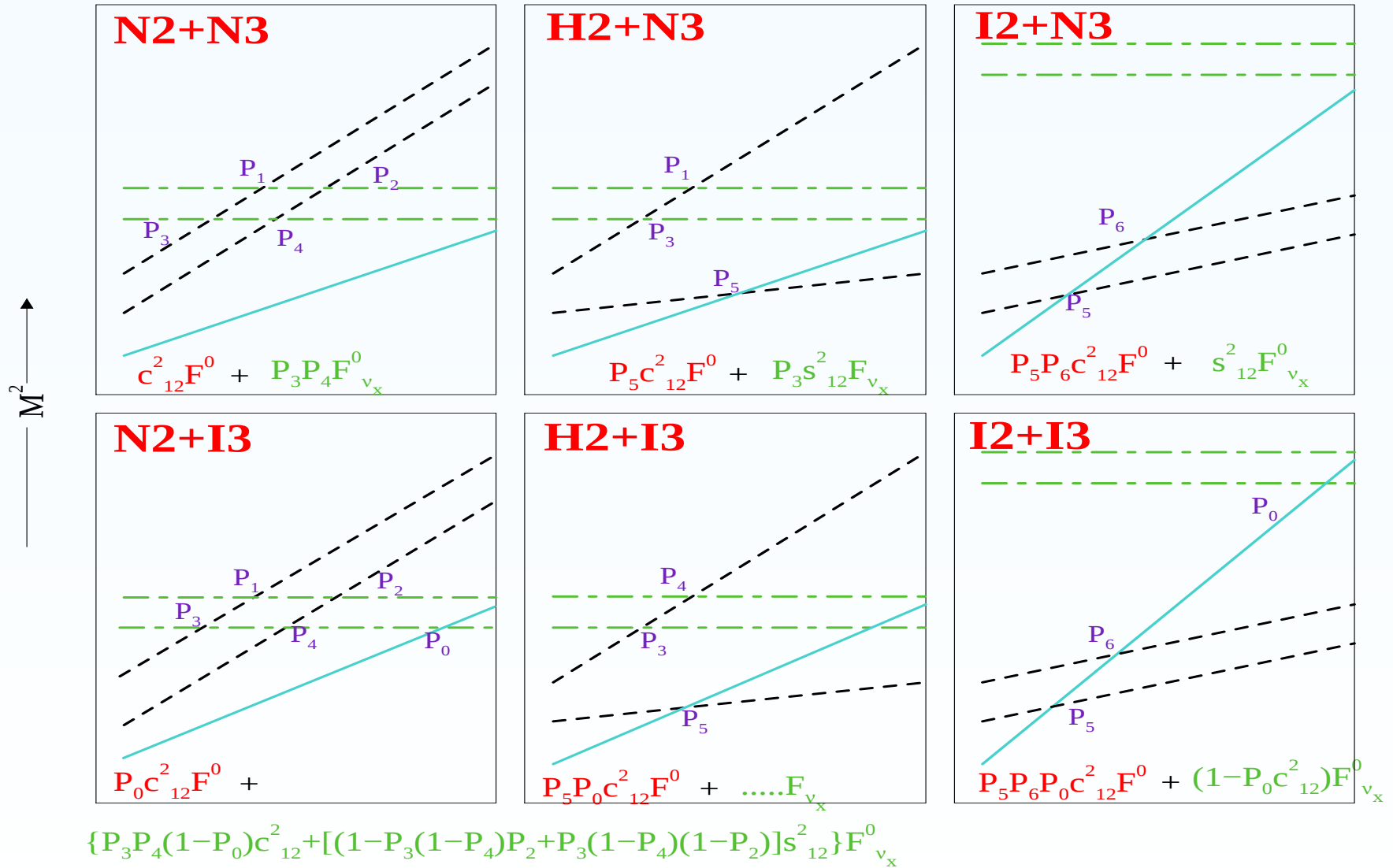
✓ $\Delta m_{21}^2 > 0$ at more than 6σ from solar neutrino data

Survival Probability of SN $\bar{\nu}_e$ in (3+2) Scenario



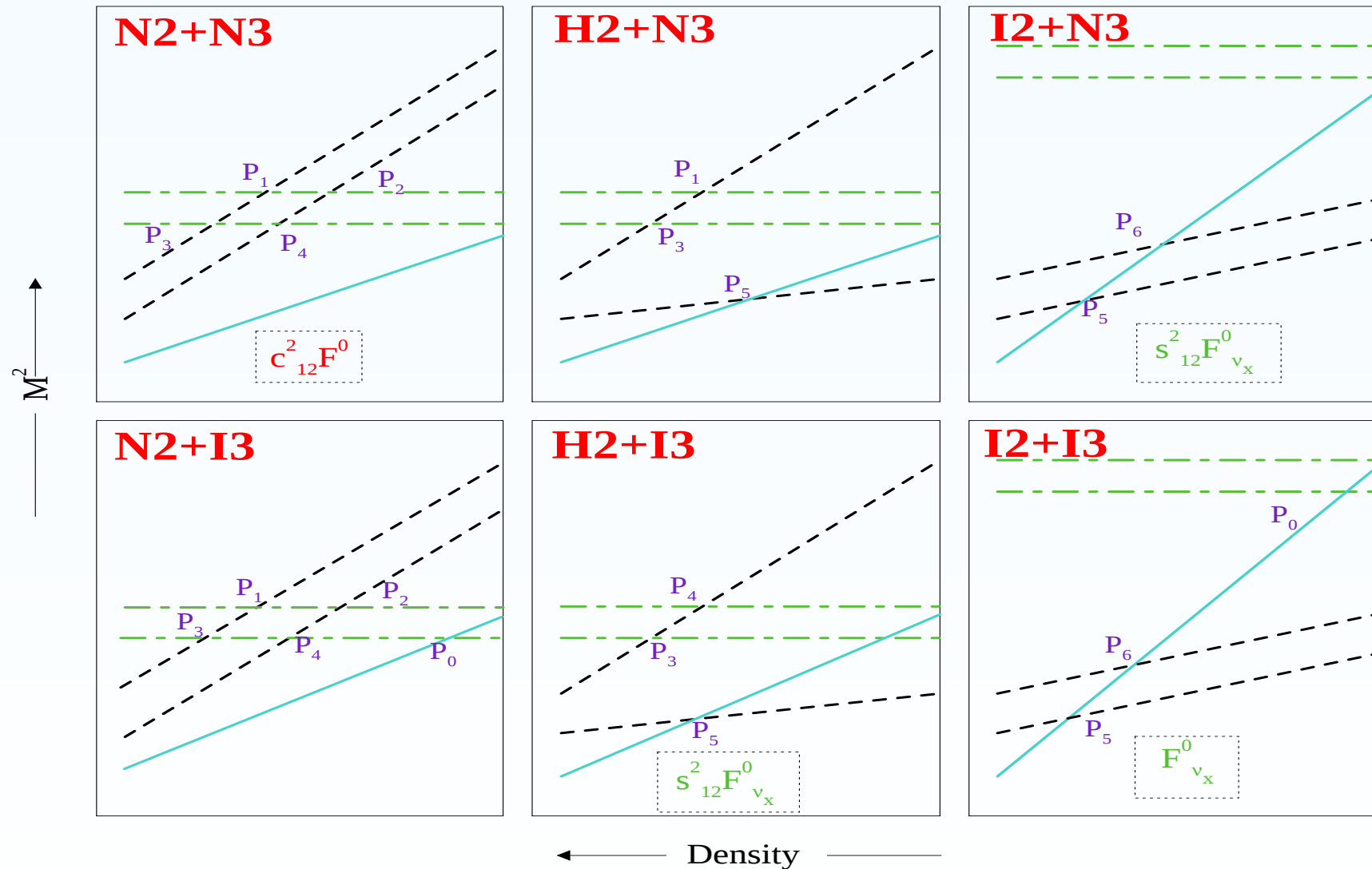
✓ Multiple resonances involving the steriles

Survival Probability of SN $\bar{\nu}_e$ in (3+2) Scenario



✓ Two terms in F_e : (i) **remainder of $\bar{\nu}_e$** (ii) **conversion from ν_x**

Survival Probability of SN $\bar{\nu}_e$ in (3+2) Scenario



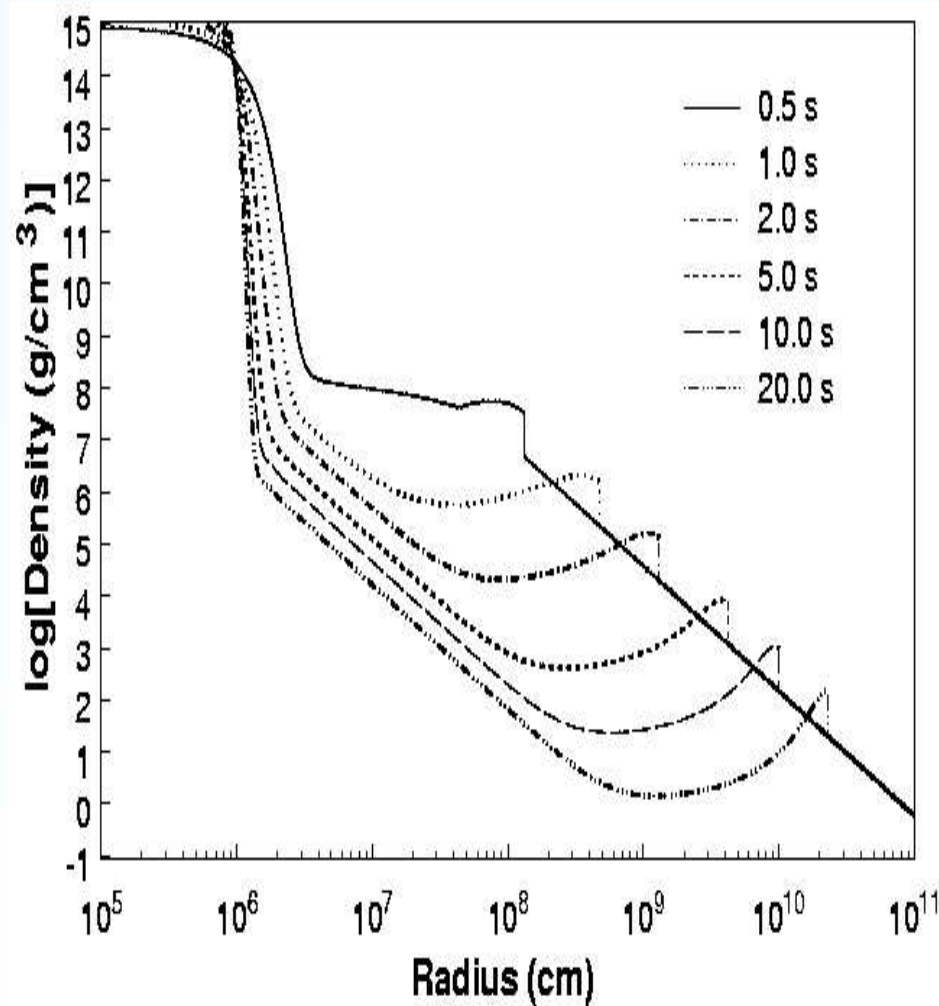
✓ The resonances are all adiabatic

Shock Wave and Level Crossing



Shock Wave and Level Crossing

FORWARD SHOCK

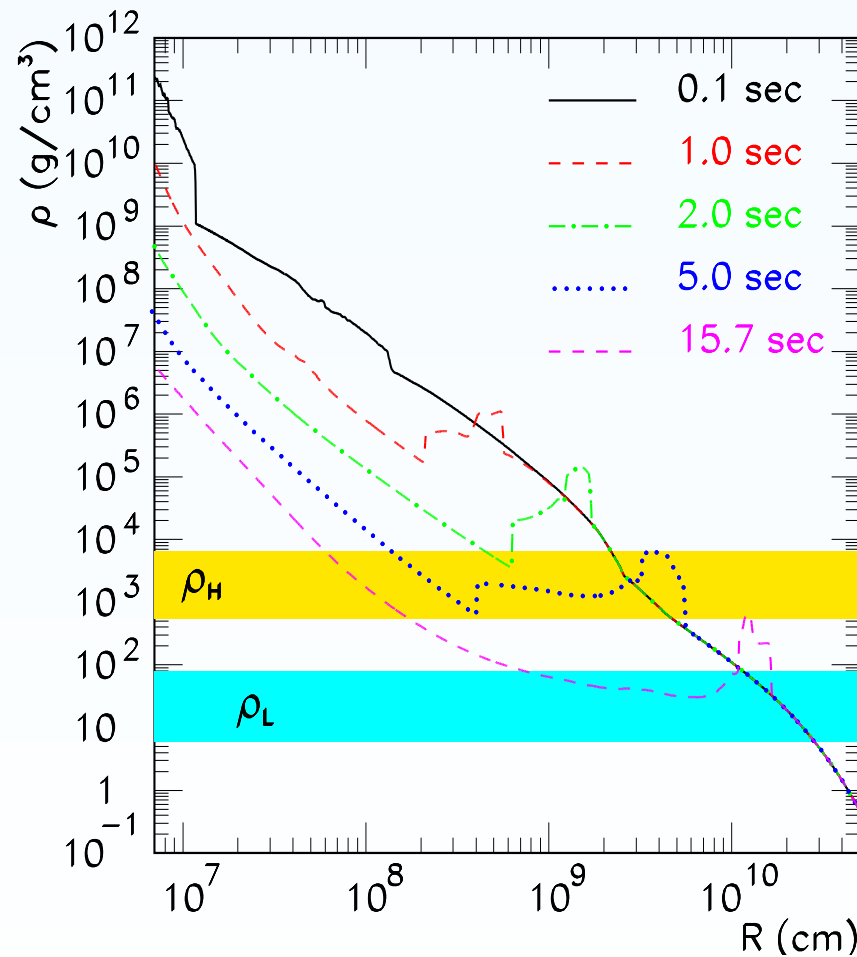


- ✓ Density profile shows a sharp jump at the shock front
- ✓ This sharp change in density causes:
- ✓ More than 1 resonance driven for a given Δm^2
- ✓ The additional resonances are non-adiabatic

Schirato and Fuller, astro-ph/0205390

Shock Wave and Level Crossing

REVERSE SHOCK



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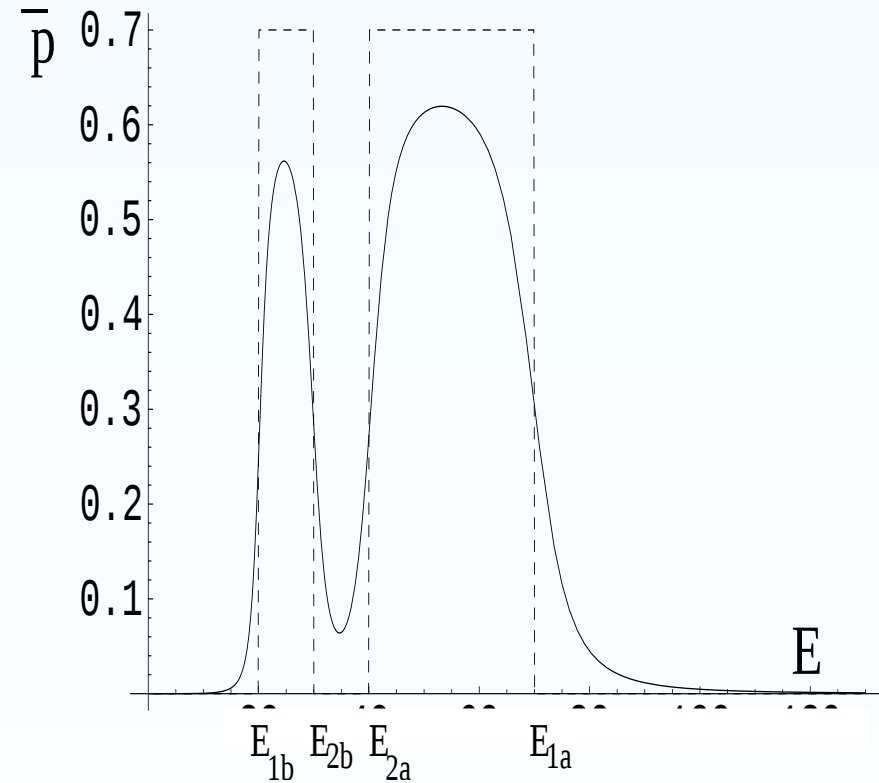
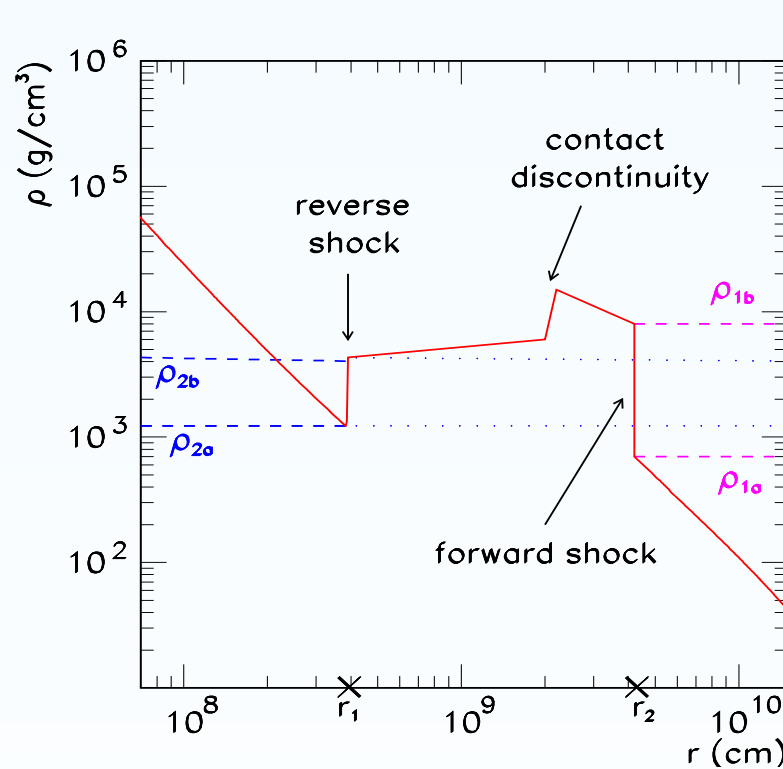
✓ More than 1 resonance driven for a given Δm^2

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Tomas, Kachelriess, Raffelt, Dighe, Janka, Scheck, hep-ph/0407132

✓ Reverse Shock is a generic feature of all SN models

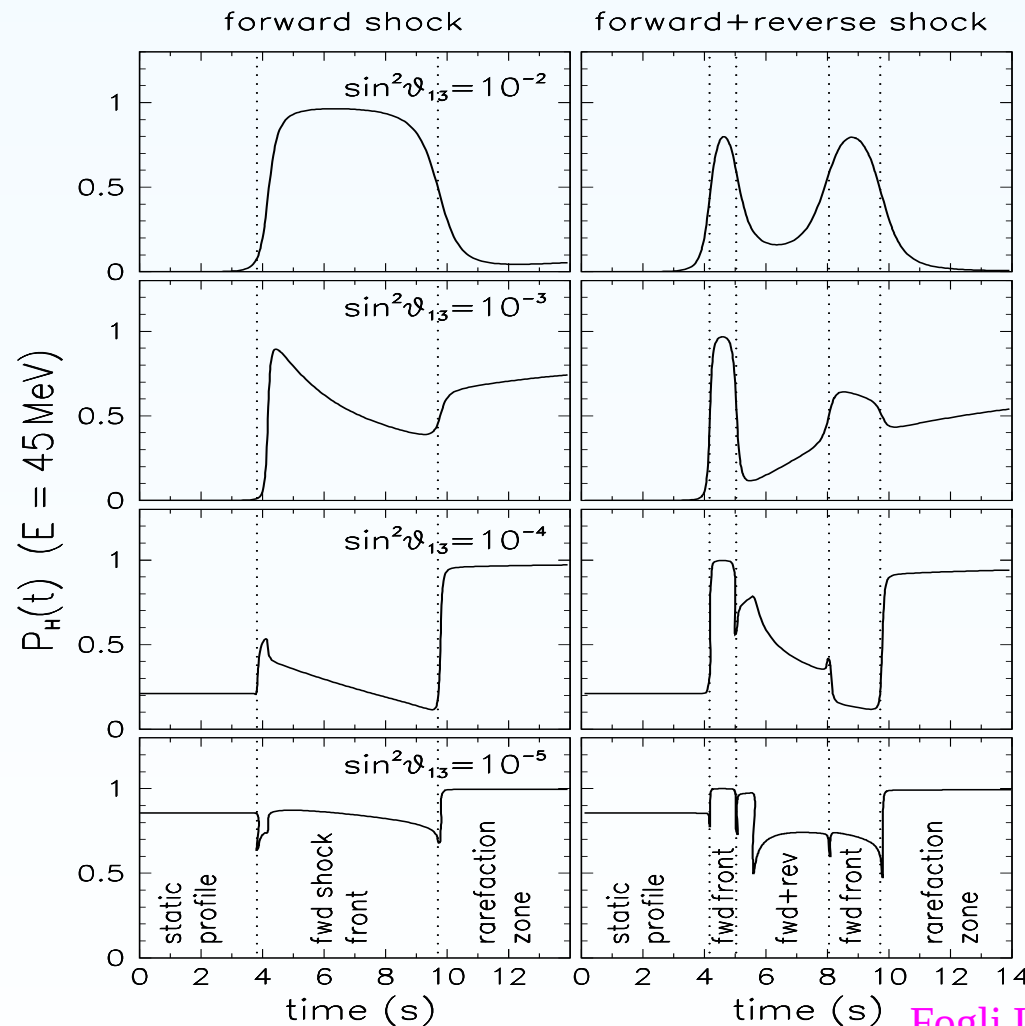
Impact of Shock on the Survival Probability



Tomas, Kachelreiss, Raffelt, Dighe, Janka, Scheck, hep-ph/0407132

$$P_J \approx \exp \left(-\pi \frac{\Delta m^2 \sin^2 \theta}{E} \left| \frac{d \ln N_e}{dx} \right|^{-1}_{x=x_{\text{res}}} \right)$$

Impact of Shock on the Survival Probability



✓ Sudden broad “dip” in E and t is the signature of the forward shock

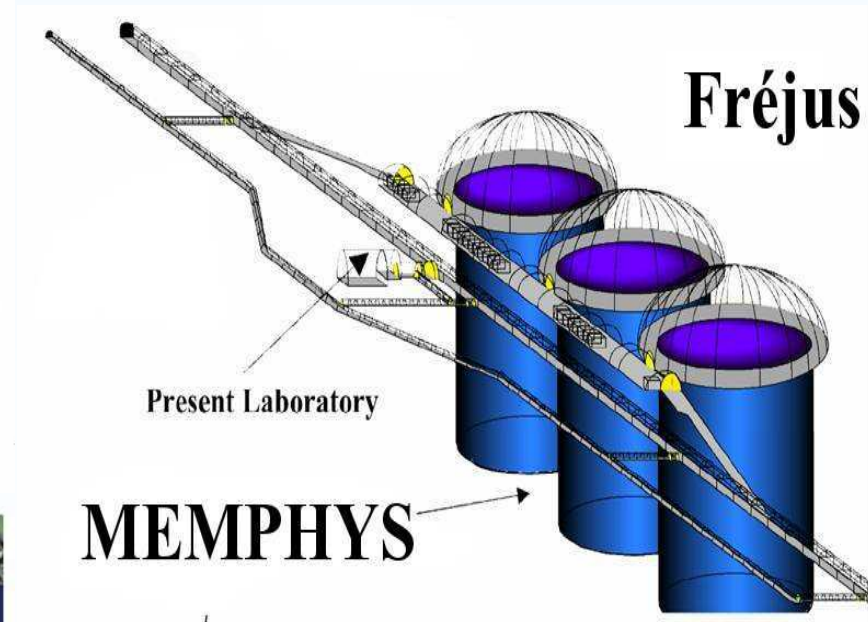
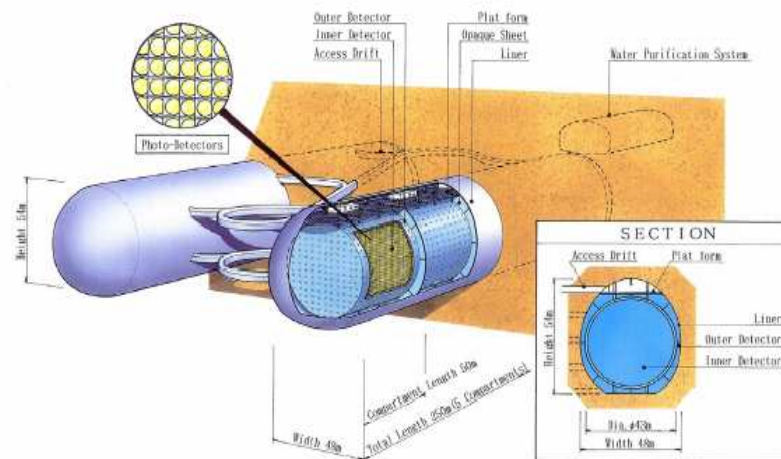
✓ “Double Dip” in E and t is the signature of the forward and reverse shock

Fogli, Lisi, Mirizzi, Montanino, hep-ph/0412046

✓ In principle, position of the dips can be mapped to the density profile of the SN and hence to its shock dynamics

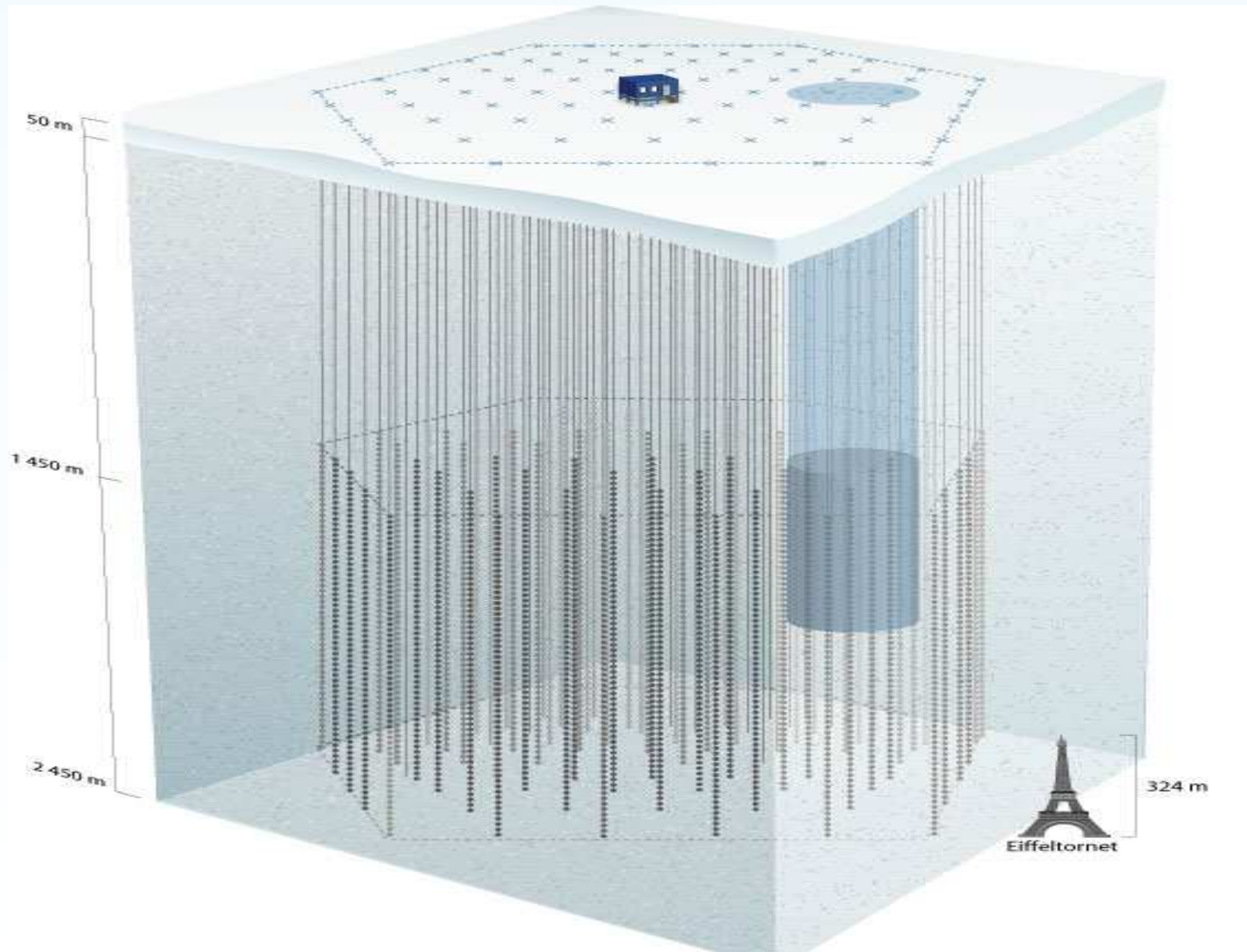
SN Neutrinos in Megaton Water Detectors

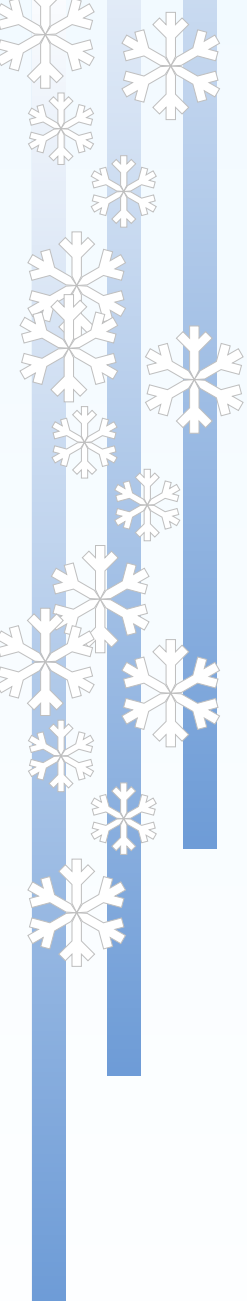
~ 1 Mton water Cherenkov detector at Kamioka



✓ Will measure both E and time

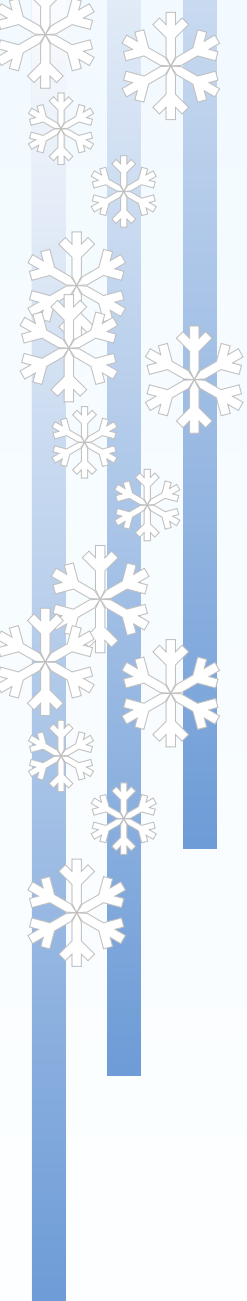
IceCube as a Supernova Detector





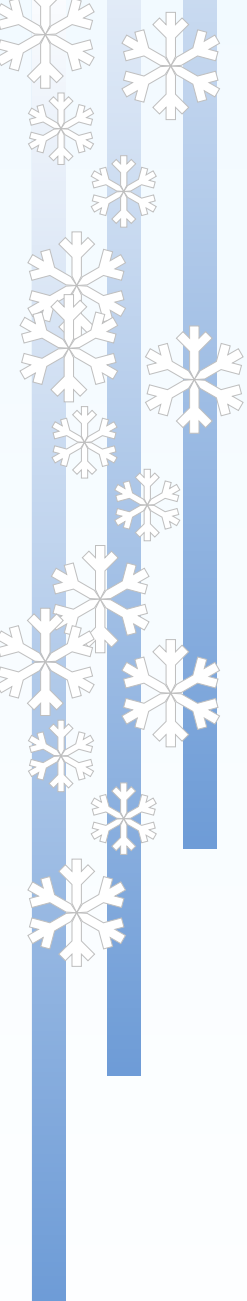
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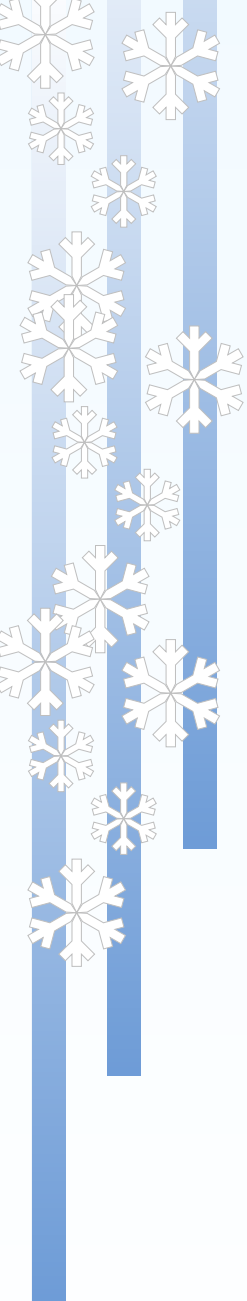
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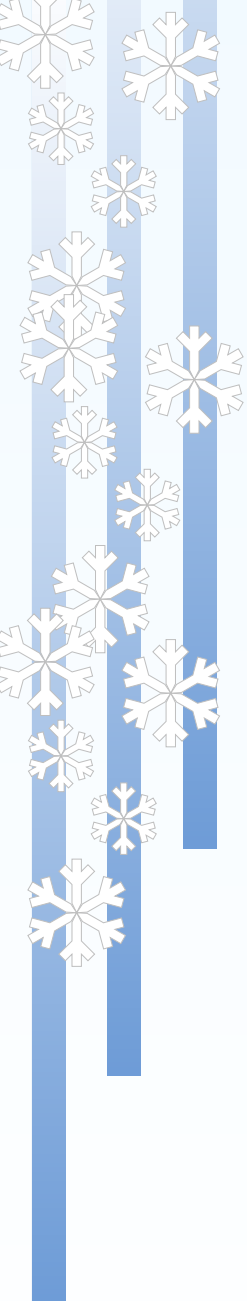
Halzen, Jacobson, Zas (1995)

Dighe, Keil, Raffelt (2003)



IceCube as a Supernova Detector

- ✓ Measuring the E of the ν events might not be possible

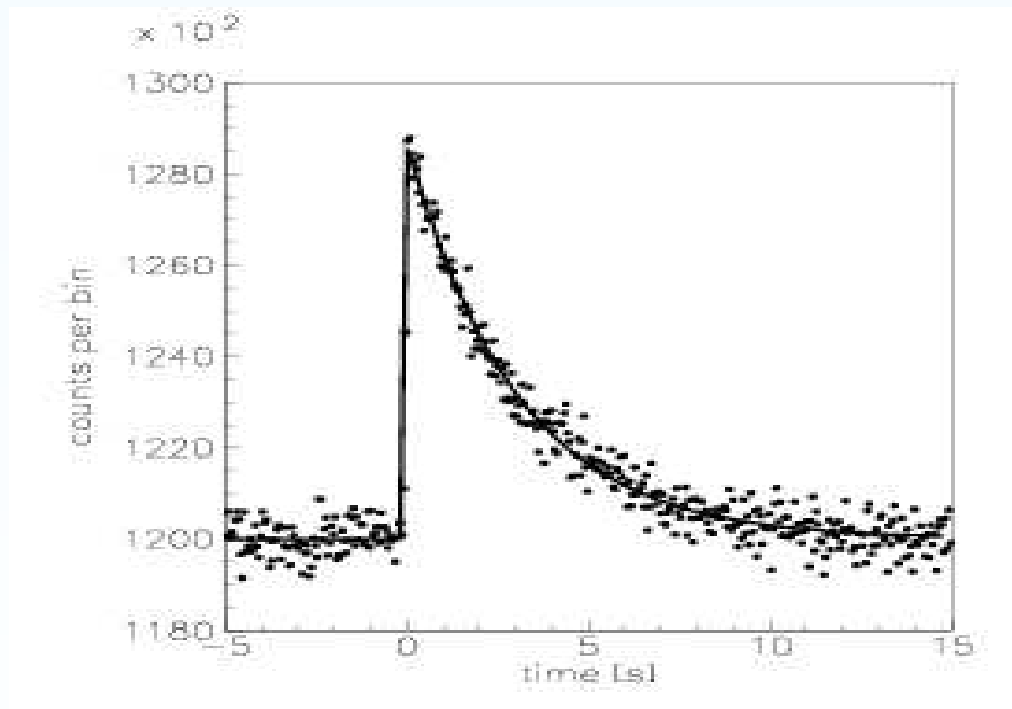


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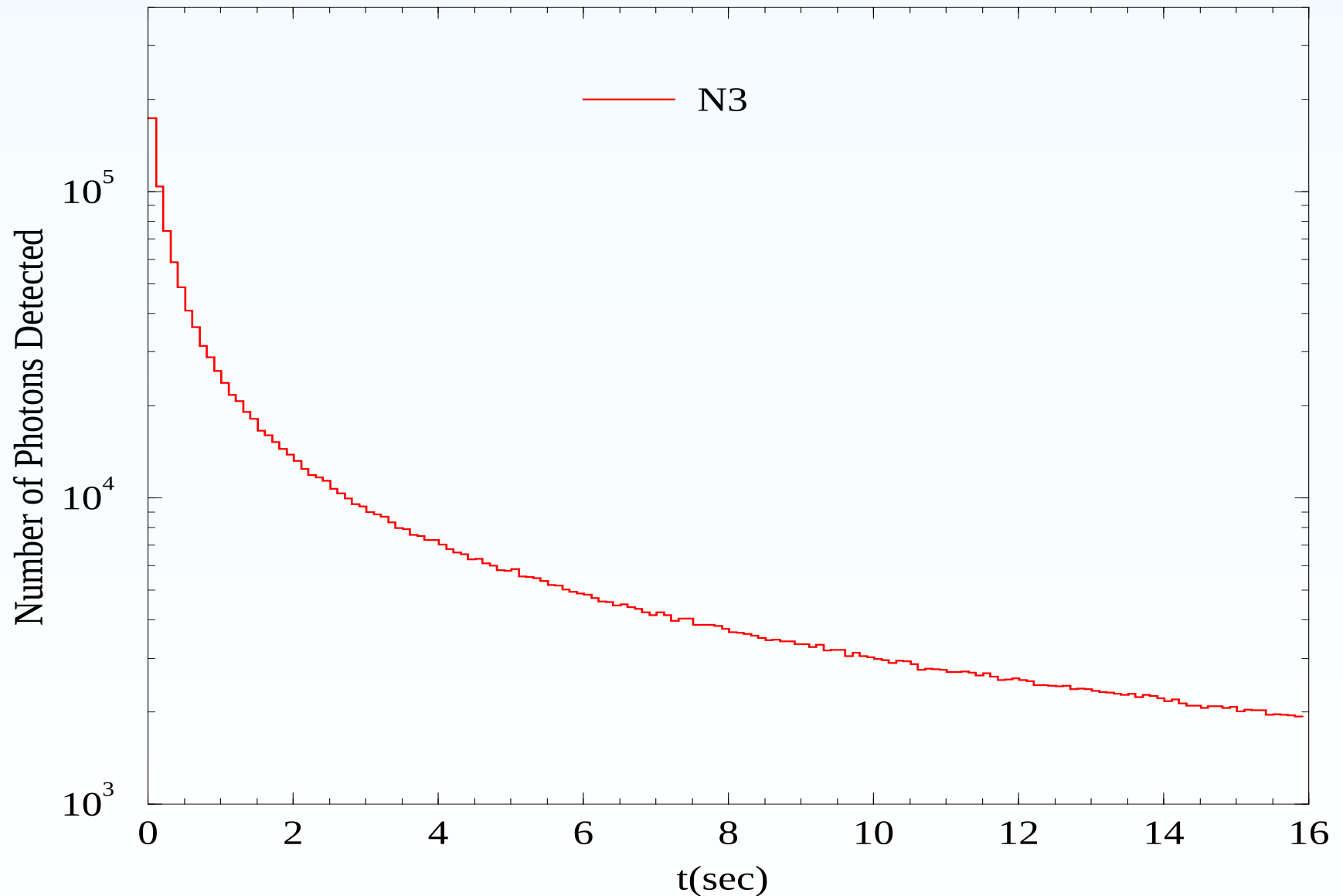
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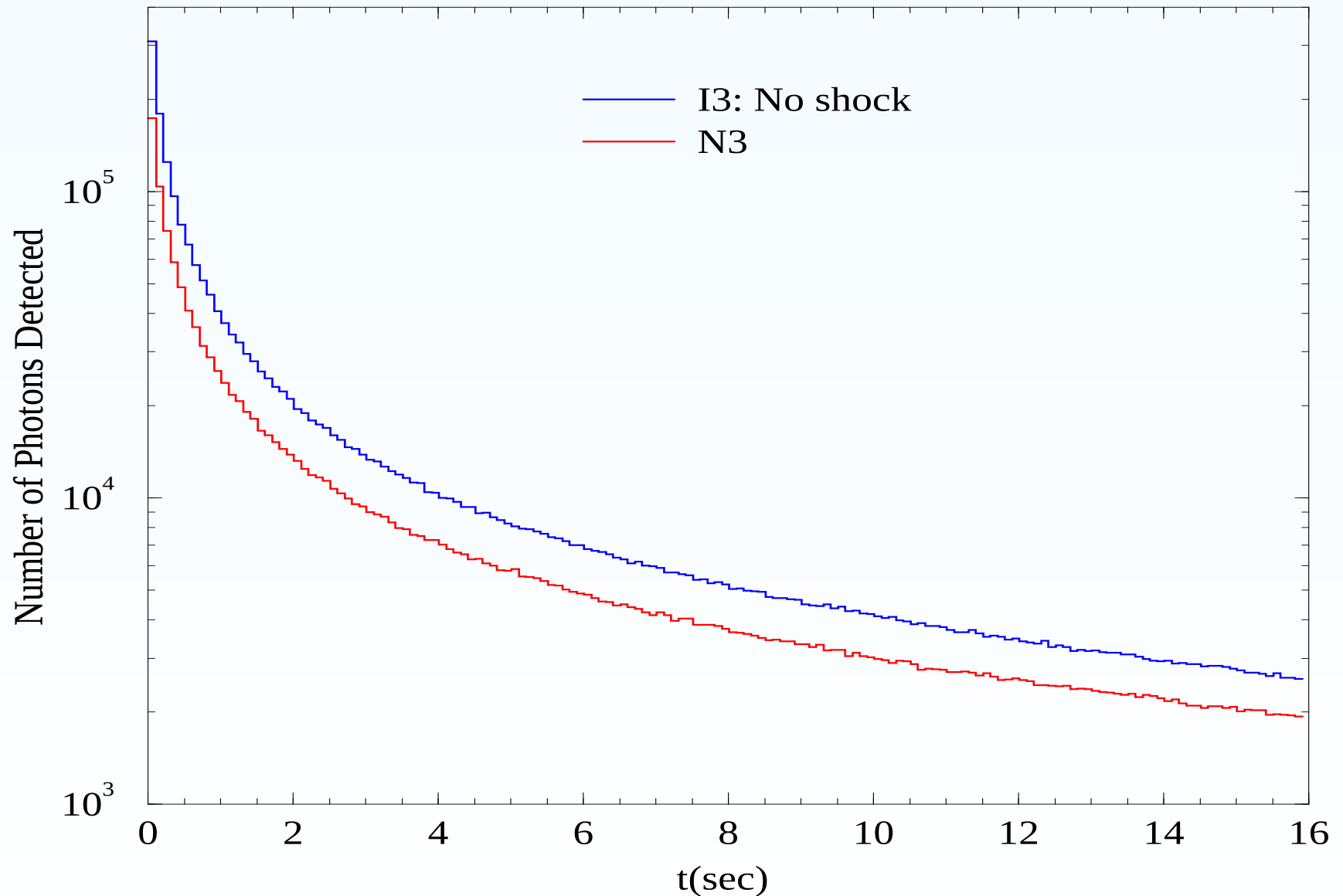


IceCube Collaboration

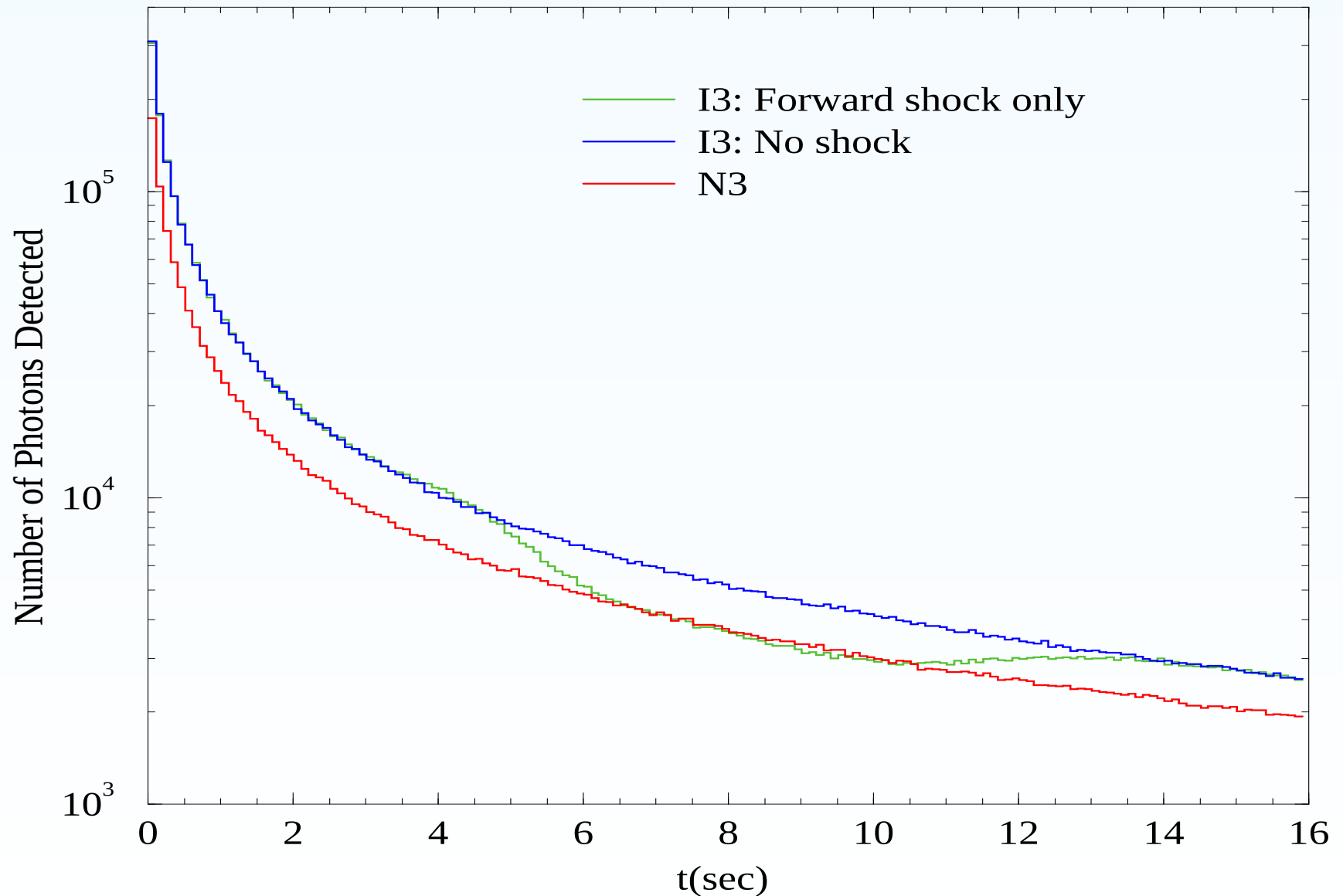
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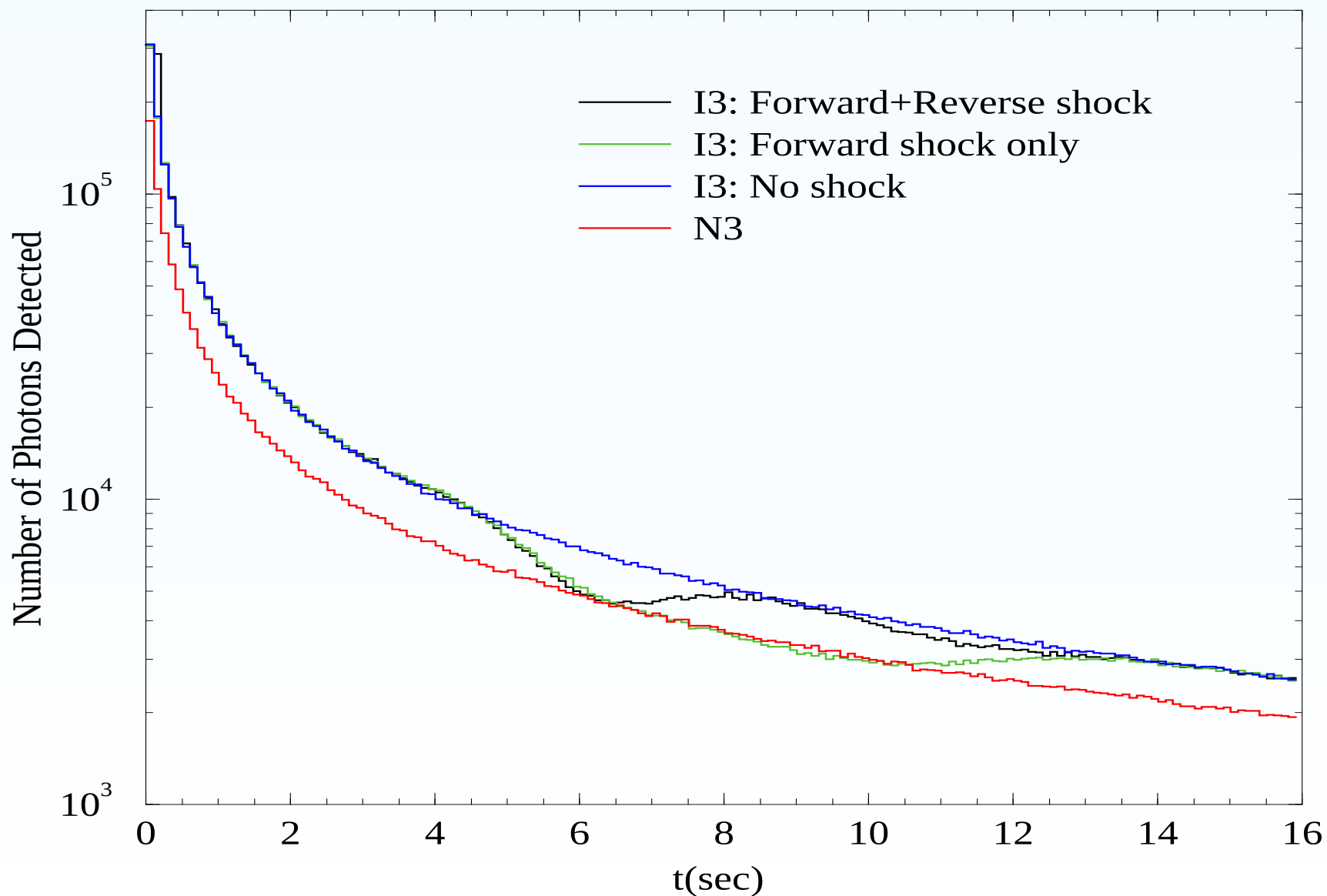
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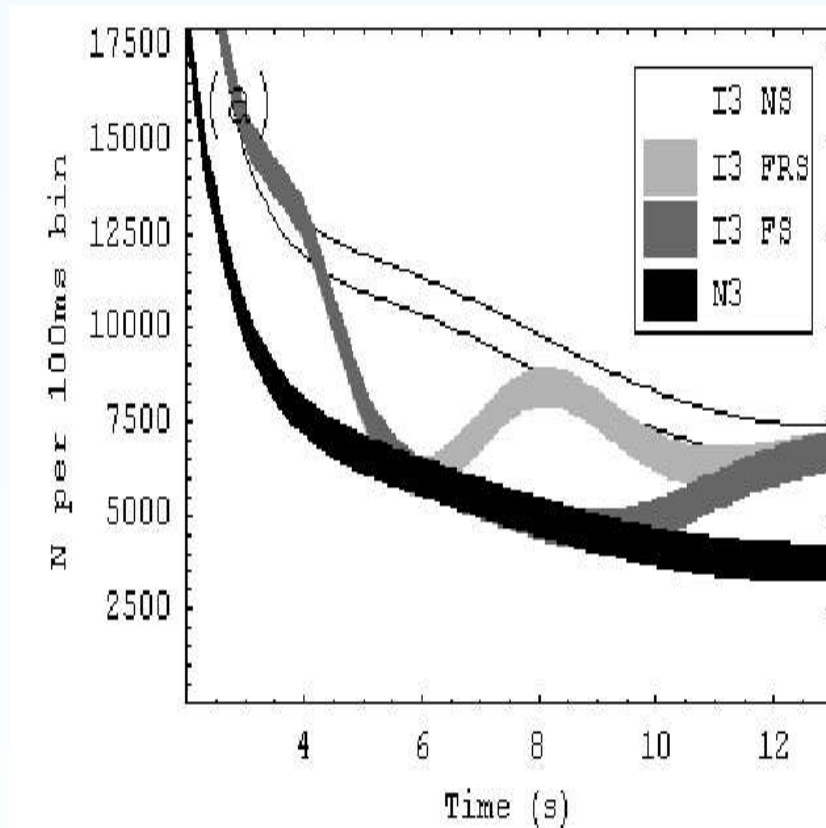
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SN ν Signature in IceCube for Three Neutrinos

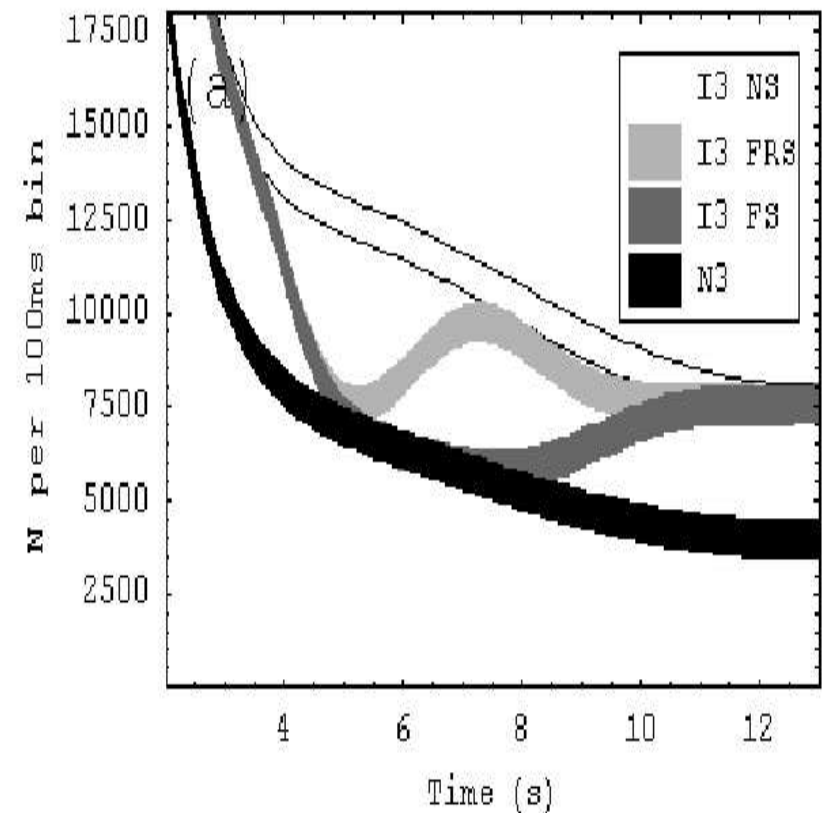


SN ν Signature in IceCube for Three Neutrinos



$$\sqrt{\langle E_{\bar{\nu}_e} \rangle} = 15, \langle E_{\nu_x} \rangle = 24$$

$$\sqrt{\frac{\Phi_{\bar{\nu}_e}^0}{\Phi_{\nu_x}^0}} = 1.6$$

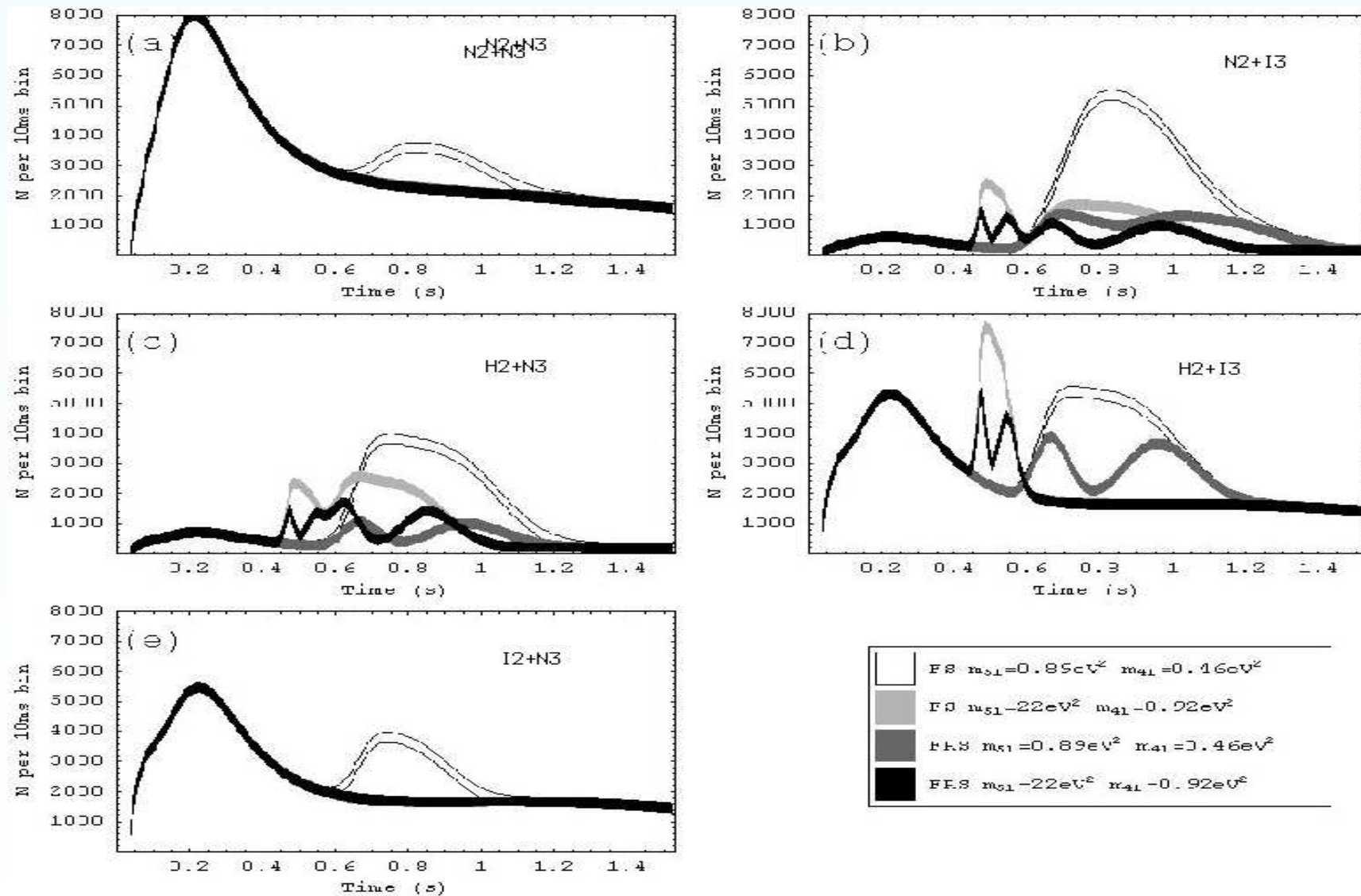


$$\sqrt{\langle E_{\bar{\nu}_e} \rangle} = 15, \langle E_{\nu_x} \rangle = 18$$

$$\sqrt{\frac{\Phi_{\bar{\nu}_e}^0}{\Phi_{\nu_x}^0}} = 0.8$$

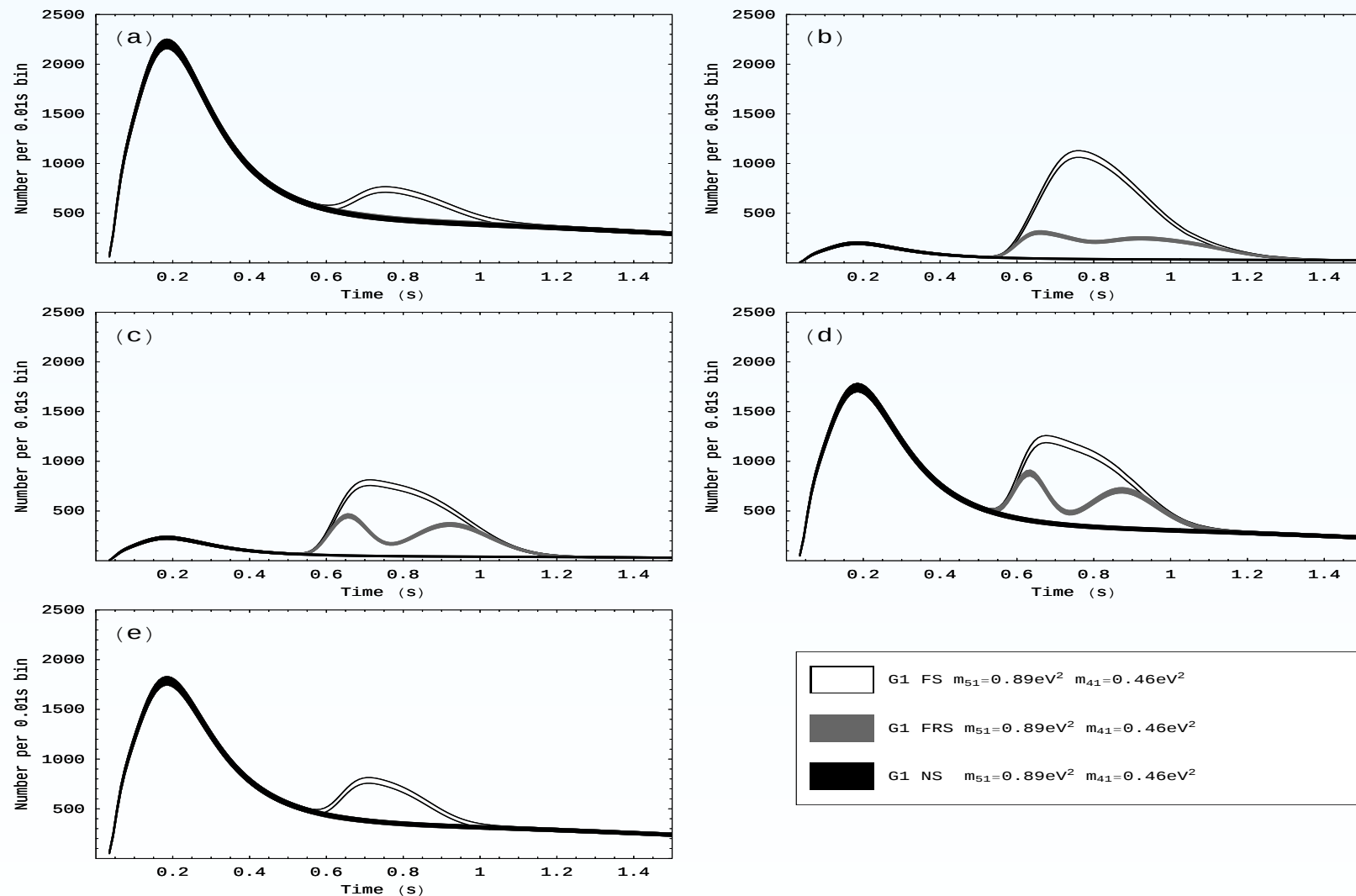
SC, Harries, Ross, hep-ph/0605255

SN ν Signature in IceCube for Five Neutrinos



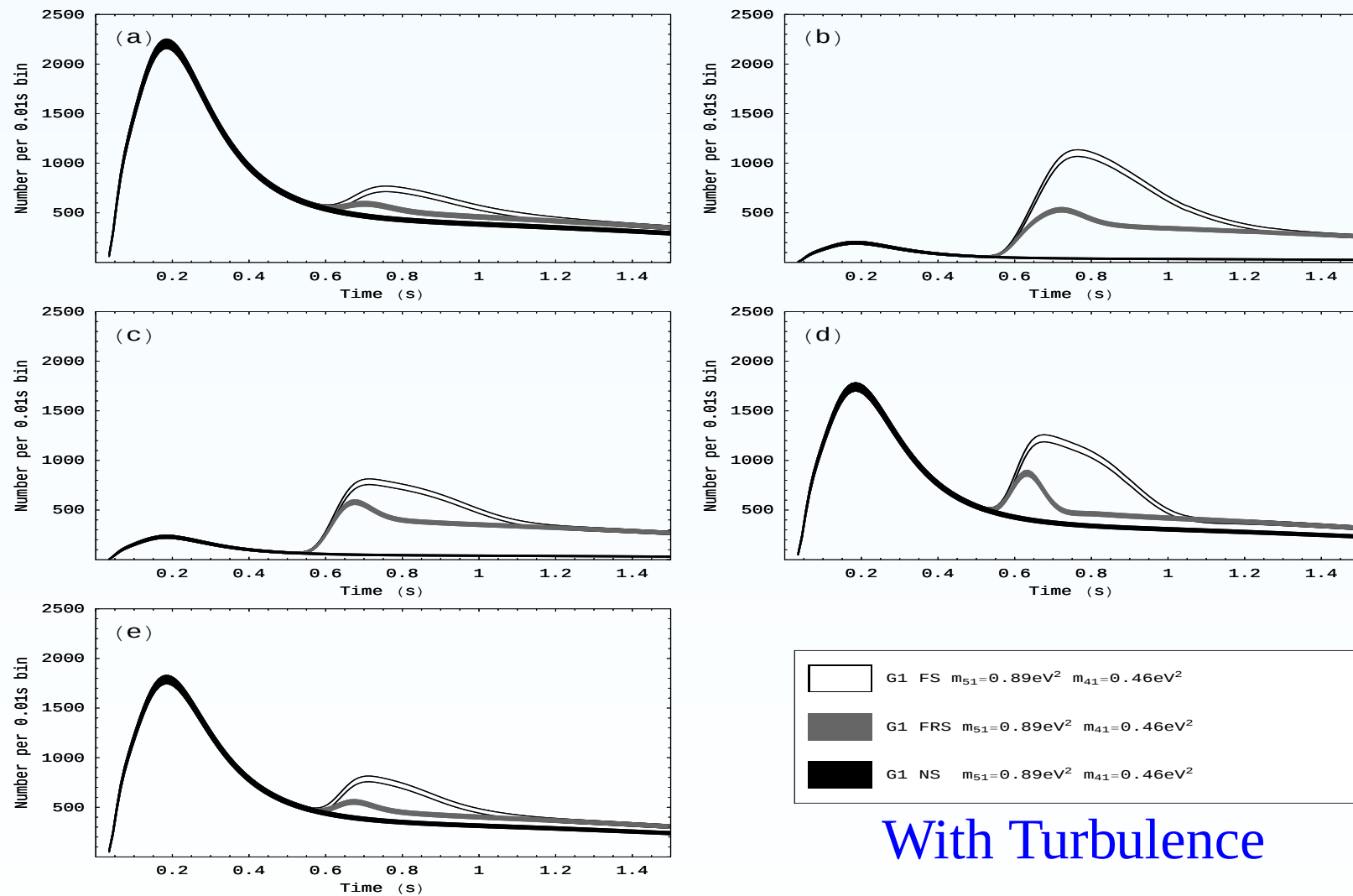
✓ Shock effects at **very early times** signal sterile neutrinos

SN ν Signature in “Hyper-K” for Five Neutrinos



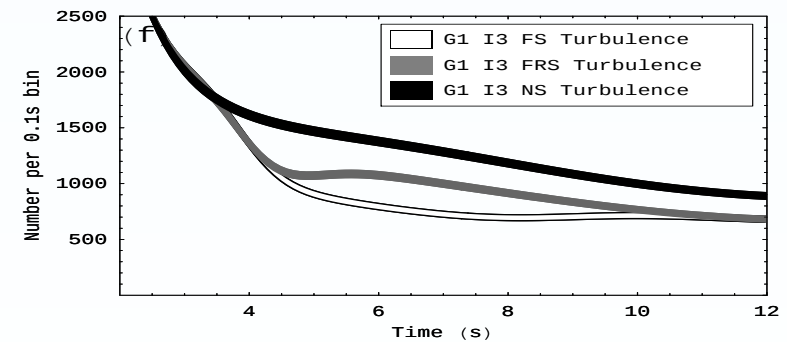
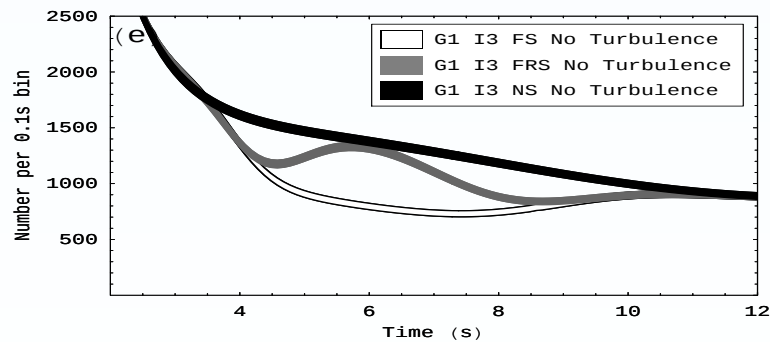
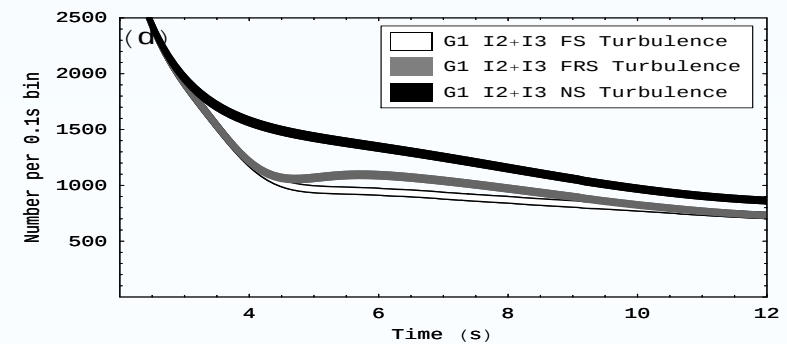
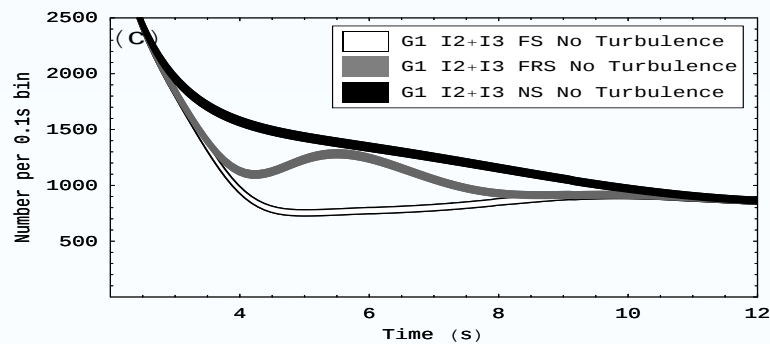
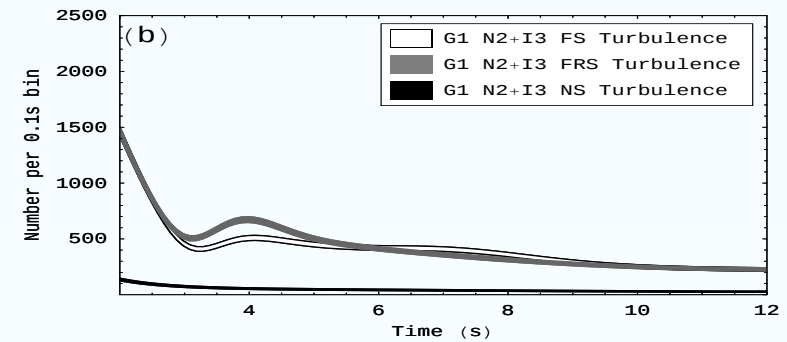
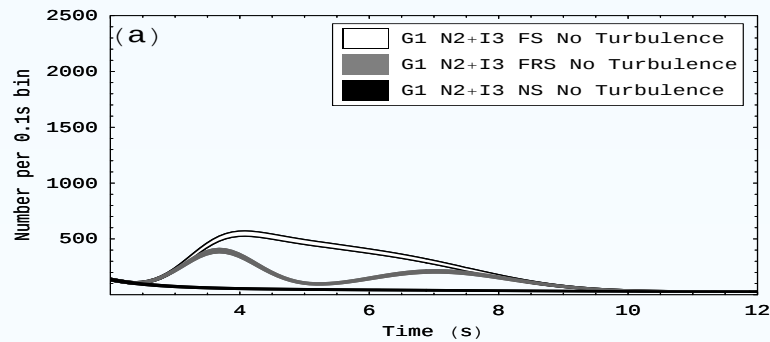
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SN ν Signature in “Hyper-K” for Five Neutrinos



✓ Shock effects at **very early times** signal sterile neutrinos

SN ν Signature in “Hyper-K” for Five Neutrinos



✓ Shock effects at late times

Comparing early and late time signal

Spectrum	Early Times		Late Times	
	Events	Shock Effect	Events	Shock Effects
N3	L	χ	L	χ
I3	L	χ	L	$\checkmark$
N2+N3	L	$\checkmark$	S	χ
N2+I3	S	$\checkmark$	S	$\checkmark$
H2+N3	S	$\checkmark$	S	χ
H2+I3	L	$\checkmark$	S	χ
I2+N3	L	$\checkmark$	S	χ
I2+I3	L	χ	L	$\checkmark$

$\checkmark$ L $\Rightarrow$ Large events; S $\Rightarrow$ Small events

Comparing early and late time signal

Spectrum	Early Times		Late Times	
	Events	Shock Effect	Events	Shock Effects
N3	L	χ	L	χ
I3	L	χ	L	$\checkmark$
N2+N3	L	$\checkmark$	S	χ
N2+I3	S	$\checkmark$	S	$\checkmark$
H2+N3	S	$\checkmark$	S	χ
H2+I3	L	$\checkmark$	S	χ
I2+N3	L	$\checkmark$	S	χ
I2+I3	L	χ	L	$\checkmark$

$\checkmark$ All sterile cases (except I2+I3) have shock effects at early times

Comparing early and late time signal

Spectrum	Early Times		Late Times	
	Events	Shock Effect	Events	Shock Effects
N3	L	χ	L	χ
I3	L	χ	L	$\checkmark$
N2+N3	L	$\checkmark$	S	χ
N2+I3	S	$\checkmark$	S	$\checkmark$
H2+N3	S	$\checkmark$	S	χ
H2+I3	L	$\checkmark$	S	χ
I2+N3	L	$\checkmark$	S	χ
I2+I3	L	χ	L	$\checkmark$

$\checkmark$ **N3** has a unique signature and hence can be identified

Comparing early and late time signal

Spectrum	Early Times		Late Times	
	Events	Shock Effect	Events	Shock Effects
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N2+I3	S	$\checkmark$	S	$\checkmark$
H2+N3	S	$\checkmark$	S	χ
H2+I3	L	$\checkmark$	S	χ
I2+N3	L	$\checkmark$	S	χ
I2+I3	L	χ	L	$\checkmark$

$\checkmark$ N2+I3 has a unique signature and hence can be identified

Comparing early and late time signal

Spectrum	Early Times		Late Times	
	Events	Shock Effect	Events	Shock Effects
N3	L	χ	L	χ
I3	L	χ	L	$\checkmark$
N2+N3	L	$\checkmark$	S	χ
N2+I3	S	$\checkmark$	S	$\checkmark$
H2+N3	S	$\checkmark$	S	χ
H2+I3	L	$\checkmark$	S	χ
I2+N3	L	$\checkmark$	S	χ
I2+I3	L	χ	L	$\checkmark$

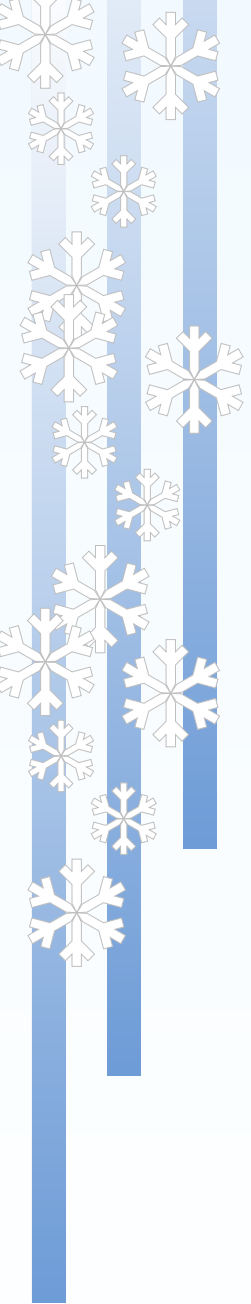
$\checkmark$ H2+N3 has a unique signature and hence can be identified

Comparing early and late time signal

Spectrum	Early Times		Late Times	
	Events	Shock Effect	Events	Shock Effects
N3	L	×	L	×
I3	L	×	L	✓
N2+N3	L	✓	S	×
N2+I3	S	✓	S	✓
H2+N3	S	✓	S	×
H2+I3	L	✓	S	×
I2+N3	L	✓	S	×
I2+I3	L	×	L	✓

✓(I3, I2+I3) are similar & (N2+N3, H2+I3, I2+N3) are similar
 ✓Model dependent analysis can differentiate between them

Conclusions



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- ✓ Existence of sterile neutrinos can be tested *model independently* using shock effects in SN neutrinos – even when $\langle E_{\bar{\nu}_e} \rangle \approx \langle E_{\bar{\nu}_\mu} \rangle$

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Thankyou!