

On the critical end point of QCD

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1. **The Taylor expansion**: non-linear quark number susceptibilities
2. **Optimizing the computation**: the Steiner problem
3. **Digression**: simple phenomena at T_c concerning NLS
4. **Radius of convergence**: a limit on the end point

hep-lat/0412035 (PRD) hep-lat/0507023

The Taylor expansion for 2 flavours

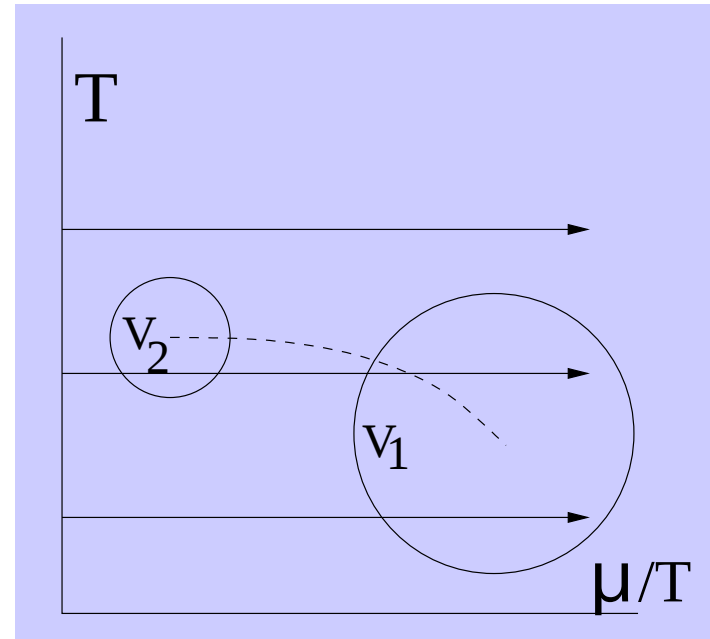
$$P(T, \mu_u, \mu_d) = \left(\frac{T}{V} \right) \log Z(T, \mu_u, \mu_d)$$

$$P(T, \mu_u, \mu_d) = P(T, 0, 0) + \sum_{n_u, n_d} \chi_{n_u, n_d} \frac{\mu_u^{n_u}}{n_u!} \frac{\mu_d^{n_d}}{n_d!}$$

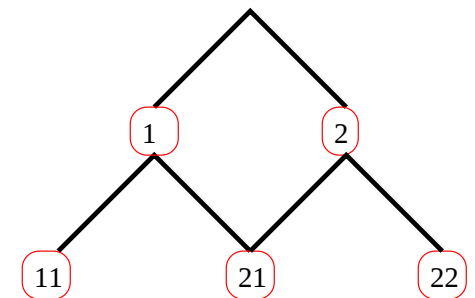
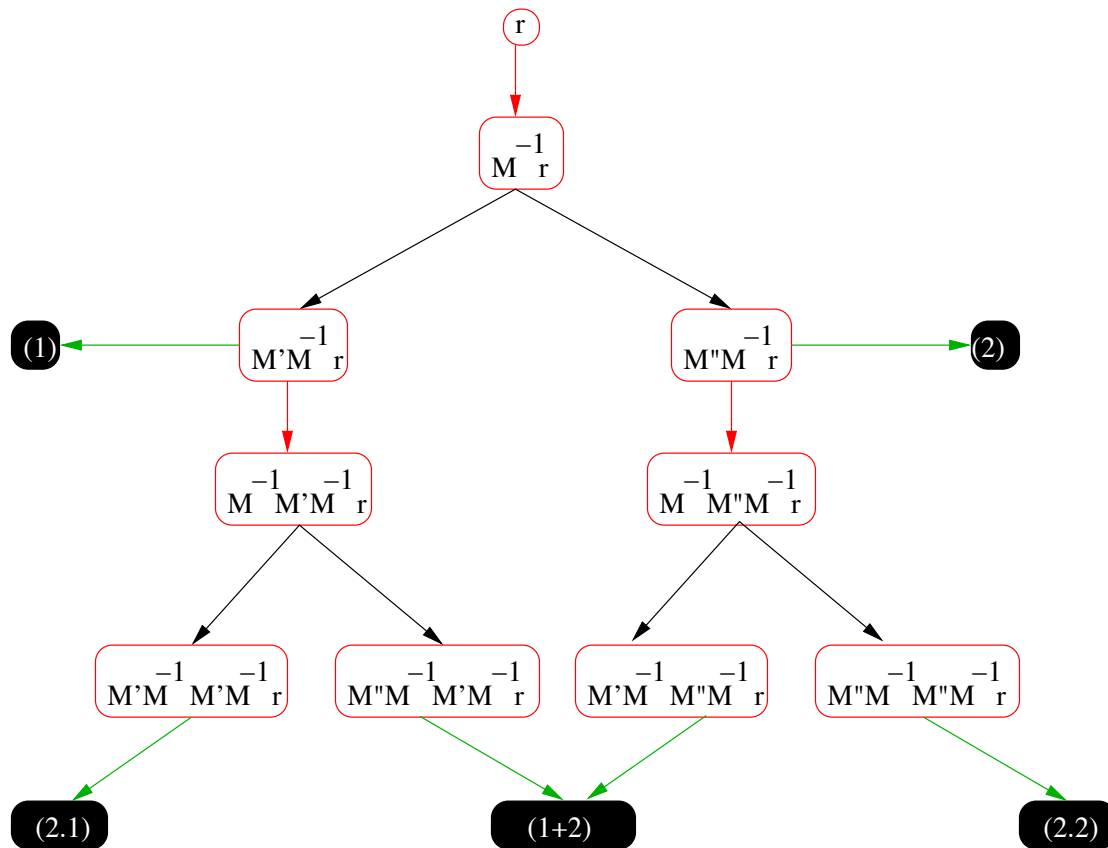
$m_u = m_d$ implies that $\chi_{n_u, n_d} = \chi_{n_d, n_u}$,
for any $\mu_u = \mu_d$. One QNS is

$$\chi_{20}(T, \mu_B) = \left. \frac{\partial^2 P(T, \mu_u, \mu_d)}{\partial \mu_u^2} \right|_{\mu_u = \mu_d = \mu_B/3}$$

$\chi_{20}(T^E, \mu_B^E)$ diverges in the infinite
volume limit: pseudo critical behaviour
at finite volumes.

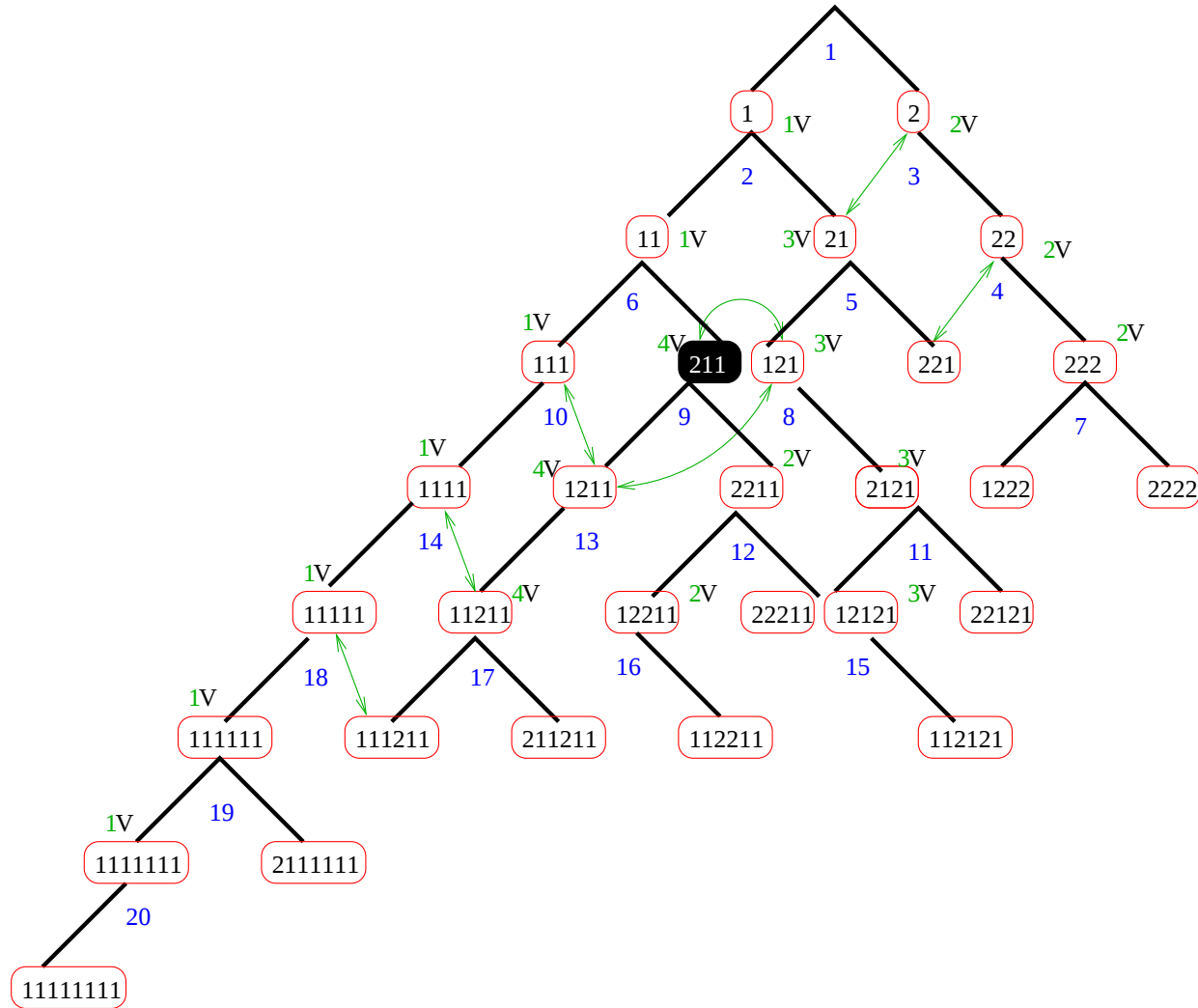


Steiner problem: efficient stochastic evaluation of traces



Charikar et al, STAN-CS-TN-97-56.

The Steiner tree for order 8 Taylor expansion

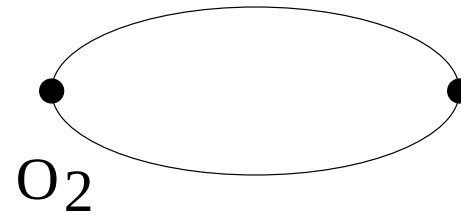
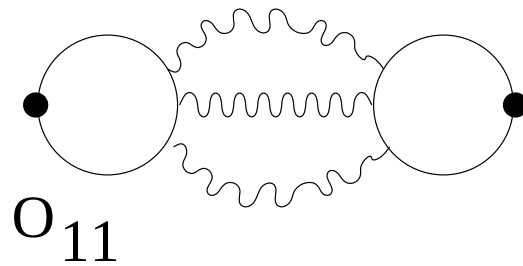
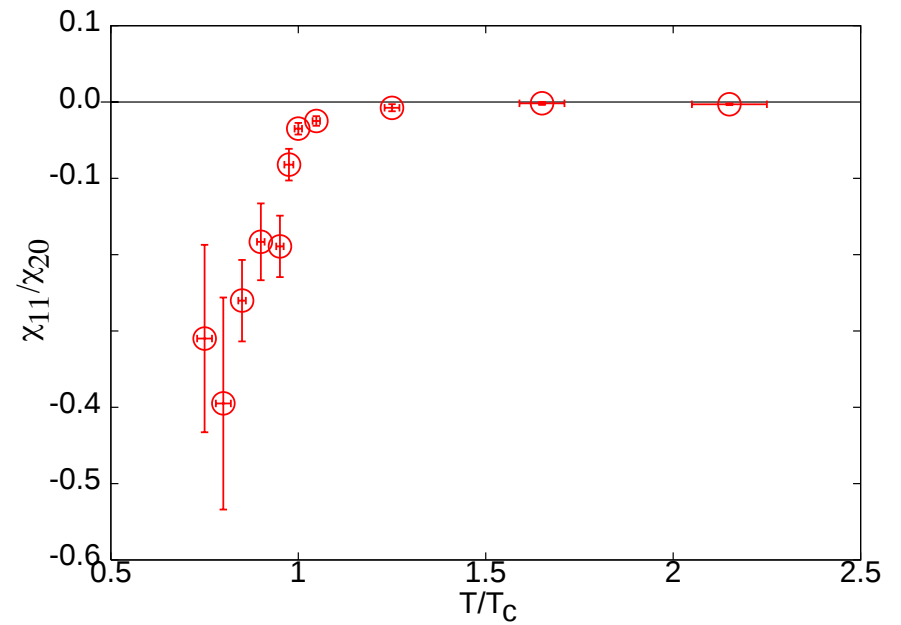
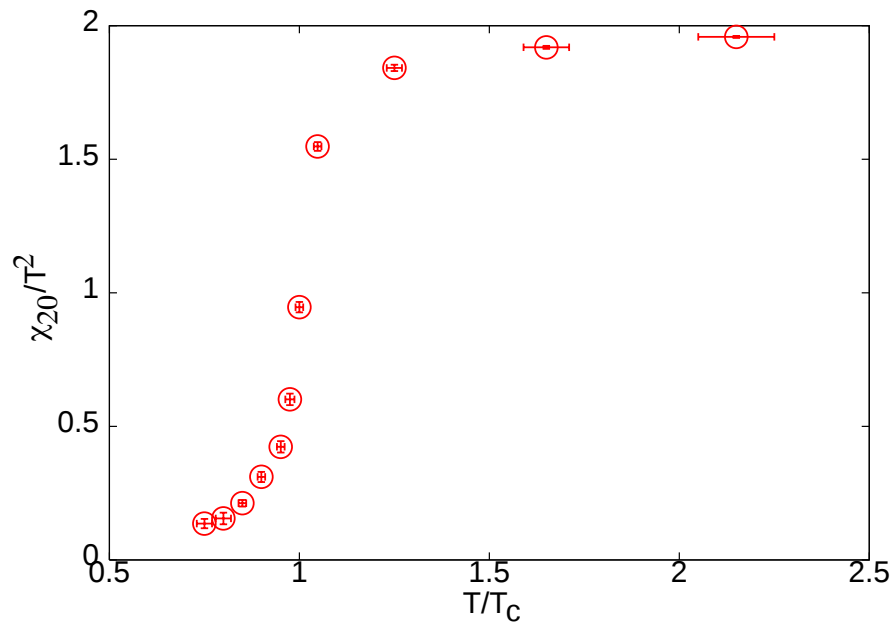


Simulation parameters

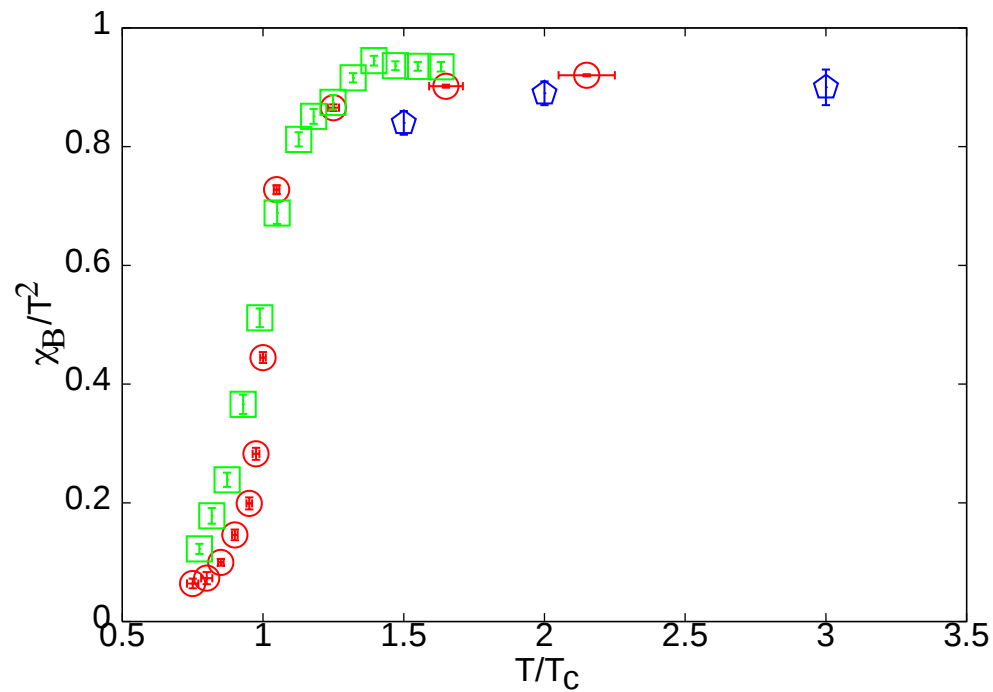
1. Taylor expansion to 8th order
2. $N_t = 4$: $m_\rho/T_c = 5.6$, $m_\pi/m_\rho = 0.3$ (Fodor Katz, de Forcrand Philipsen)
3. N_s : $3.3 \leq m_\pi V^{1/3} \leq 10.0$ (interpolates to Bielefeld-Swansea)
4. Statistics: more than $50\tau_{int}$ of configs at each T and V

$$\begin{aligned} \frac{\chi_{20}}{m_\rho^2} &= \frac{T}{m_\rho^2 V} \left[\sum_{ij} \frac{|\langle i|\gamma_0|j\rangle|^2}{\lambda_i \lambda_j} + \left(\sum_i \frac{\langle i|\gamma_0|i\rangle}{\lambda_i} \right)^2 \right] \\ &= \frac{(m_\pi/m_\rho)^2}{m_\pi^3 V} \left[\sum_{ij} |\langle i|\gamma_0|j\rangle|^2 \frac{T m_\pi}{\lambda_i \lambda_j} + \left(\sum_i \langle i|\gamma_0|i\rangle \frac{\sqrt{T m_\pi}}{\lambda_i} \right)^2 \right] \end{aligned}$$

The QNS: sign fluctuations



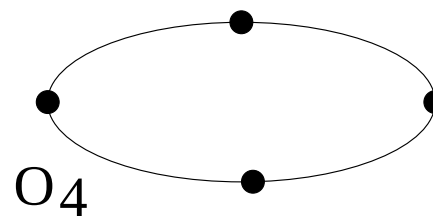
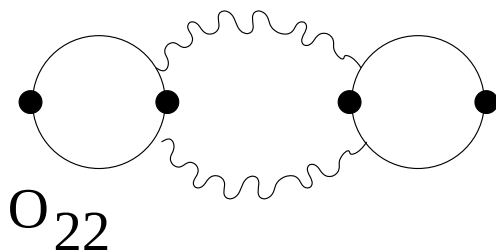
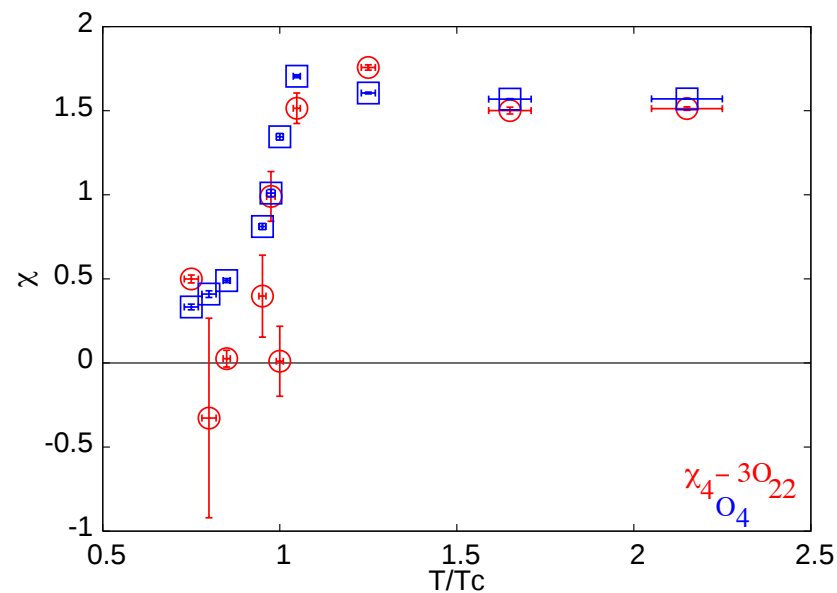
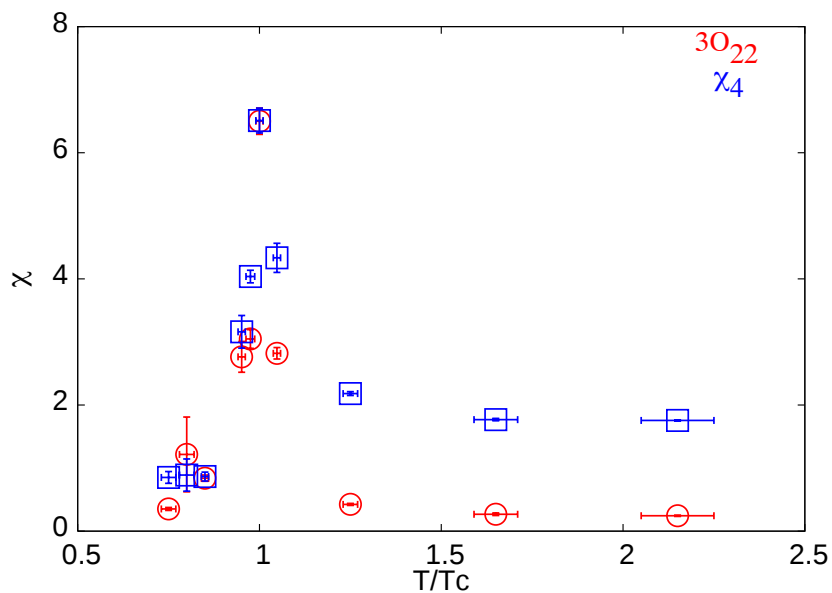
Fluctuations



Quenched, MILC $N_t = 8$, this work, scaled by quenched

Weak coupling: Blaizot et al, Mustafa et al, Vuorinen

4th order NLS



... and similar systematics at higher orders

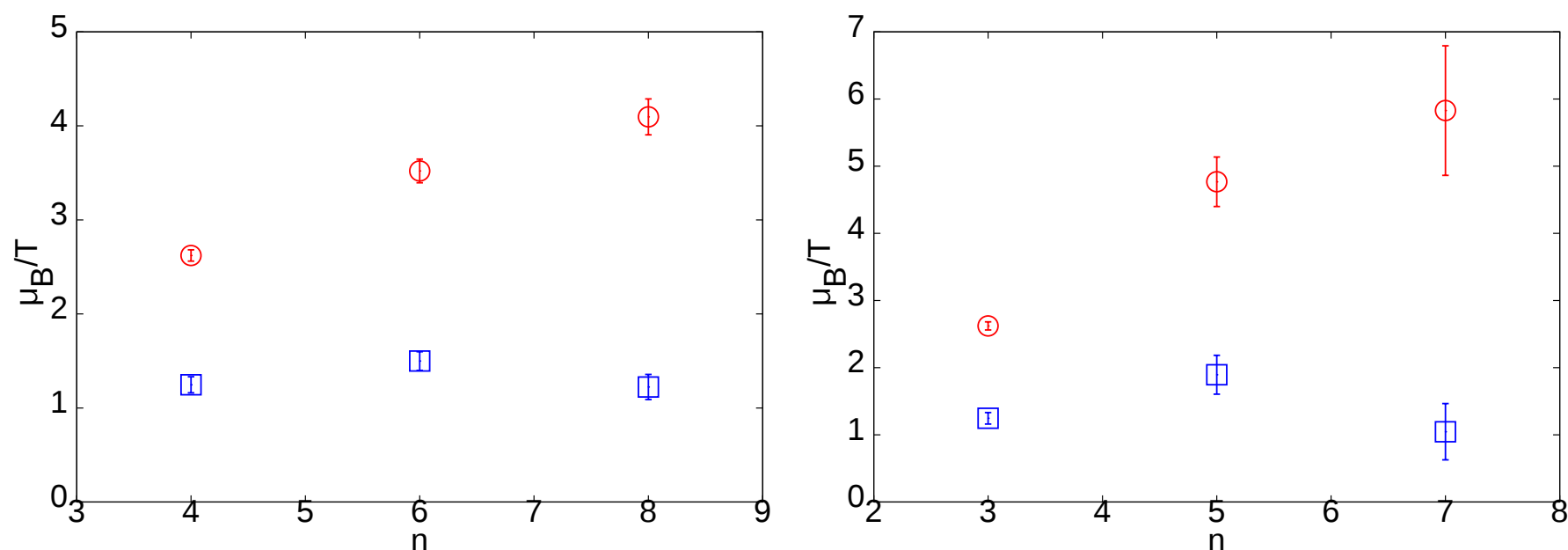
Systematics of susceptibilities

1. $T \geq 1.5T_c$ perturbative power counting roughly sets the right scale of various operators. More detailed tests, comparisons with weak-coupling theory, might be useful.
2. $T \leq 0.9T_c$ most NLS of given order are of the same order of magnitude. No special simplification.
3. **neighbourhood of T_c** Susceptibilities of the composite operator O_2 dominate and cause all the peaks that are seen. Shapes are consistent with $(T/V)\langle O_2 \rangle$ behaving as an order parameter.

Volume dependence under control at fixed mass. Mass dependence? Chiral limit? Is there an effective theory of O_2 near T_c ? What is the role of O_{11} in the chiral limit?

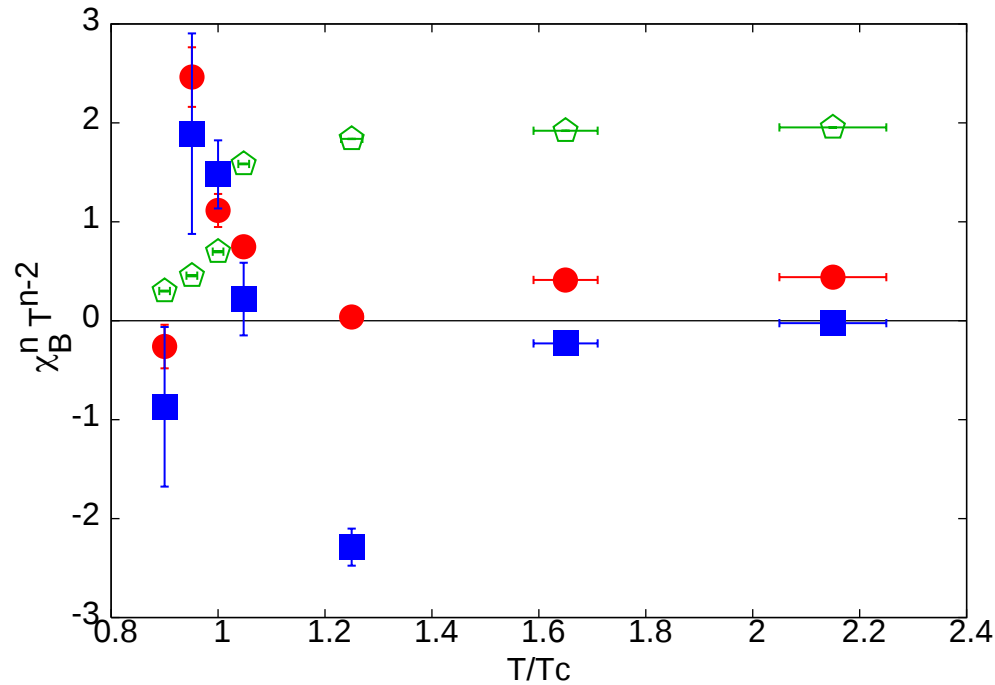
Radius of convergence

Notation: If $f(x) = \sum_n f_{2n} x^{2n}$ then $\rho_{2n} = \left| \frac{f_0}{f_{2n}} \right|^{1/2n}$ and $r_{2n} = \sqrt{\left| \frac{f_{2n}}{f_{2n+2}} \right|}$



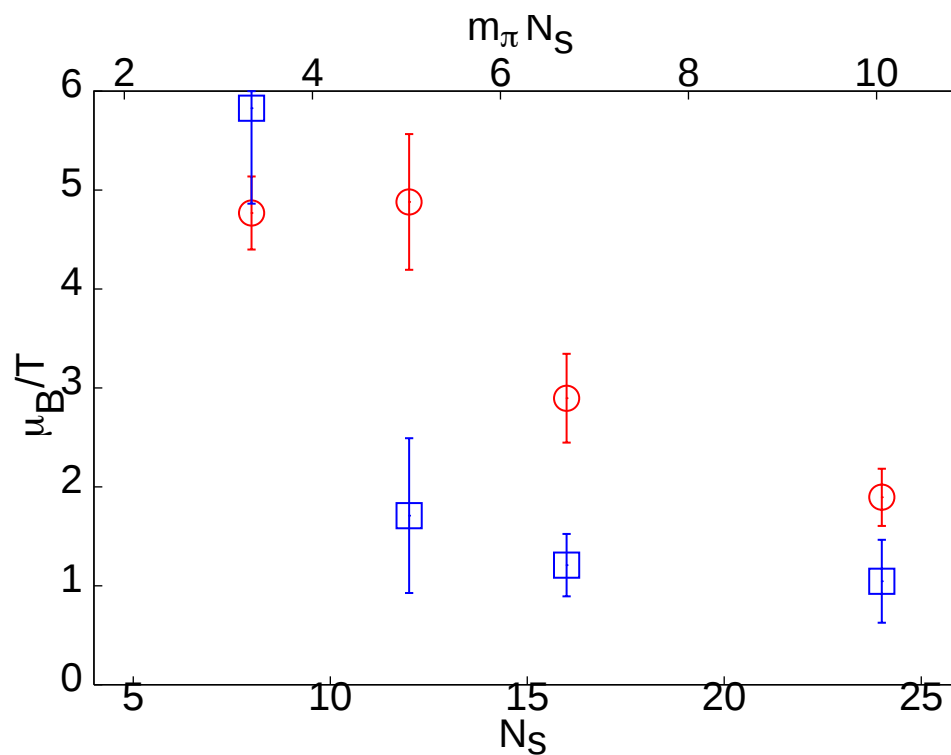
Pseudocritical behaviour: larger volumes required at larger order

Sign and convergence of subexpressions



4×24^3 lattice: χ_B^0/T^2 , $\chi_B^4/2$, $\chi_B^6 T^2/12$.

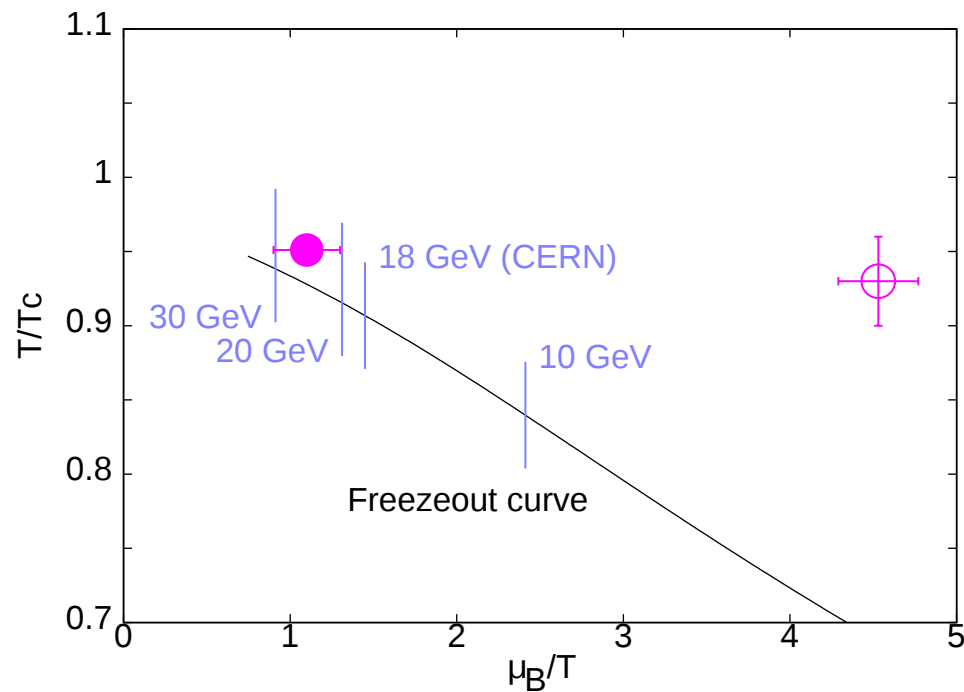
Finite size scaling



$T/T_c = 0.95$: r_4 (circles), r_6 (boxes).

FS shift in μ_B^E estimated assuming CEP is in Ising universality class.

Summary: Critical end point



Volume effects under control. Examine mass dependence and continuum limit. Use more volumes to refine FSS and find critical exponents.