

# QUARKONIA PRODUCTION AT RHIC

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Hot and Dense Matter in RHIC LHC era  
Mumbai Fev/2008

# The Importance of the Quarkonia

2285 papers with J/ $\psi$  in the title according to SPIRES

	Mass	radius
$\pi$	0.14 GeV	0.06fm
p	0.94 GeV	0.87fm
$\psi'$	3.68 GeV	0.90fm
$X_c$	3.53 GeV	0.72fm
J/ $\psi$	3.1 GeV	0.50fm
$\Upsilon$	9.5 GeV	0.28fm

$$\Lambda_{\text{QCD}} \sim 200 \text{ MeV}$$

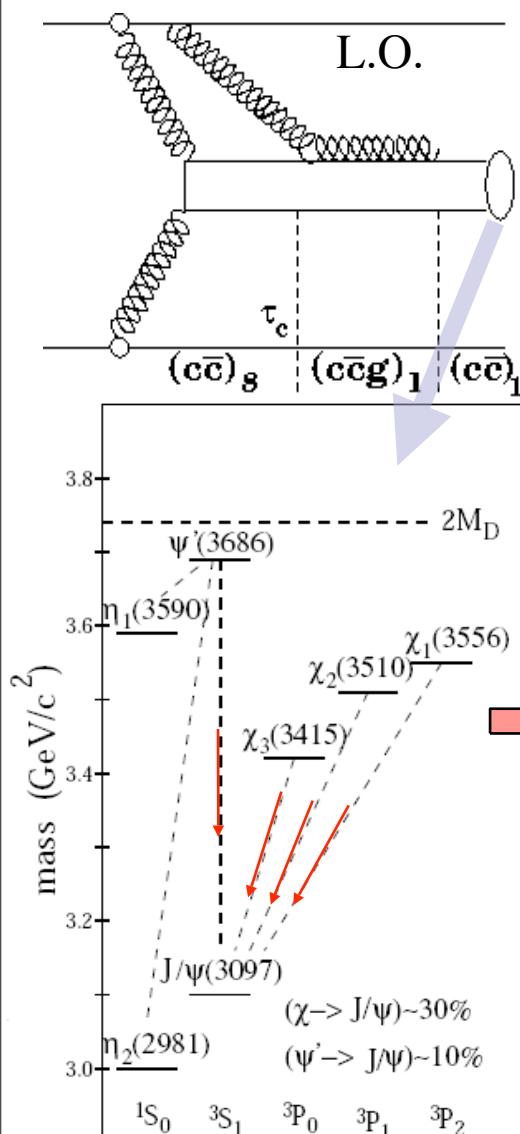
$$\alpha_s(M_Q) \ll 1$$

- weak coupling with light mesons
- small size
- strong binding

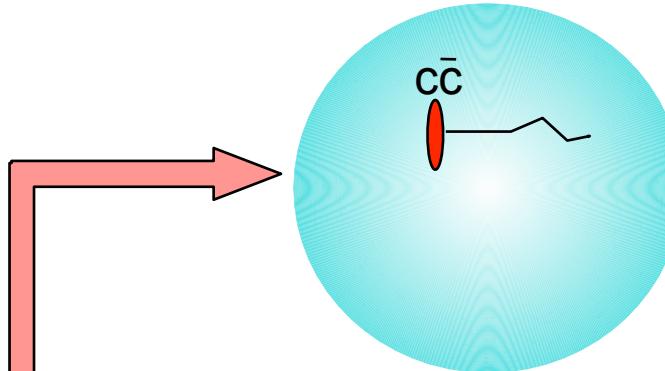
All these features make quarkonia an excellent QCD probe.

The complete quarkonia dissociation could serve as a QGP thermometer.

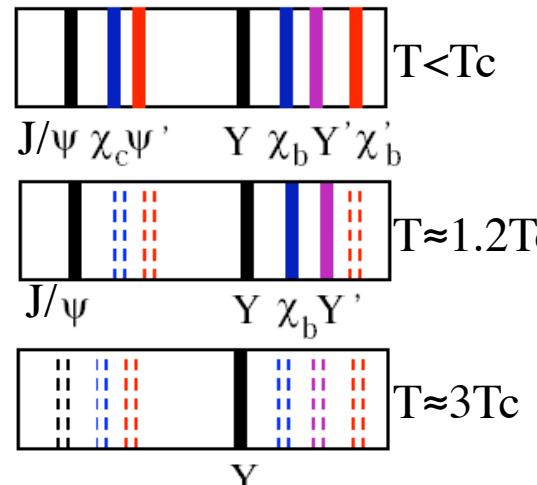
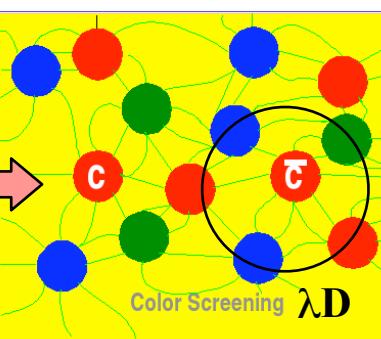
# Quarkonia Production/Suppression



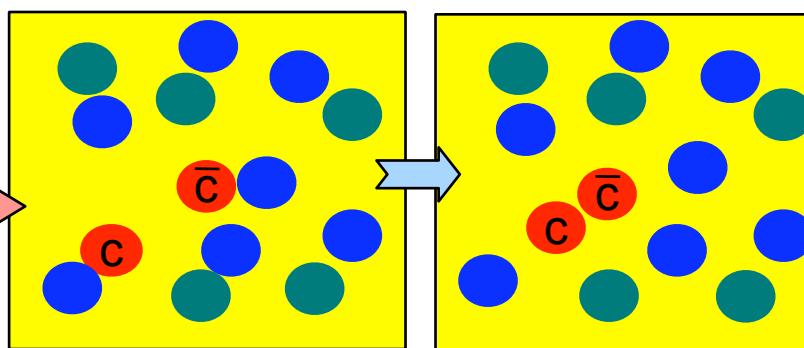
Production



- ▶ Nuclear Absorbtion
- ▶ Parton Distribution Modifications
- ▶ Comovers

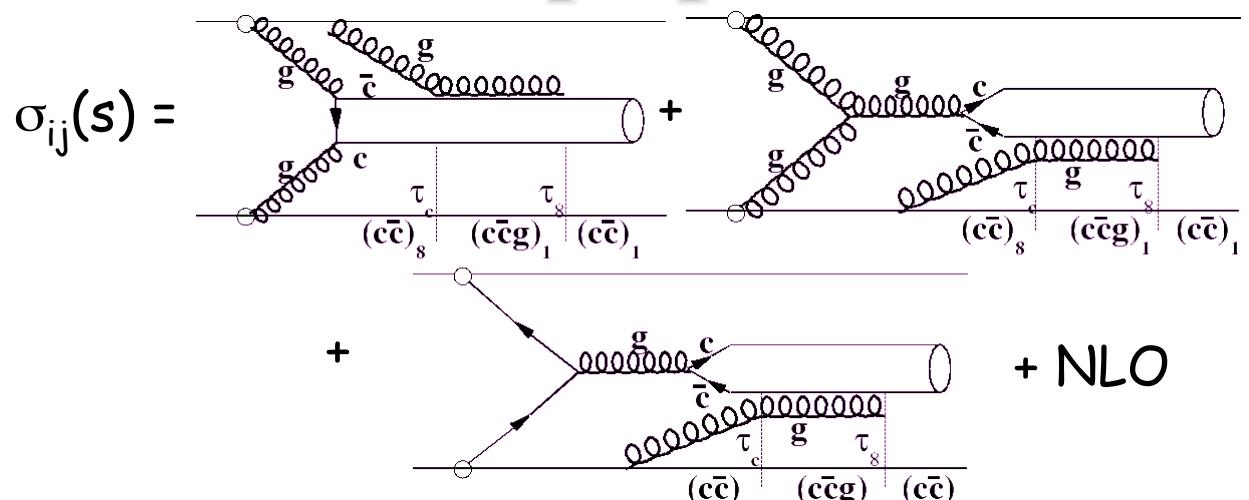
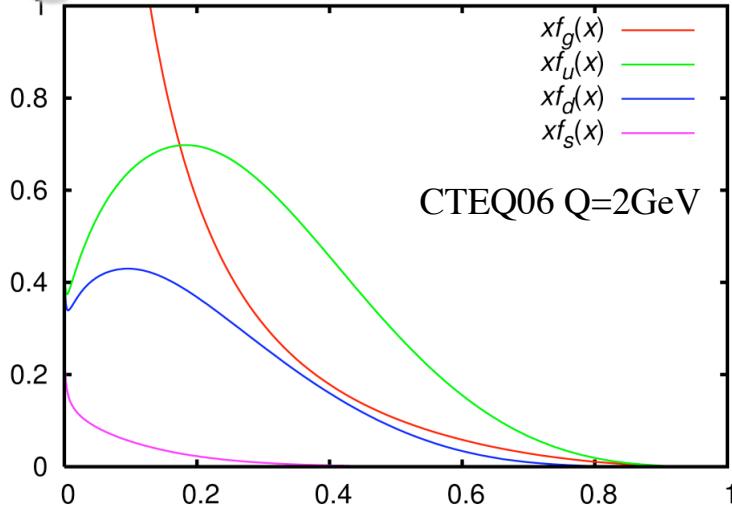


- ▶ Dissociation in deconfined medium



- ▶ Recombination or Coalescence

# Quarkonia Production in p+p collisions



$$\sigma(S_{NN}, Q^2) = \sum_{i,j=\bar{q},q,g} \int_0^1 dx_1 \int_0^1 dx_2 x f_i(x_1, Q^2) x f_j(x_2, Q^2) \sigma_{ij}(\hat{s}, Q^2) \delta(\hat{s} - x_1 x_2 S_{NN})$$

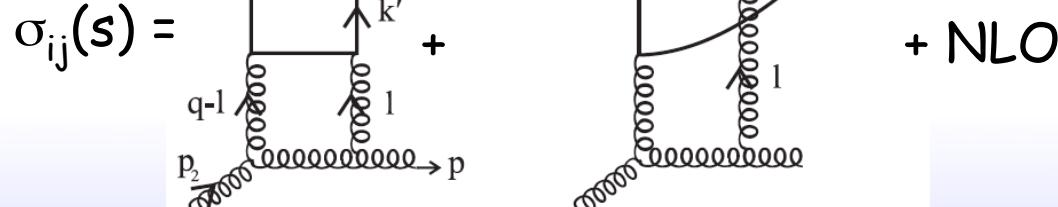
Color neutralization is a non-perturbative process. Methods to handle it :

Non-relativistic QCD (NRQCD) :  $\sigma(\hat{s}, Q^2) = \sum_n \frac{C_n(\Lambda)}{m^{d_n-4}} \langle \mathcal{O}_n^H(\Lambda) \rangle$  Formation of  $(ccg)_8$   
 PLB167 (1986) 437, PRD43 (1991) 196,  
 PRD51 (1995) 1125.

Color Evaporation Model (CEM) :  $\sigma(S_{NN}, Q^2) = F_H \int_{2m_c}^{2m_D} d\hat{s} \frac{d\sigma(\hat{s}, Q^2)}{d\hat{s}}$  data.  
 Int. J. Mod. Phys. A 10 (1995) 3043

pQCD w/ 3-gluon fusion :

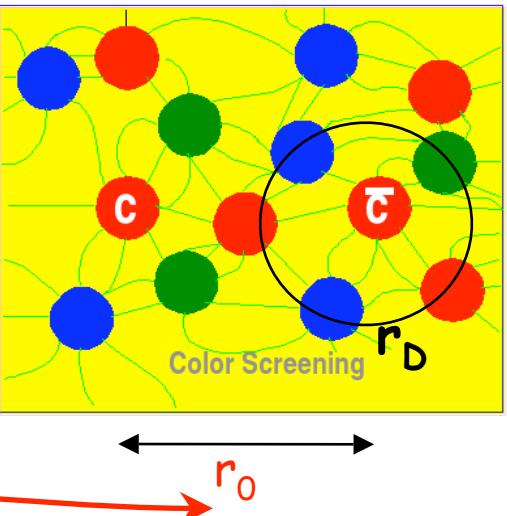
Eur. Phys. J. C 39, 163–171 (2005)



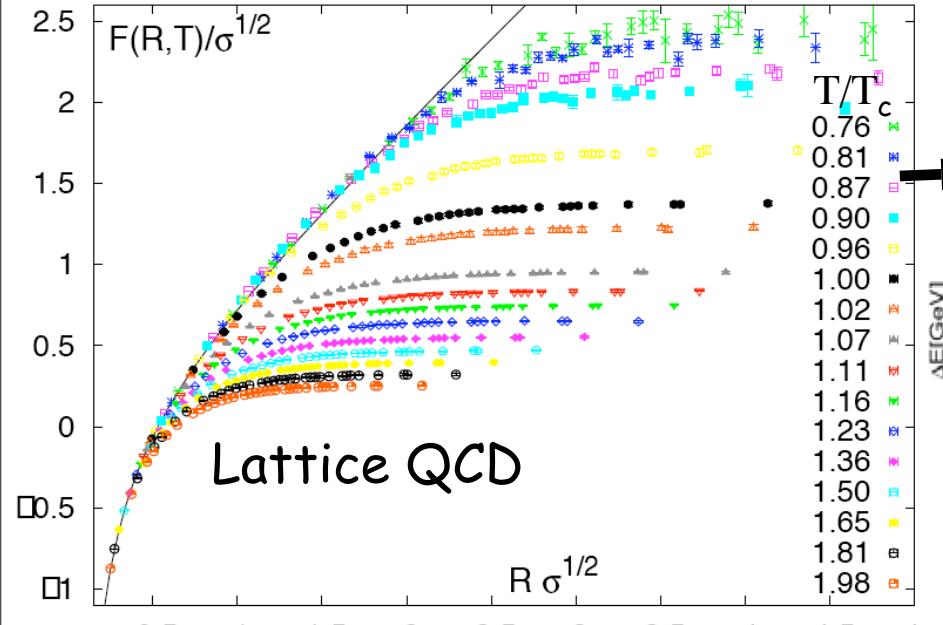
# Screening in Colored Matter

Matsui & Satz, PLB178 ('86),  
Miyamura et al., PRL57 ('86)

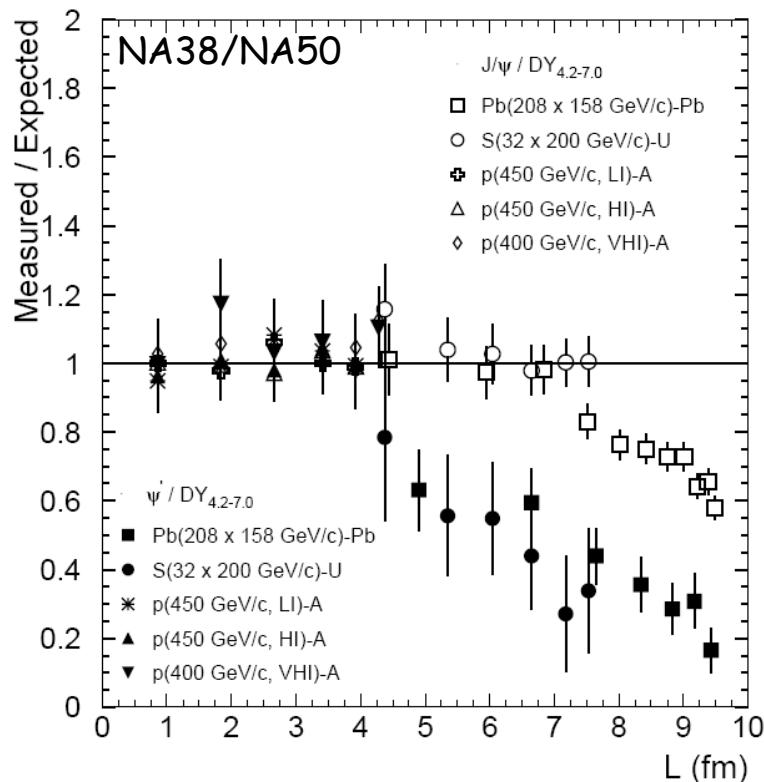
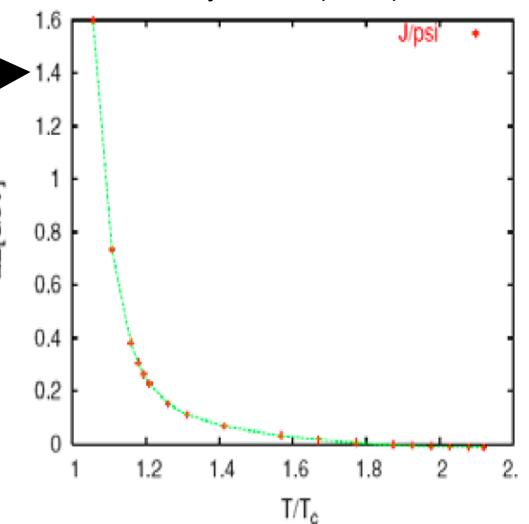
state	$J/\psi$	$\chi_c$	$\psi'$
mass [GeV]	3.10	3.53	3.68
$\Delta E$ [GeV]	0.64	0.20	0.05
$\Delta M$ [GeV]	0.02	-0.03	0.03
$r_0$ [fm]	0.50	0.72	0.90



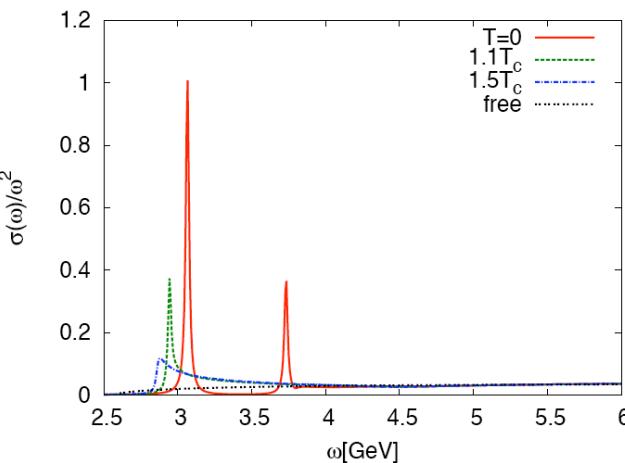
O. Kaczmarek, F. Zantow, Phys. Rev. D71 (2005) 114510.



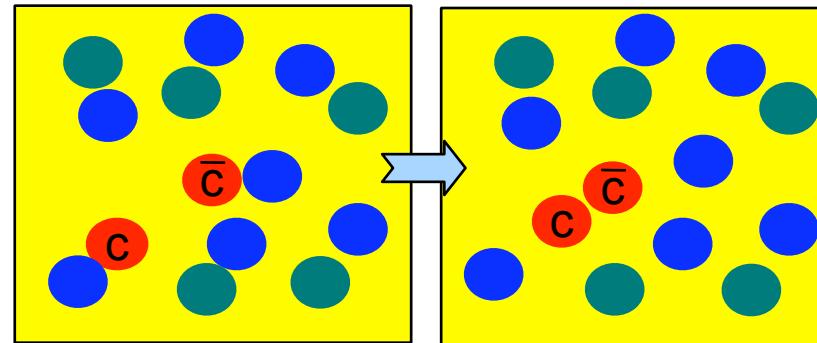
H.Satz, J. Phys. G32 (2006) R25



A.Mocsy, P.Petreczky, PRL99 (2007) 211602



# Recombination and Coalescence



$$\sigma_{c\bar{c}} = \sigma(p + p \rightarrow c\bar{c} + X) = \begin{cases} 567 \pm 57(\text{stat}) \pm 224(\text{sys}) \mu\text{b} @ \text{PHENIX} [\text{PRL97 (2006) 252002}] \\ 256 +400-146 \mu\text{b} \text{ by FONLL} \end{cases}$$

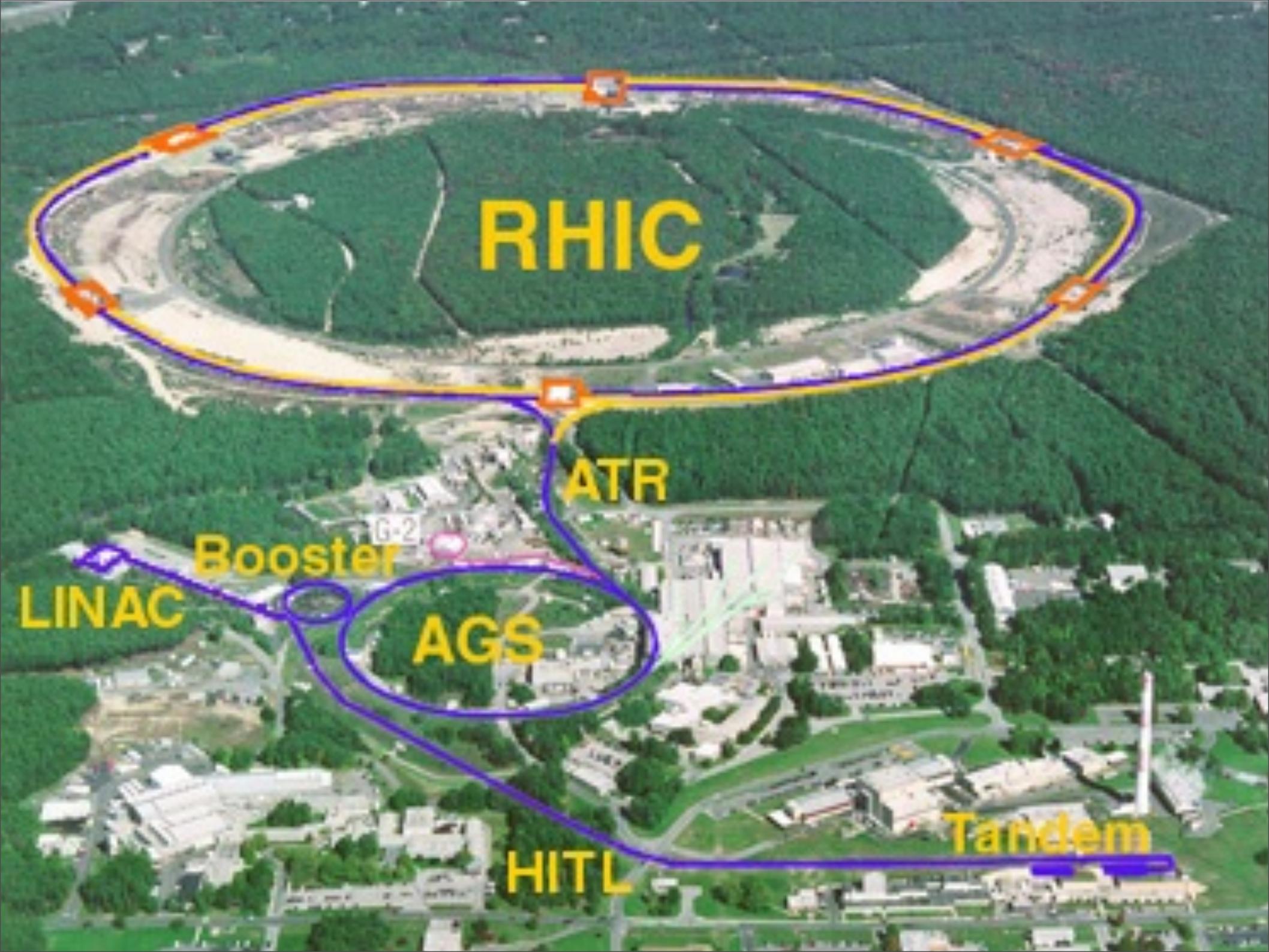
The charm cross section scales with binary scaling [PRL94 (2005) 082301]. Thus, the average number of  $c\bar{c}$  pairs is:

$$\langle N_{c\bar{c}} \rangle = N_{coll} \frac{\sigma_{c\bar{c}}^{p+p}}{\sigma_{inel}^{p+p}}$$

Using PHENIX  $\sigma_{c\bar{c}}$ ,  $\sigma_{pp} = 42 \text{ mb}$  and 1000 binary collisions in very central Au+Au collisions we find **13±7  $c\bar{c}$  pairs per event**.

Number of  $J/\psi$ s produced by recombination of charms :  $N_{J/\psi} \propto \frac{N_{c\bar{c}}^2}{N_{ch}}$

The charmonium recombination may be an important source of final  $J/\psi$  yield at RHIC.



RHIC

ATR

Booster

LINAC

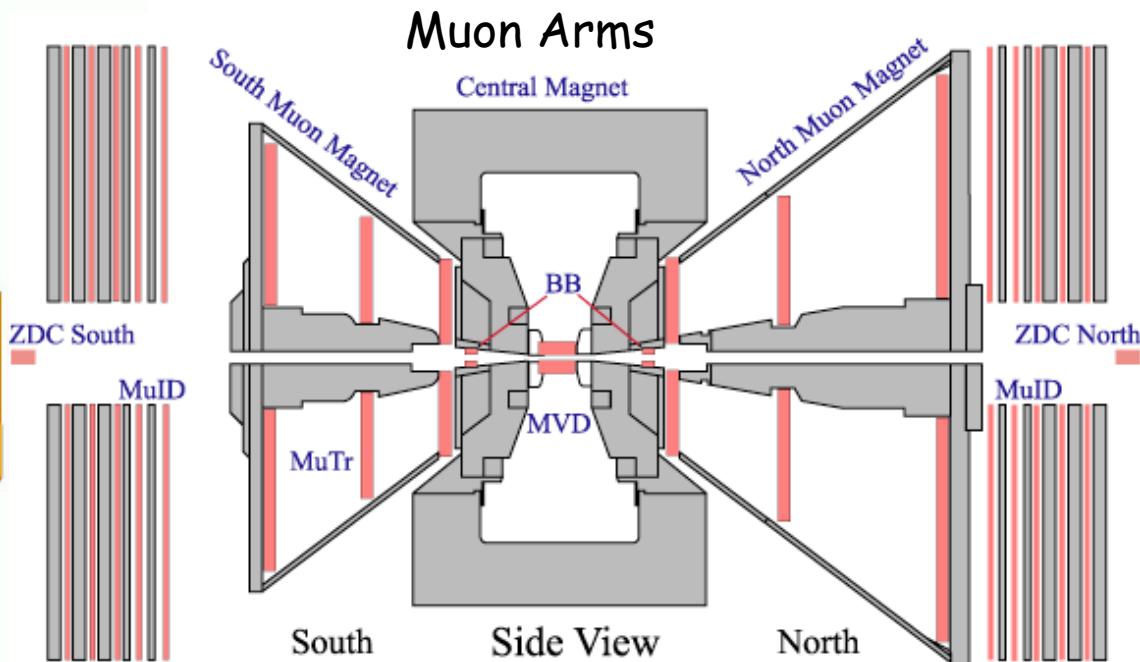
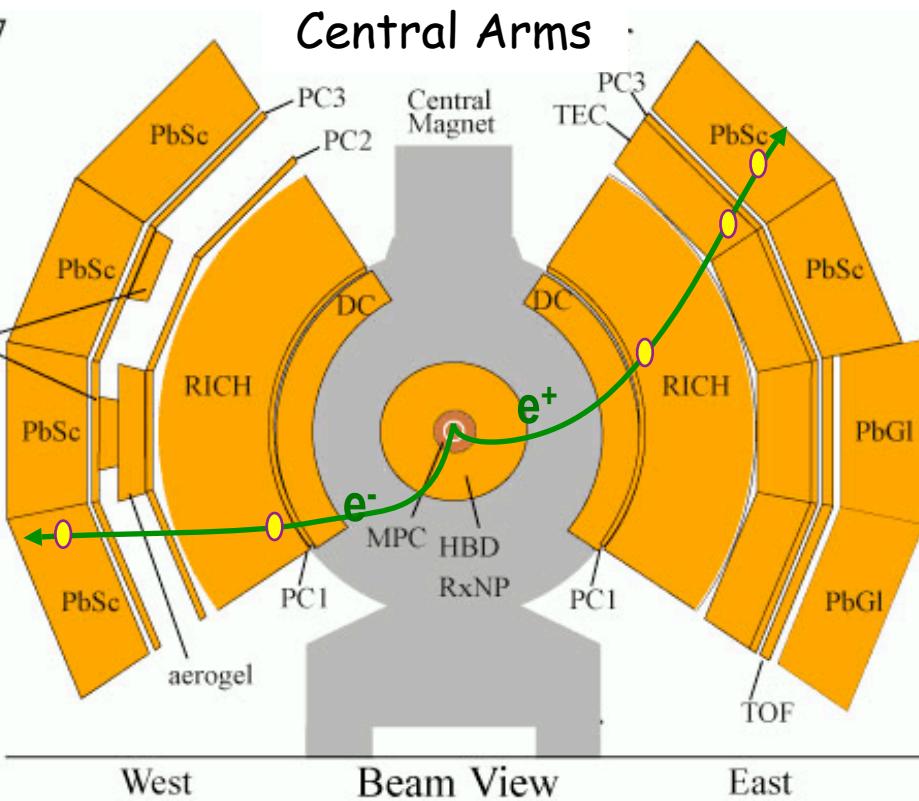
AGS

HITL

Tandem

# PHENIX Measures Electrons and Muons

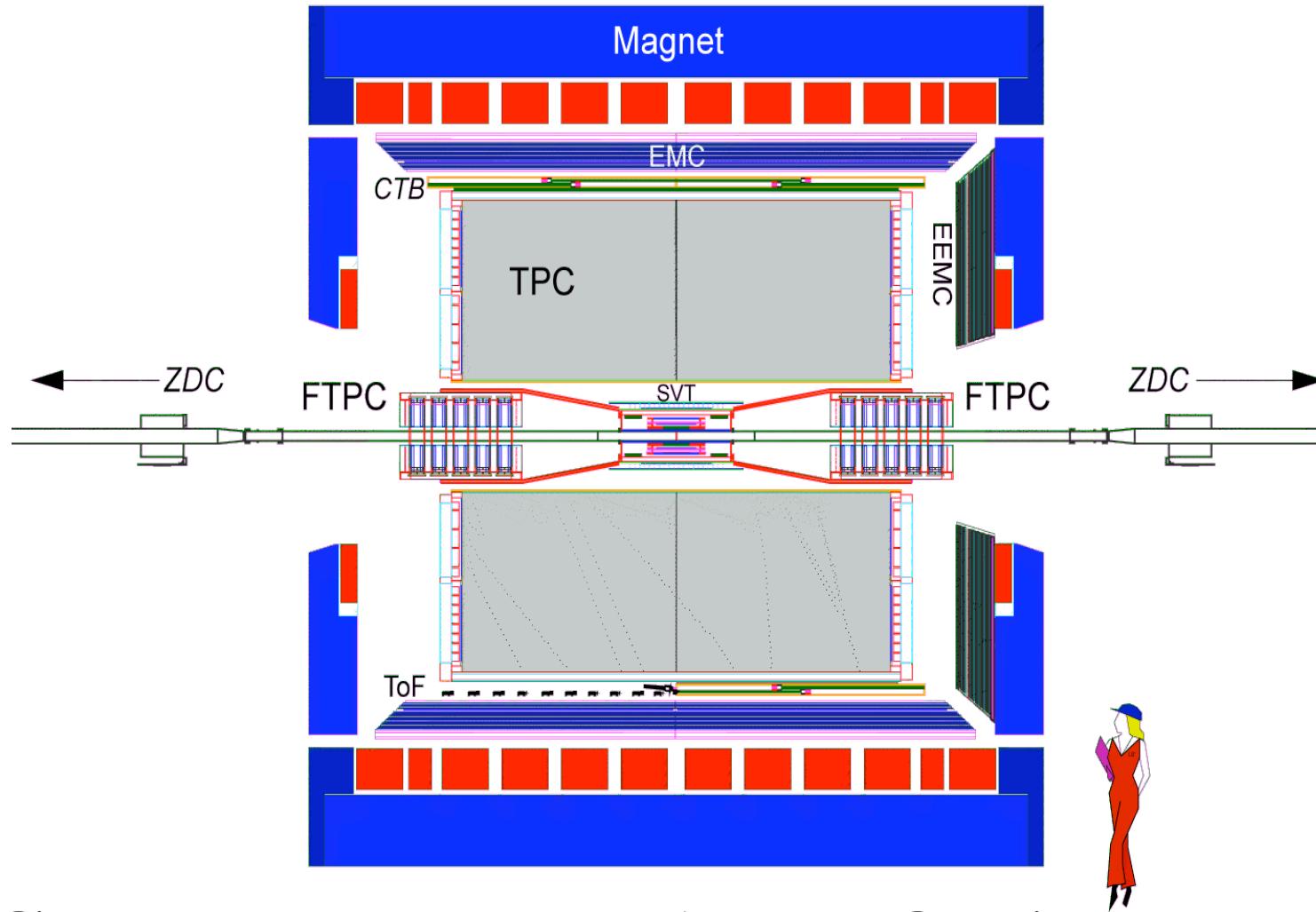
007



- $|\eta| < 0.35$
- $\Delta\phi = 2 \times \pi/2$
- eID: EmCal and RICH

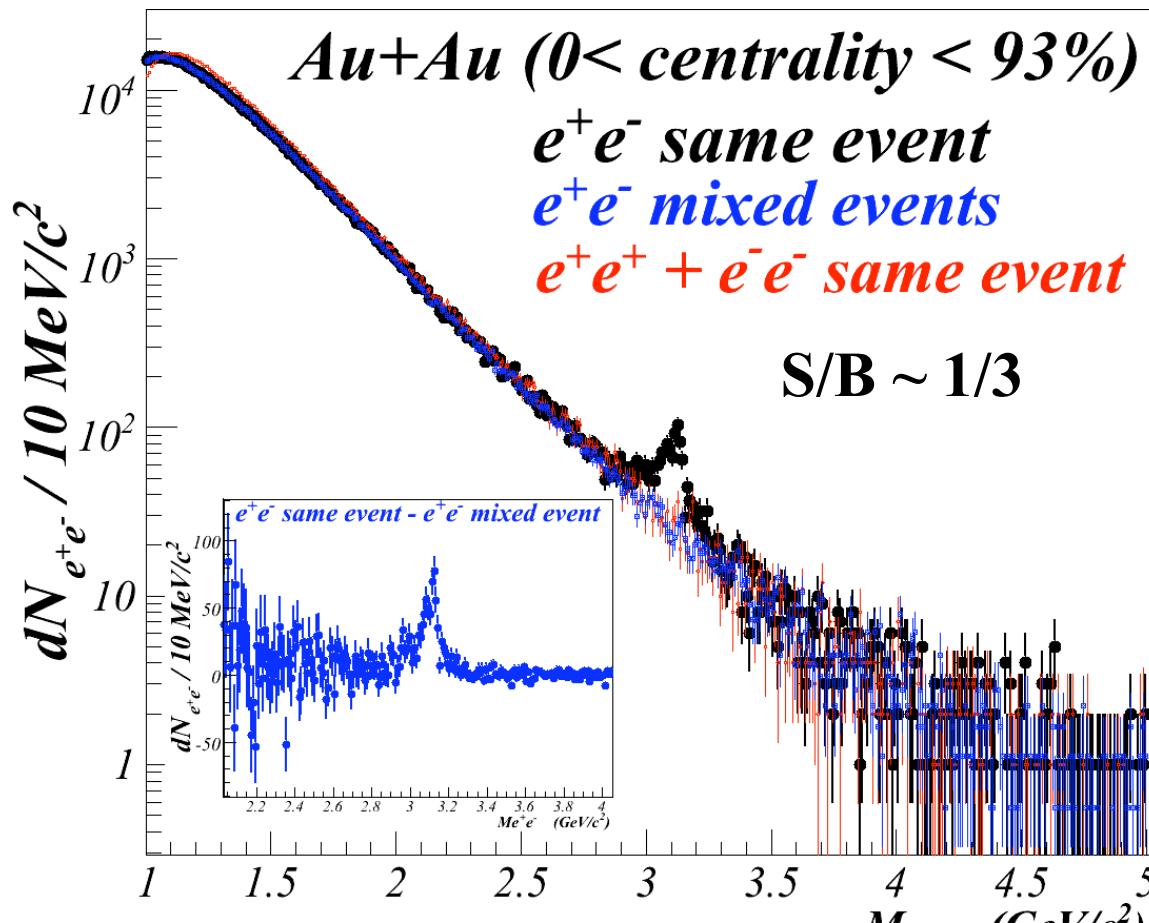
- $-1.2 > \eta > -2.2$  and  $2.4 > \eta > 1.2$
- $\Delta\phi = 2\pi$
- Muon  $|p| > 2 \text{ GeV}/c$
- muon tracking
- 4 hadron absorbers in MuID
- common accepted  $\pi/\mu \sim 10^{-4}$

# STAR Measures Electrons



- Electrons are reconstructed by TPC+EmCal
- $|\eta| < 1$ ,  $0 < \phi < 2\pi$
- Events are triggered by EmCal ( $p_T > 5 \text{ GeV}/c$ )

# Measuring J/ $\psi$ s



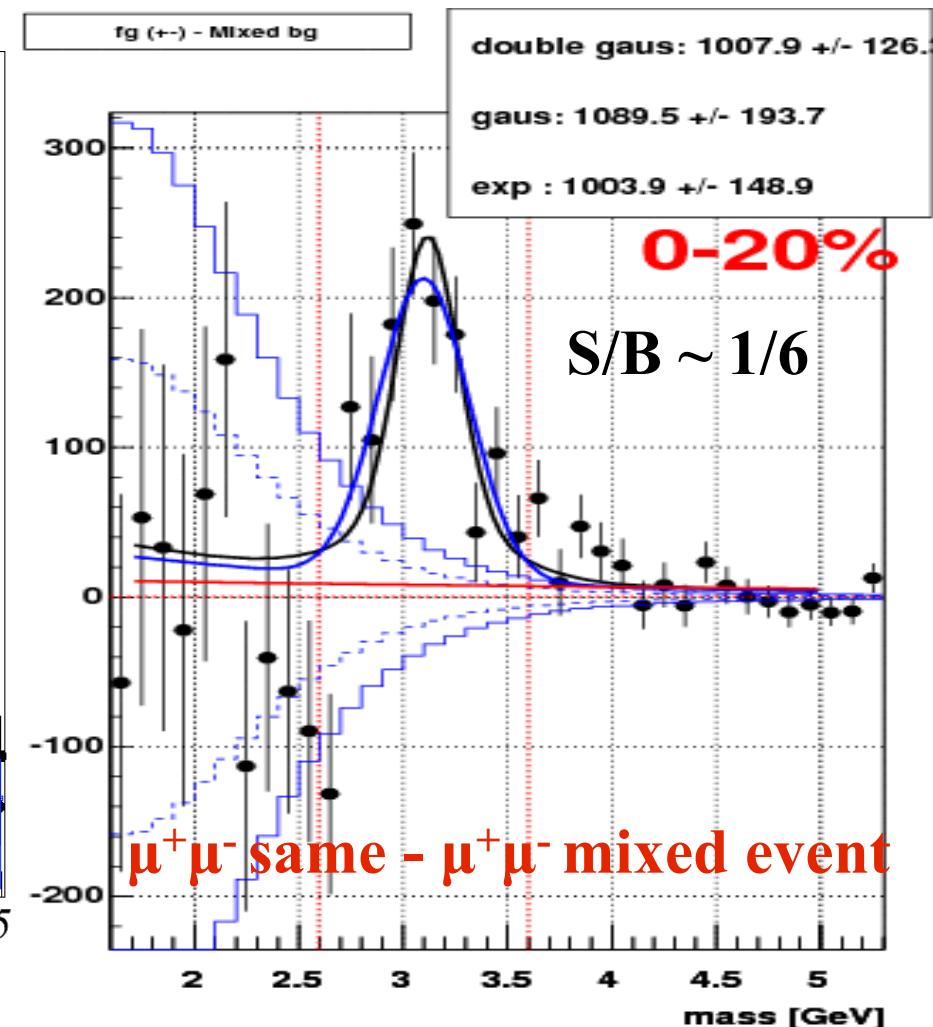
$$B \frac{dN}{dy} = \frac{N_{J/\psi}}{N_{ev} \Delta y \epsilon_{acc} \epsilon_{mult}}$$

B : branching ratio for  $e^+e^-$  or  $\mu^+\mu^-$  from PDG

$\Delta y$  : rapidity coverage

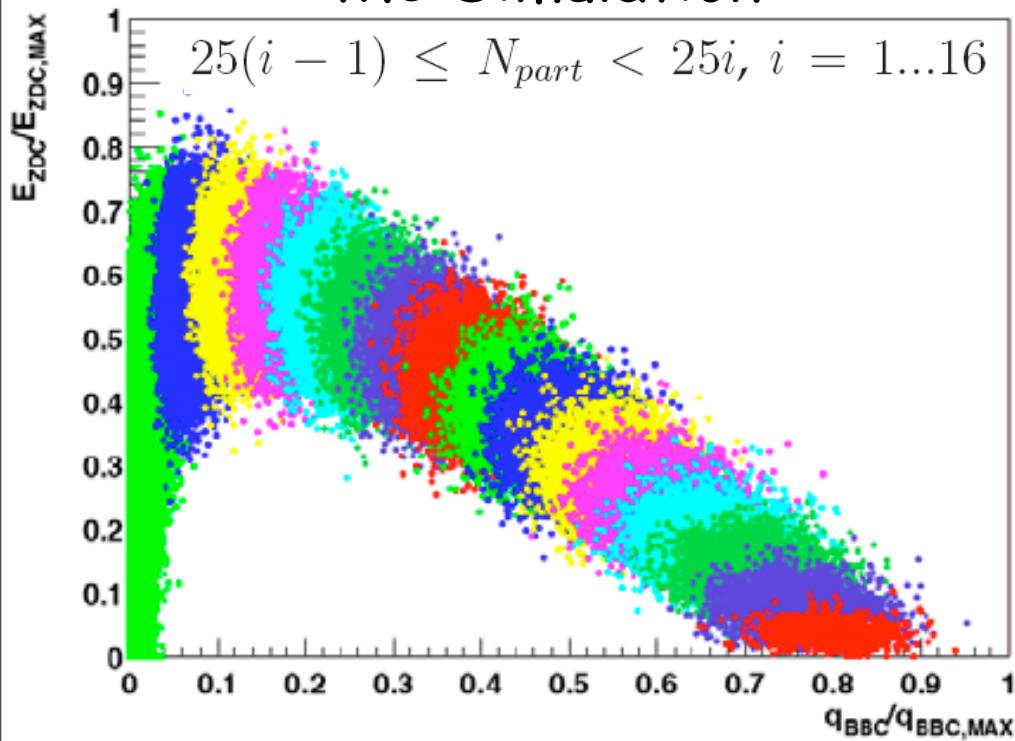
$\epsilon_{acc}$  : overall acceptance for  $J/\psi$  in each arm from simulation

$\epsilon_{mult}$  : efficiency dependence with particle multiplicity.

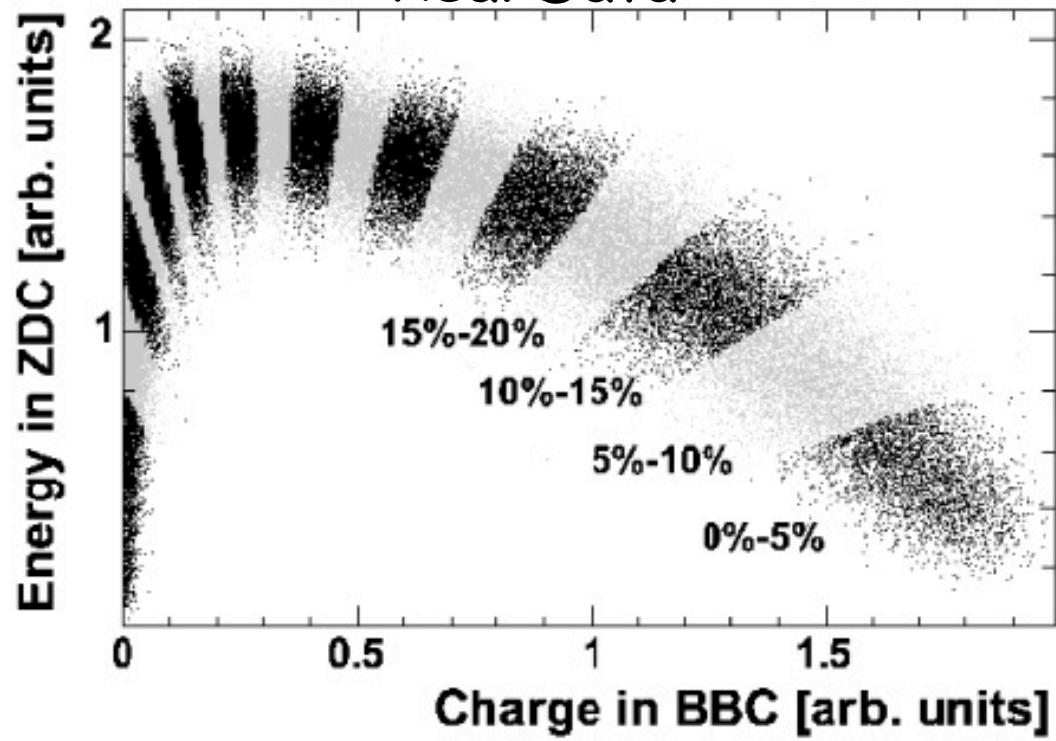


# Determining Centrality

MC Simulation

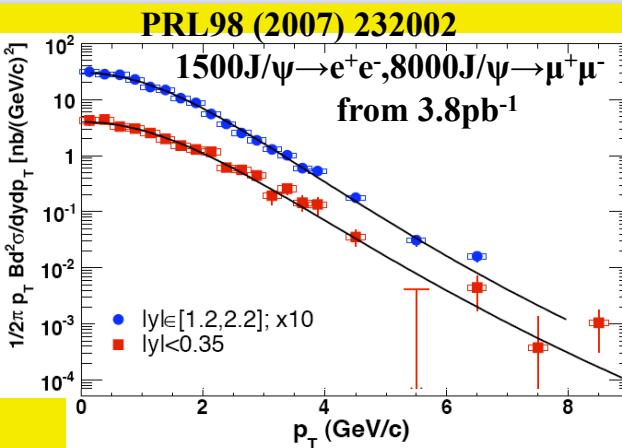
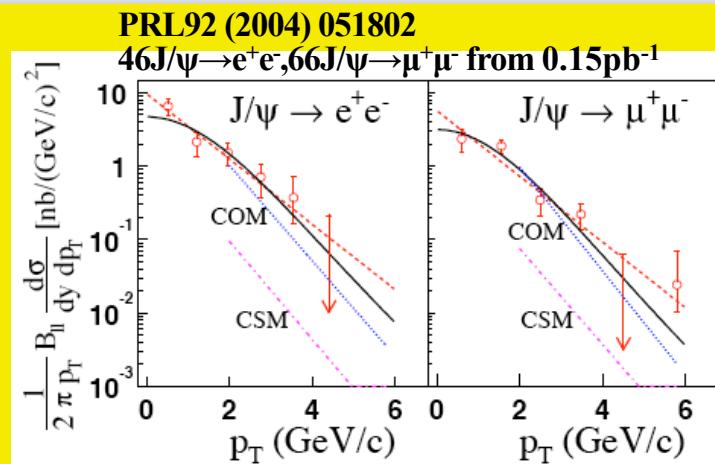


Real Data

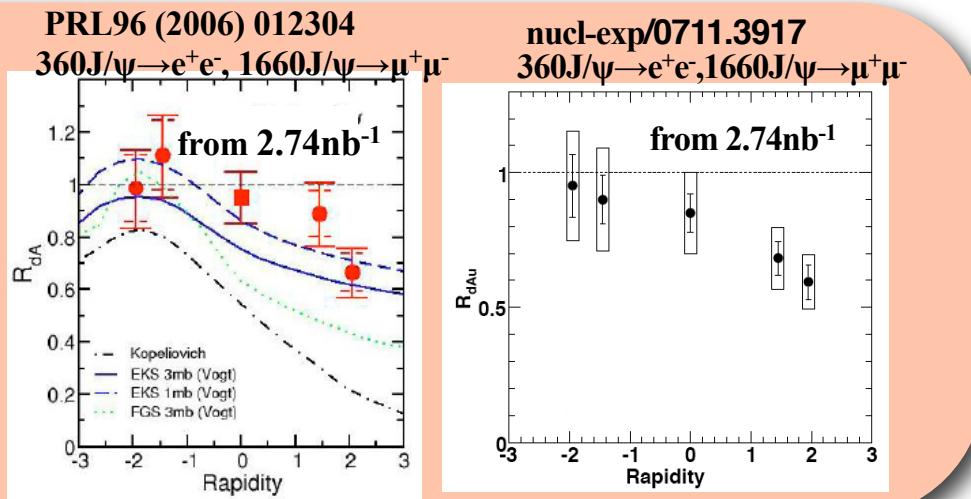


Centrality determined by **Zero Degree Calorimeter**( $|\eta|>6$ ) and **Beam Counter** ( $3.0<|\eta|<3.9$ )

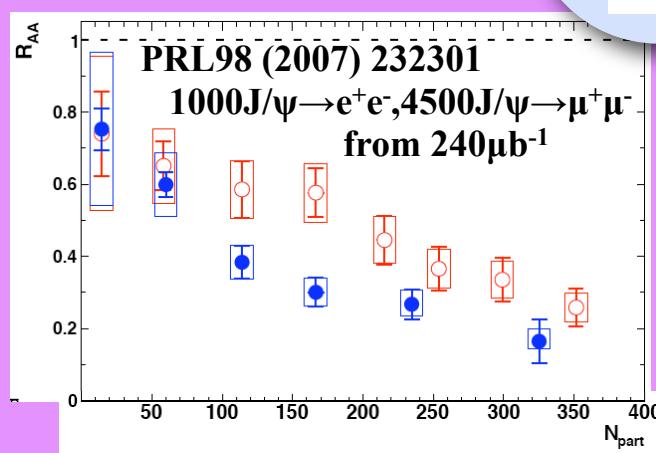
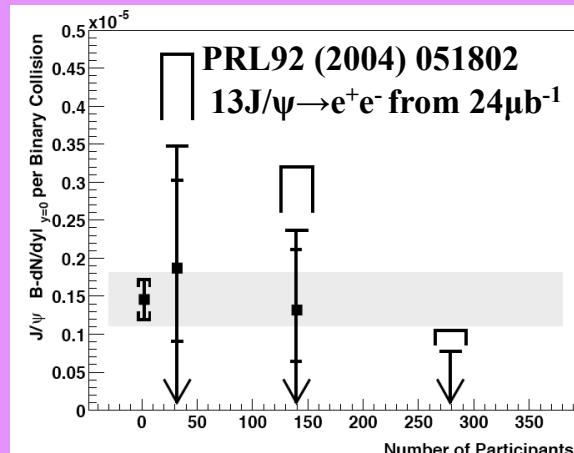
**p+p**



**d+Au**

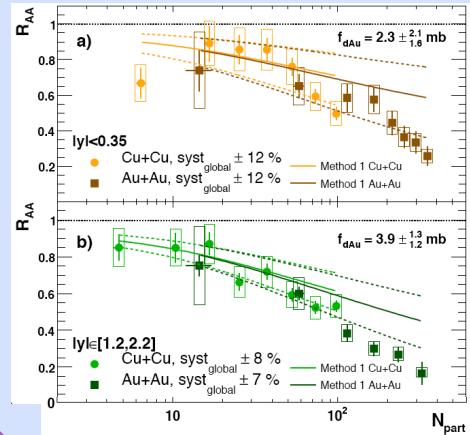


**Au+Au**

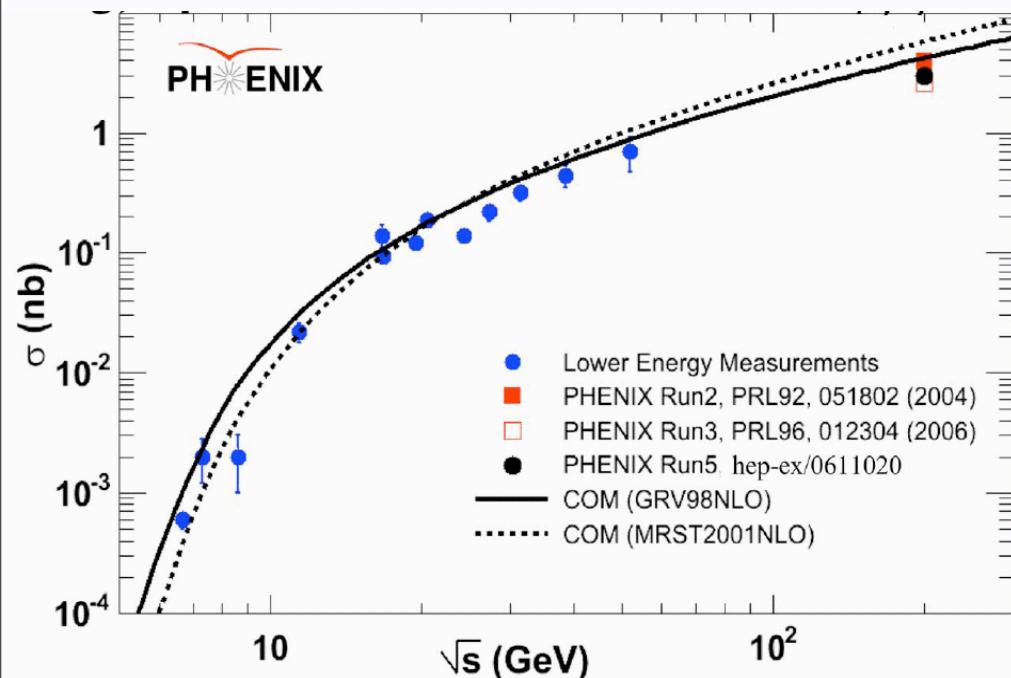


**Cu+Cu**

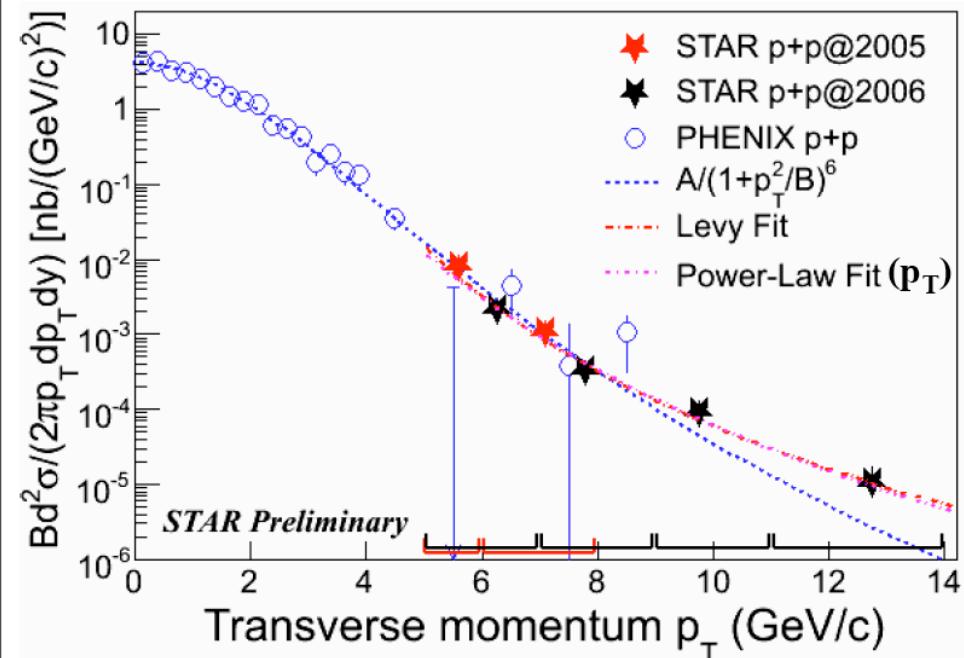
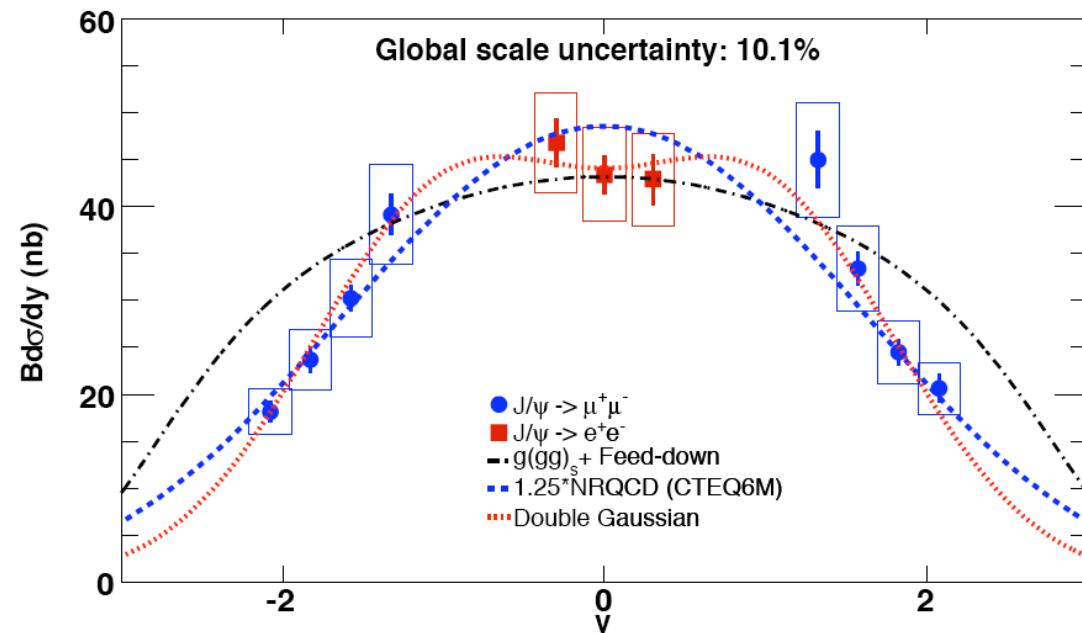
**nucl-exp/0801.0220**  
 $2000J/\psi \rightarrow e^+e^-$ ,  $9000J/\psi \rightarrow \mu^+\mu^-$  from  $2.1\text{nb}^{-1}$



**p+p Data**



$p+p \rightarrow J/\psi$

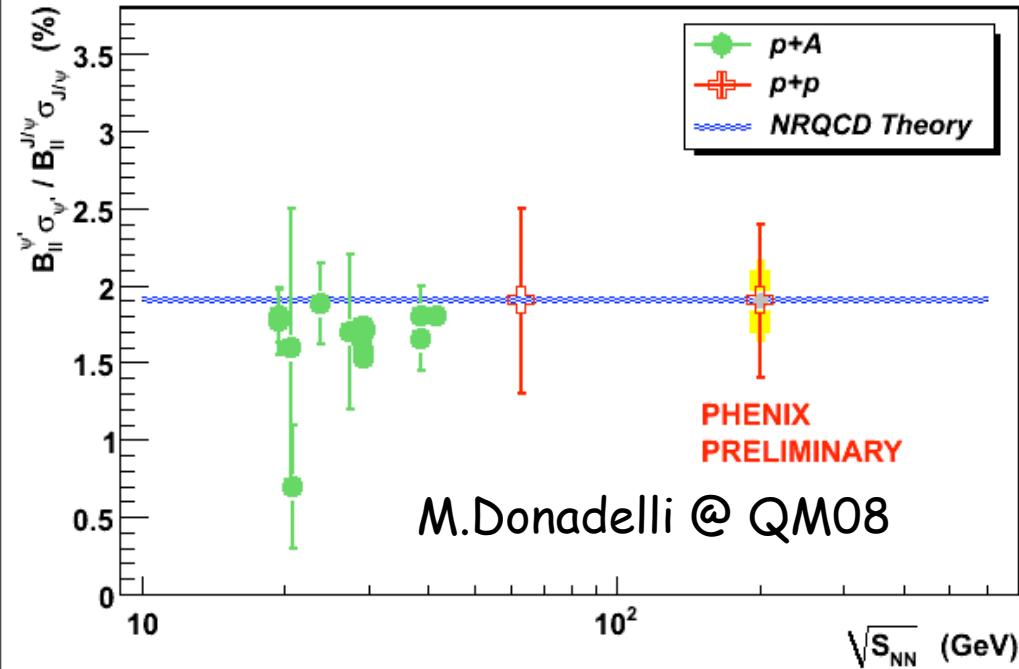


- Inclusive measurement: includes feed down
- Different models fit to the p+p data returning a total cross section

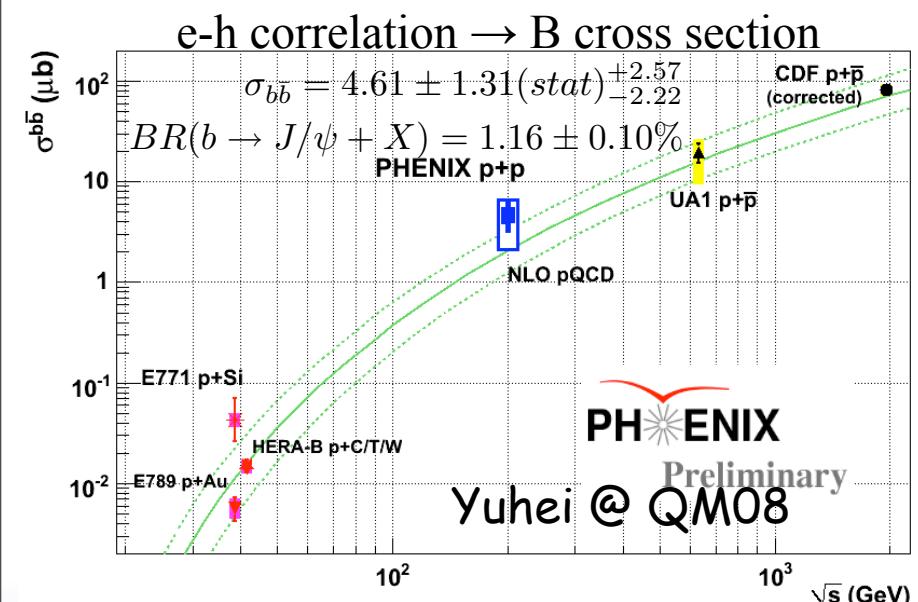
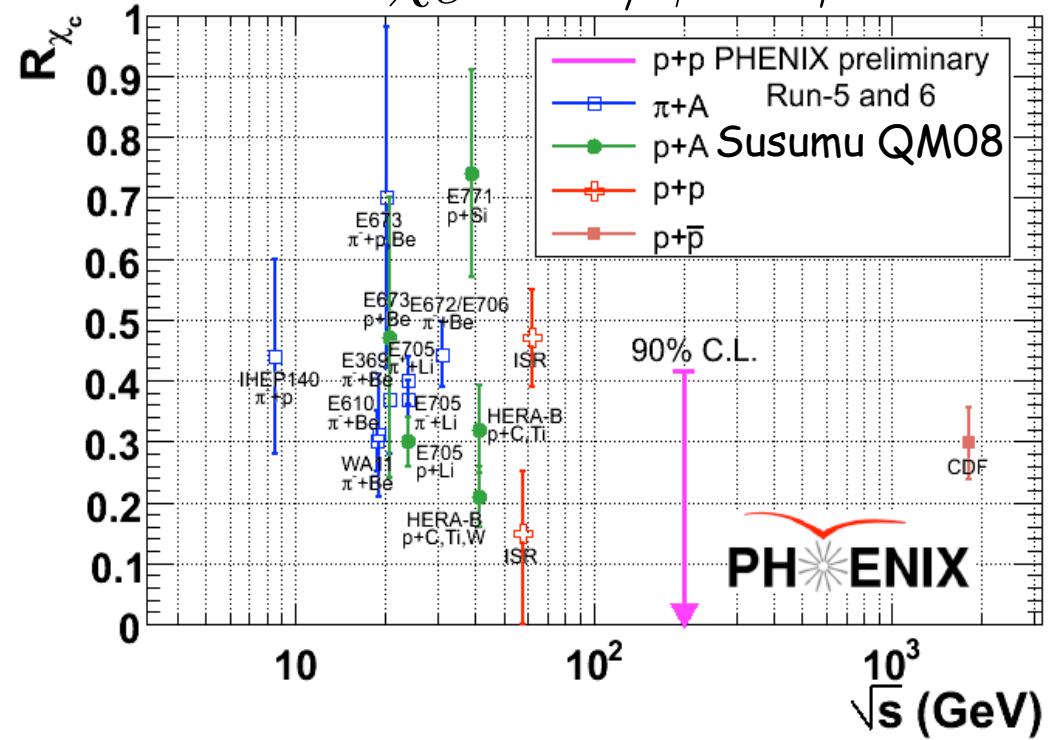
$$\text{BR}_{||} \sigma_{pp}(p+p \rightarrow J/\psi + X) = 178 \pm 3(\text{stat}) \pm 53 (\text{sys}) \pm 18 (\text{norm}) \text{ nb}$$

# Feed Down Contributions in p+p

$$\psi' \rightarrow e^+ e^-$$



$$\chi_c \rightarrow J/\psi + \gamma$$



decay	PHENIX	theory
$\psi' \rightarrow J/\psi$	$0.086 \pm 0.025$	$0.08^*$
$\chi_c \rightarrow J/\psi$	$< 0.42$ (90% CL)	$0.30^*$
$B \rightarrow J/\psi$	$0.036^{+0.025}_{-0.023}$	$0.02 \pm 0.01^{**}$

\*Digal et al., Phys. Rev. D 64 (2001) 094015

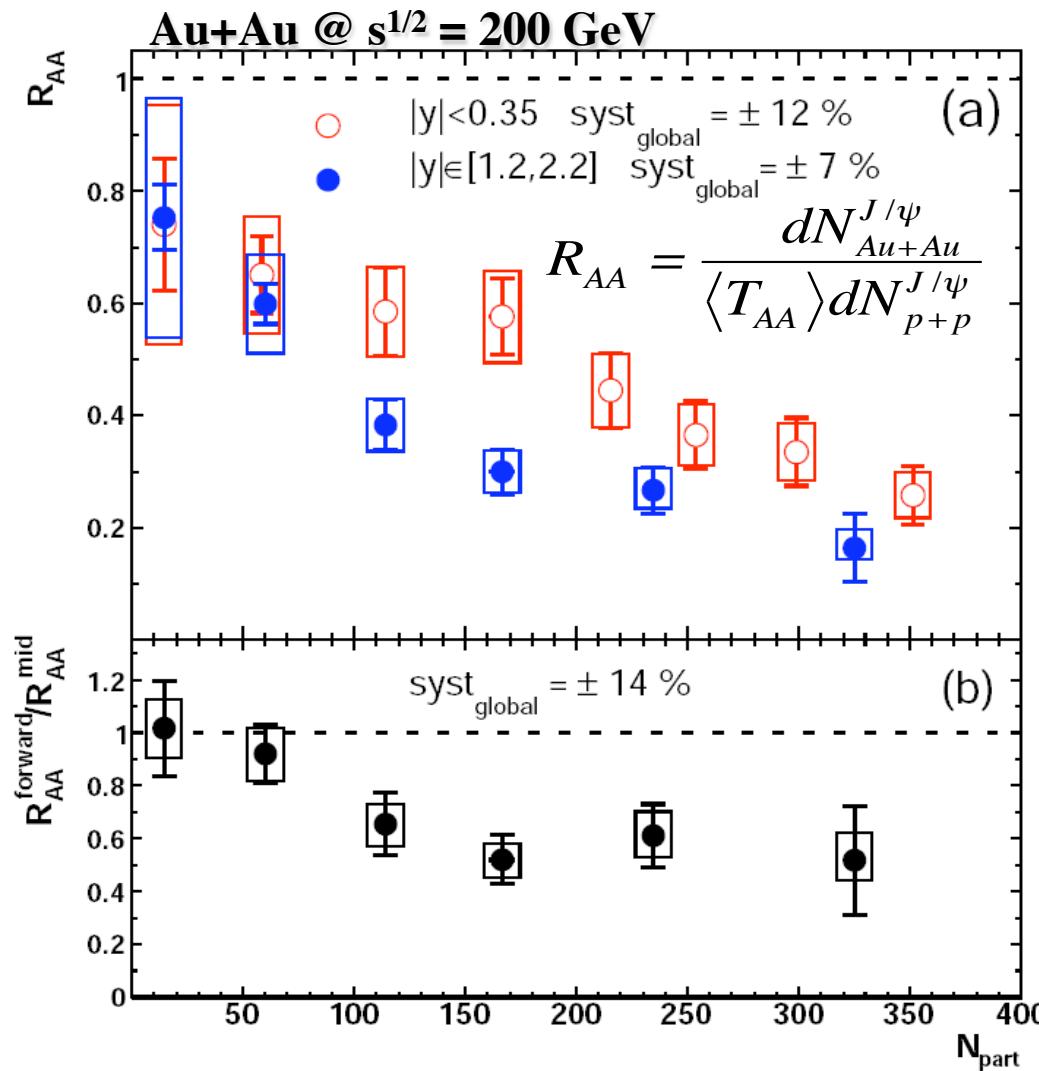
\*\* NLO pQCD extrapolation

**A+A Data**

# Au+Au → J/ψ

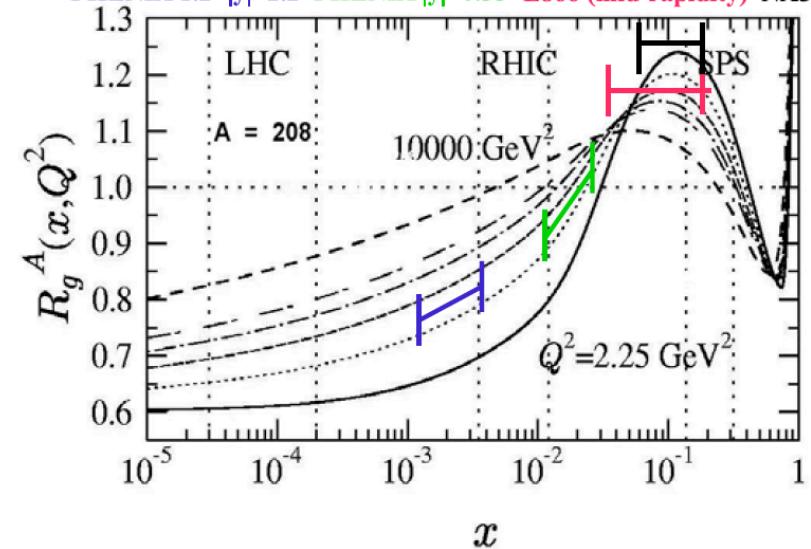
Stronger suppression at forward rapidities:

- $\sigma_{\text{breakup}}(|y| > 1.2) > \sigma_{\text{breakup}}(|y| \sim 0)$  ?
- $R_g(|y| > 1.2) < R_g(|y| \sim 0)$  ?
- regeneration at mid-rapidity ?

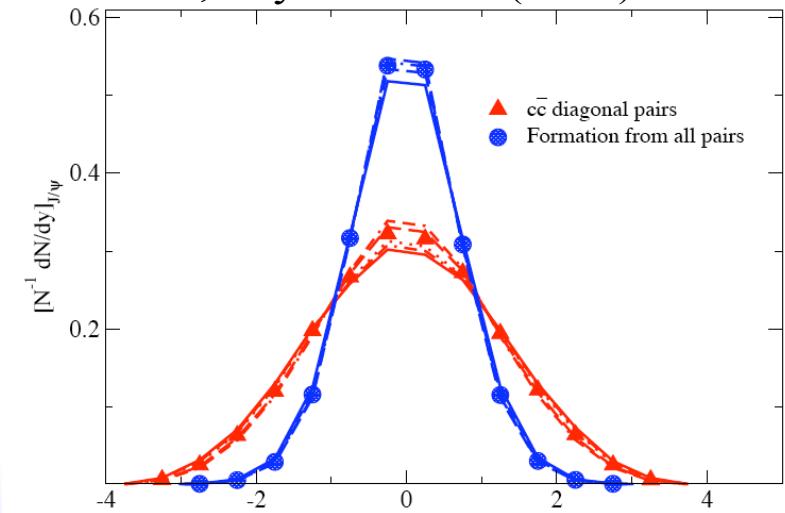


$$R_g(x, Q^2) = \frac{x_g^{\text{nucleus}}(x, Q^2)}{x_g^{\text{nucleon}}(x, Q^2)}$$

K.J.Eskola, V.J.Kolhinen, R. Vogt, *Nucl.Phys.* A696, 729  
PHENIX  $1.2 < |y| < 2.2$  PHENIX  $|y| < 0.35$  E866 (mid-rapidity) NA50

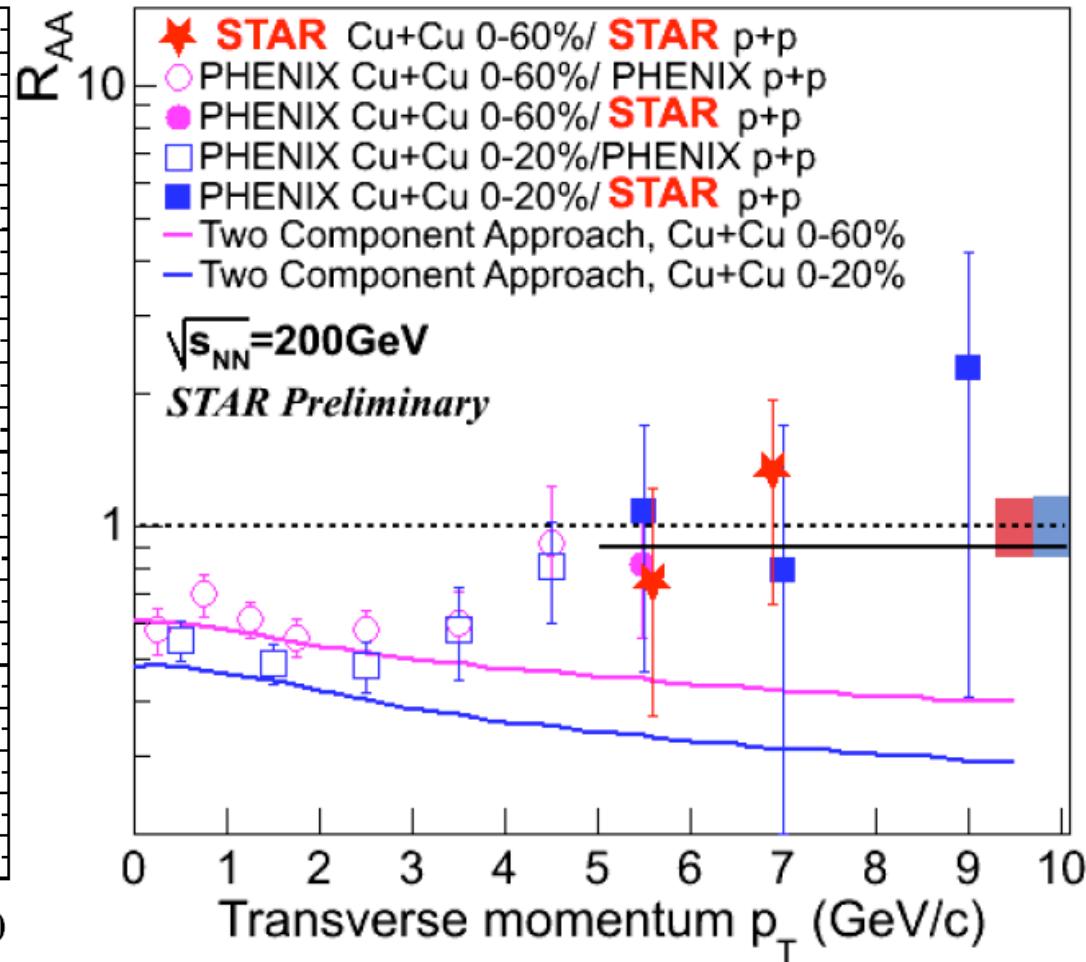
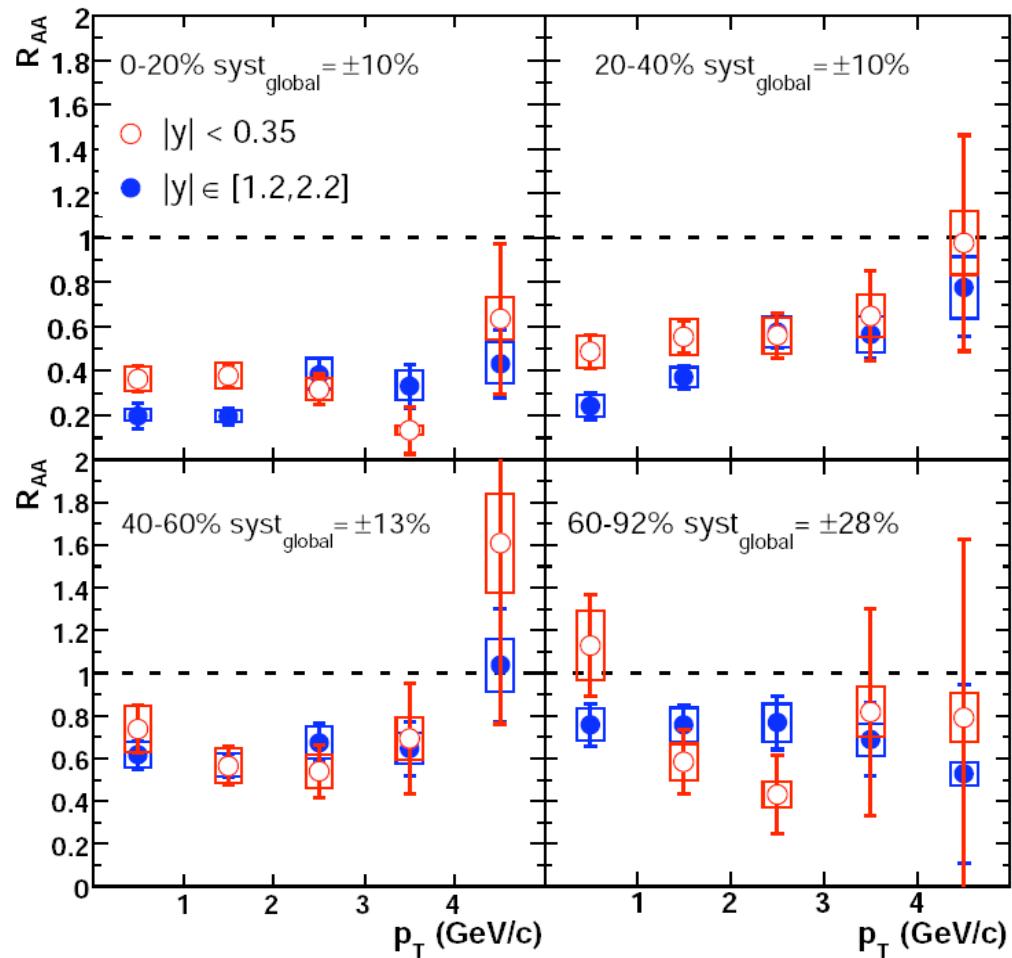


R.L.Tews, Phys.Rev.C73 (2006) 014904



# R<sub>AA</sub> VS p<sub>T</sub>

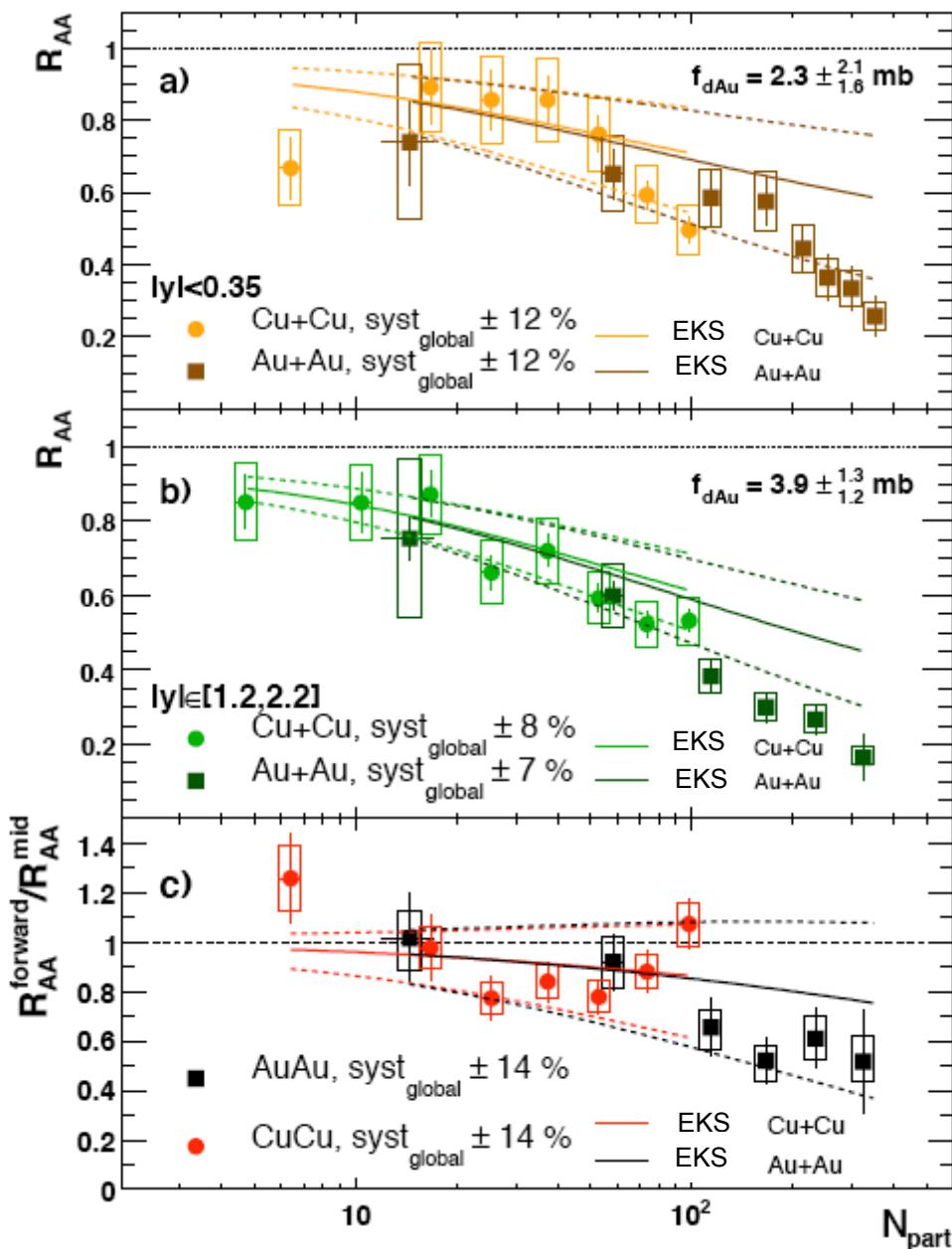
Au+Au @ s<sup>1/2</sup> = 200 GeV



No strong messages from this side.

More data coming soon to have a complete  $R_{AA} \times \text{centrality} \times p_T$  picture.

# Extrapolation of d+Au cold nuclear matter effects to Cu+Cu and Au+Au results



$f_{dAu}$  = breakup cross section

arxiv:0801.0220

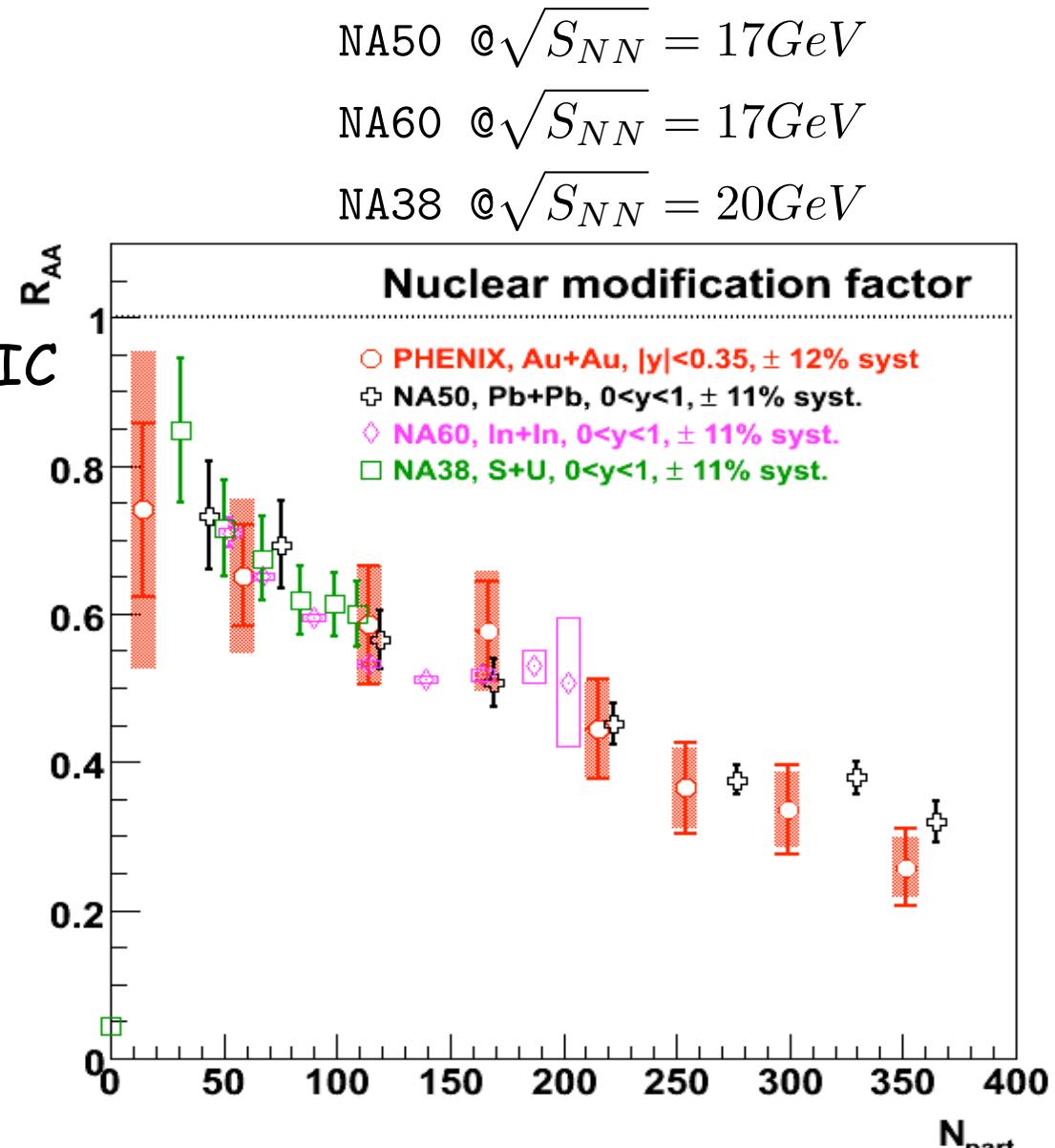
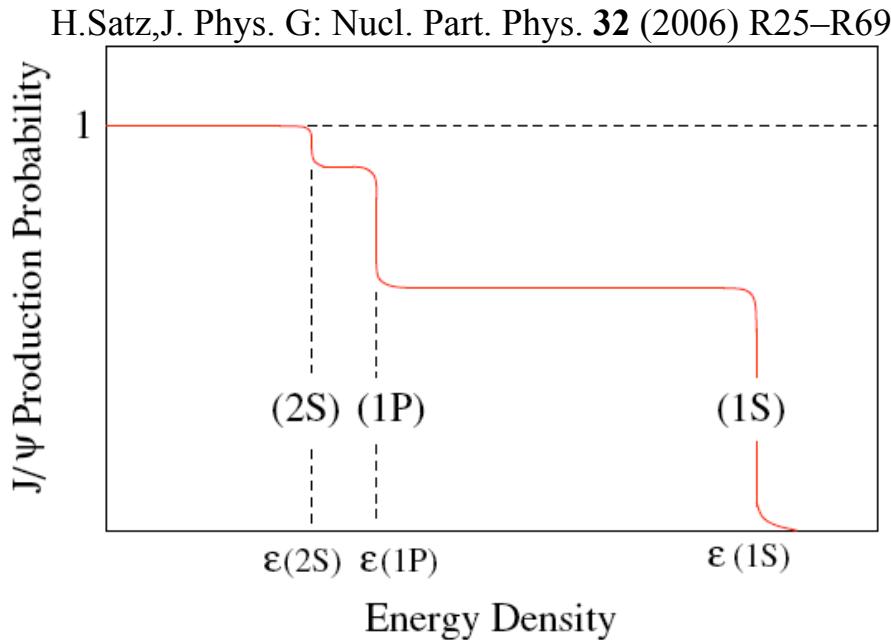
- there is a large room for CNM effects
- anomalous suppression is considerable at least at forward rapidity
- but the CNM is not enough constrained yet for a serious calculation of anomalous suppression  $R_{AA}/\text{CNM}$
- see Raphael's talk

# Comparison with SPS

Same suppression as found in SPS

Caveat:

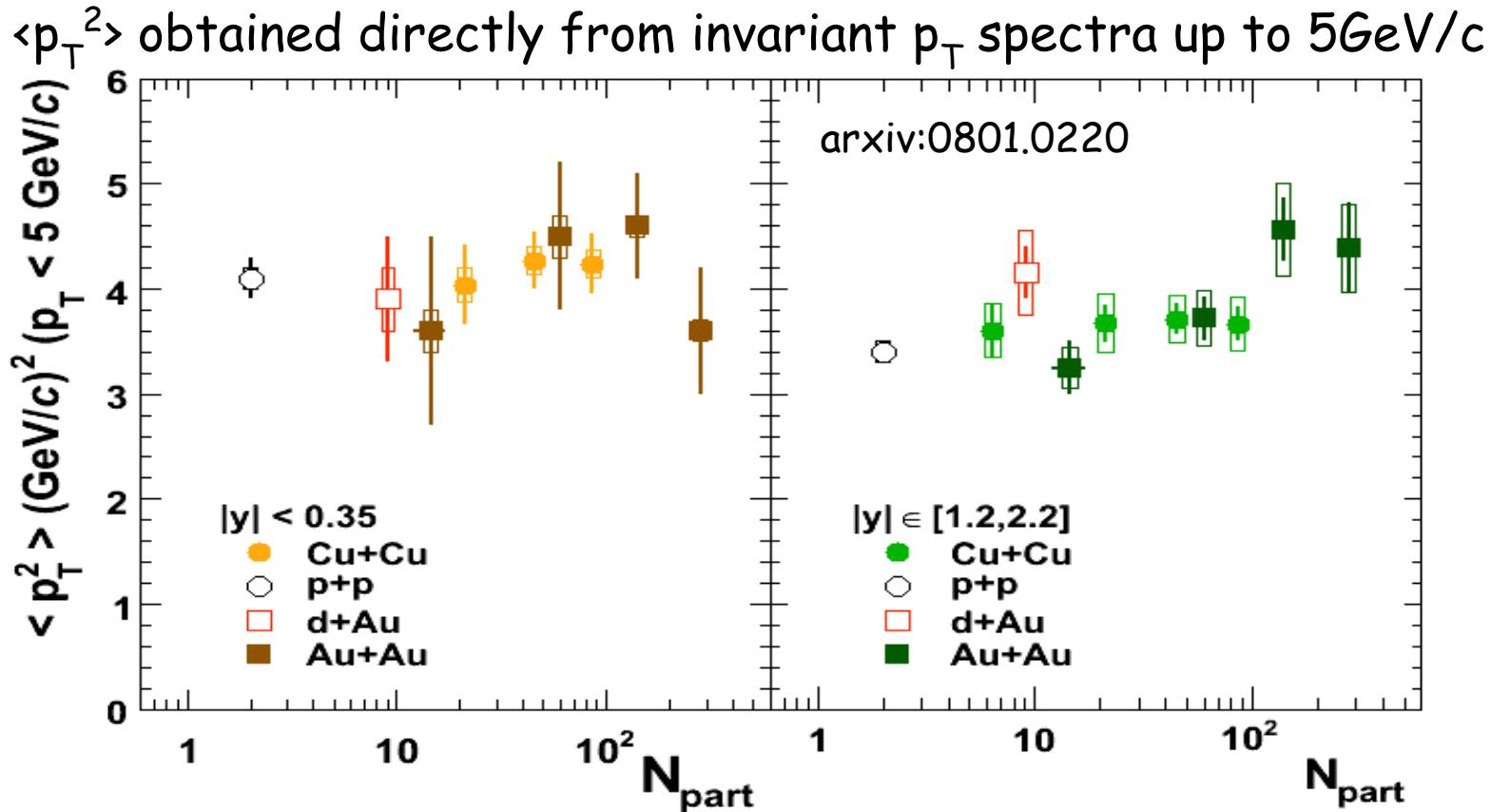
- different CNM
- stronger QGP suppression at RHIC
- recombination at RHIC
- melting of different feed down sources of J/ $\psi$  ?



Needs  $R_{AA}(\chi_c)$ ,  $R_{AA}(\psi')$  measurements.

# Recombination/Coalescence Probe I

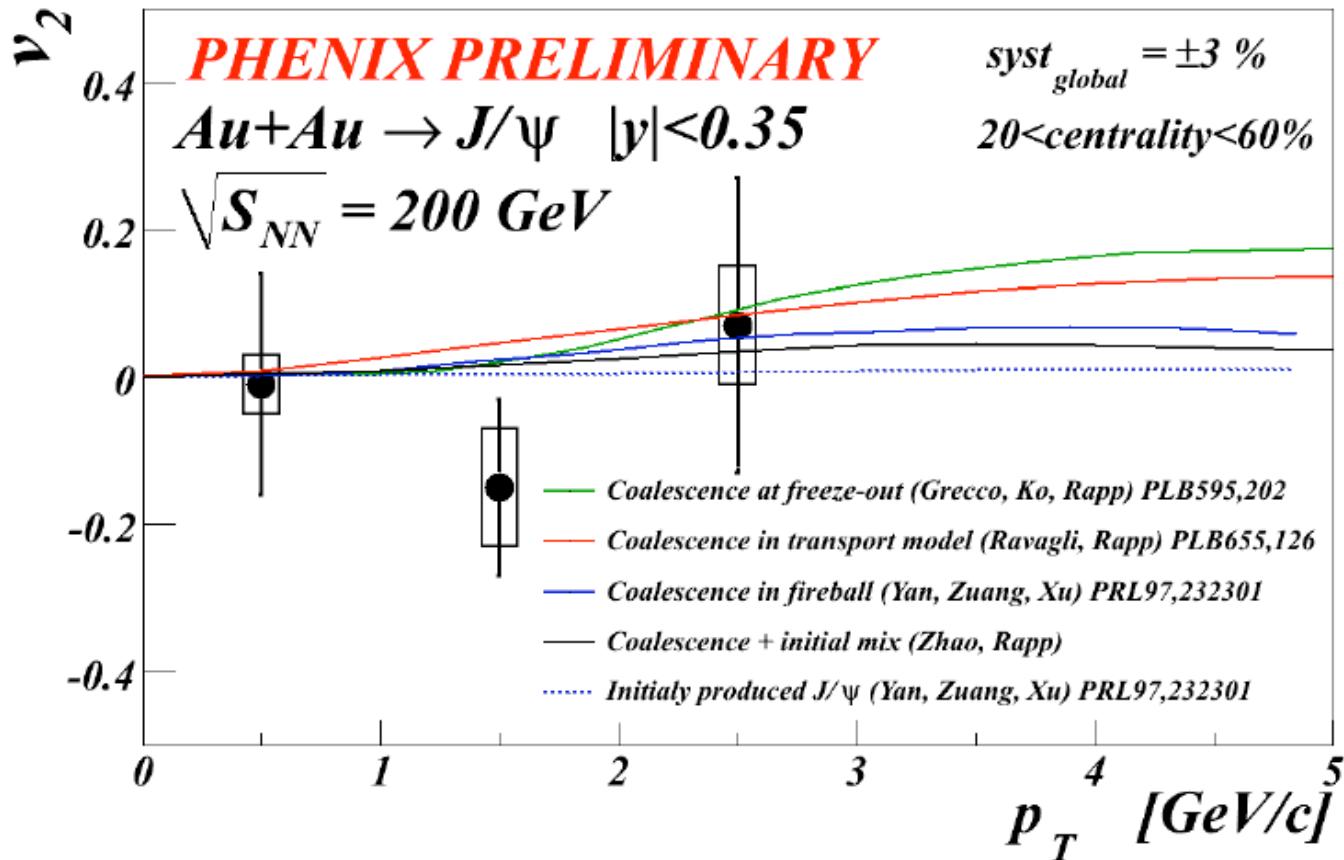
- Cronin effect broaden the  $p_T$  of primary J/ψs
- J/ψs from recombined charms are not affected by Cronin effect and  $\langle p_T^2 \rangle$  should be flat



- $p_T$  broadening observed at forward rapidity but not at mid-rapidity
- but there is no broadening in d+Au data neither at mid-rapidity
- again, we need more d+Au to confirm this trend

# Recombination/Coalescence Probe II

- primary J/ $\psi$ s has isotropic behavior
- open charms flow (see A. Dion's presentation)
- J/ $\psi$ s from these charms will show similar flow



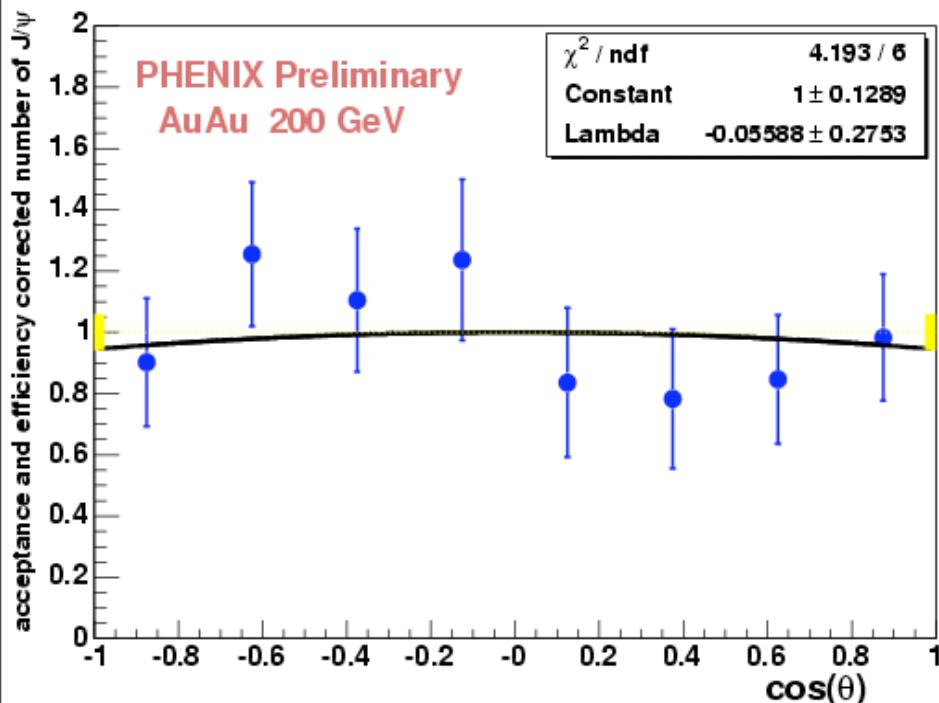
- partial data from last Au+Au run doesn't provide any hint for now
- be patient, remember the first PHENIX J/ $\psi$  R<sub>AA</sub> measurement?

# (Non-)Perturbative Gauge

## Polarization measurement

$$N(\theta) = A[1 + \lambda \cdot \cos^2 \theta]$$

$\theta \equiv$  angle btw.  $e^+$  and  $J/\psi$  directions  
in its rest frame

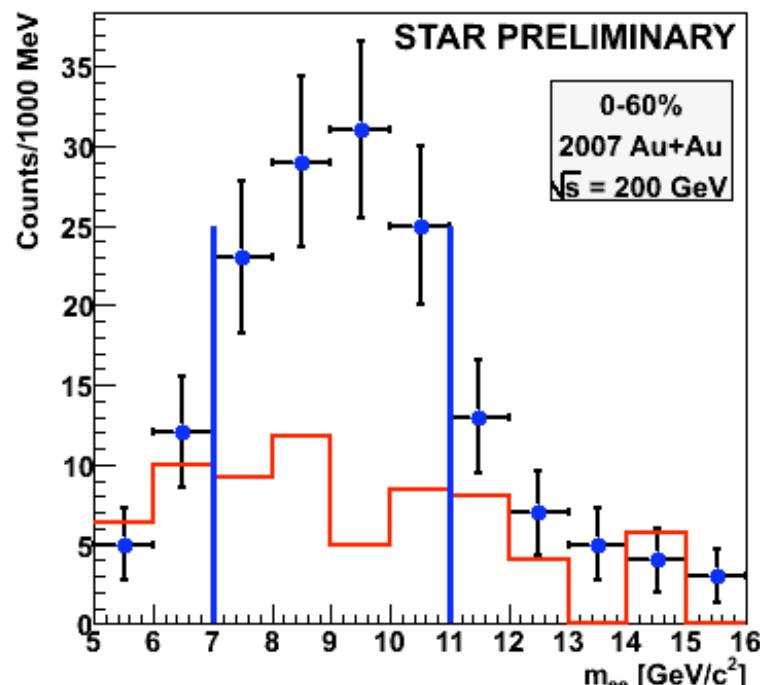
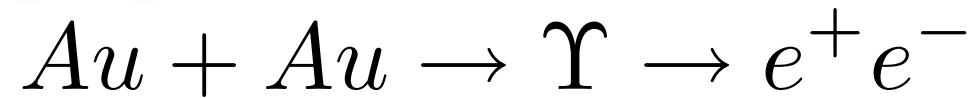
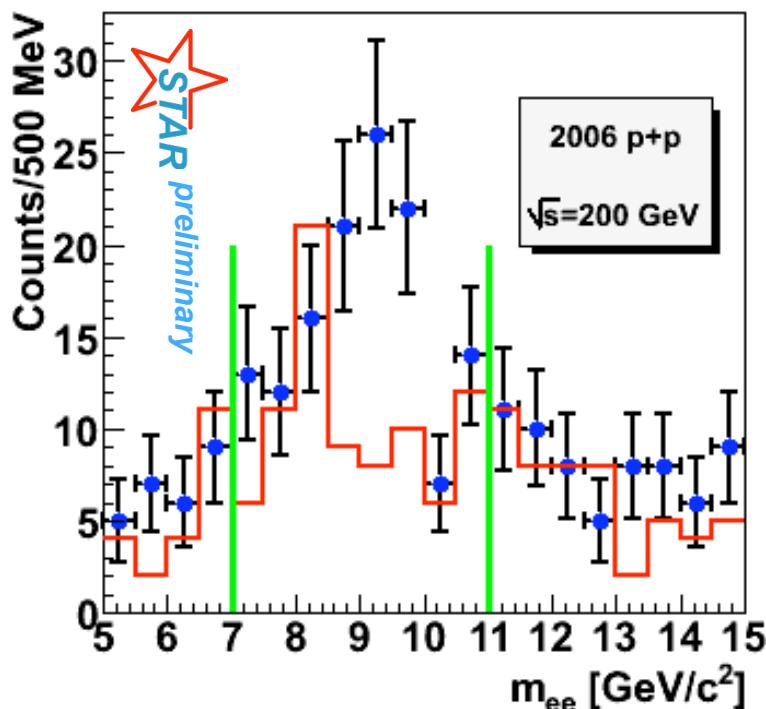
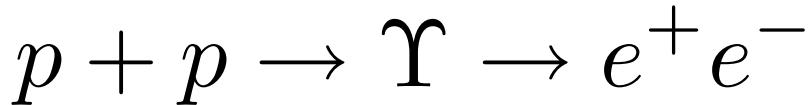


Polarization consistent with zero in  
the current  $Au+Au$  PHENIX  
preliminary result

B. L. Ioffe, D. E. Kharzeev, PRC68, 061902 (2003)

- non-perturbative effects in  $J/\psi$  production leads to unpolarization
- perturbative behavior leads to a polarized  $J/\psi$
- deconfined matter wash out non-perturbative effects in  $J/\psi$  and result in non zero polarization for  $p_T \sim 0$  in  $Au+Au$  collisions
- one needs precise measurement of polarization evolution in  $Au+Au$ ,  $d+Au$  and  $p+p$

# $\Upsilon$ Measurement in p+p and Au+Au

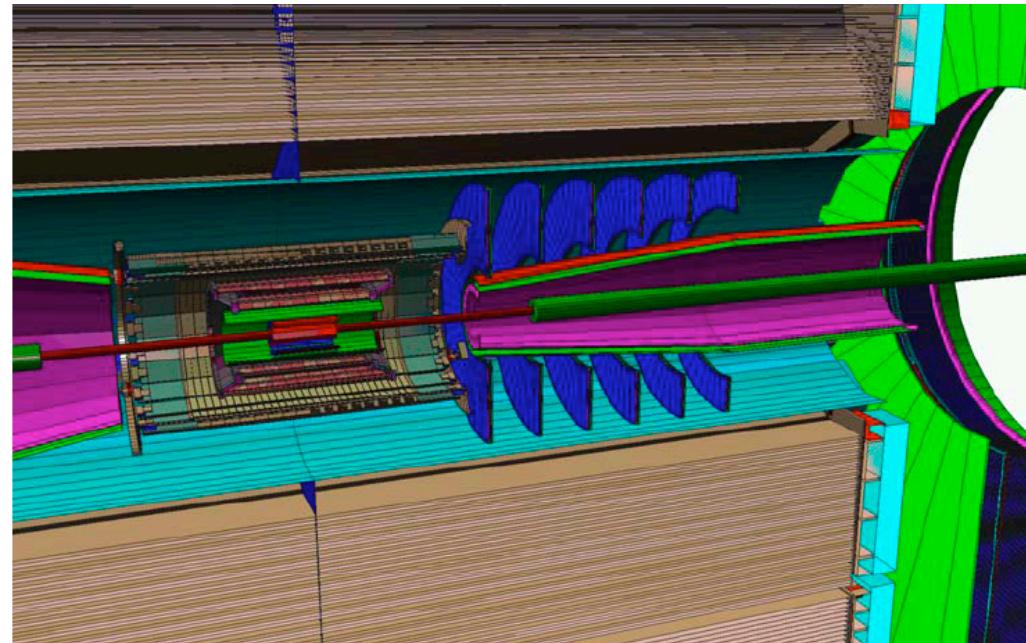


- First measurement of  $\Upsilon$ s in A+A collision so far
- nuclear modification matter much easier to interpret
  - less shadowing (CNM)
  - negligible recombination
- nevertheless, dissociation occur only for  $T/T_c > 3$  for  $\Upsilon$ s and  $\sim 1.6$  for  $\Upsilon'$

# Detectors Upgrades

PHENIX

STAR



## Silicon Vertex Detector

- will improve S/B
- precise measurement of  $B \rightarrow J/\psi$

## Nose Cone Calorimeter

- large coverage for  $\gamma$
- larger acceptance for  $\chi_c \rightarrow J/\psi + \gamma$

## Forward Tracker

- measures  $J/\psi$  at forward rapidity

## Compact Muon Detector

- measurement of  $\Upsilon \rightarrow \mu^+ \mu^-$

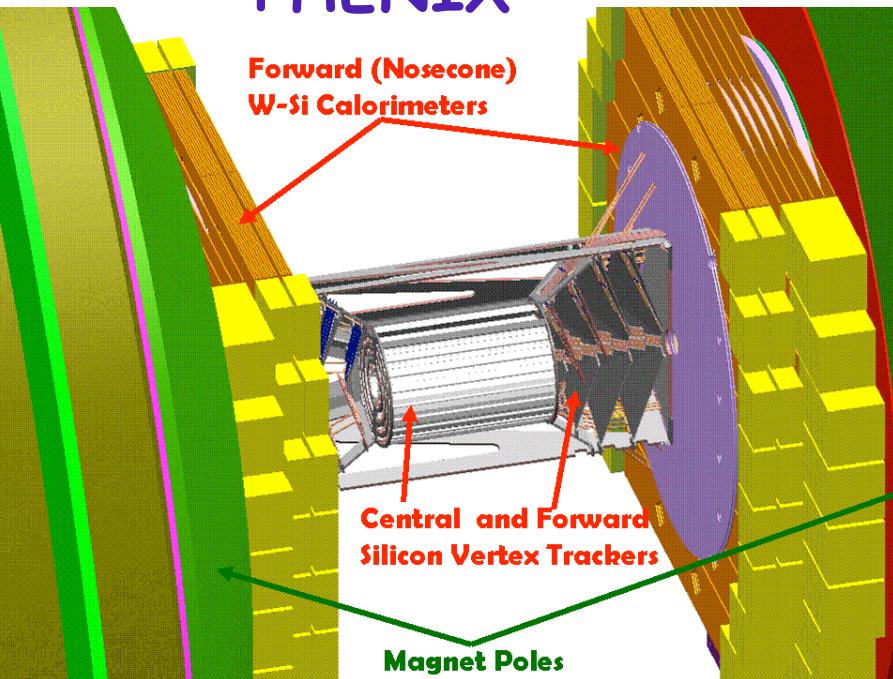
## DAQ upgrade

- take advantage of higher luminosities

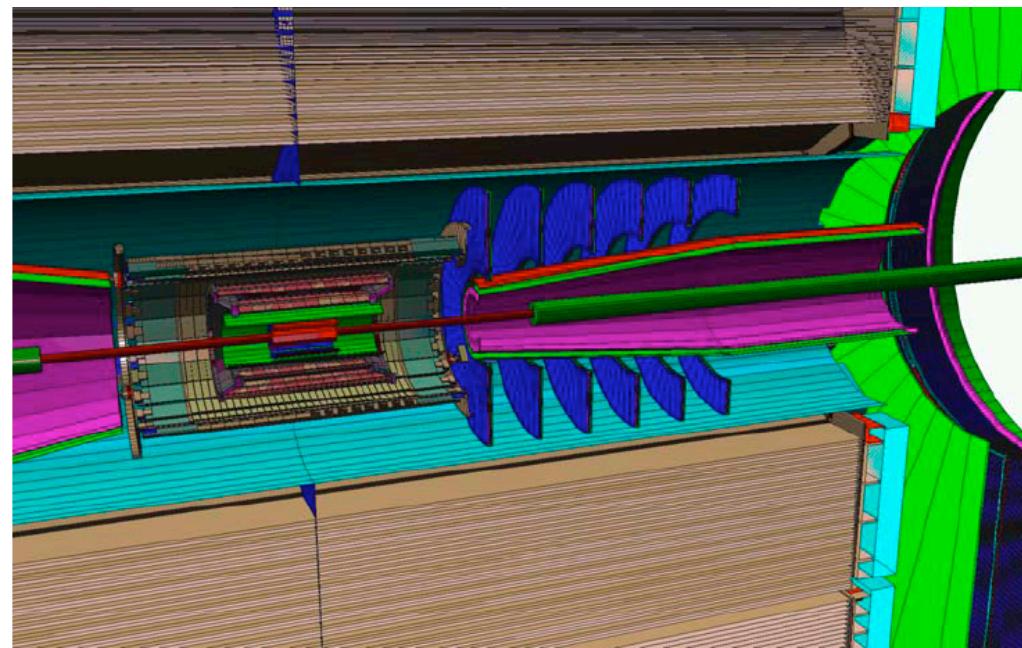
Stochastic cooling at RHIC will provide an order of magnitude higher luminosity in few years

# Detectors Upgrades

## PHENIX



## STAR



### Silicon Vertex Detector

- will improve S/B
- precise measurement of  $B \rightarrow J/\psi$

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### Forward Tracker

- measures  $J/\psi$  at forward rapidity
- Compact Muon Detector**
- measurement of  $\Upsilon \rightarrow \mu^+ \mu^-$
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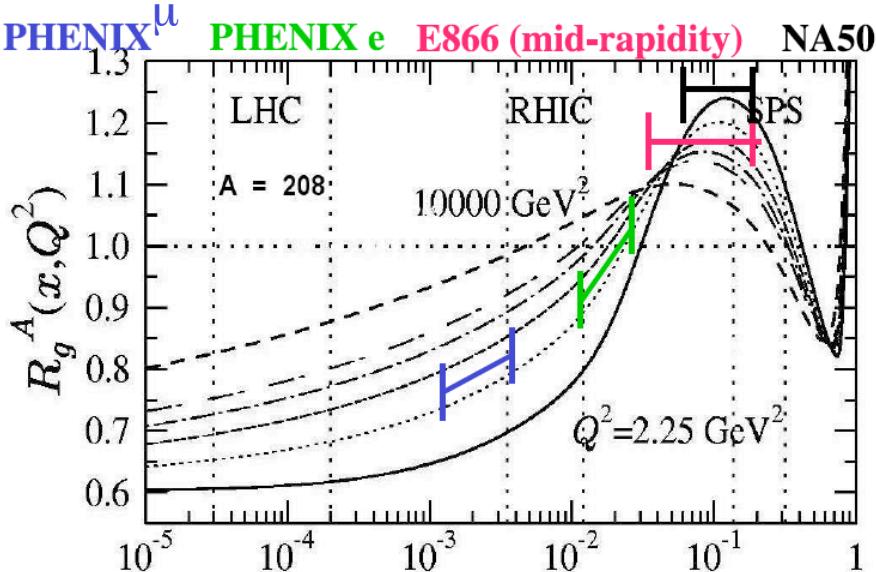
# Current Status and Outlook

- a big progress in the lattice QCD was made to use quarkonia suppression as a hot and dense matter thermometer
- in the experimental side the suppression has been measured but we are still learning its components: CNM, recombination, feed down
- $R_{AA}$  measurement is not sufficient to separate the suppression components
- other measurements which are on the way:
  - extended momentum distributions
  - elliptic flow
  - polarization
- suppression of excited states needs detector upgrades
- $\Upsilon$  suppression measurement seems feasible in a short term

# BACKUP SLIDES

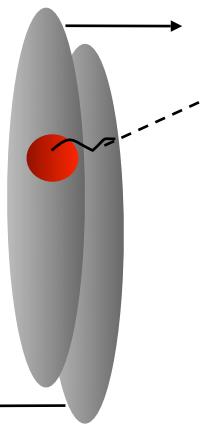
# Cold Nuclear Effects

PDF modification in nucleus from DGLAP evolution equations fits to DIS and Drell-Yan data.[Nucl.Phys.B535(1998)351]



Quarkonia formation doesn't scale with number of collisions due to :

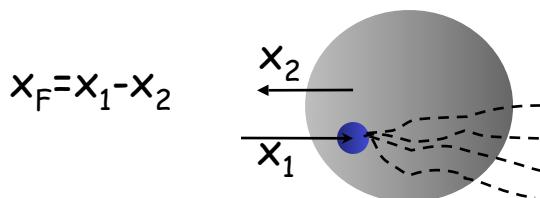
- ▶ different parton distribution
- ▶ dissociation with hadrons ( $J/\psi + \pi^- \rightarrow D^+ + D^-$ )
- ▶ energy loss



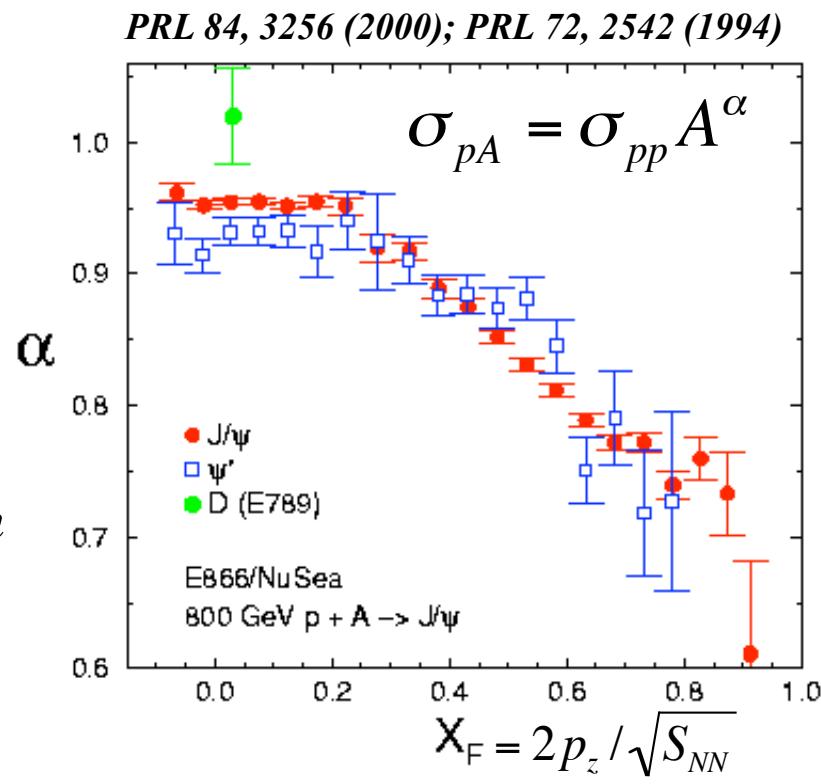
$(ccg)_8$  formation  $\tau_c \approx \hbar / 2m_c \approx 0.07 \text{ fm}/c$

Color neutralization :  $\tau_8 \approx \hbar / \sqrt{2m_c \Lambda_{QCD}} \approx 0.25 \text{ fm}/c$

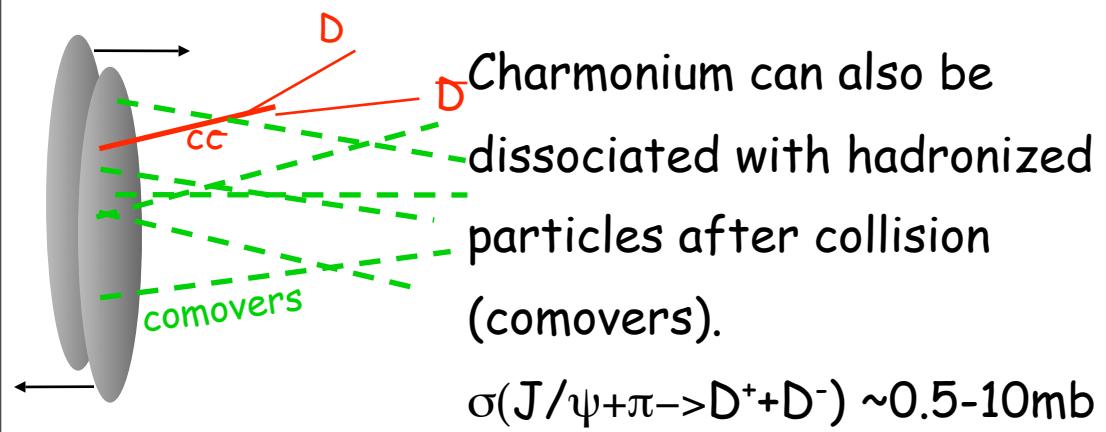
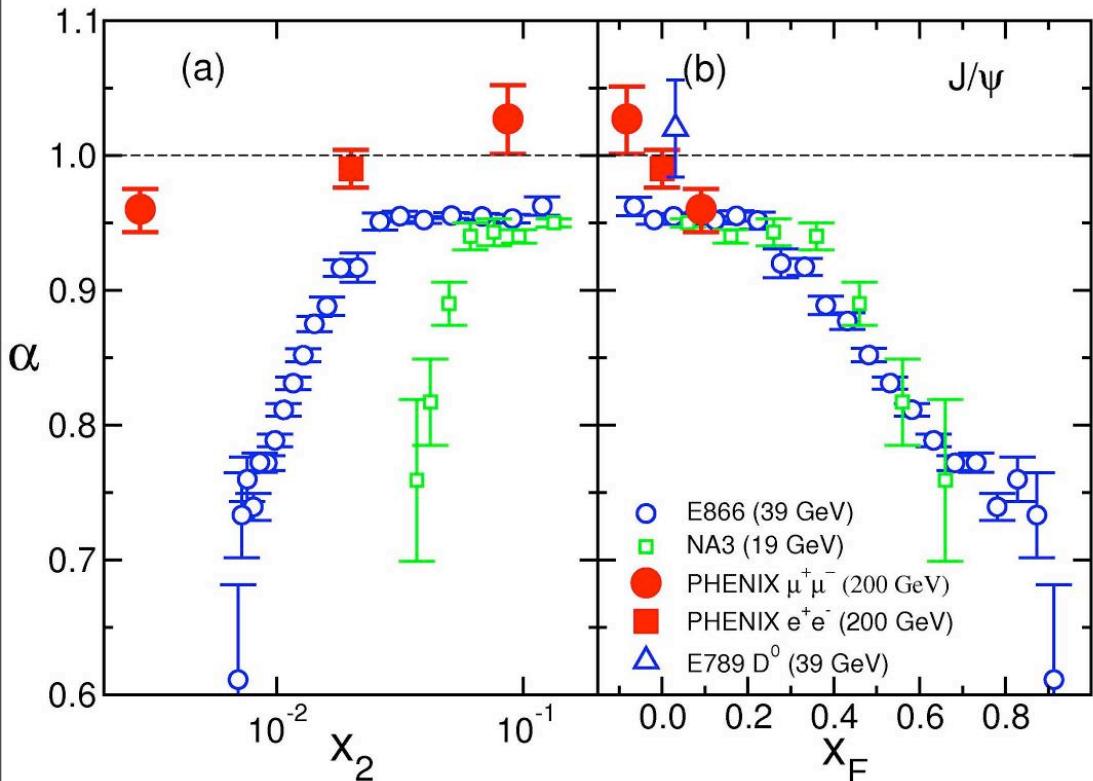
Path @ RHIC :  $d_8(x_F \sim 0) \approx 26 \text{ fm}$      $R_{Au} = 6.38 \text{ fm}$



Most of Cold Nuclear Effects  
probed with p+Au or d+Au

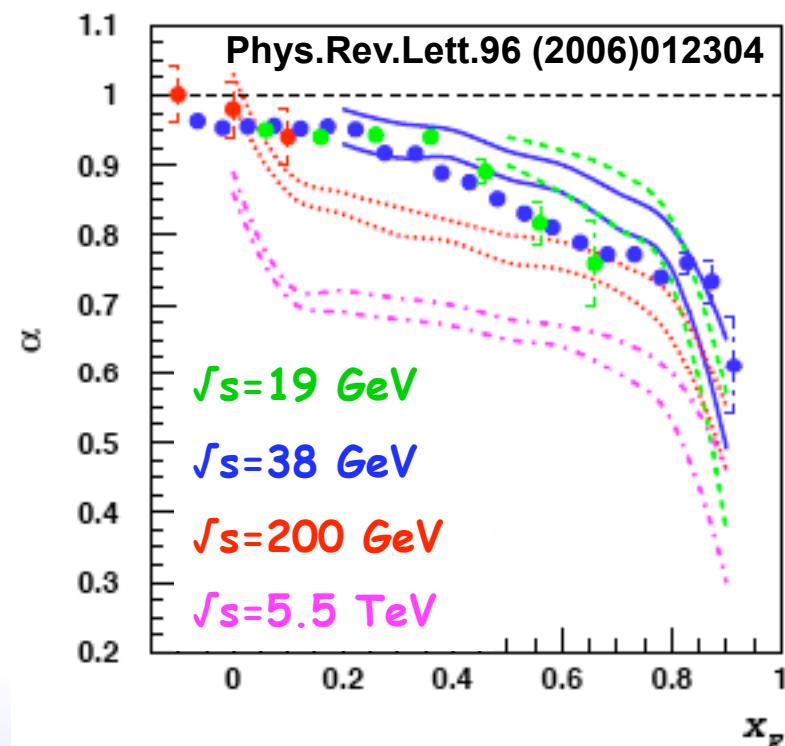
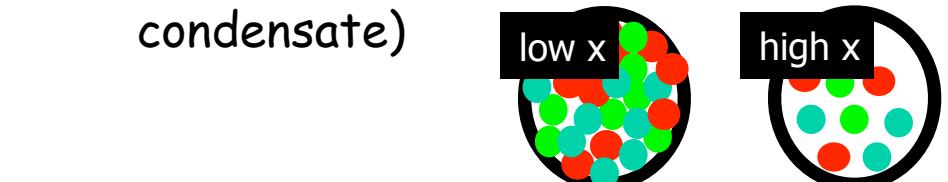


# Cold Nuclear Effects II

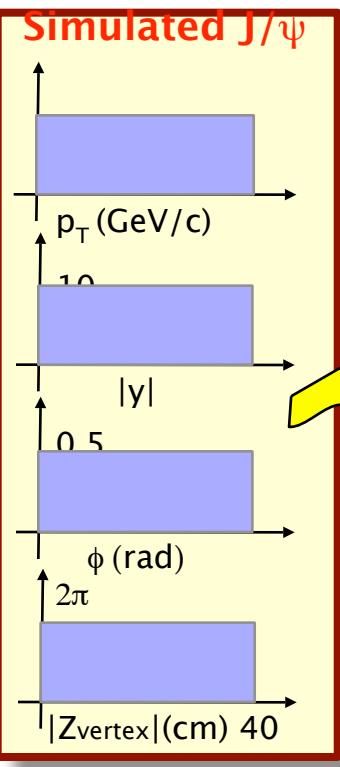


Absorption scales with  $x_F$ , but not in  $x_2$  (nucleus frame). Possible explanation :

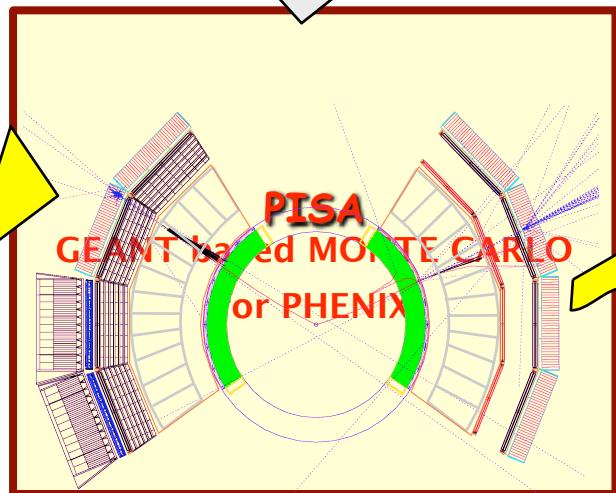
- limiting fragmentation caused by saturation of gluons (color glass condensate)



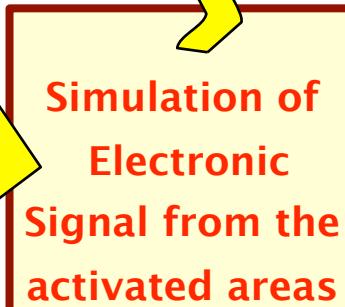
# Detector Efficiency Estimation



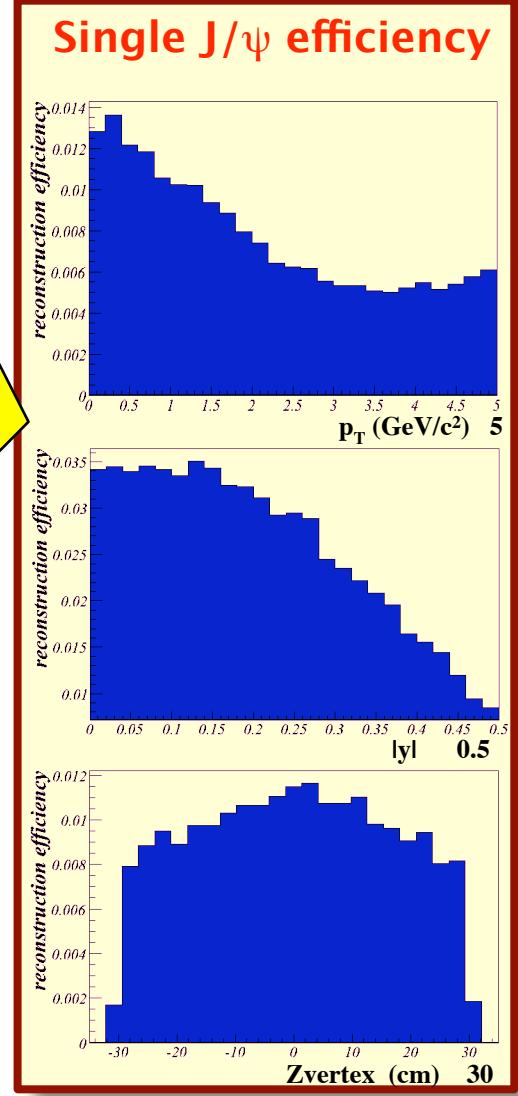
- Geometry
- Material composition



- Tuning calibration to match  $\gamma \rightarrow e^+e^-$  response



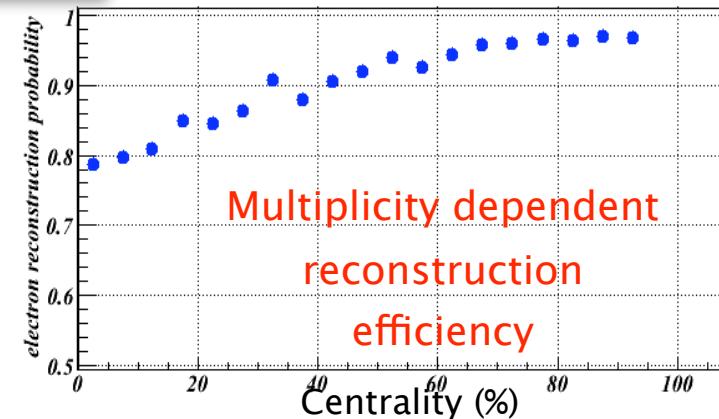
- Dead areas
- Gains



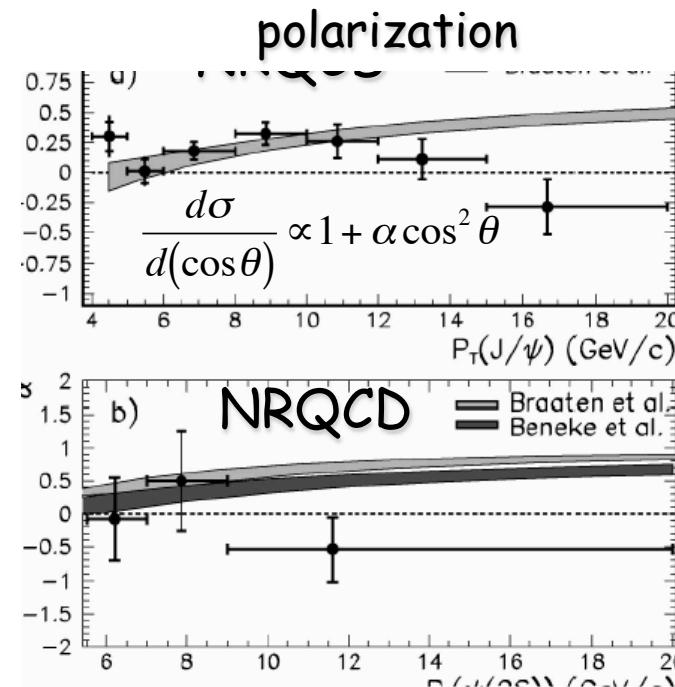
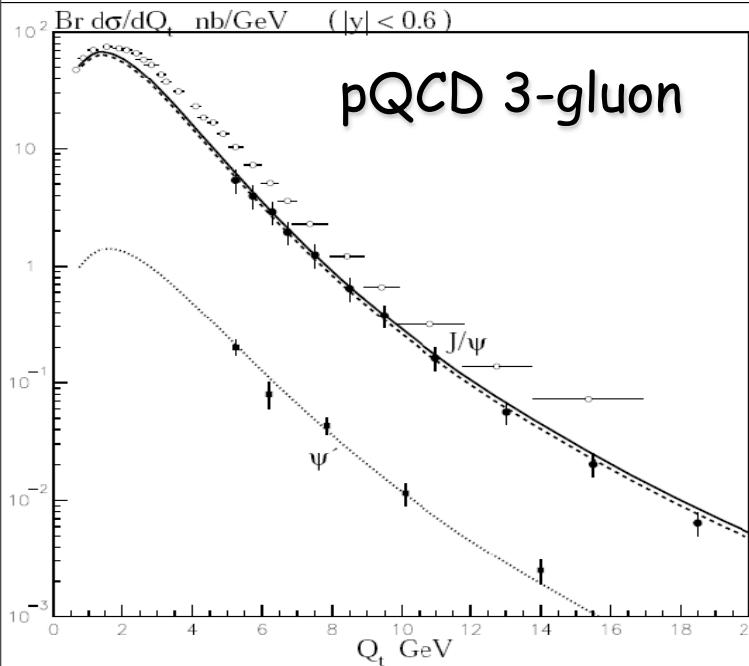
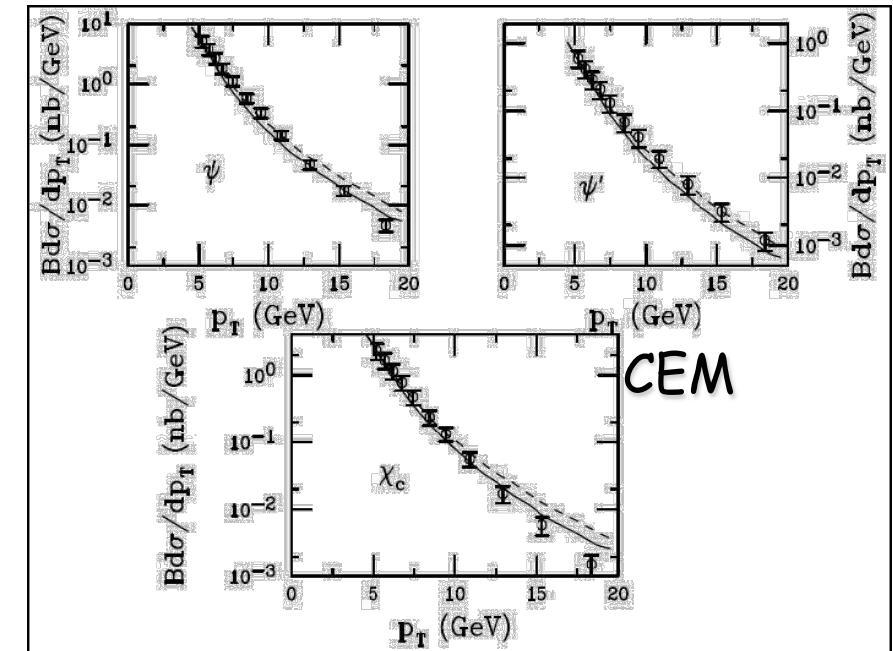
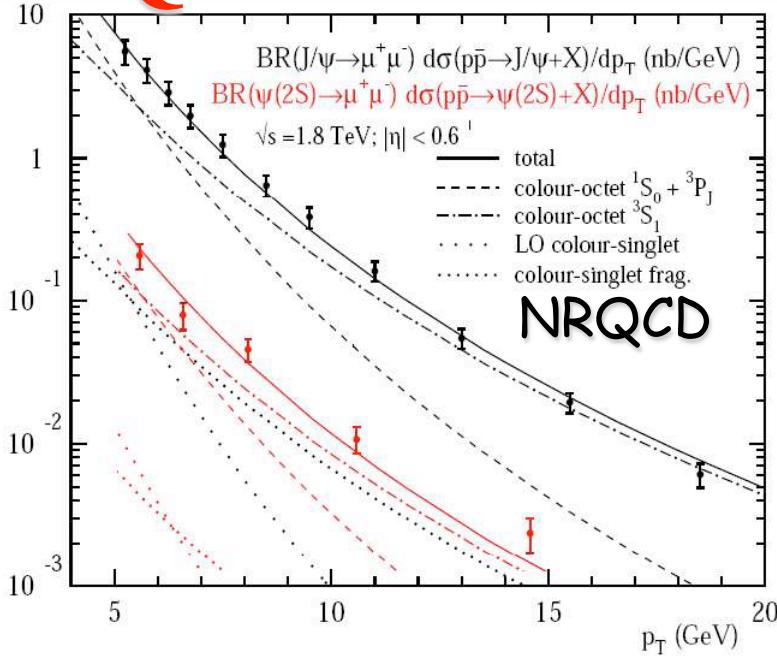
**Minimum Bias Real Data**

**Raw data embedding**

**Event reconstruction**



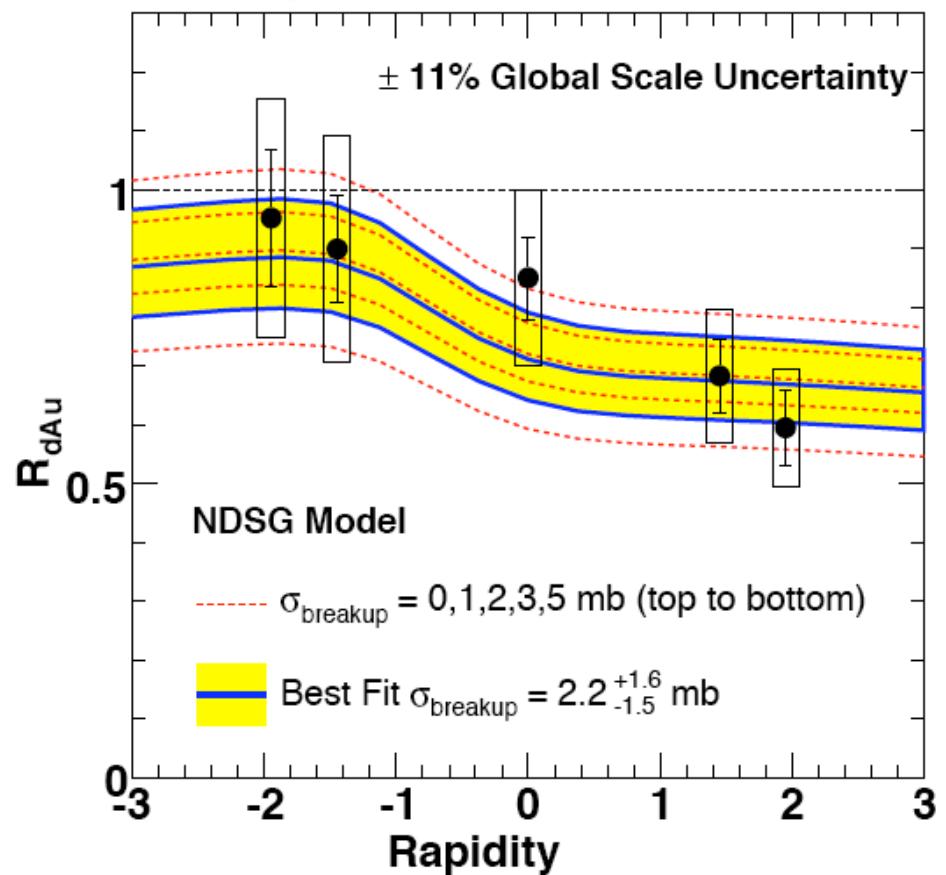
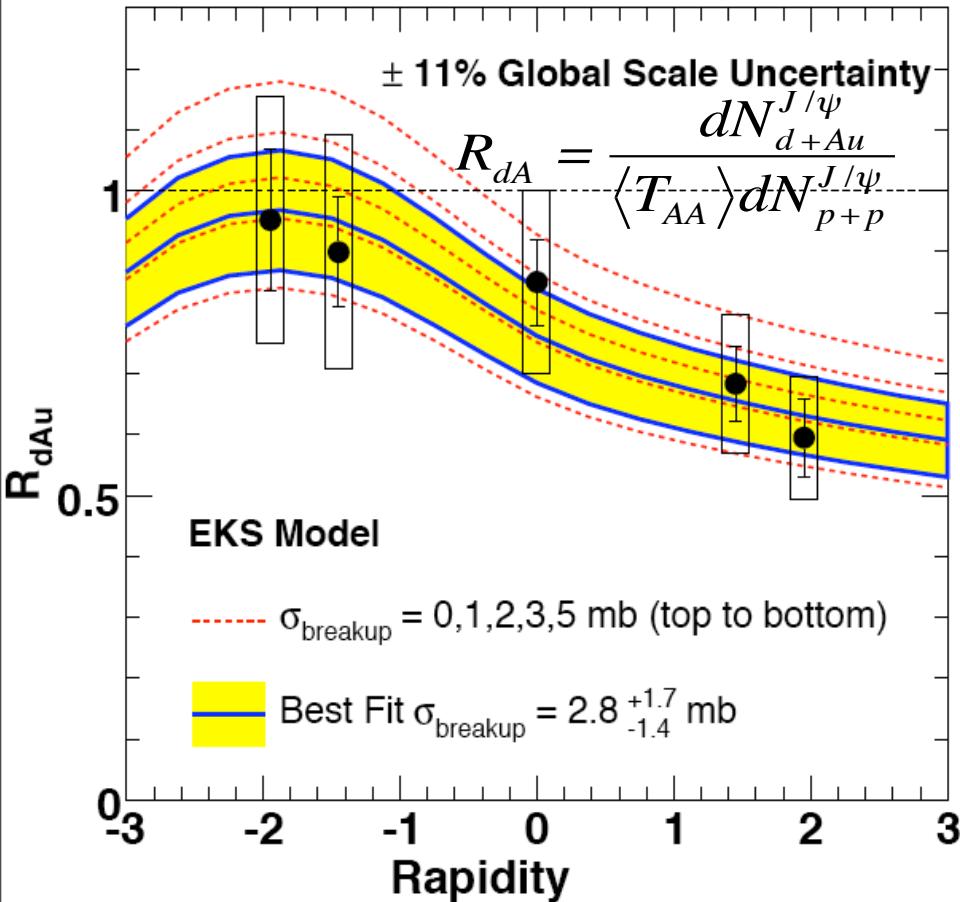
# Quarkonia Production in Tevatron



All models can describe Tevatron p+p-> $J/\psi + X$  cross section @  $s^{1/2} = 1.8 \text{ TeV}$

Only pQCD predict longitudinal polarization for high  $p_T$

# d+Au → J/ψ

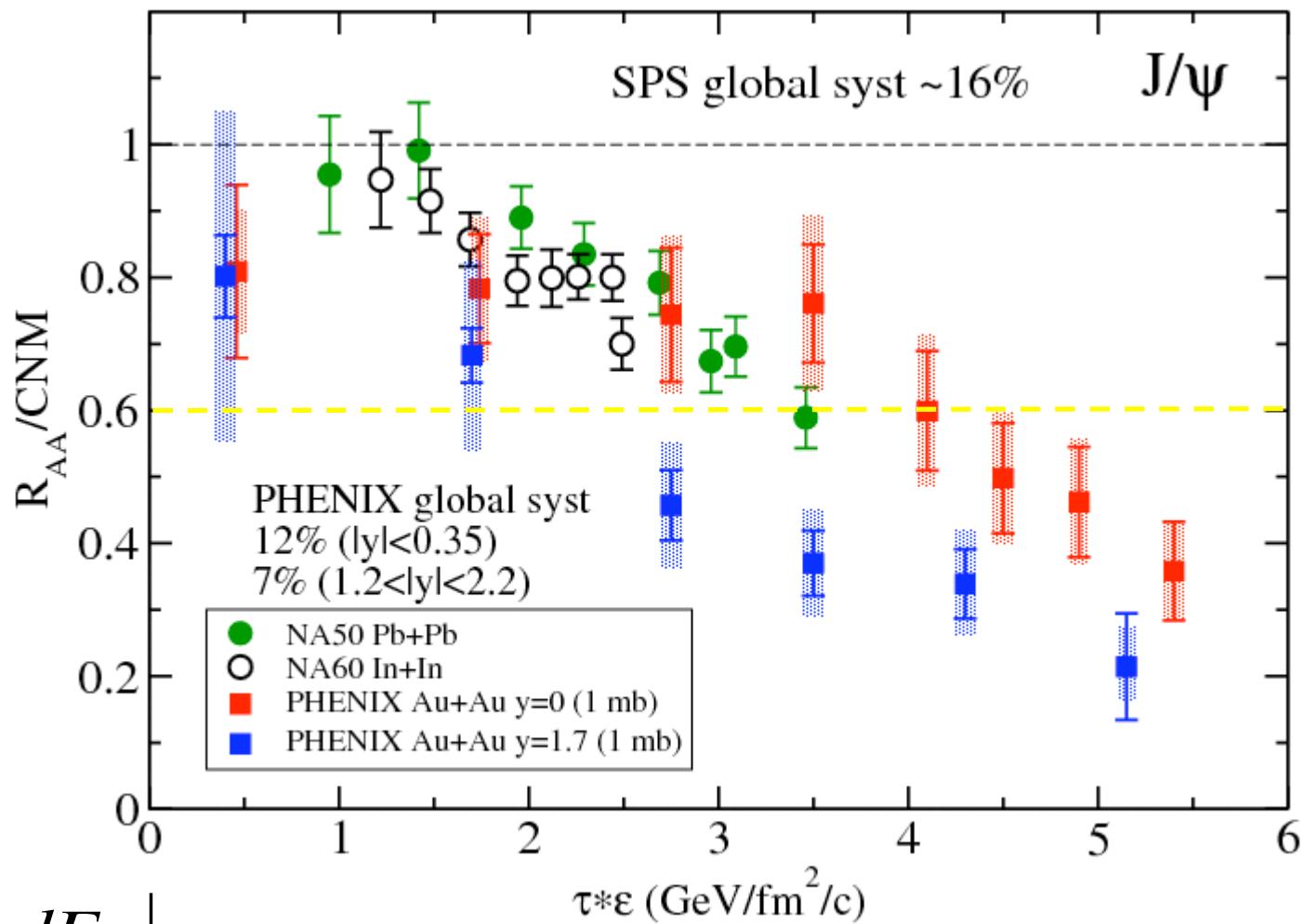


**EKS:** K. J. Eskola et al., Nucl. Phys. A696, 729 (2001)

**NDSG:** D. deFlorian and R. Sassot, Phys. Rev. D 69, 074028(2004).

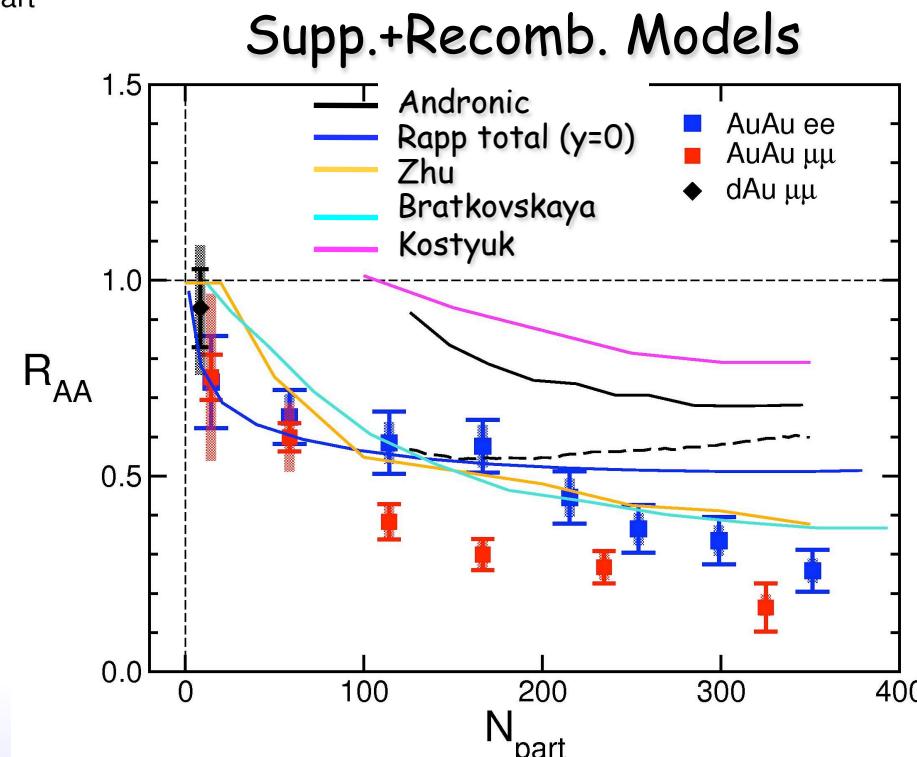
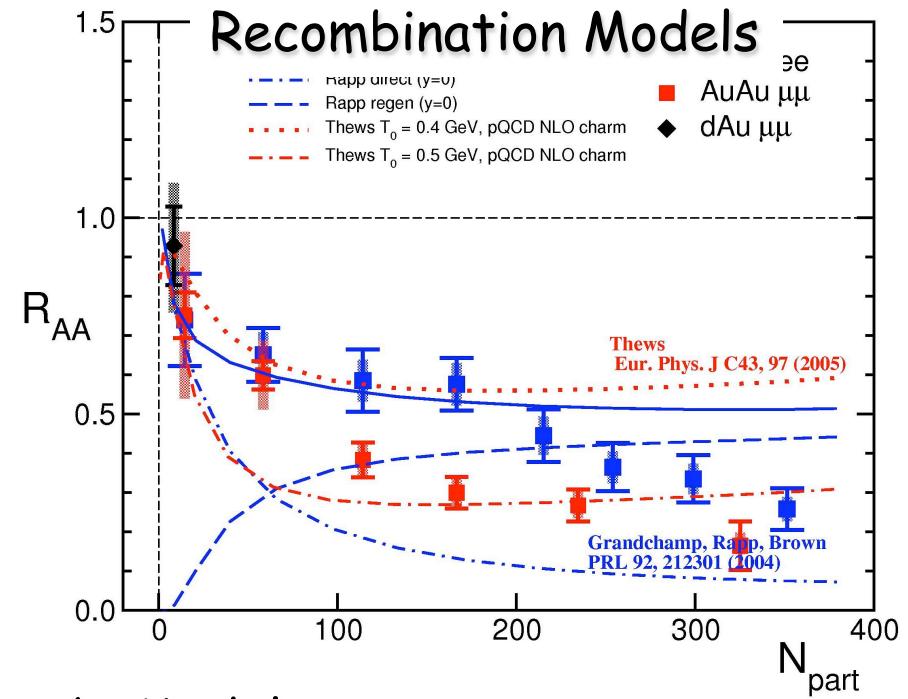
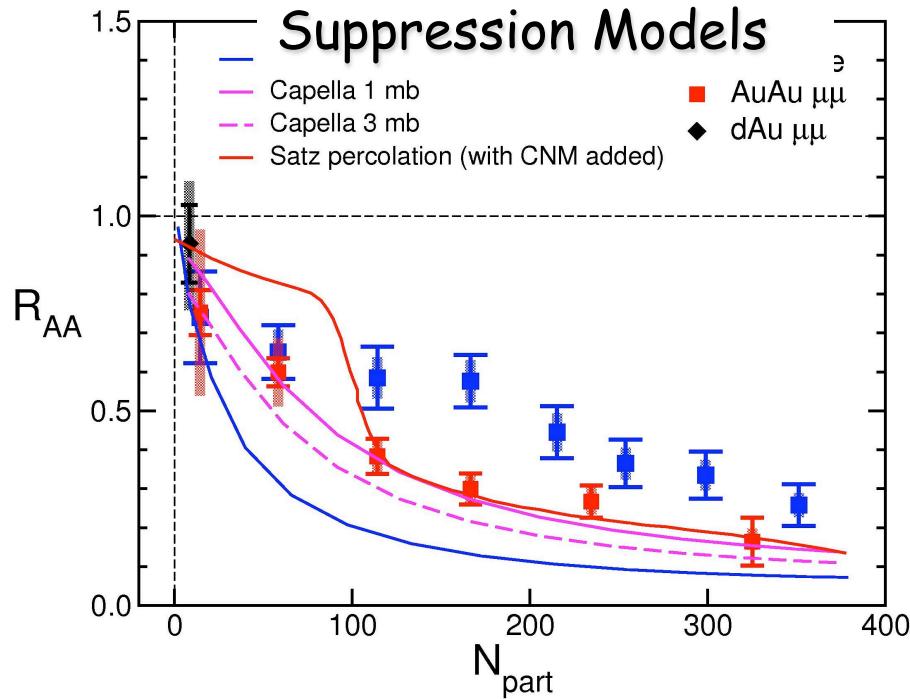
Modest absorption at RHIC energies.  
Brand new d+Au data with 30x more statistics will constrain CNM models.

# Sequential Screening



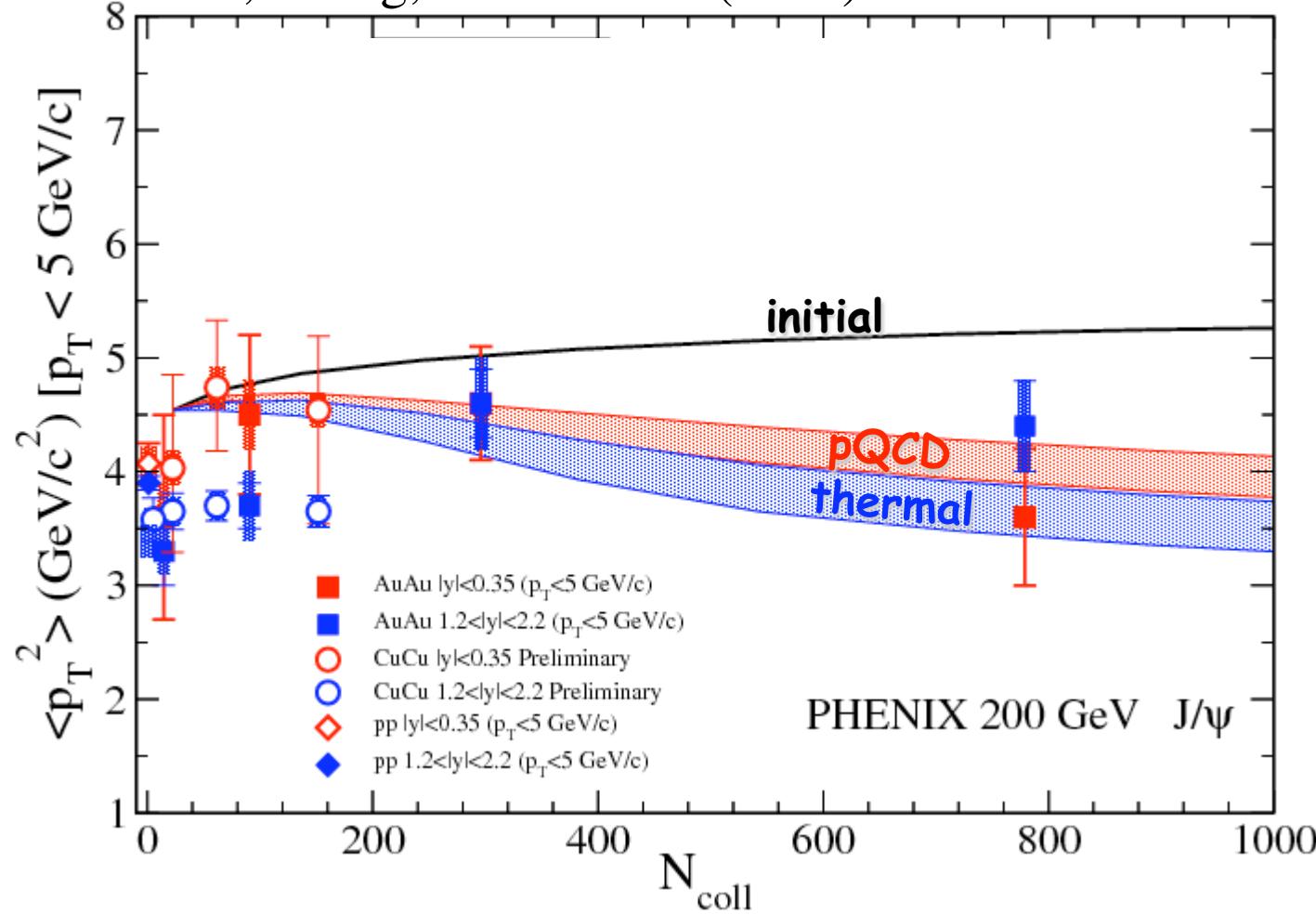
$$\varepsilon = \frac{1}{\tau_0 \pi R^2} \left. \frac{dE_T}{dy} \right|_{y=0}$$

# Models to Au+Au collisions



$\langle p_T^2 \rangle$

Yan, Zhang, Xu - PRL97 (2006) 232301.



$\langle p_T^2 \rangle$  is obtained directly to invariant  $p_T$  spectra for different centralities.  
Only  $p_T < 5$  GeV/c was accounted for in this calculation.