

# Fluctuations of conserved quantities

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TIFR  
From Strings to LHC II

December 20, 2007

# 1 Introduction

# 2 Fluctuations

# 3 Critical end point

# 4 Linkage

# 5 Summary

# General conditions in heavy-ion collisions

Fluctuations  
of conserved  
quantities

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to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

- In a heavy-ion collision, almost all quantum numbers pass through without collision: valence quarks are fast particles, QCD is asymptotically free.
- When  $\sqrt{S}$  is large, the central rapidity region contains significant energy with almost no flavour quantum numbers: due essentially to soft gluon and sea quark component of hadron wavefunction.
- This fireball expands and cools: strong evidence for collective (final state) effects in measurements of elliptic flow and jet quenching.
- Is there a QCD plasma formed in this region? A quantitative question: what are the properties of the plasma? Was this cup of water ever brought to boil?

# Ensembles

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of conserved  
quantities

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Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

- **Event ensembles at colliders:** defined through various selection criteria. Values of conserved quantities can fluctuate from one event to another
- **Thermodynamic ensembles:** defined by letting certain conserved quantities fluctuate (canonical) while keeping others fixed (micro-canonical).
- **Mapping between the two:** The system is a small part of a big (ion-ion) collider event, the heat-bath is the remainder. Is the event a thermostat? Need experimental check.

# Is a collider event a thermostat?

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From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

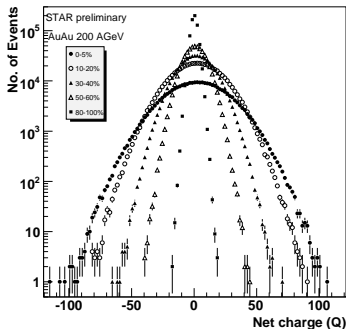
Summary

- **Energy average over events:**  $m_T$  stands in for energy, and has an exponential distribution. No other measurement of temperature performed yet. Jets at LHC could measure temperature.
- **Other conserved quantities:** If there are spatial gradients of densities, then not a thermostat. No evidence for a spatial gradient. Diffusion time could be much smaller than the lifetime of fireball. (Lattice, toy models: some “solvable” using AdS/CFT)

# STAR Collaboration

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Distribution of net charge is Gaussian ( $p_T < 1$  GeV, central rapidity slice).

Nayak (STAR) nucl-ex/0608021

# Susceptibilities

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of conserved  
quantities

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TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

If  $P(T, \mu, \dots)$  is the pressure, then a susceptibility is

$$\chi = \frac{\partial n}{\partial \mu}, \quad \text{where} \quad n = \frac{\partial P}{\partial \mu}$$

Gottlieb, Liu, Toussaint, Renken, Sugar, Phys.Rev.Lett. 59, 2247, 1987.

For each conserved charge one chemical potential:  $\mu_B, \mu_Q, \mu_S$  etc.. Two kinds of susceptibilities

$$\chi_i = \frac{\partial^2 P}{\partial \mu_i^2} \qquad \chi_{ij} = \frac{\partial^2 P}{\partial \mu_i \partial \mu_j}$$

diagonal and off-diagonal. Higher derivatives are called non-linear susceptibilities. Gavai and SG, Phys.Rev. D68, 034506, 2003.

# Diagonal susceptibilities

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of conserved  
quantities

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From Strings  
to LHC II

Outline

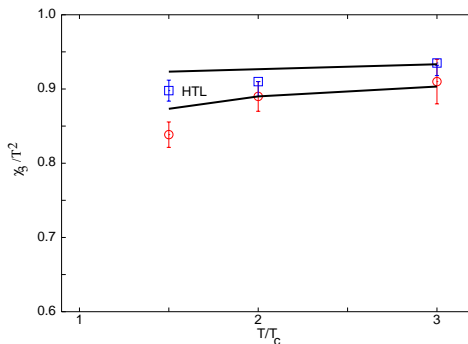
Introduction

Fluctuations

Critical end  
point

Linkage

Summary



In the high temperature phase of QCD the diagonal susceptibilities (but not the off-diagonal) agree with a resummed weak coupling computation.

Blaizot, Iancu, Rebhan, Phys.Lett. B523, 143, 2001

Gavai and SG Phys.Rev. D67, 034501, 2003



# Thermodynamic fluctuations

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$$P(Q) \propto \exp\left(-\frac{Q^2}{2VT\chi_Q}\right), \quad \text{so} \quad \langle \Delta Q^2 \rangle = VT\chi_Q$$

Bias free experimental measurement of  $\chi_Q$  possible.

Asakawa, Heinz and Muller, Phys.Rev.Lett. 85, 2072, 2000;

Jeon and Koch, Phys.Rev.Lett. 85, 2076, 2000.

$$P(B, Q) \propto \exp\left(-\frac{Q^2}{2VT\chi_Q} - \frac{B^2}{2VT\chi_B} - \frac{QB}{VT\chi_{BQ}}\right)$$

Variances determined by diagonal susceptibilities, covariances by off-diagonal susceptibilities.

# What do fluctuations measure?

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From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

## In thermodynamics

Fluctuations of energy measures specific heat ( $c_v$ ), of volume measures compressibility ( $\kappa$ ), of charge measures susceptibility ( $\chi$ ). Fluctuations are usually Gaussian.

# What do fluctuations measure?

## Fluctuations of conserved quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

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## At colliders

Many sources of fluctuations: incomplete characterization of event, beam energy fluctuations, ion charge fluctuations, volume, detector noise: all are Gaussian. Hard to disentangle from thermodynamics without proper event characterization.

# Baby and bathwater

## Fluctuations of conserved quantities

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From Strings  
to LHC II

Outline

Introduction

**Fluctuations**

Critical end  
point

Linkage

Summary

- One popular analysis trick is to construct variables which remove all Gaussian fluctuations. This removes all thermodynamical information as well.

# Baby and bathwater

## Fluctuations of conserved quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

- One popular analysis trick is to construct variables which remove all Gaussian fluctuations. This removes all thermodynamical information as well.
- Useful in the context of multi-particle hadron dynamics where one wants to understand correlations between particles.

# Baby and bathwater

## Fluctuations of conserved quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

- One popular analysis trick is to construct variables which remove all Gaussian fluctuations. This removes all thermodynamical information as well.
- Useful in the context of multi-particle hadron dynamics where one wants to understand correlations between particles.
- No use in thermodynamics, where correlations are irrelevant: no matter what the dynamics is (almost) always get a Gaussian. Physics is in the width of the Gaussian.

# Baby and bathwater

## Fluctuations of conserved quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

### Outline

### Introduction

### Fluctuations

### Critical end point

### Linkage

### Summary

- One popular analysis trick is to construct variables which remove all Gaussian fluctuations. This removes all thermodynamical information as well.
- Useful in the context of multi-particle hadron dynamics where one wants to understand correlations between particles.
- No use in thermodynamics, where correlations are irrelevant: no matter what the dynamics is (almost) always get a Gaussian. Physics is in the width of the Gaussian.
- Either find non-Gaussian thermodynamics or eliminate dependence on uncontrolled variables.

# Use non-Gaussian fluctuations



# Phase diagram of QCD

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quantities

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Outline

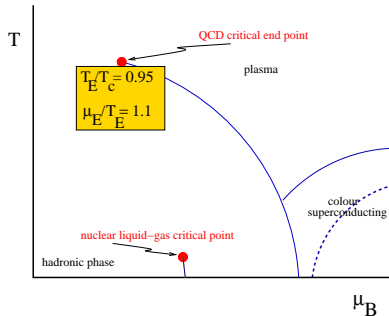
Introduction

Fluctuations

Critical end  
point

Linkage

Summary



At a critical point fluctuations are non-gaussian.

Stephanov, Rajagopal, Shuryak, Phys.Rev. D60, 114028, 1999

# Sign problem

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of conserved  
quantities

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TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

## Monte Carlo

$\text{Det } D$  is real if  $QDQ^\dagger = D^\dagger$  for unitary  $Q$ .  $Q = \gamma_5$  for the Dirac operator (free and minimally coupled to a gauge field). Hence Monte Carlo evaluation of partition function possible.

# Sign problem

Fluctuations  
of conserved  
quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

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# Sign problem

Fluctuations  
of conserved  
quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

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## Taylor expansion

Phase diagram explored using Taylor expansion of free energy around  $\mu = 0$ . Taylor coefficients are non-linear susceptibilities. Taylor series fails to converge at the phase boundary: **non-gaussian fluctuations**. Gavai and SG, Phys.Rev. D71, 114014, 2005

# Non-Gaussianity

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quantities

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From Strings  
to LHC II

Outline

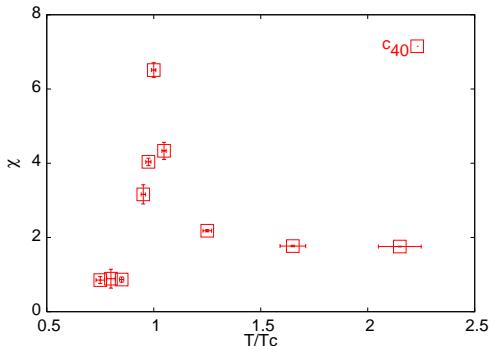
Introduction

Fluctuations

Critical end  
point

Linkage

Summary



Gvai, SG, Phys.Rev. D72, 054006, 2005

Strong peak in the fourth order susceptibility near  $T_c$  signals strongly non-Gaussian fluctuations. Residual non-Gaussian behaviour in the high-temperature phase: true even for an ideal gas.

# Eliminate uncontrolled variables

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

# Equation of state of hot matter

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of conserved  
quantities

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TIFR  
From Strings  
to LHC II

Outline

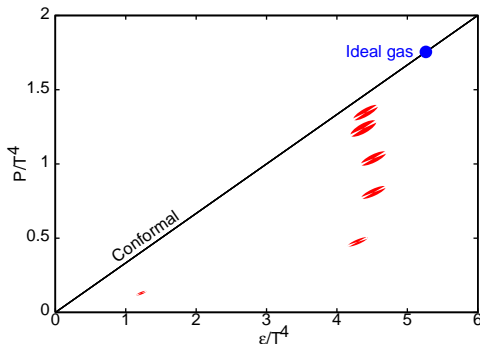
Introduction

Fluctuations

Critical end  
point

Linkage

Summary



Gavai, SG, Mukherjee, Phys.Rev. D71, 074013, 2005

The EOS of an ideal gas says how many degrees of freedom there are. In an interacting gas?

# Linkage

Fluctuations  
of conserved  
quantities

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TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

- **Linkage** between two quantum numbers:

$$C_{KL/L} = \frac{\langle KL \rangle - \langle K \rangle \langle L \rangle}{\langle L^2 \rangle - \langle L \rangle^2} = \frac{\chi_{KL}}{\chi_L}$$

- **Example:**  $C_{BS/S}$  is zero at low temperatures, because the lightest strange particle is the K, which has no baryon number. It is  $-1/3$  for quarks because  $S = 1$  goes with  $B = -1/3$ .
- **Robust:** systematic lattice effects cancel out in the ratio, systematic experimental effects also expected to reduced tremendously in the ratio: factors of VT cancel.



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Fluctuations  
of conserved  
quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

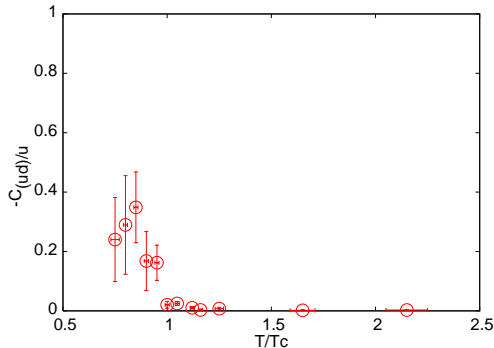
Introduction

Fluctuations

Critical end  
point

Linkage

Summary



Flavours are decoupled from each other: up and down shown,  
others similar. Gavai and SG, Phys.Rev. D73, 014004, 2006

# Quarks in matter

Fluctuations  
of conserved  
quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

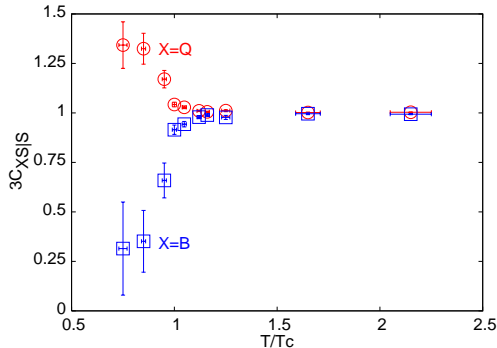
Introduction

Fluctuations

Critical end  
point

Linkage

Summary



Strangeness is carried by quarks. Similarly for up and down flavours. Gavai and SG, Phys.Rev. D73, 014004, 2006

# At the LHC

## Fluctuations of conserved quantities

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TIFR  
From Strings  
to LHC II

## Outline

## Introduction

## Fluctuations

## Critical end point

## Linkage

## Summary

- Measurement of linkage possible in principle since the factors of  $VT$  in two fluctuation measures cancel out. Need to characterize event volume and centrality in greater detail. (Possible: STAR preliminary)
- Influence of hadronization mechanism remains to be studied. Flow measurements ( $v_2$ ) throw light on this: quark coalescence dynamics seems to determine  $v_2$ .
- Hadronization will definitely disrupt certain linkages: for example,  $C_{BS|S}$ . However, if quark coalescence works then  $C_{QS|S}$  will not be disrupted. Such off-equilibrium linkages between flavour quantum numbers can be frozen in at the freezeout time.

# Summary

## Fluctuations of conserved quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

## Non-Gaussian fluctuations

Non-Gaussian fluctuations are visible near a phase transition. The QCD critical point can be studied through the energy ( $\sqrt{S}$ ) dependence of measures such as the kurtosis (**Binder cumulant**) of the event-to-event baryon number fluctuations.

# Summary

Fluctuations  
of conserved  
quantities

Sourendu  
Gupta  
TIFR  
From Strings  
to LHC II

Outline

Introduction

Fluctuations

Critical end  
point

Linkage

Summary

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## Linkage between fluctuations

Some of the uncertainties in going from observed fluctuations to QCD predictions cancel out in a robust measure called **Linkage**. Some linkages seem to persist during the hadronization of the fireball and can be directly used to look for quark-like quasiparticles in the plasma.