

Cluster algorithms for frustrated Ising models

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Ingredients

- ▶ Cluster constructions
- ▶ (Dual) “loop-like” (dimer) representations
- ▶ Directed worm constructions

Particularly appropriate setting...

Frustration and entropic interactions

- ▶ Frustrated magnets: Large degeneracy of minimum energy configurations
- ▶ At $T \ll J$: system samples minimally frustrated subspace
(Or falls out of equilibrium...)
- ▶ Fluctuations generate entropic interactions

Order by disorder:

- ▶ Low temperature physics dominated by entropic interactions
- ▶ Characteristic signatures in structure factor
- ▶ More dramatic cases: Order-by-thermal/quantum disorder

Example

- ▶ $S = 1$ easy-axis triangular lattice antiferromagnet with large single-ion anisotropy

$$H = \sum_{\langle rr' \rangle} \vec{S}_r \cdot \vec{S}_{r'} - D \sum_r (S_r^z)^2$$

- ▶ Reduction to: $S = 1/2$

$$H_{XXZ} = J_z \sum_{\langle rr' \rangle} \sigma_r^z \sigma_{r'}^z - J_\perp \sum_{\langle rr' \rangle} (\sigma_r^+ \sigma_{r'}^- + h.c.)$$

(KD & Senthil 2006)

- ▶ QMC of H_{XXZ} using SSE
(Heidarian & KD 2005, Melko *et.al.* 2005, Wessel & Troyer 2005, Boninsegni & Prokofev 2005 ...)

Key challenge for algorithm

- ▶ Needs to “know” structure of minimally frustrated landscape
Within stochastic Series Expansion (SSE): Directed-loop
updates must be able to move system within minimally frustrated
subspace
- ▶ Solution: Cluster decomposition of H_{XXZ} in triangle Hamiltonians
(Kim & Gross, KD & Heidarian 2004)
- ▶ advantage: incorporates $\Delta E = 0$ changes & $\Delta E \neq 0$ on equal
footing

Another example

- ▶ $S \geq 3/2$ easy axis antiferromagnets on triangular and Kagome lattices
- ▶ Reduction to frustrated Ising model with further-neighbour/multi-spin interactions

$$H_{\text{Ising}} = J \sum_{\langle ij \rangle} \sigma_i \sigma_j + J' \sum_{\langle\langle ij \rangle\rangle} \sigma_i \sigma_j + \dots$$

Challenge for algorithm

- ▶ At $T = 0$, formulate as interacting dimer model and use dimer worm algorithm of Alet *et. al.*.
(Sen *et. al.*, 2008, 09)
- ▶ Difficulty at small nonzero temperature: Need to include higher-energy configurations with correct weight in efficient way
Standard (Wolff-inspired) cluster constructions don't work so well with frustration

Coddington & Han 1994, Zhang and Yang 1994

More recent attempt

- ▶ Worm construction (Wang,Sterck & Melko 2012)
Uses dual geometric worm algorithm (*a la*
Hitchcock,Sorenson,Alet 2004)
- ▶ **Works when $T = 0$ limit is dual to non-interacting dimers**
- ▶ Involves rejection of significant fraction of worms

Third example

- ▶ Triangular lattice Ising antiferromagnet in a transverse field

$$H_{\text{TFIM}} = J \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z - \Gamma \sum_i \sigma_i^x + \dots$$

- ▶ Quantum cluster algorithm available in SSE representation
(Sandvik 2003)

Clusters reduce to variant of Swendsen-Wang clusters in $\Gamma = 0$

limit \rightarrow frustrated J_{ij} again leads to problems(?)

Need to “teach” algorithm physics of minimally frustrated configurations(?)

In this talk...

- ▶ Quantum cluster construction for *frustrated* TFIM
- ▶ Cluster algorithm for frustrated two-dimensional H_{Ising} with up to third neighbour interactions

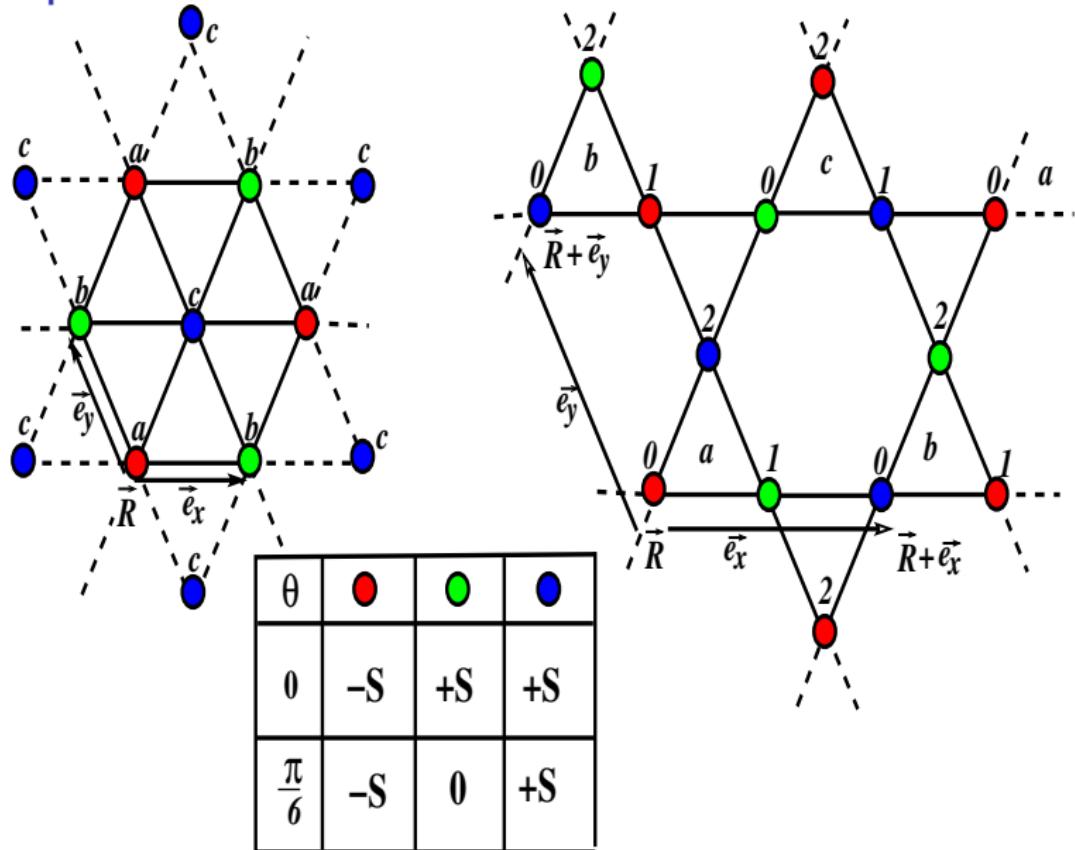
SSE for frustrated TFIM

- ▶ Example: Transverse field Ising antiferromagnet on triangular lattice (also with further neighbour (J_2, J_3) couplings...)
- ▶ Interesting physics questions

Thermodynamic signature of two-step melting of three-sublattice order

Transition from plaquette to columnar three-sublattice order

Order parameter



For triangular lattice: $\Psi = \sum_r e^{i\mathbf{Q} \cdot \mathbf{r}} \sigma_r^z$

Columnar vs Plaquette type orders

$$\Psi = |\Psi| e^{i\theta}$$

$\theta = 2\pi m/6$: Columnar three-sublattice order ($m = 0, 1, 2 \dots 6$)

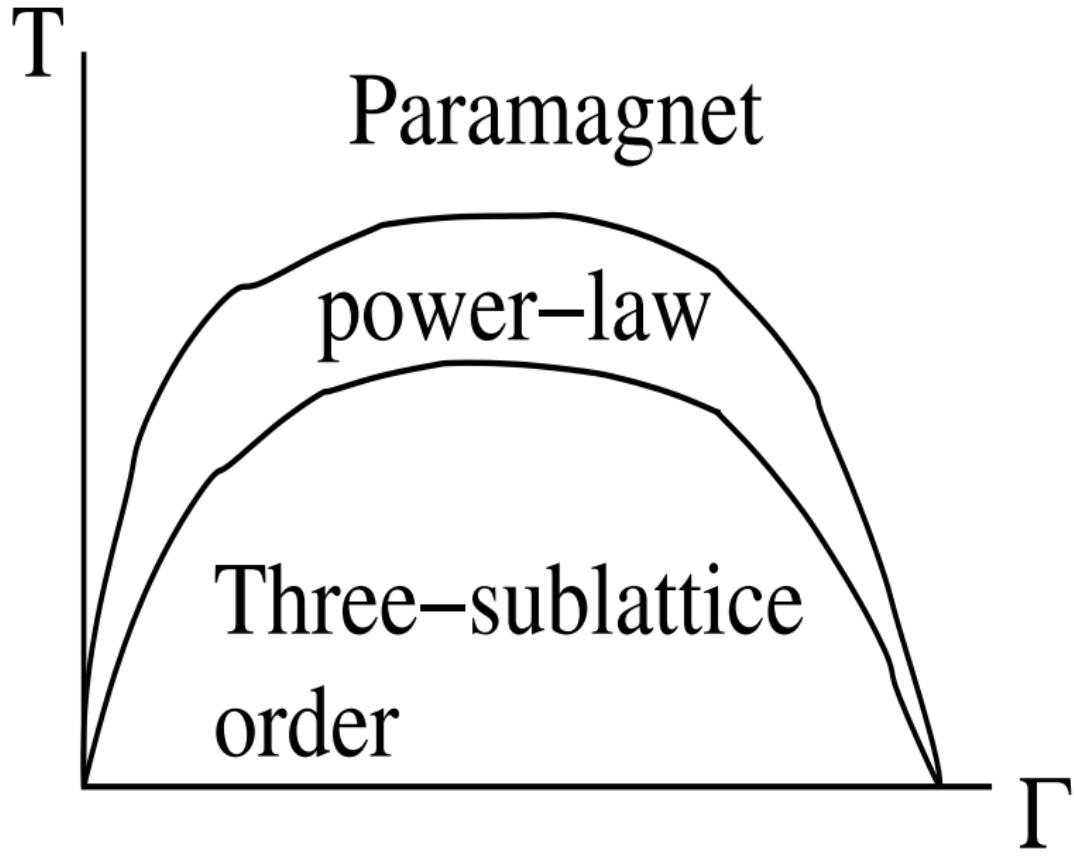
$\theta = (2m + 1)\pi/6$: Plaquette three-sublattice order ($m = 0, 1, 2 \dots 6$)

In ordered state: θ pinned to these values

Columnar phase is ferrimagnetic $m \propto \cos(3\theta)$

In power-law phase: θ has gaussian fluctuations with no pinning

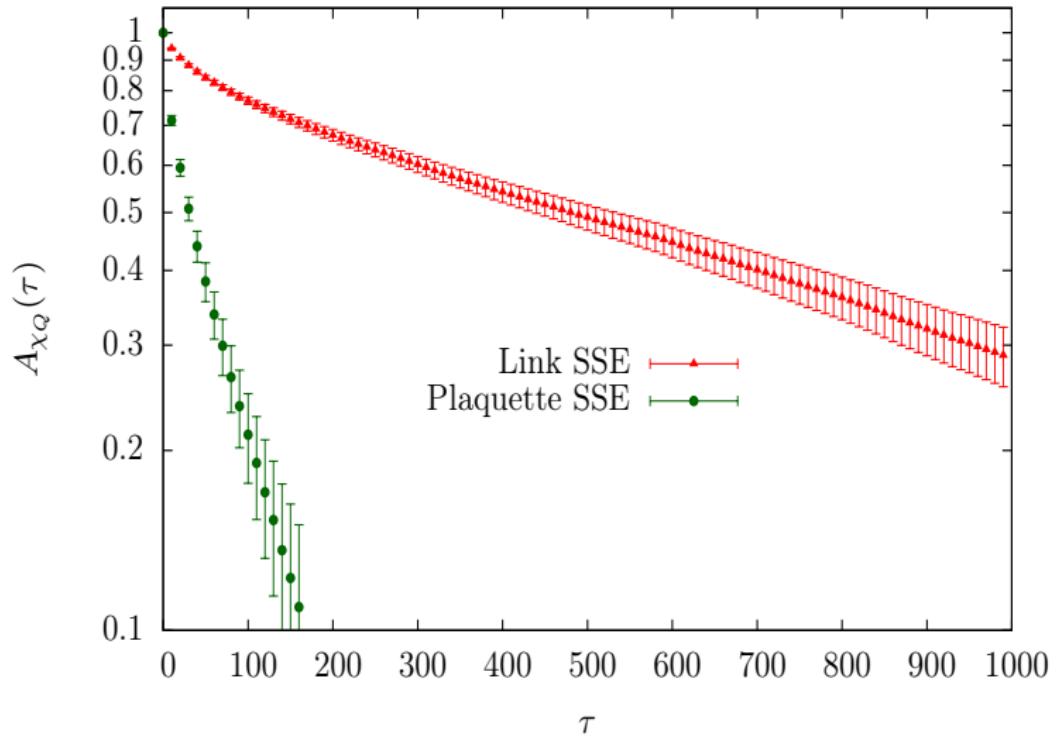
Physics with $J_2 = 0$



(Isakov & Moessner 2003)

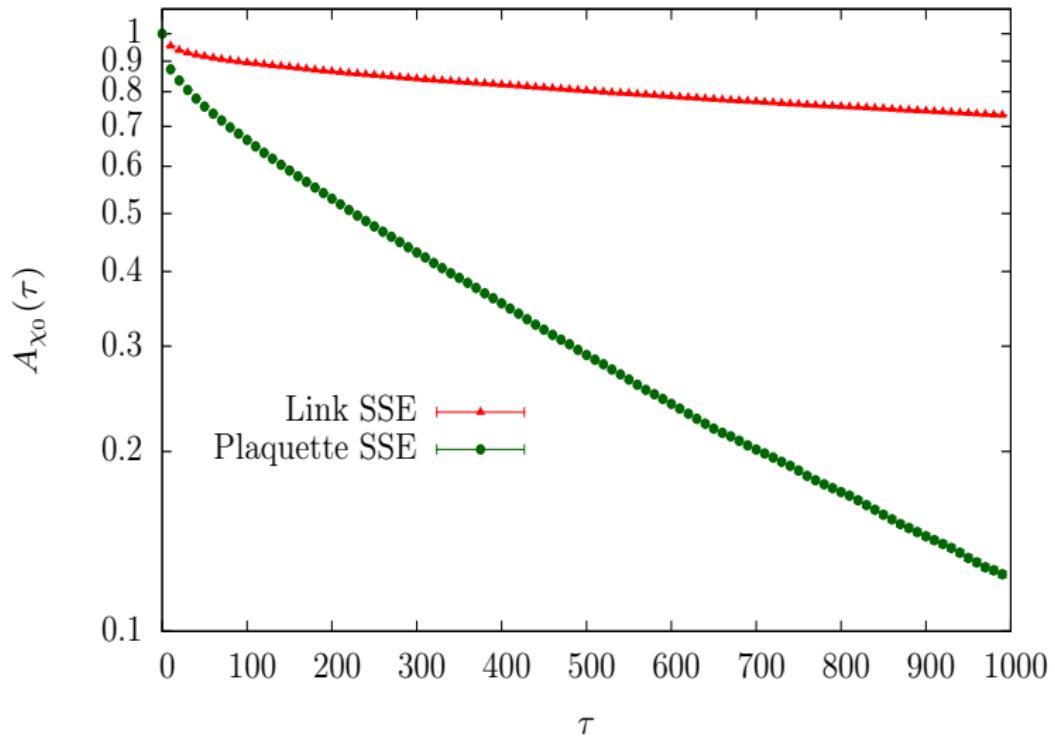
New ideas needed?

$$L = 48, \Gamma = 0.8, T = 0.1, J_1 = 1.0, J_2 = 0.0$$



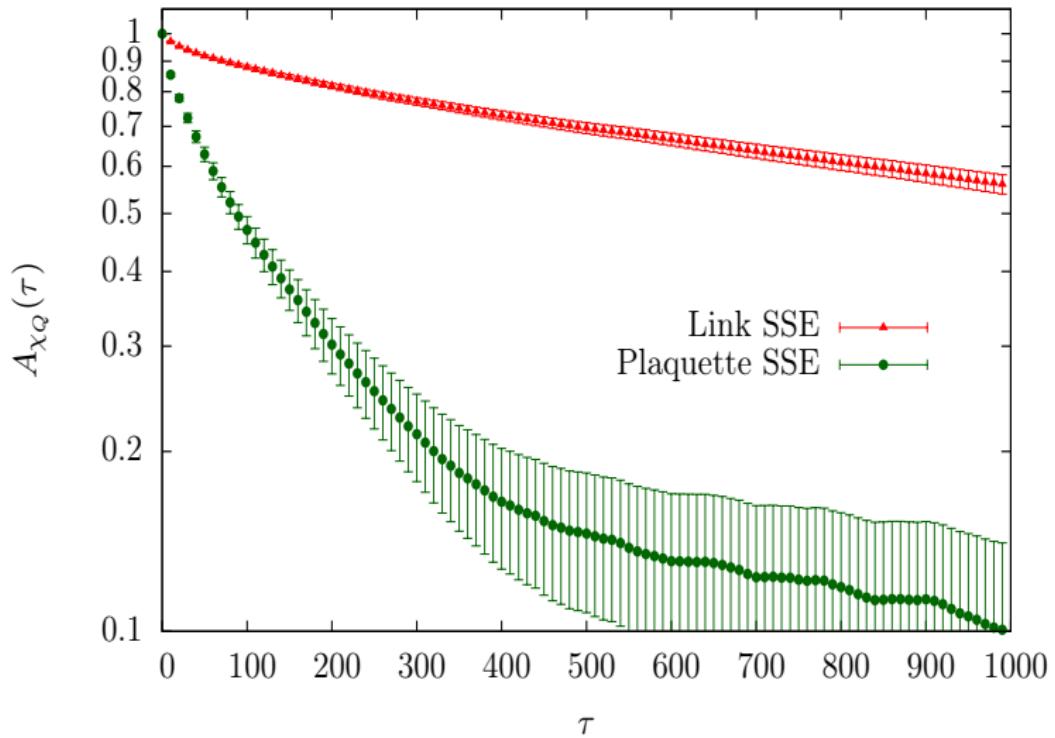
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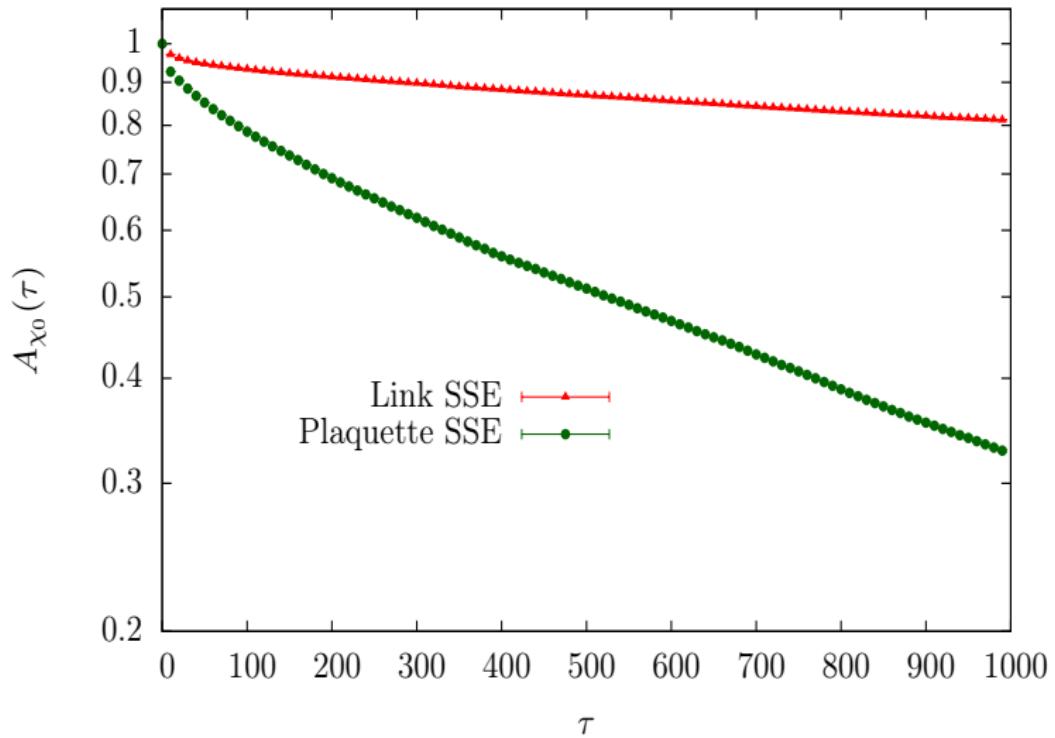
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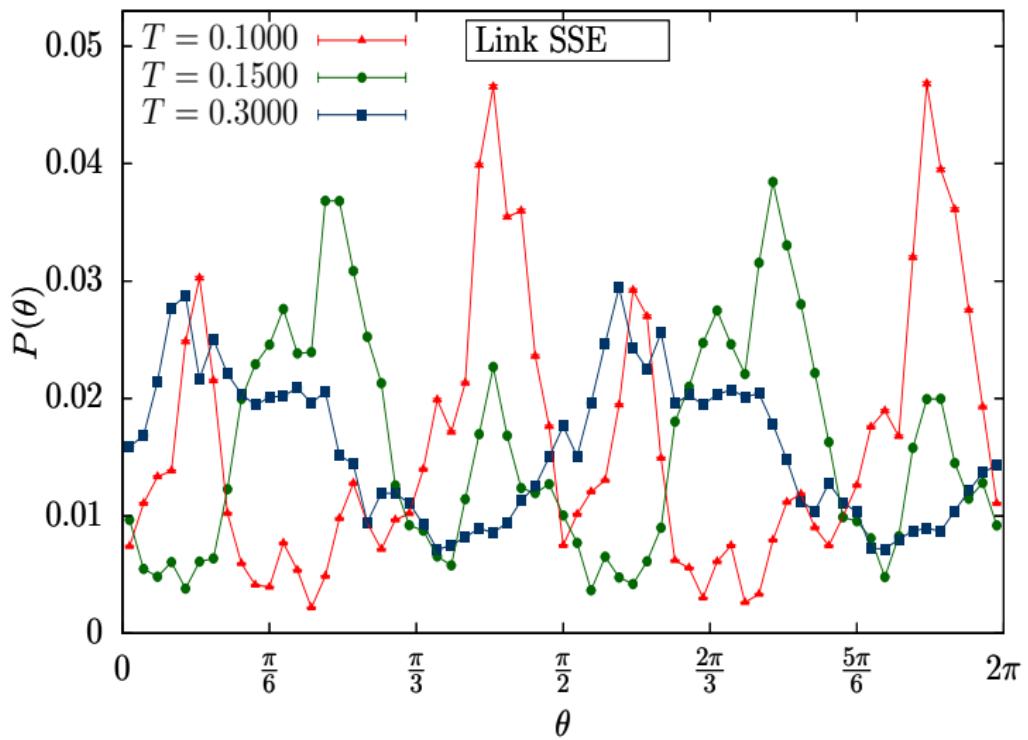
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Approach

$$H_{\text{TFIM}} = \sum_{\Delta} H_{\Delta} + \sum_{\text{link}} H_{\text{link}} + \sum_{\text{sites}} H_{\text{sites}}$$

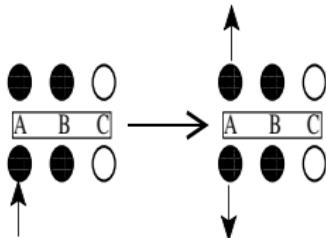
H_{Δ} : Triangle decomposition of all antiferromagnetic couplings

H_{Link} : Bond decomposition of all ferromagnetic couplings

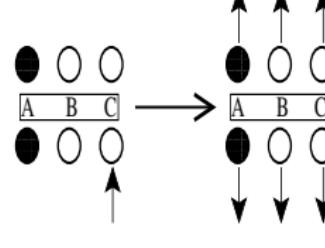
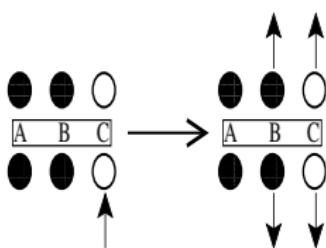
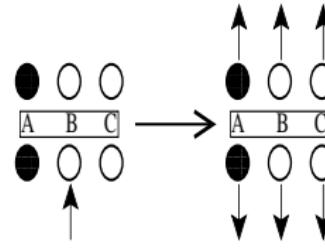
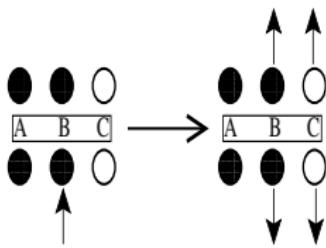
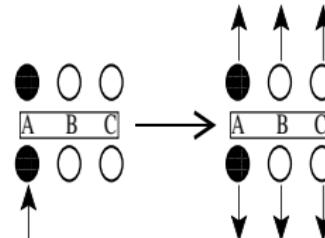
H_{sites} : site decomposition of transverse field term

Quantum-cluster construction for frustrated TFIM

A-Majority site

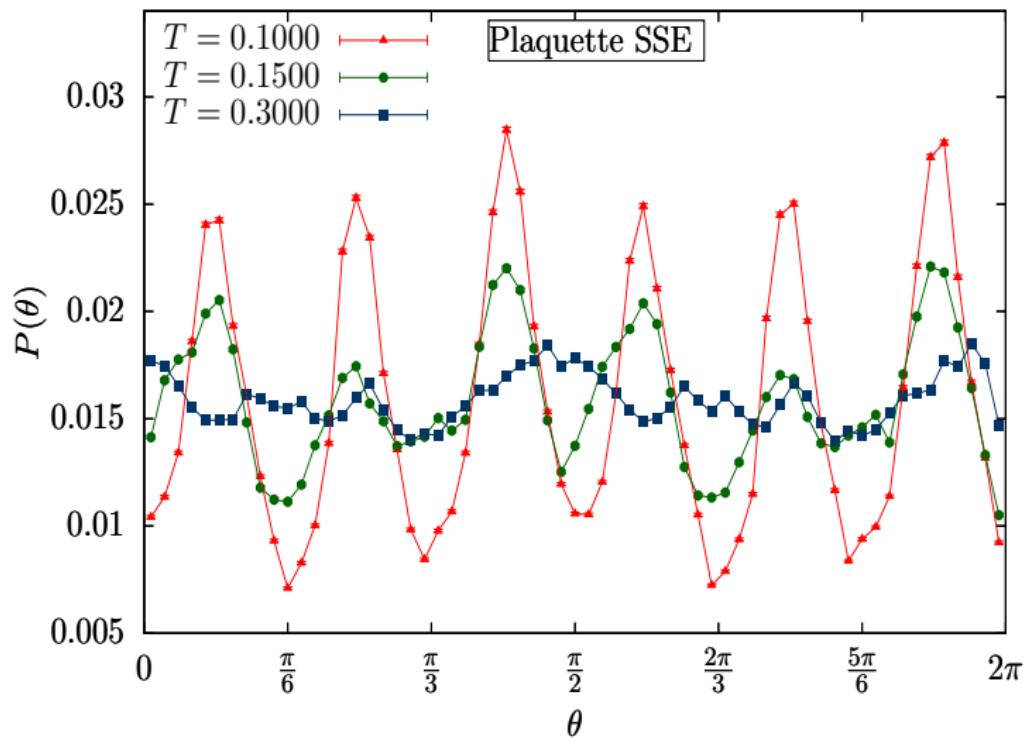


A-Minority Site

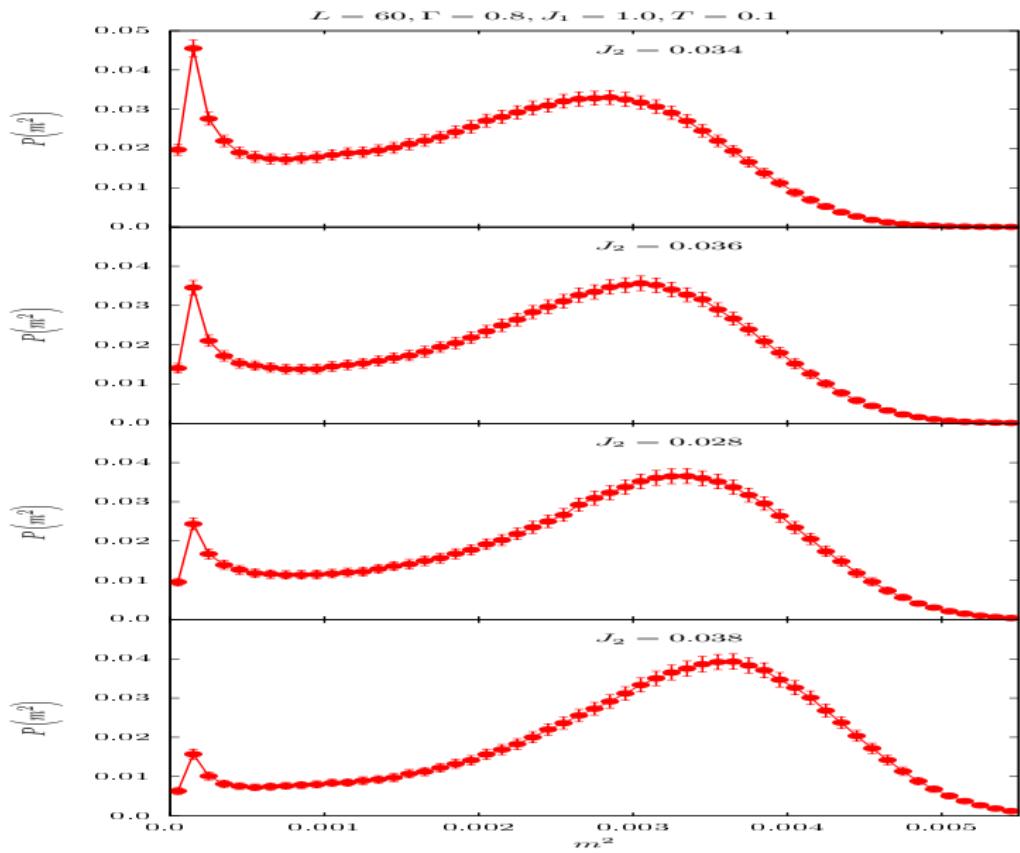


Improvement

$$L = 48, \Gamma = 0.8, J_1 = 1.0, J_2 = 0.0$$

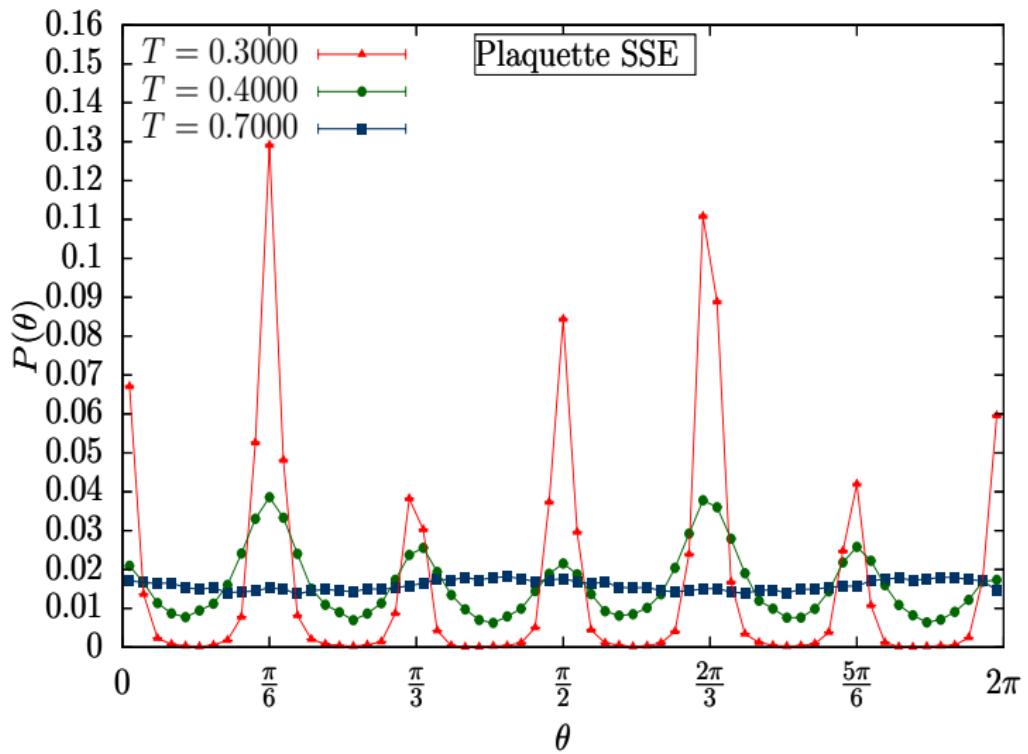


Going into columnar phase



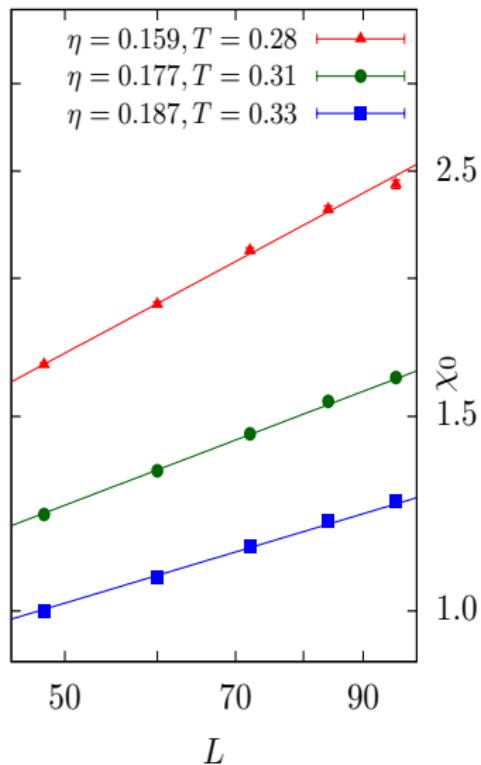
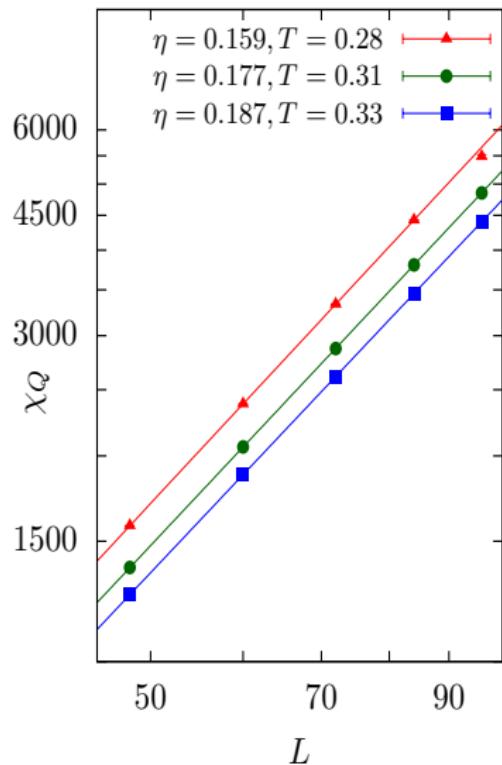
In columnar phase

$L = 72, \Gamma = 0.8, J_1 = 1.0, J_2 = 0.1$

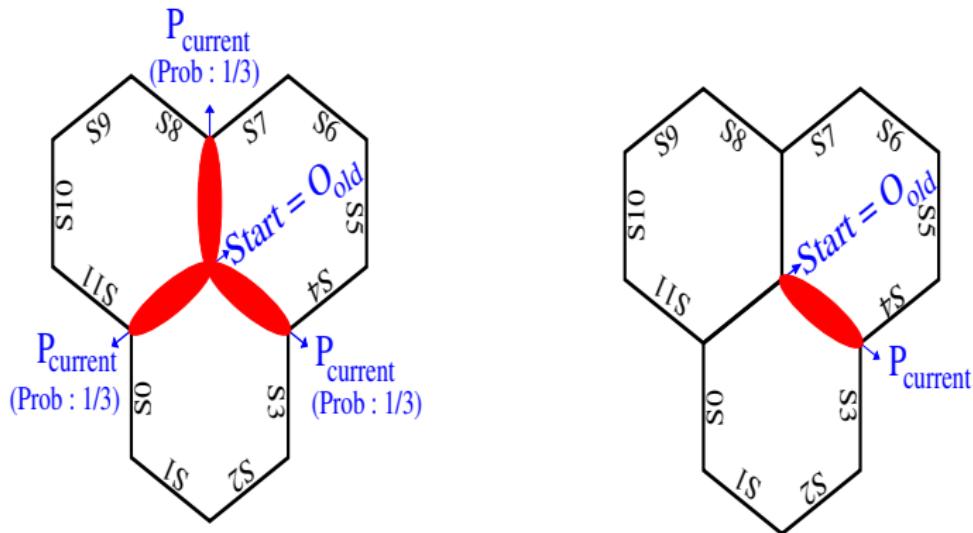


Divergent ferromagnetic susceptibility of antiferromagnet

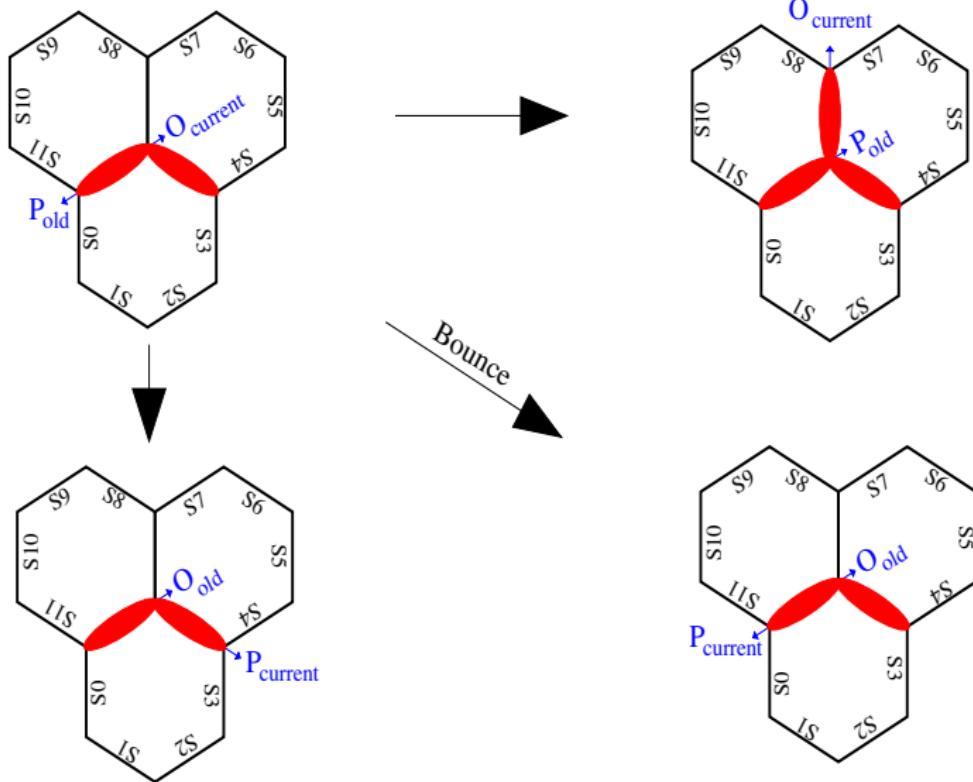
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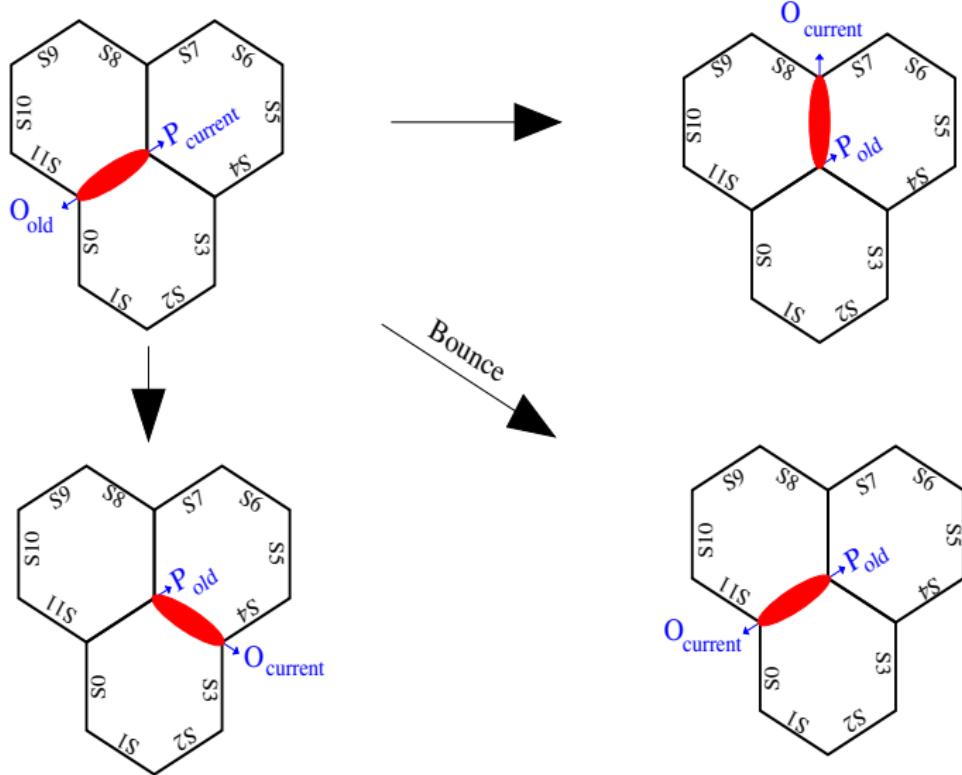
Algorithm for H_{Ising}



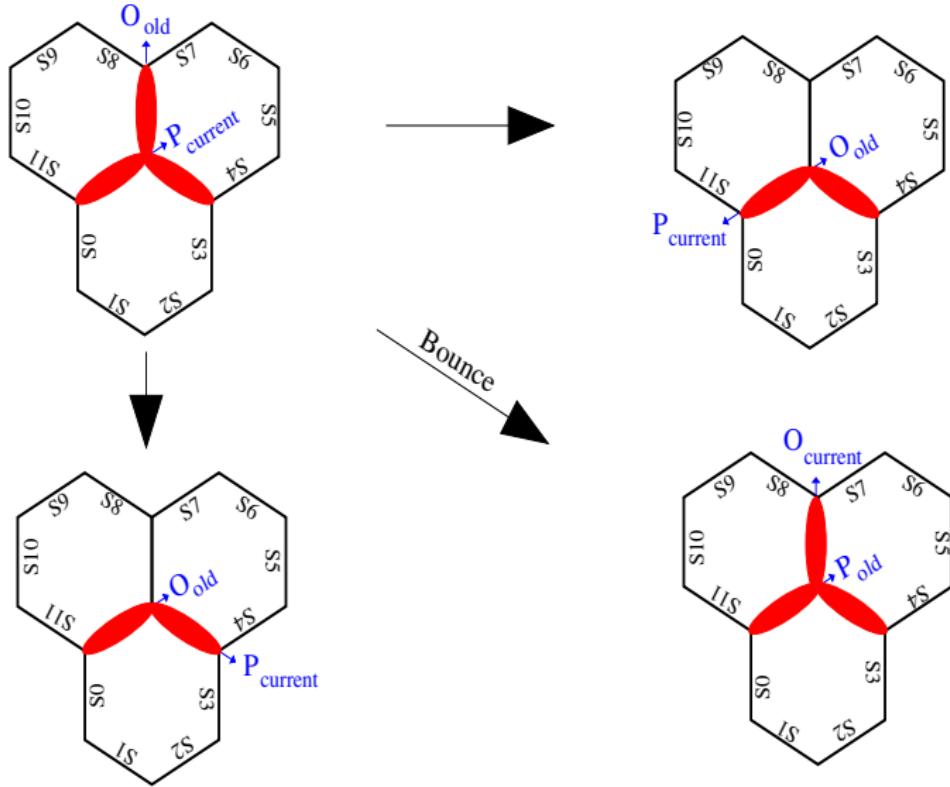
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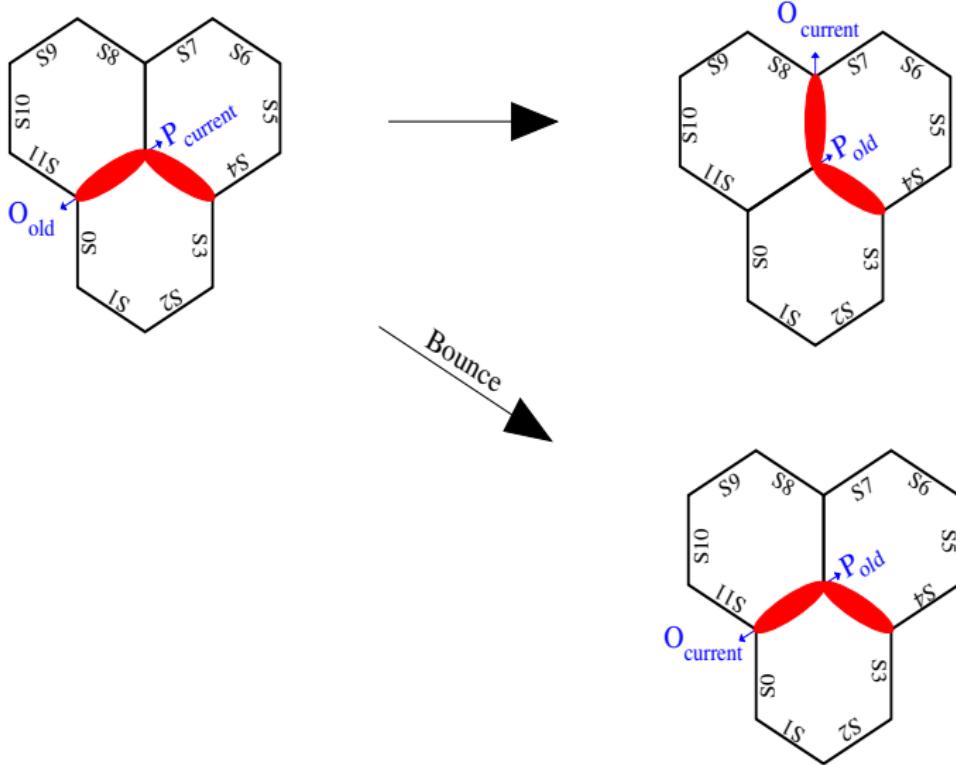
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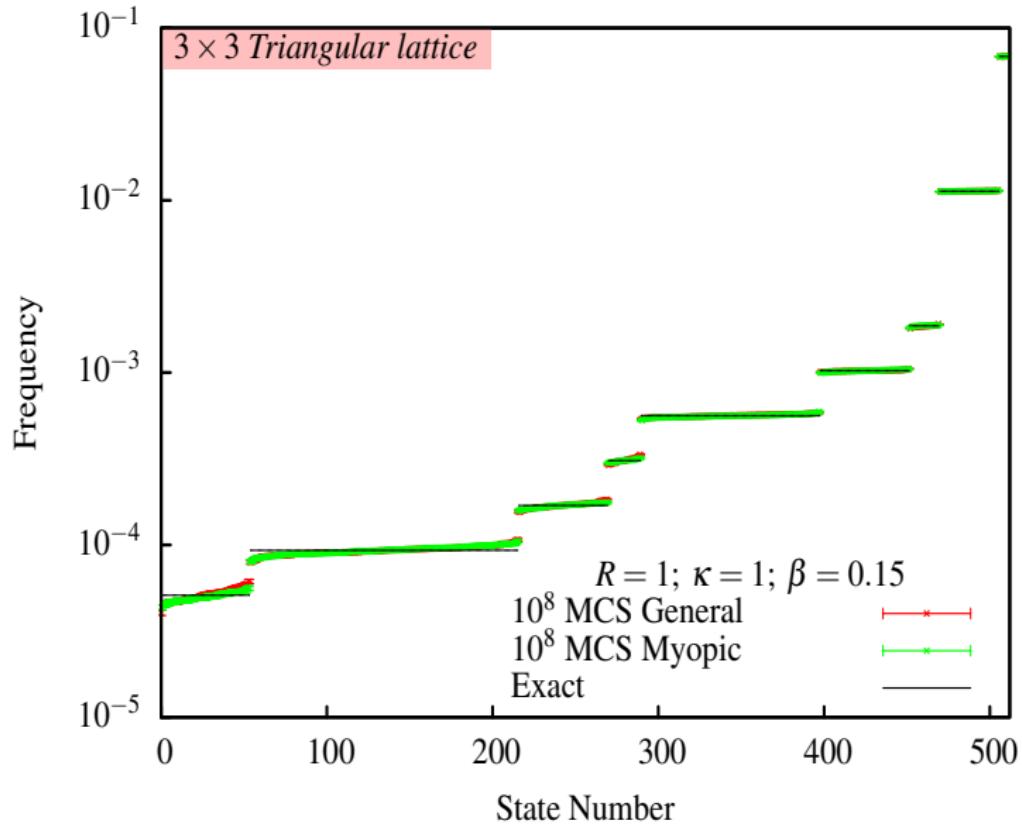
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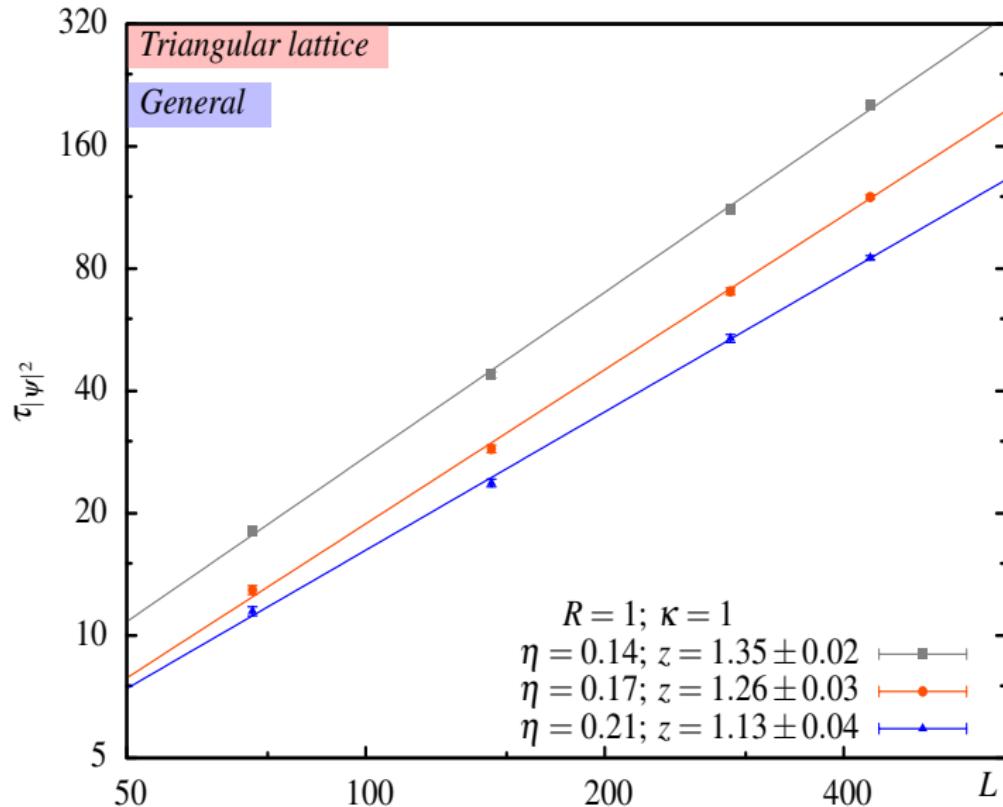
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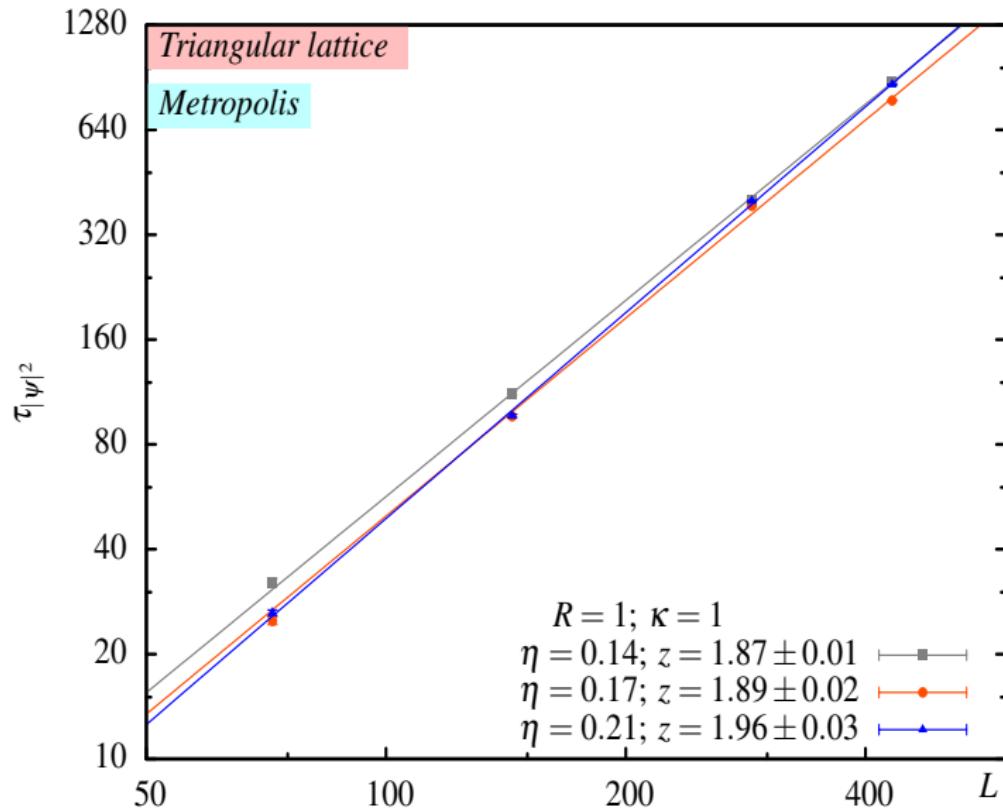
Test against exact enumeration



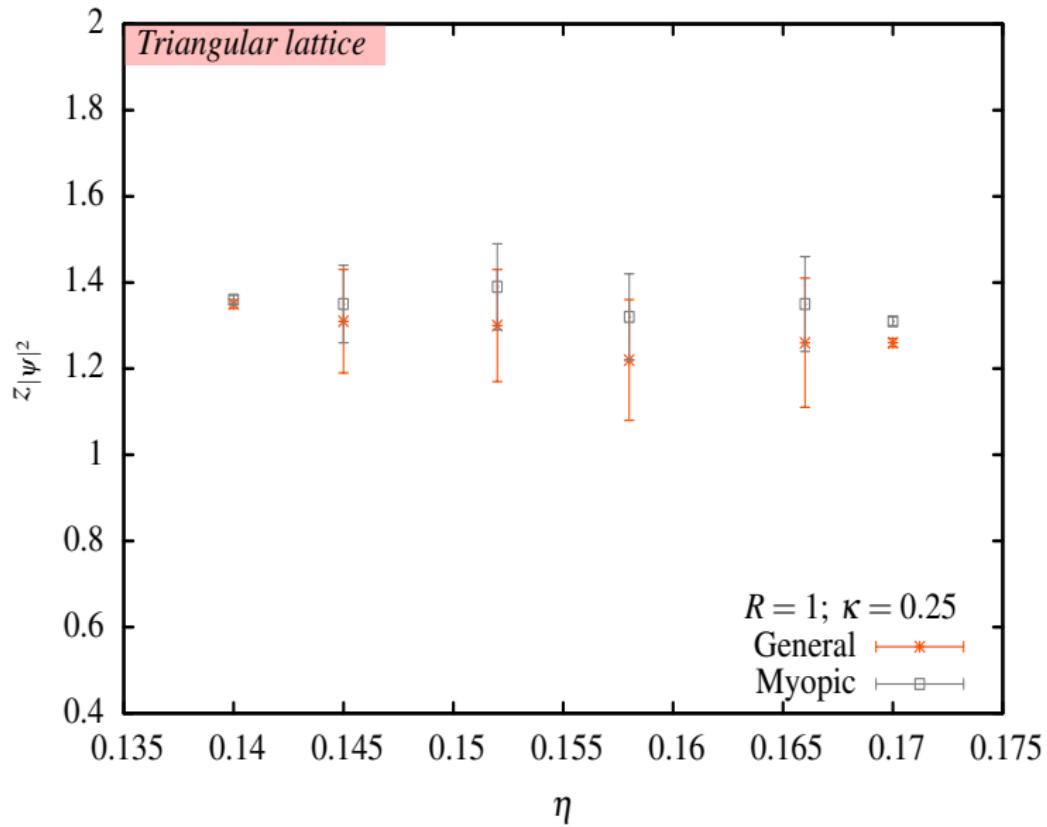
In power-law three-sublattice ordered phase



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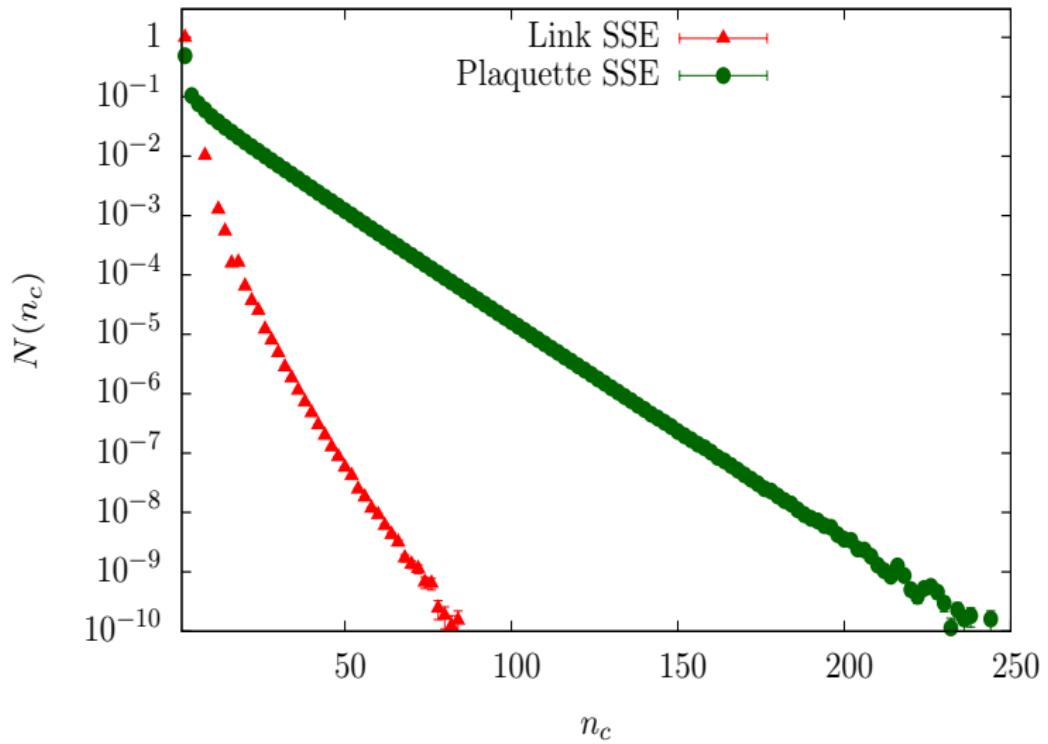


Acknowledgements

- ▶ Collaborators:
Sounak Biswas & Geet Ghanshyam TIFR
- ▶ Computational resources: funded by DTP TIFR

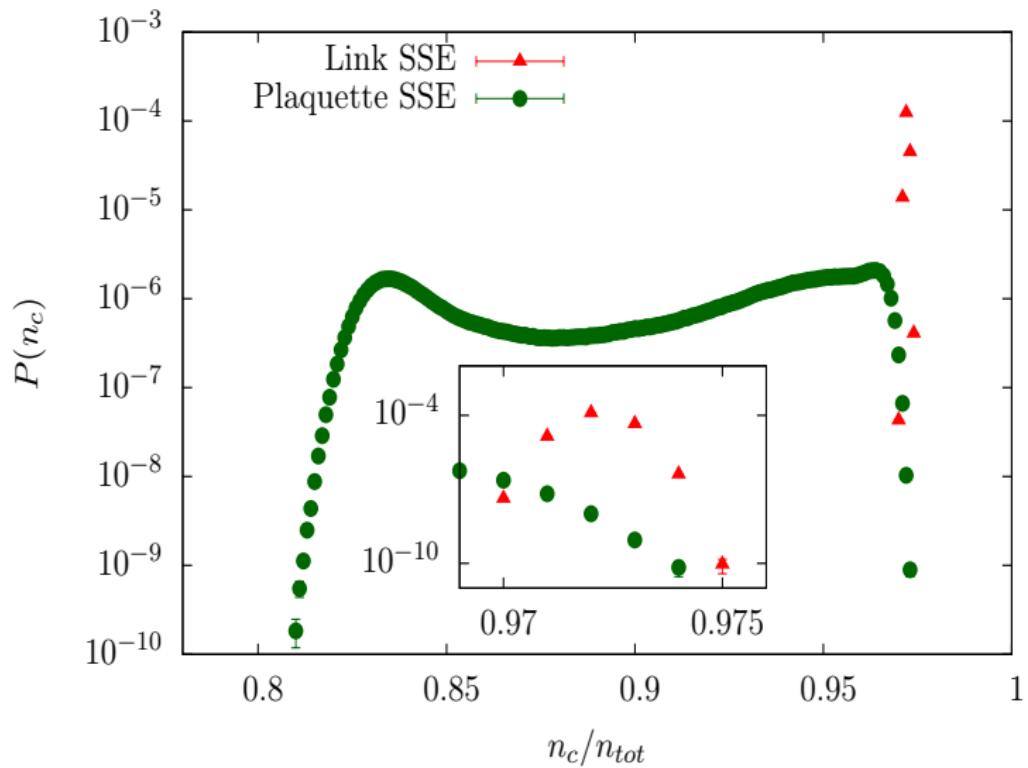
Cluster size distribution

$$L = 48, \Gamma = 0.8, T = 0.1, J_1 = 1.0, J_2 = 0.0$$



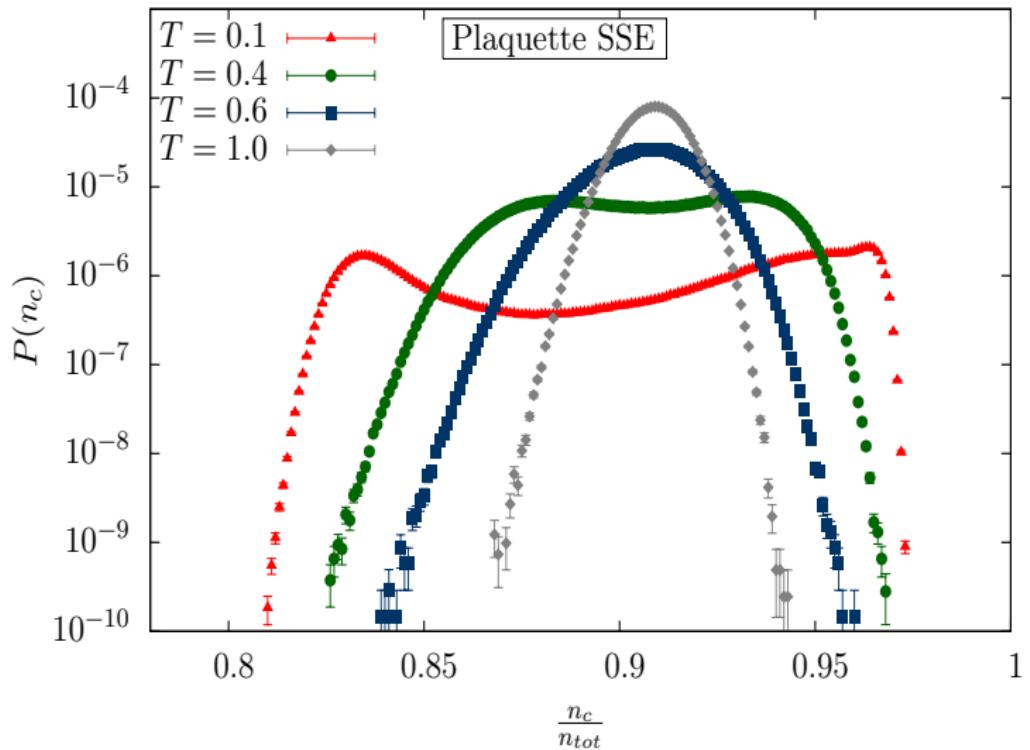
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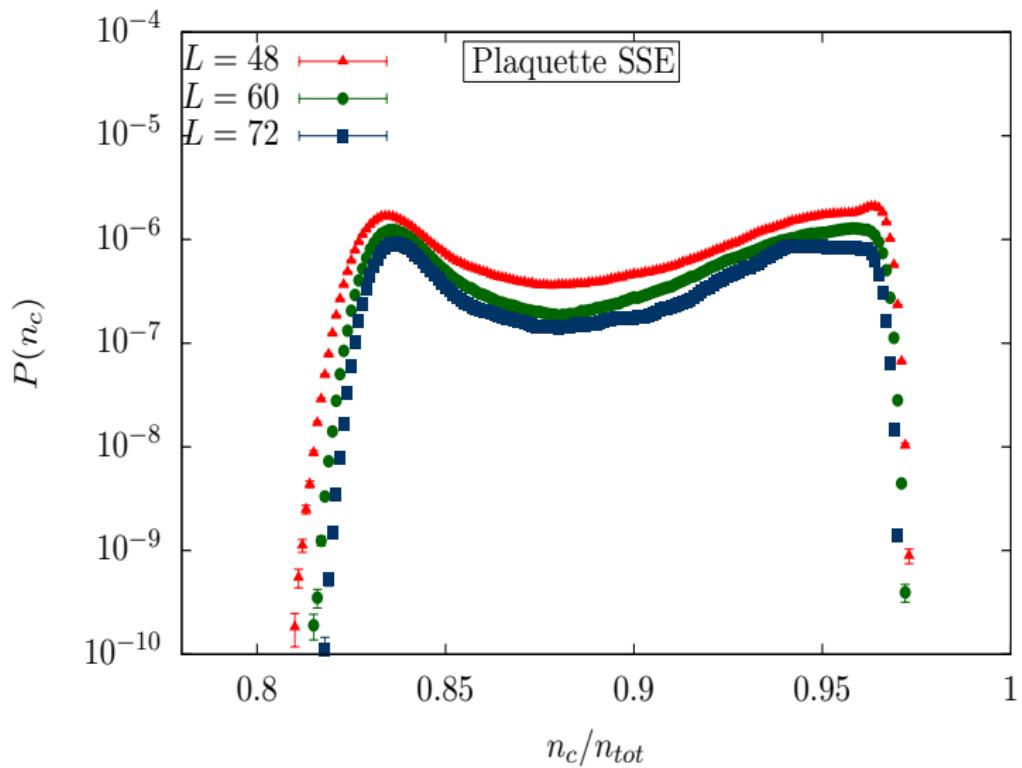
Cluster size distribution

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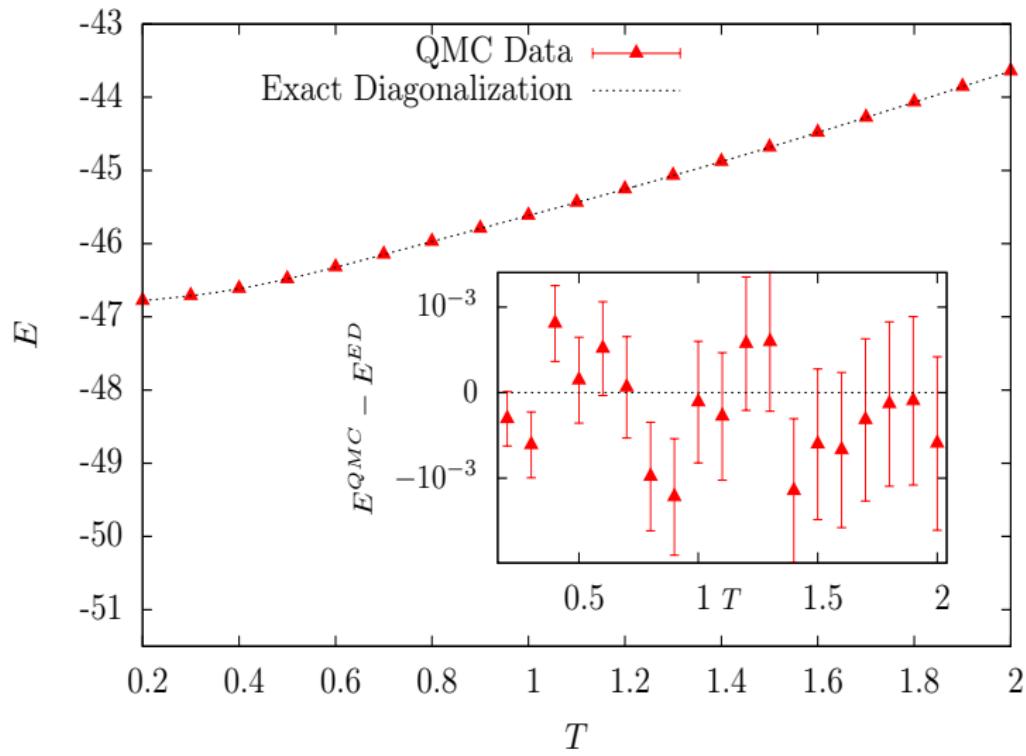
Cluster size distribution

$$T = 1.0, \Gamma = 0.8, J_1 = 1.0, J_2 = 0.0$$



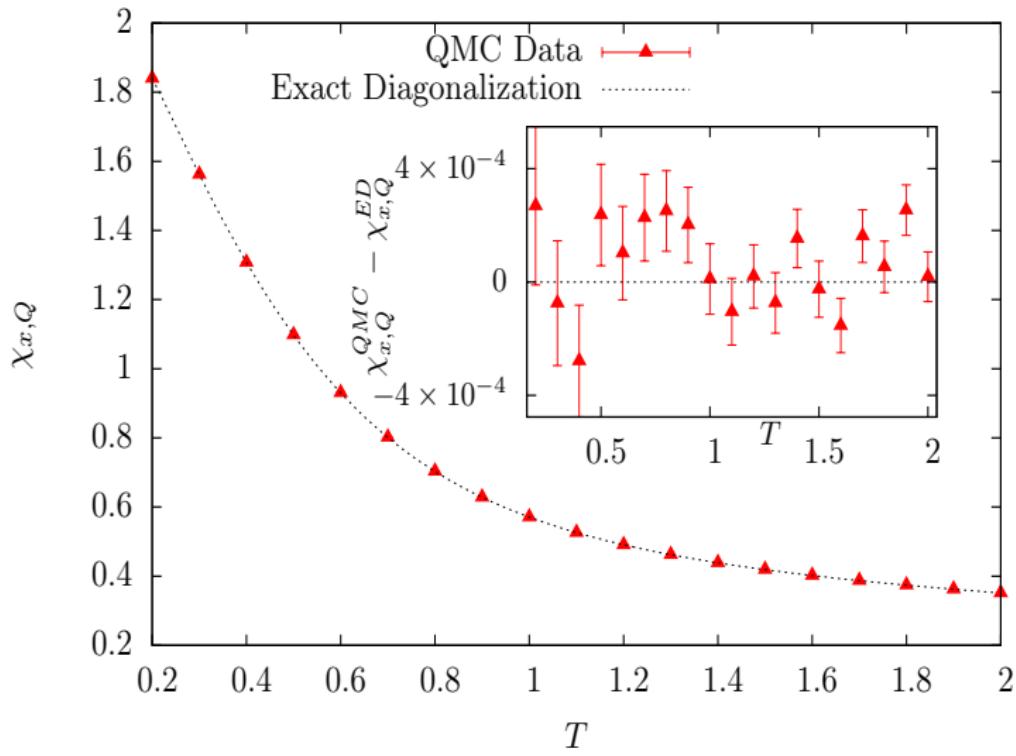
ED tests

$$L = 3, J_1 = 1.0, J_2 = 0.0$$



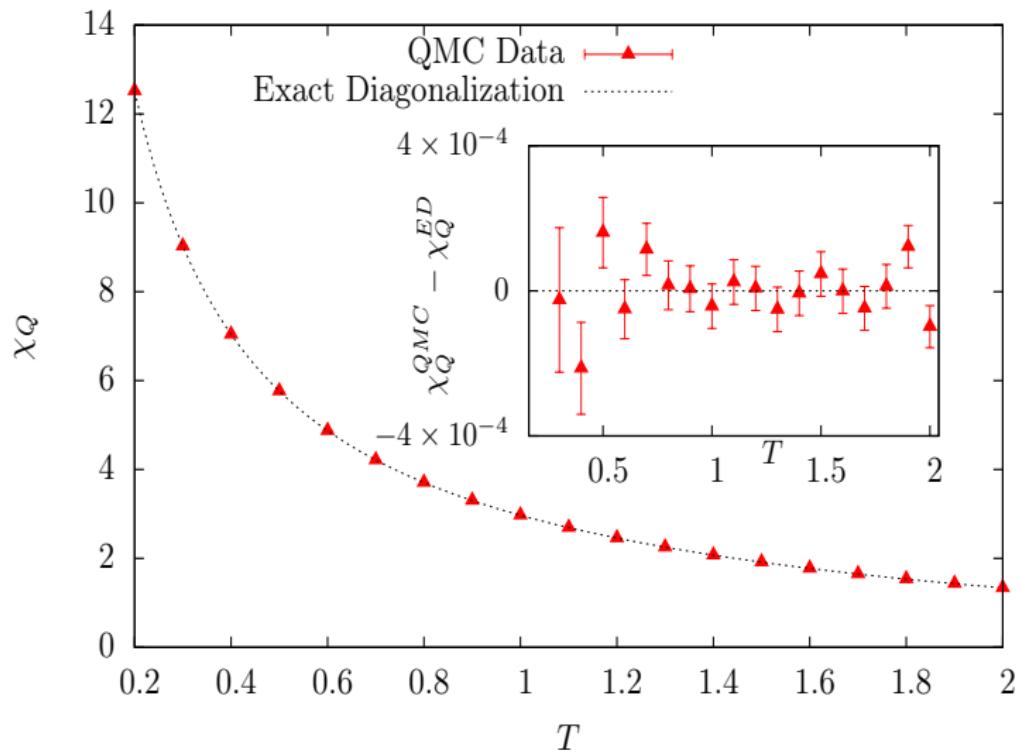
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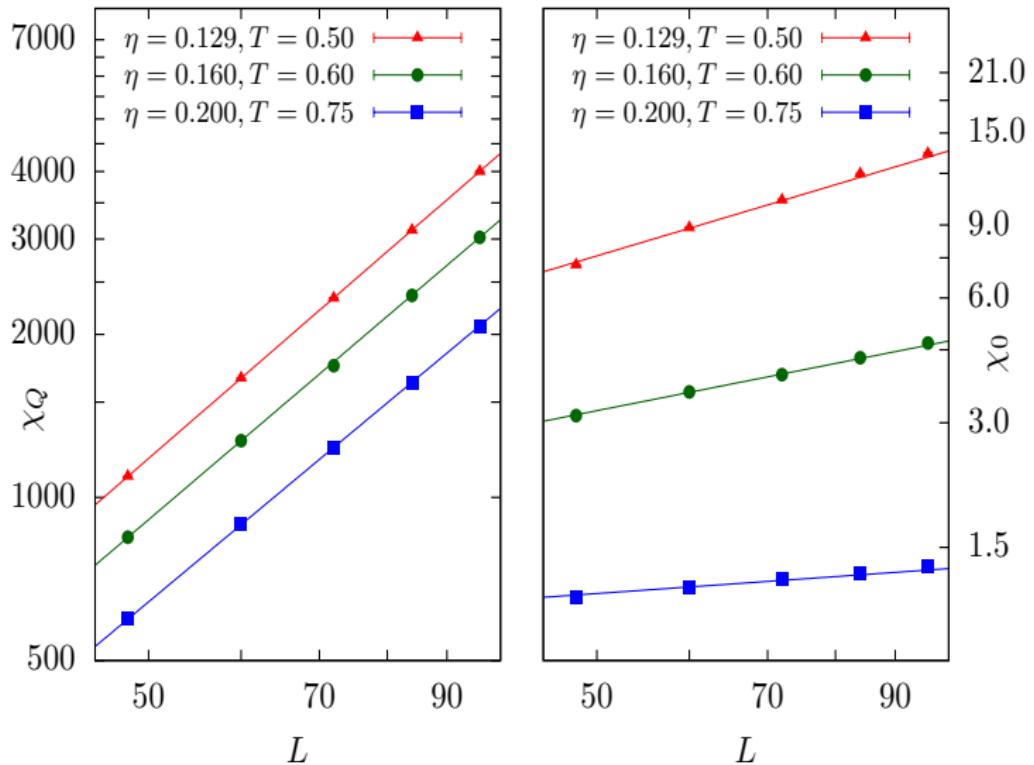
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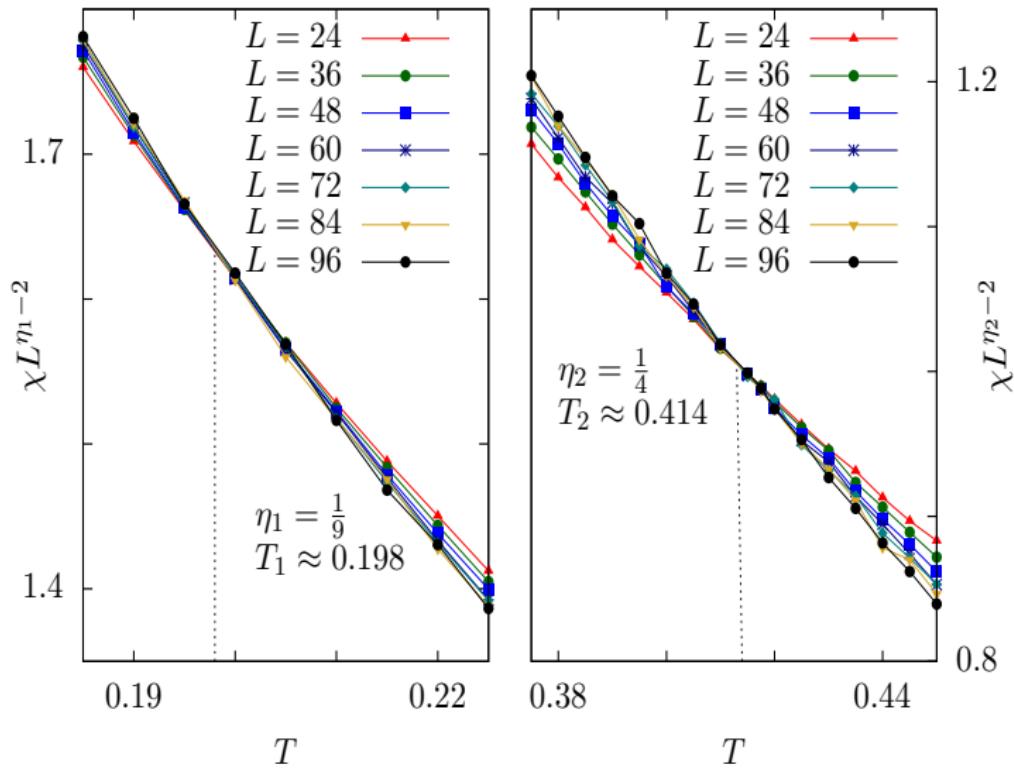
Thermodynamic signature of melting of columnar order

$$\Gamma = 0.8, J_1 = 1.0, J_2 = 0.1$$



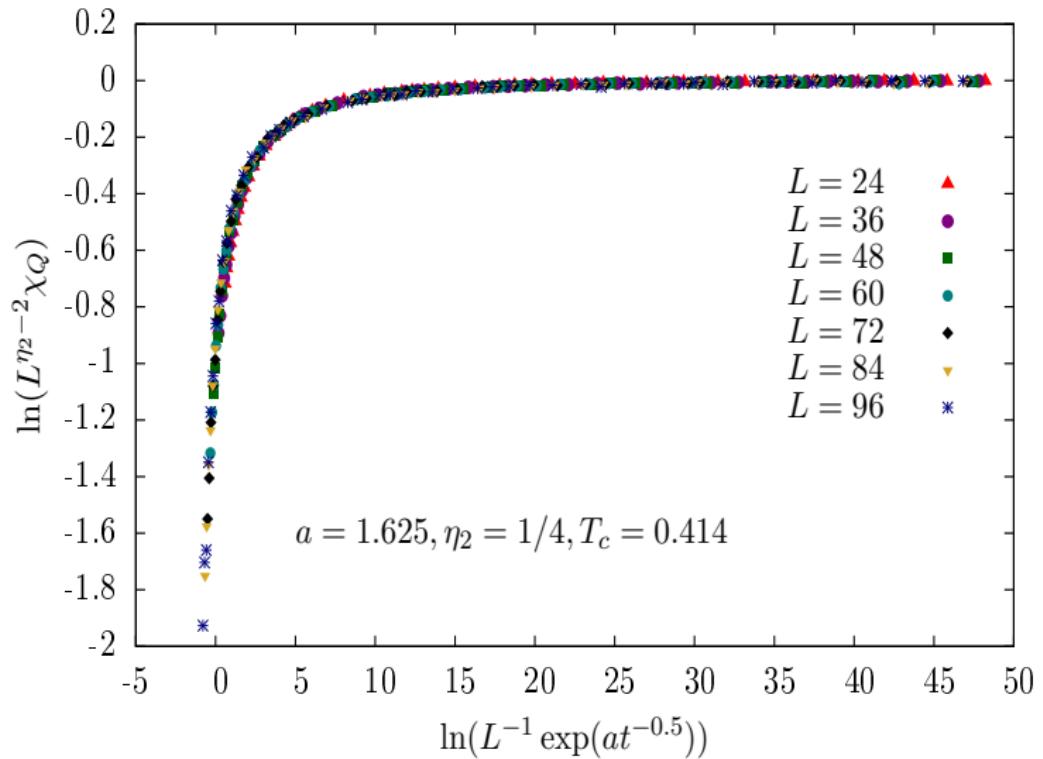
Scaling for two step transition

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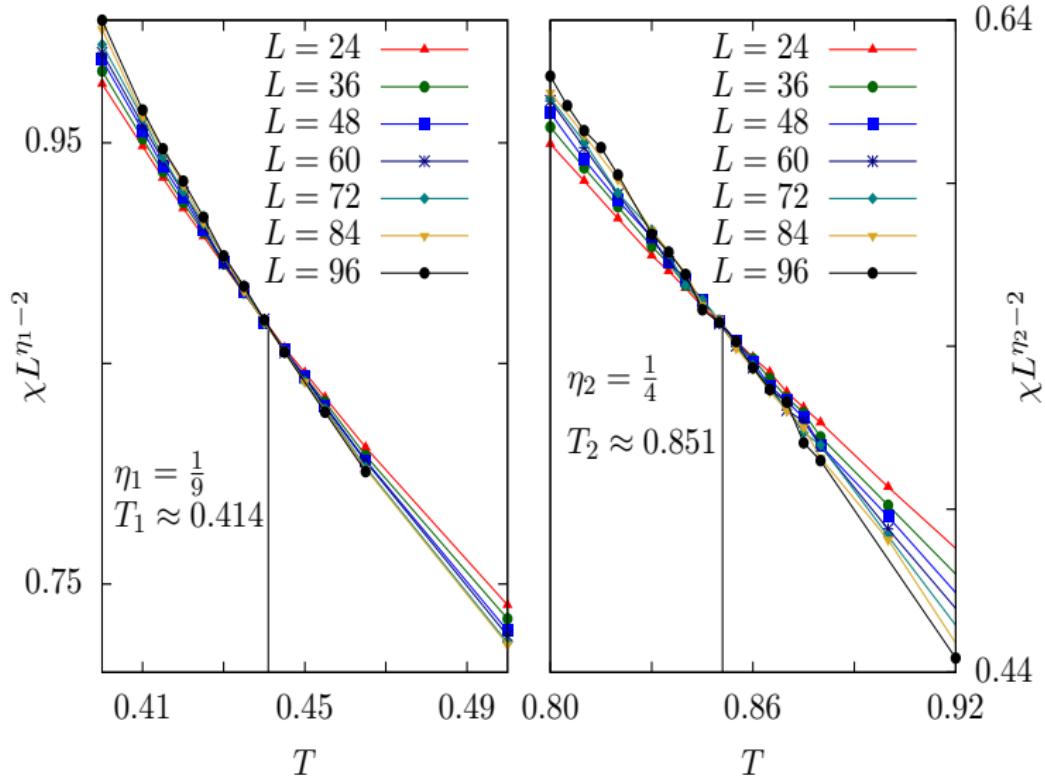
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