

Quark Recombination and Elliptic flow

Subrata Pal

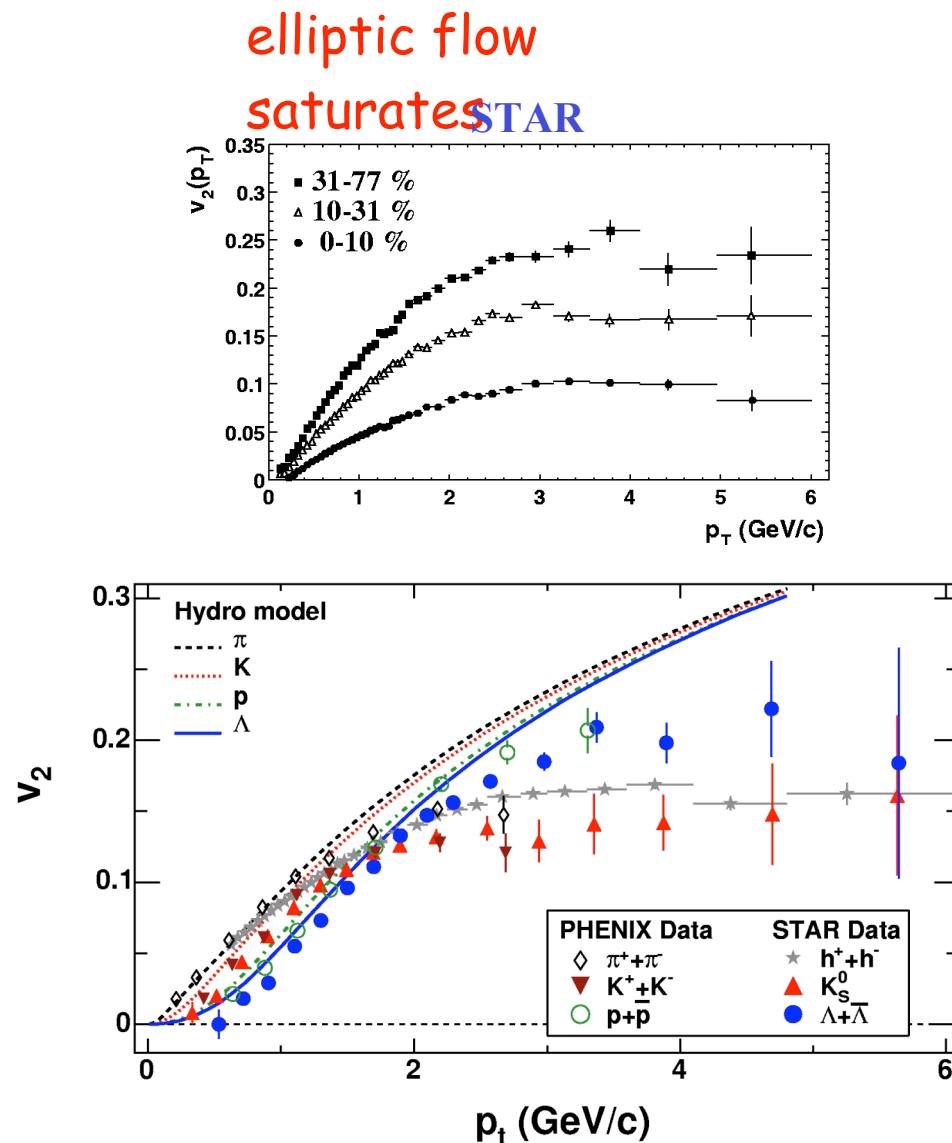
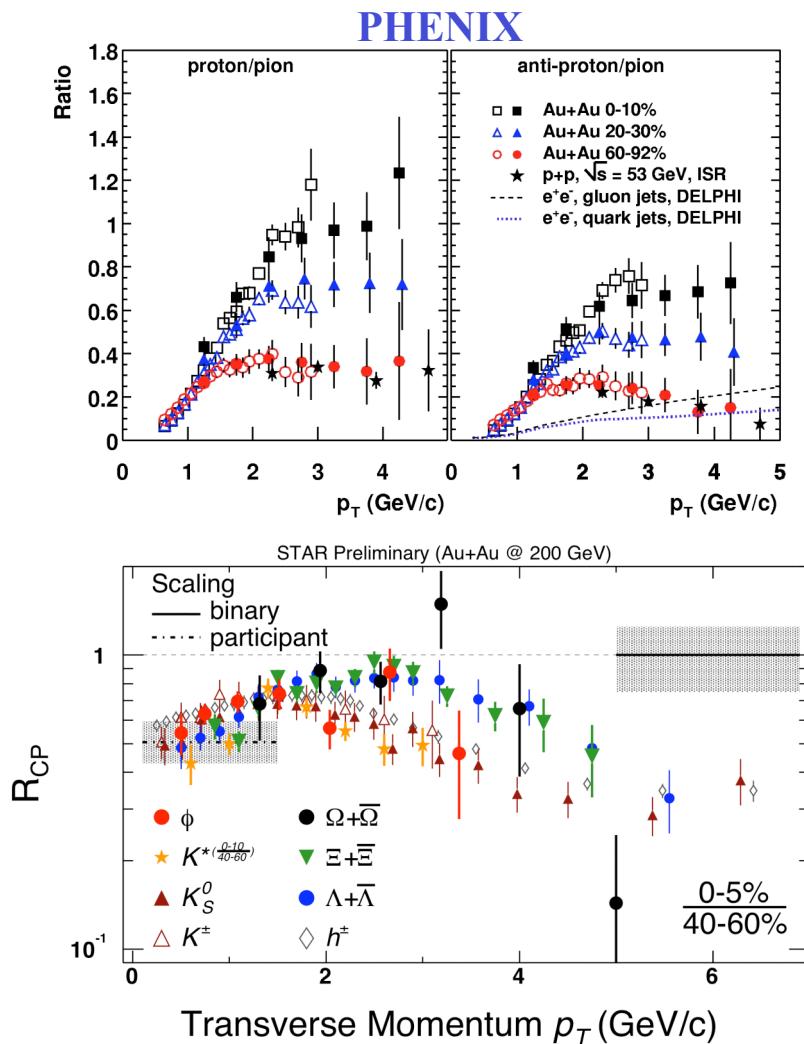
Tata Institute of Fundamental Research

Outline

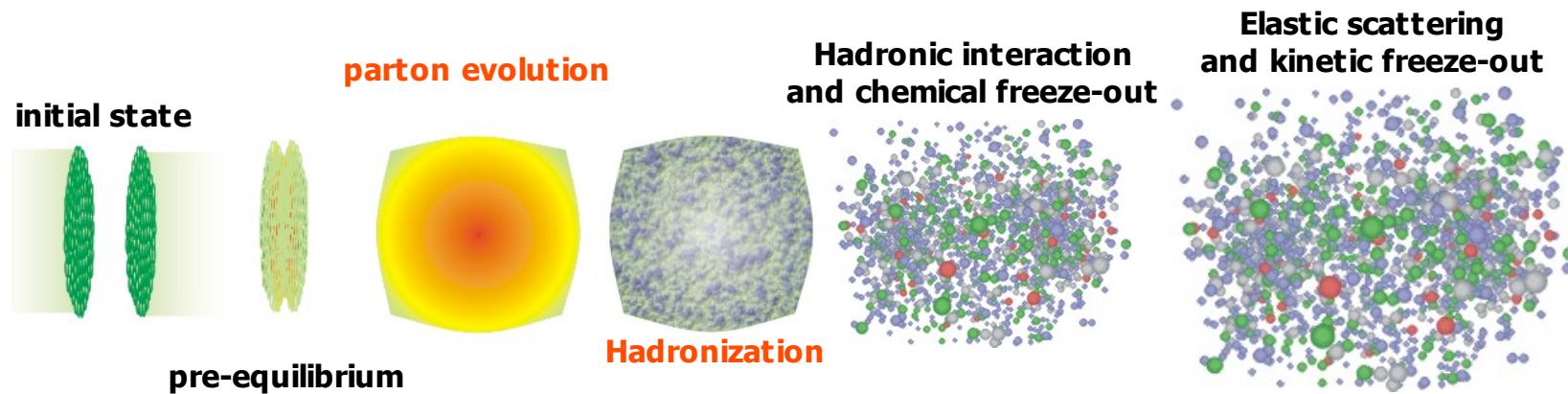
- **Quark Recombination**
 - motivation
 - simple quark recombination formula
 - recombination versus fragmentation
- **Elliptic flow**
 - elliptic flow at RHIC
 - model prediction of elliptic flow
 - quark number scaling of elliptic flow
- **Caveats of quark recombination formalism**
 - phase space density vs volume anisotropy
 - model calculations of quark number scaling
- **Summary**

Two RHIC puzzles

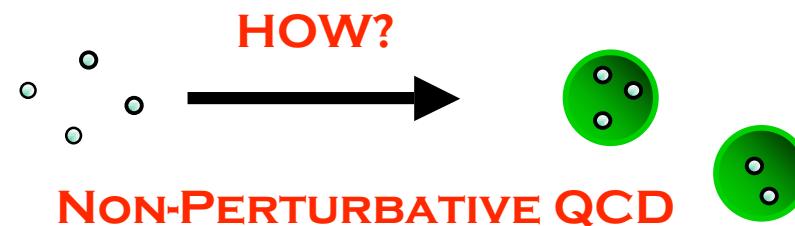
Baryon non suppression: $p/\pi \sim 1$



Hadronization problem in HI collision



- ❖ *Free* partons produced in high energy collisions: $e^+ + e^-$, $e + p$, $p + p$, $A + A$
- ❖ Partons in the final state have to be converted into hadrons (confinement)



- ❖ Can we learn about the parton phase from a theory of hadronization?

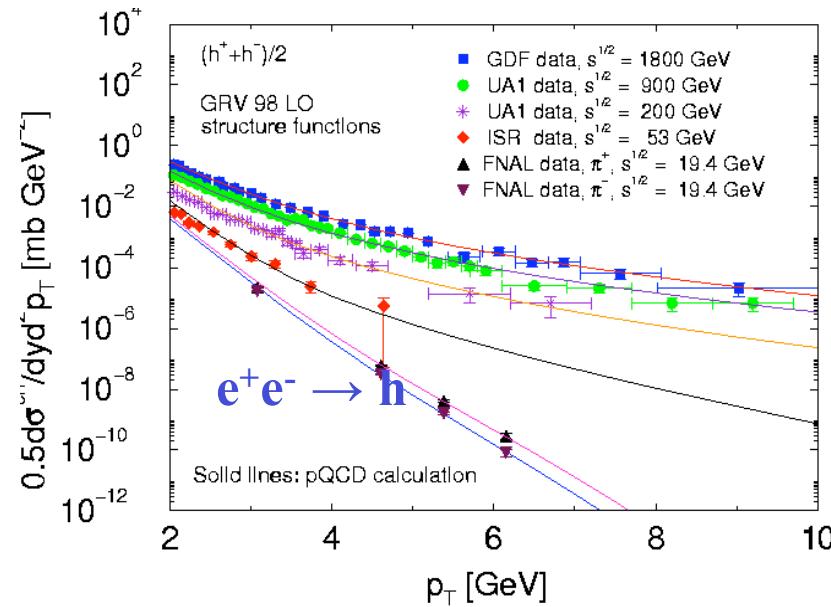
Independent fragmentation in p+p colln

Inclusive hadron distribution – calculable in pQCD

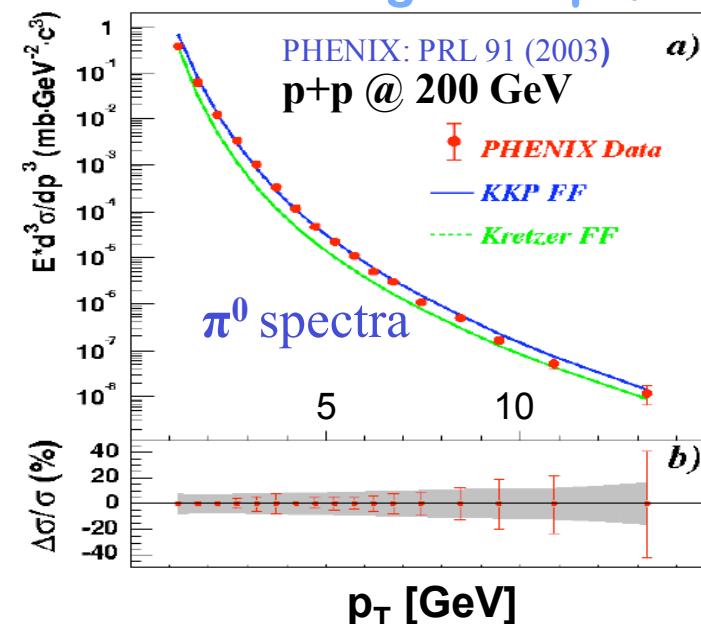
$$dN_h = f_{a/N}(x_a, Q^2) \otimes f_{b/N}(x_b, Q^2) \otimes d\sigma_{ab \rightarrow cX} \otimes D_{c \rightarrow h}(z_c, Q^2)$$

Parton distribution fns
 Perturbative cross section
 Fragmentation fn

Leading order pQCD phenomenology



Next-to-Leading order pQCD

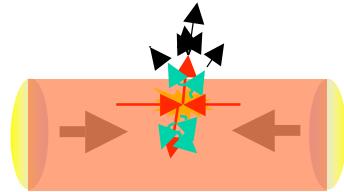


Fragmentation function D works for $p+p$ collision at RHIC

Indep. fragmentation in A+A collision

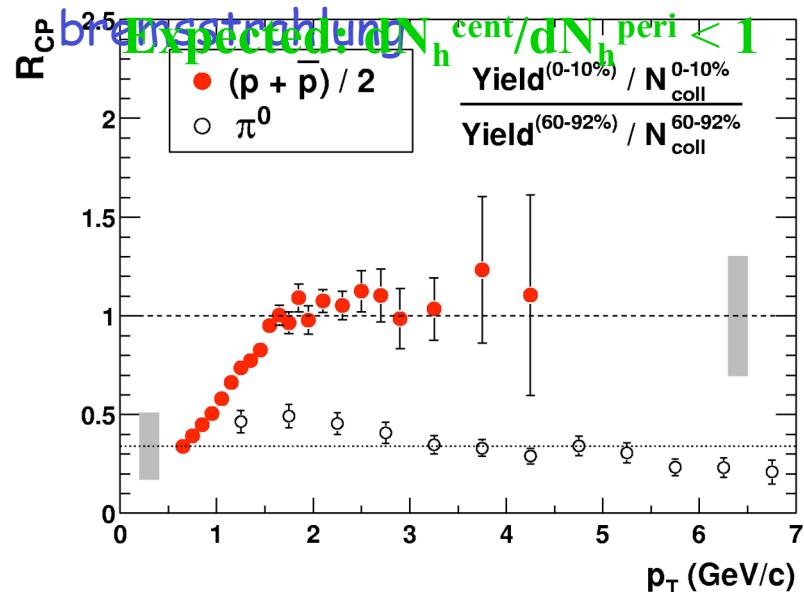
A+A collision: *inclusive hadron distribution from pQCD*

$$dN_h = f_{a/A}(x_a, Q^2) \otimes f_{b/A}(x_b, Q^2) \otimes d\sigma_{ab \rightarrow cX} \otimes \Delta E \otimes D_{c \rightarrow h}(z_c, Q^2)$$



Gyulassy, Levai,
Vitev NPB594 ('01)

Jet quenching: pQCD energy loss
 ΔE of partons by gluon



PHENIX: Phys. Rev. Lett. 91 (2003)

$$\Delta E^{(1)} \approx \frac{C_R \alpha_s}{4} \frac{\mu^2 L^2}{\lambda_g} \text{Log} \frac{2E}{\mu^2(L)L} + \dots ,$$

- Static medium

$$\Delta E^{(1)} \approx \frac{9\pi C_R \alpha_s^3}{4} L \frac{1}{A_\perp} \frac{dN^g}{dy} \text{Log} \frac{2E}{\mu^2(L)L} + \dots ,$$

- 1+1D Bjorken

A+A collision with energy loss

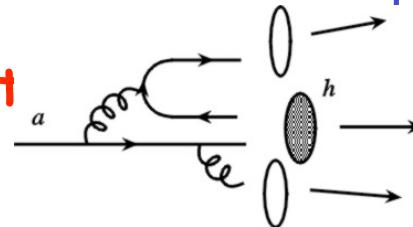
works for pions (suppression) but
fails for baryons (non-suppression)

- Failure of fragmentation function

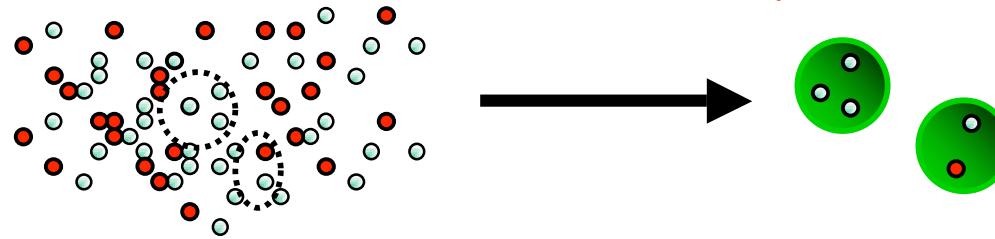
Quark Recombination

- Collision of elementary particles or for $A+A$ at $p_T > 6 \text{ GeV}/c$

Jet fragmentation is dominant

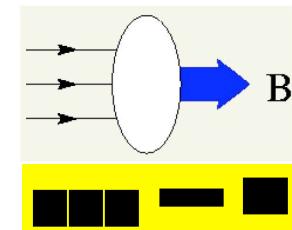
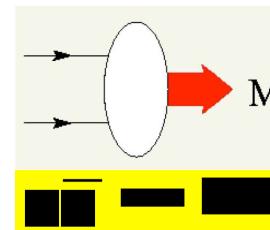


- Heavy ion collisions: large parton density, multi-parton processes possible, additional hadronization channels via **parton recombination**



- parton ReCo assumption:

Consider lowest order Fock state, i.e.
valence constituent (massive) quarks



Simple recombination formula

- Based on $n+p \rightarrow d$

Recombination: $q\bar{q} \rightarrow meson, \quad qqq \rightarrow baryon \quad$ Convolution of Wigner functions

$$\frac{dN_M(p)}{d^3 p} = g_M \int \left(\prod_{i=1,2} d^3 x_i d^3 p_i \right) f_\alpha(p_1, x_1) f_\beta(p_2, x_2) W_M(x_1 - x_2, p_1 - p_2) \delta^3(p - p_1 - p_2)$$

$$\frac{dN_B(p)}{d^3 p} = g_B \int \left(\prod_{i=1,2,3} d^3 x_i d^3 p_i \right) f_\alpha(p_1, x_1) f_\beta(p_2, x_2) f_\gamma(p_3, x_3) W_B(x_{12}, x_{13}, p_{12}, p_{13}) \delta^3(p - \sum p_i)$$

Hadron yield Space-time Quark distribution fn Hadron function

□ Basic assumptions:

- uncorrelated quark distribution function
- weak binding
- sudden (instantaneous) hadronization

Das & Hwa, PLB 68 ('77)

Biro et al, PLB 347 ('97)

Voloshin NPA 517 ('03)

Greco, Ko, Levai, PRC 68 ('03)

Fries, Muller, Nonaka, Bass, PRC68 ('03)

Molnar & Voloshin, PRL 91 ('03)

Two approaches to ReCo:

- (i) parametrize f_q , (ii) Compute f_q in a dynamical model

Recombination vs fragmentation

Exponential: $f_q \sim e^{-p_T/T}$

$$N_{\text{frag}} = f \otimes D \sim e^{-P_T/\langle z \rangle T} \langle D \rangle$$

$$N_{\text{reco}} = f_\alpha \otimes W_M \otimes f_\beta \sim e^{-P_T/T}$$

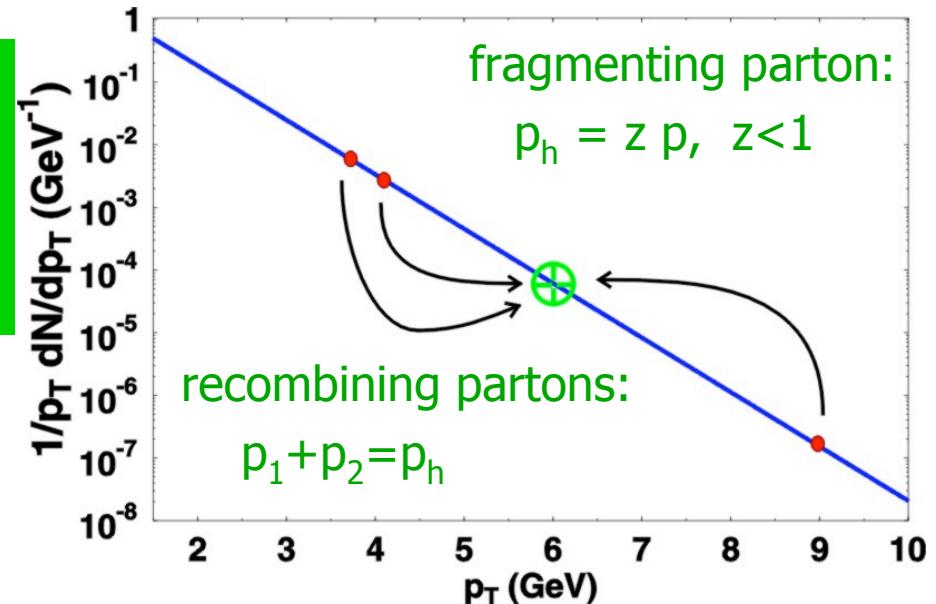
Recombination wins

Power law: $f_q \sim p_T^{-\alpha}$

$$N_{\text{frag}} \sim P_T^{-\alpha}$$

$$N_{\text{reco}} \sim P_T^{-2\alpha}$$

Fragmentation wins



Baryon: $N_B(p_T) = N_q(p_T/3)$

Meson: $N_M(p_T) = N_q(p_T/2)$

- Baryon number dominates meson

Recombination competes with fragmentation and wins at $p_T < 6$ GeV

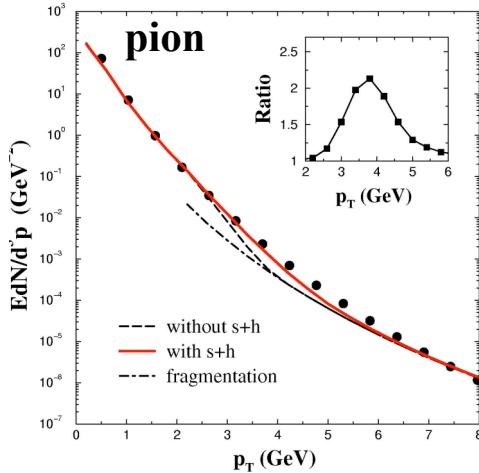
ReCo by source parametrization

One approach to obtain f_q is via “source parametrization”:

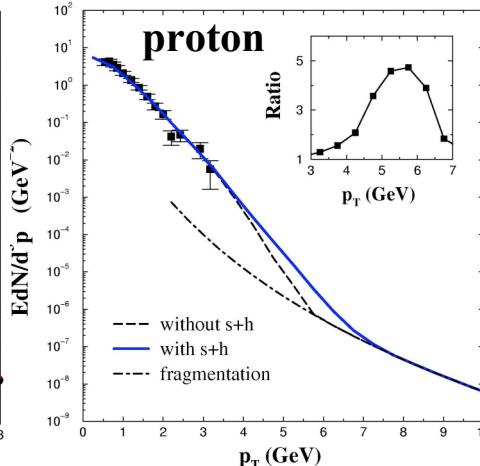
- Hard spectra : LO pQCD minijets with jet quenching and indep. fragmentation
 - Soft spectra ($p_T < 2$ GeV): (anti)quark thermalized plasma with radial flow
- Total hadron yield $dN_h = dN_h^{coal}(f_q \otimes f_{\bar{q}} \otimes W_h) + dN_h^{frag}(d\sigma \otimes E_{loss} \otimes D)$

TAMU group: $Tq=170$ MeV, $v_T=0.5$

Greco, Ko, Levai, PRC 68 ('03)

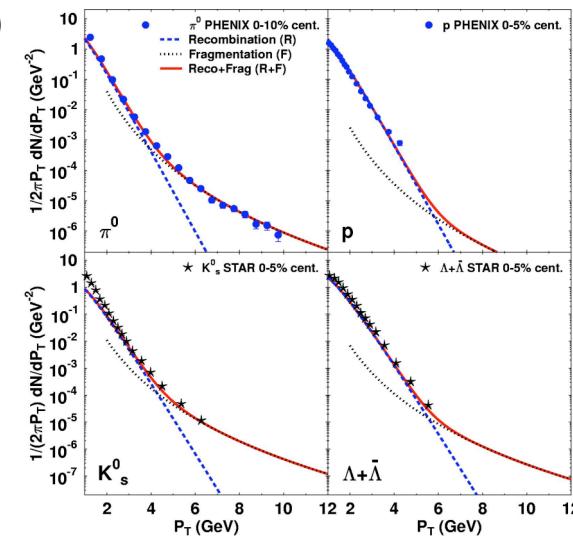


Au+Au @ 200AGeV (central)



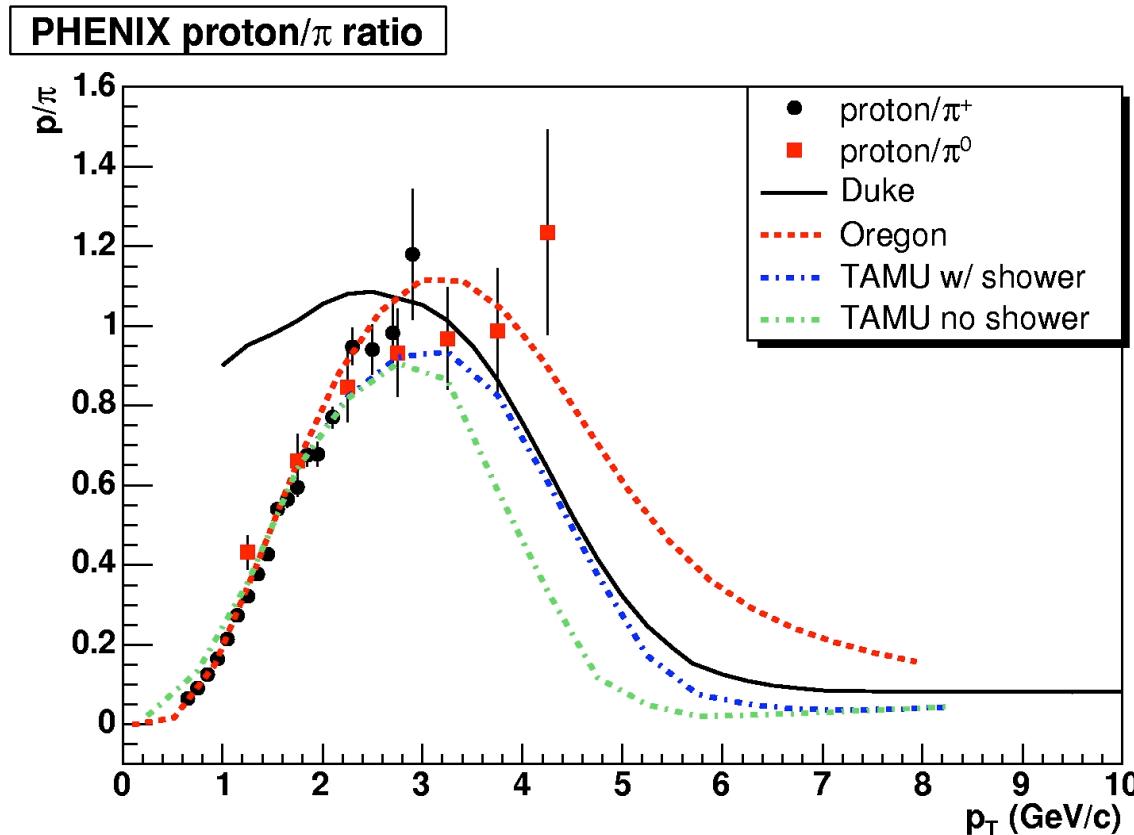
Duke group: $Tq=175$ MeV, $v_T=0.55$

Fries, Muller, Nonaka, Bass, PRC 68 ('03)



ReCo dominates up to $p_T \sim 4$ GeV (meson) & to ~ 6 GeV (baryon)

Baryon/Meson anomaly in ReCo

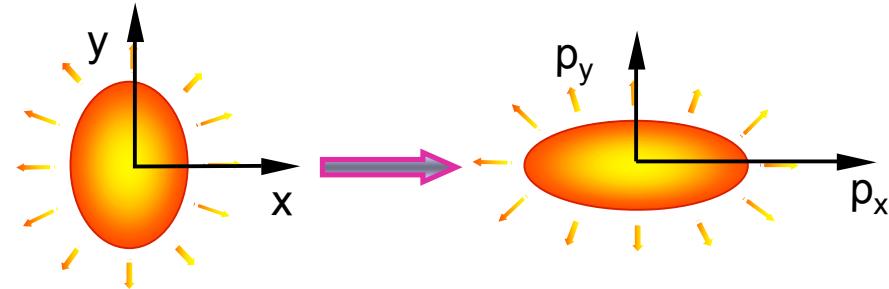
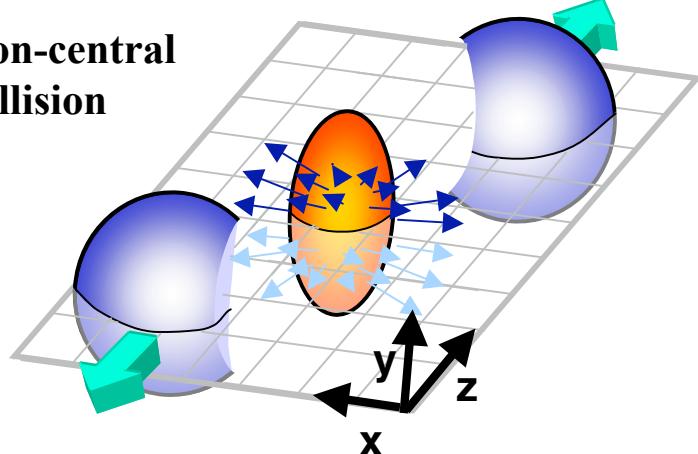


$p_T < \sim 2$ GeV recombination via soft partons dominates:

Mesons up to $2x(\sim 2)$ GeV, Baryons up to $3x(\sim 2)$ GeV

Elliptic flow

Non-central
collision



Origin of $v_2 \neq 0$: coordinate-space anisotropy ($b > 0$) and re-interaction

Anisotropic flows v_n are Fourier coefficients in azimuthal distribution of particles with respect to reaction plane

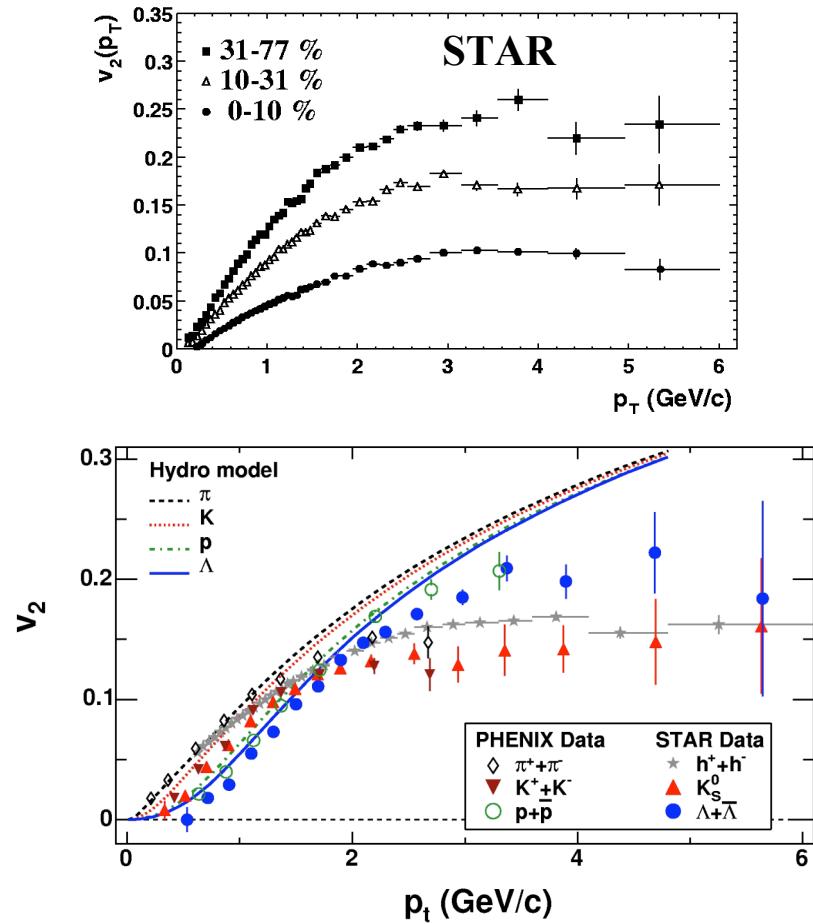
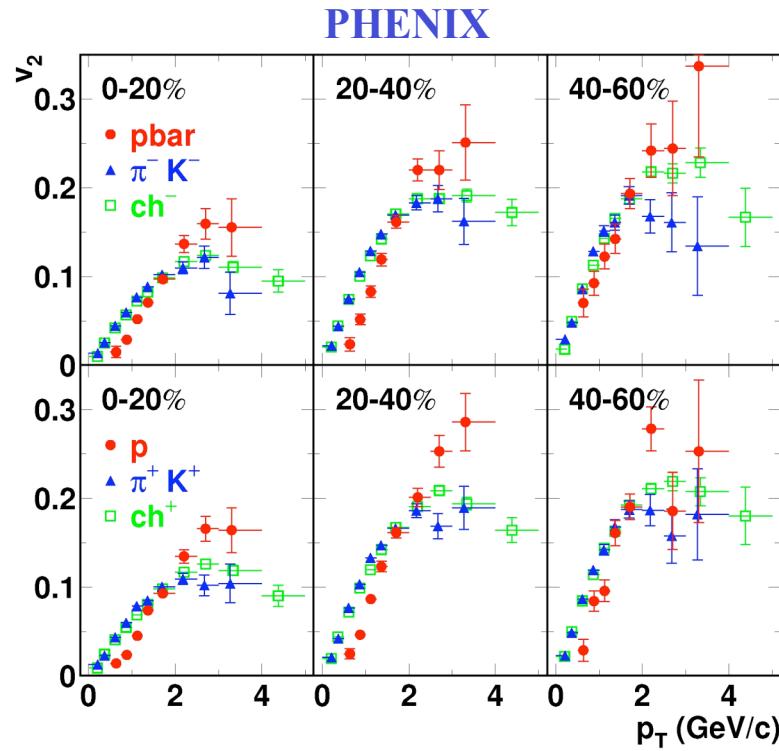
$$\frac{dN}{d\varphi dX} = \frac{1}{2\pi} \frac{dN}{dX} \left[1 + 2 \sum_n v_n(X) \cos(n\varphi) \right]$$

Ollitrault PRD 46 ('92)

X: particle species and event, eg: centrality, transverse momentum

Elliptic flow $\Rightarrow v_2(p_T) \equiv \langle \cos(2\varphi) \rangle_{p_T} = \frac{\int d\varphi \cos(2\varphi) (dN/dp_T dy d\varphi)}{\int d\varphi (dN/dp_T dy d\varphi)} = \left\langle \frac{p_x^2 - p_y^2}{p_x^2 + p_y^2} \right\rangle$

Elliptic flow at RHIC



- Large elliptic flow for all hadron species
- v_2 saturation at different p_T for hadrons
- At $p_T < 2$ GeV, heavier hadrons smaller v_2
- At $p_T > 2$ GeV, heavier hadrons larger v_2

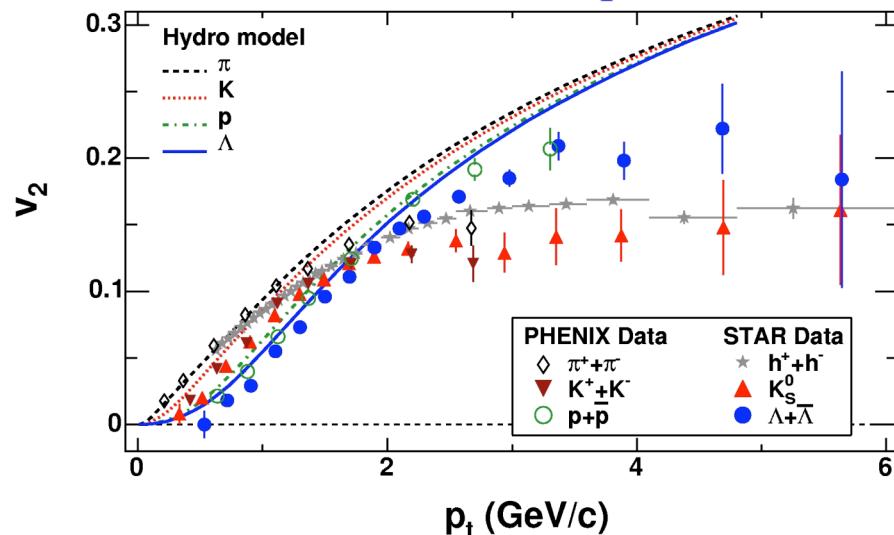
Ideal hydrodynamics & elliptic flow

Basic assumption: local thermodynamic equilibrium

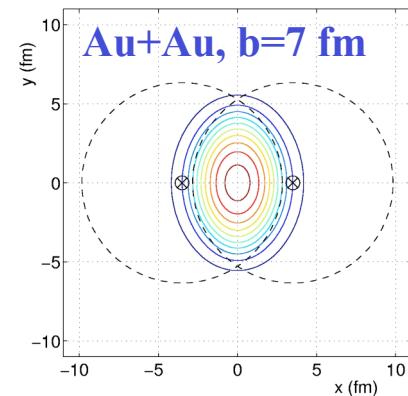
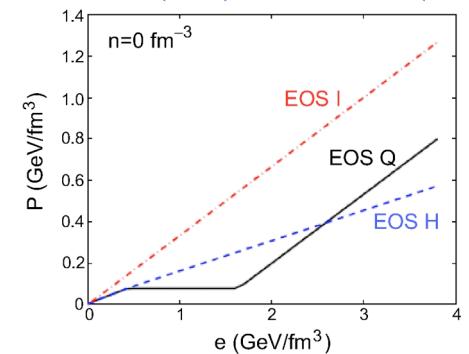
Evolution equations are local conservation of energy-momentum and net baryon number

$$\partial_\mu T^{\mu\nu}(x) = 0, \quad \partial_\mu N^\mu(x) = 0$$

5 eqns, 6 unknowns (ε, n, p, u)
requires an EOS: $p = p(\varepsilon, n)$



Kolb et al, PRC62 ('00),
NPA696 ('01), PLB500 ('01)



- ❖ Hydro explains v_2 systematics with a **QGP EOS** for $p_T < 2$ GeV
- ❖ **Mass ordering** at $p_T < 2$ GeV - massive hadrons smaller v_2
- ❖ mass ordering **violated** at large p_T and v_2 saturates in data

Elliptic flow from recombination

If all partons have same momentum distribution

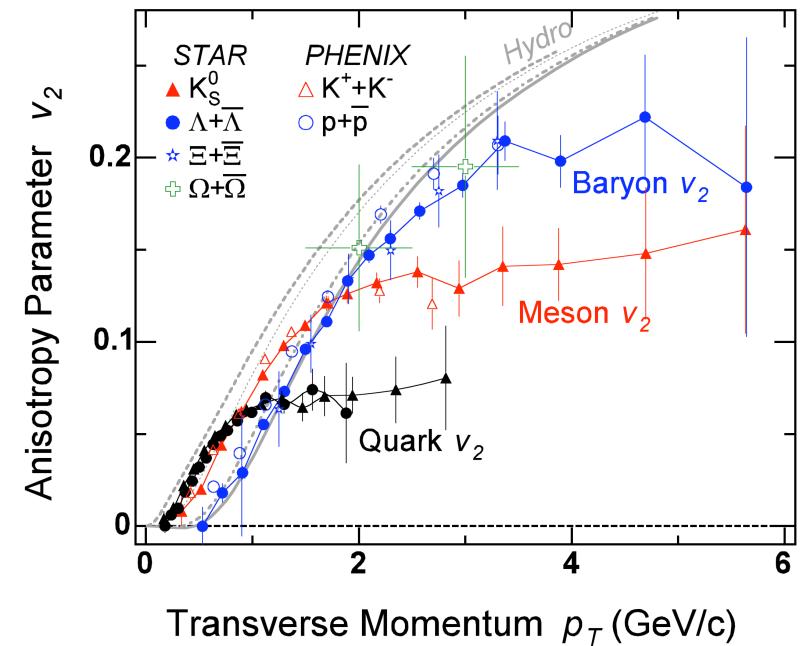
$$\frac{dN_M(p_T)}{d^3 p_T} \propto \left[\frac{dN_q(p_T/2)}{d^3 p_T} \right]^2, \quad \frac{dN_B(p_T)}{d^3 p_T} \propto \left[\frac{dN_q(p_T/3)}{d^3 p_T} \right]^3$$

$$v_2^{(M)}(p_T) \approx 2v_2^{(q)}(p_T/2), \quad v_2^{(B)}(p_T) \approx 3v_2^{(q)}(p_T/3)$$

At low $p_T < 2$ GeV: $v_2^{(B)}(p_T) < v_2^{(M)}(p_T) < v_2^{(q)}(p_T)$
 → mass ordering

At high $p_T > 2$ GeV: $v_2^{(B)}(p_T) > v_2^{(M)}(p_T) > v_2^{(q)}(p_T)$
 → hadron flow amplified

Amplification: 3x for baryons
 2x for mesons
 → 50% larger baryon flow



➤ By measuring v_2 for various hadrons, quark flow can be extracted

Quark number scaling of v_2

$$v_2^{(H)}(p_T) \approx n v_2^{(q)}(p_T/n), \quad n = 2, 3 \text{ for mesons, baryons}$$

$$v_2^{(H)}(p_T)/n \rightarrow \text{universal}$$

If $v_2^{(q)} \neq v_2^{(s)}$ (light vs. strange)

Mesons :

$$v_2^{(\partial)}(p_T) \approx 2v_2^{(q)}(p_T/2)$$

$$v_2^{(K)}(p_T) \approx v_2^{(q)}(p_T/2) + 2v_2^{(s)}(p_T/2)$$

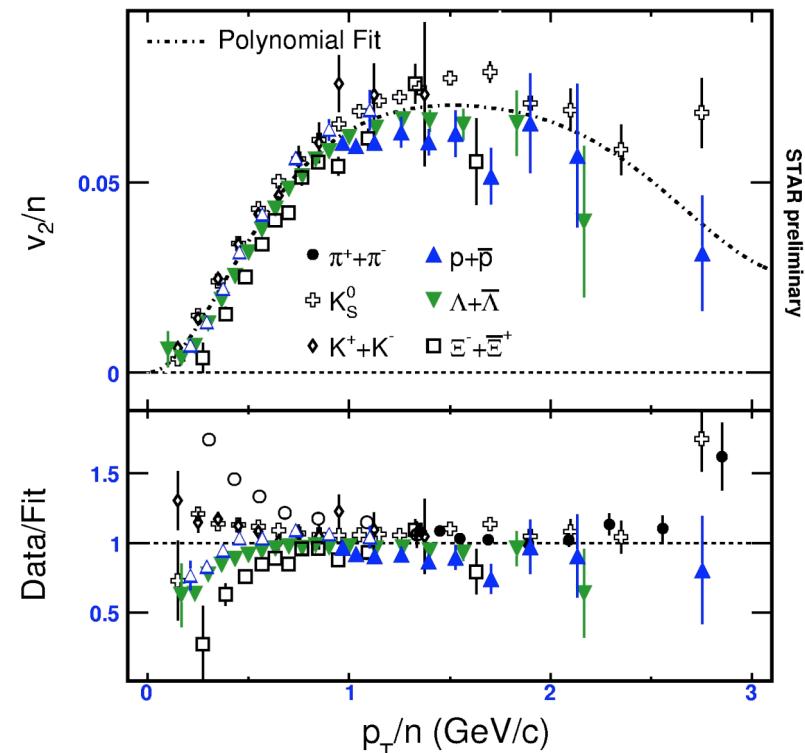
Baryons :

$$v_2^{(p)}(p_T) \approx 3v_2^{(q)}(p_T/3)$$

$$v_2^{(\Lambda,\Sigma)}(p_T) \approx 2v_2^{(q)}(p_T/3) + v_2^{(s)}(p_T/3)$$

$$v_2^{(\Xi)}(p_T) \approx v_2^{(q)}(p_T/3) + 2v_2^{(s)}(p_T/3)$$

$$v_2^{(\Omega)}(p_T) \approx 3v_2^{(q)}(p_T/3)$$



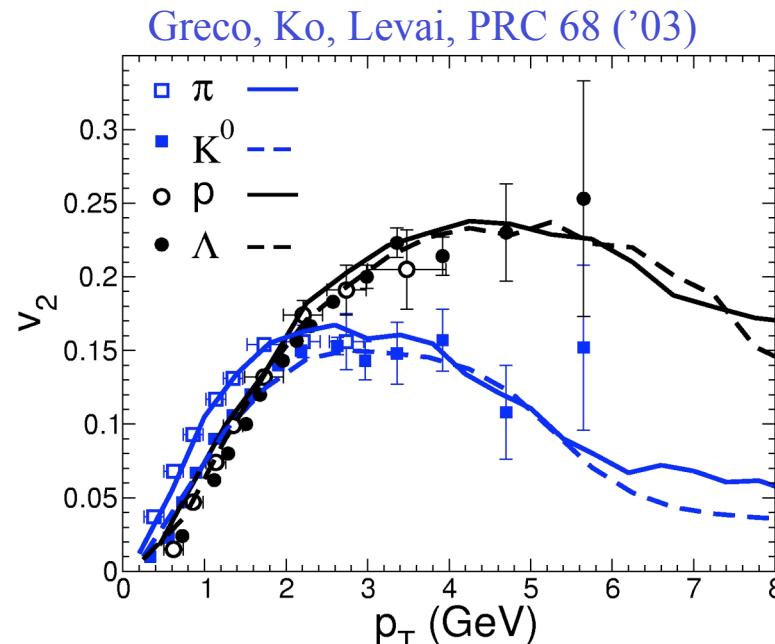
- Coalescence prediction approx. verified
- RHIC data shows $v_2^{(q)} = v_2^{(s)}$

v_2 from source parametrization

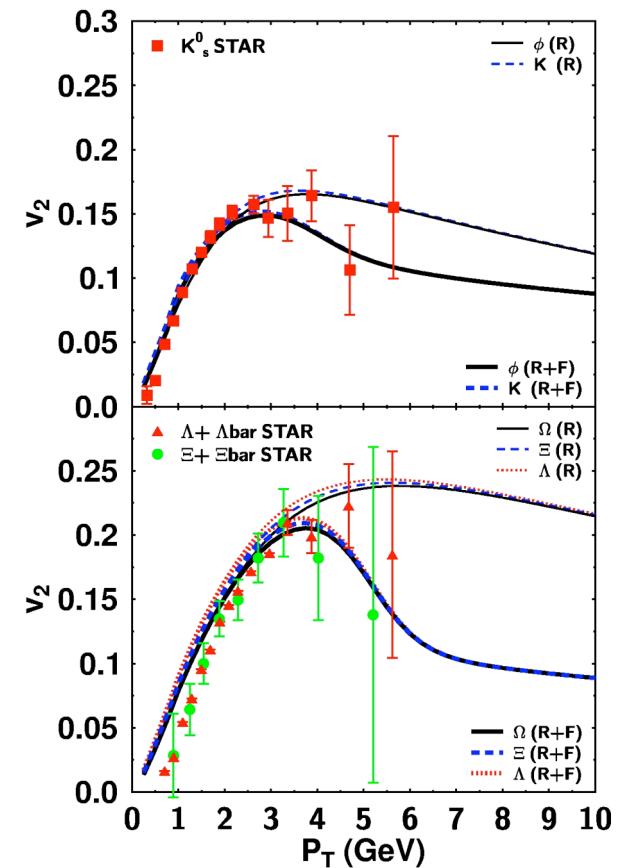
Quark transverse momentum anisotropy included:

$$\frac{dN_q}{d^2 p_T} = \frac{dN_q}{p_T dp_T} [1 + 2v_2 \cos(2\phi)]$$

→ Quark elliptic flow extracted from pion v_2



Nonaka, Fries, Bass, PLB 583 ('04)



Covariant transport model

- Interpolates between ideal hydro $\lambda = 0$ and free streaming $\lambda = \infty$

MPC: kinetic theory of parton gas (Boltzmann)

$$p^\mu \partial_\mu f_i(x, p) = S_i(x, p) + \sum_{jkl} F_{ij \rightarrow kl}^{2 \rightarrow 2}[f_i, f_j, f_k, f_l]$$

- Minijet initial conditions in $S_i(x, p)$
- Elastic $2 \rightarrow 2$ cross section
- Debye screened $2 \rightarrow 2$ pQCD $\sigma \sim 1/(t - \mu^2)^2$
- 1 parton \rightarrow 1 hadron hadronization

Require large opacities:

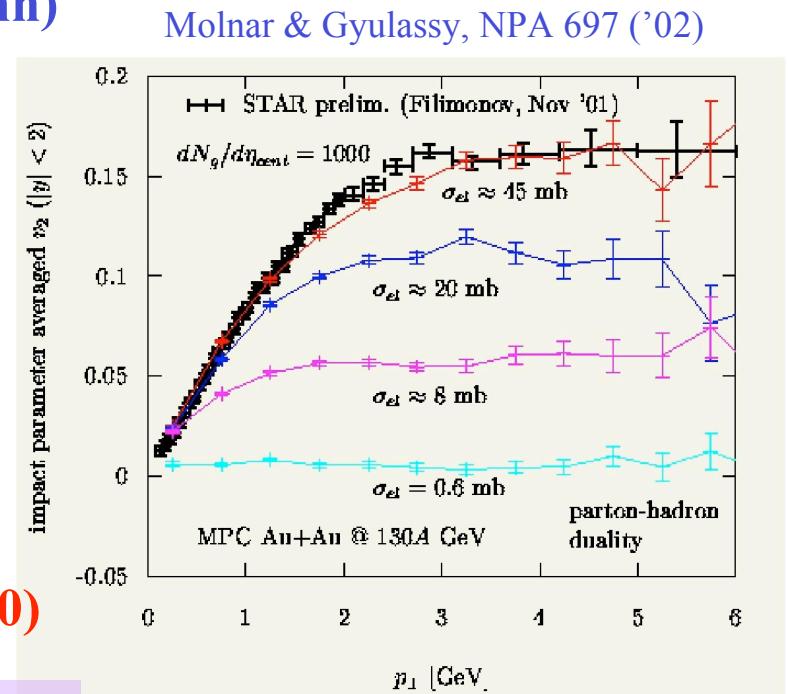
$$\sigma_{\text{el}} \times dN_g/d\eta \approx 45000 \text{ mb} \gg \text{pQCD } (3 \text{ mb} \times 1000)$$

Solution to opacity puzzle: Recombination

ReCo dynamics: f_q , $f_{\bar{q}}$ and hypersurface

Elliptic flow fit with ReCo requires

smaller $\sigma_{\text{el}} \times dN_g/d\eta$ ($b=0$) $\approx 7000-15000$ mb

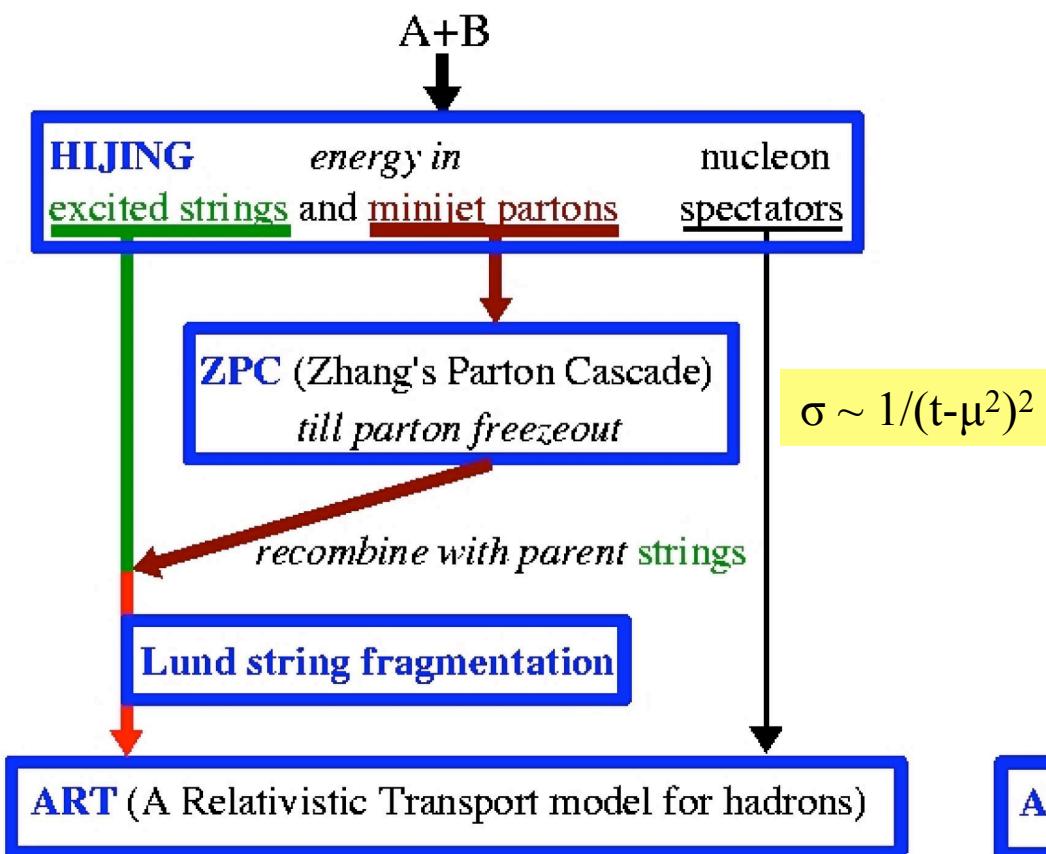


Molnar & Voloshin, PRL 91 ('03)

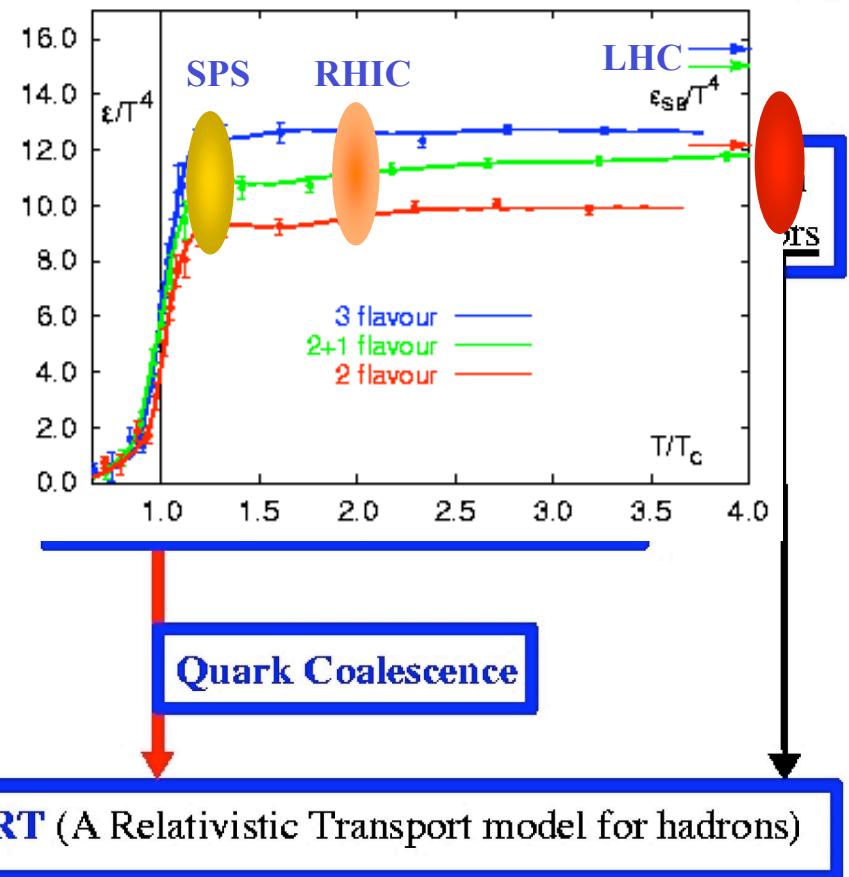
A MultiPhase Transport model (AMPT)

Lin, Ko, Li, Zhang, S.P, nucl-th/0411110

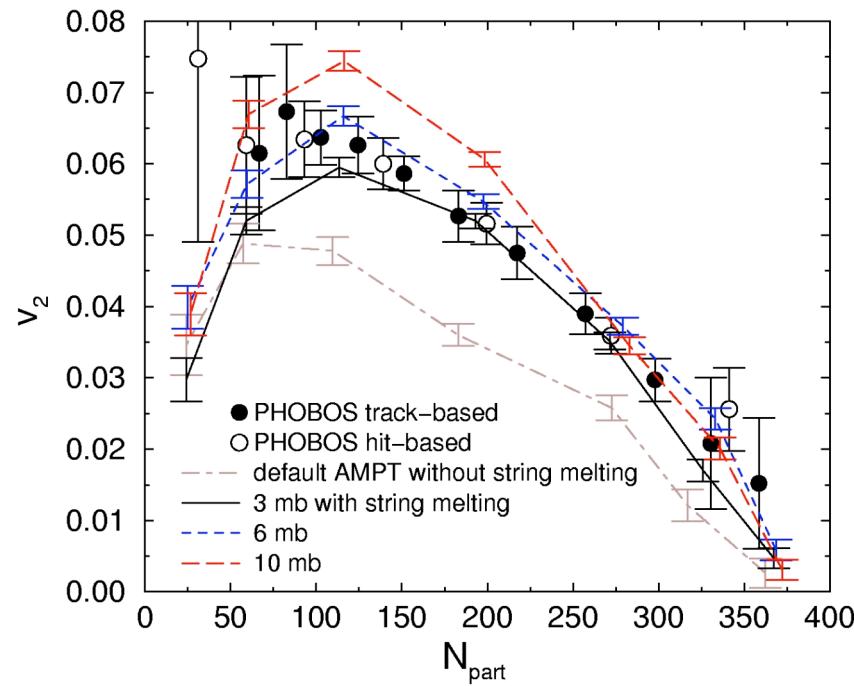
Structure of the default AMPT model



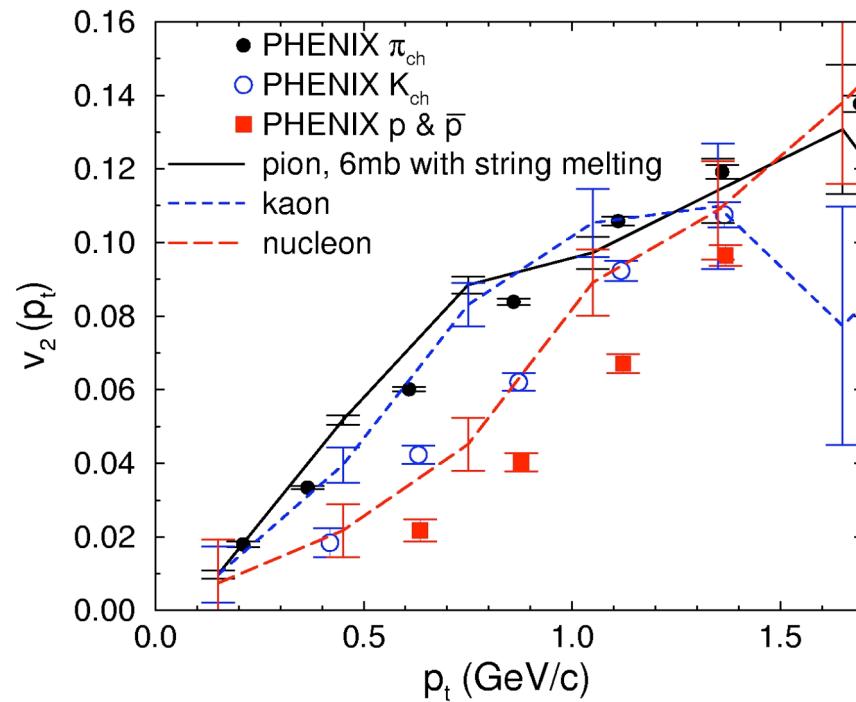
Structure of Karsch, NPA698 (2002) in melting



v_2 in AMPT model



Au+Au @ 200A GeV (min bias)



- String melting enhances parton density
- v_2 well explained with recombination & string melting
- v_2 sensitive to parton elastic $\sigma_{q\bar{q} \rightarrow q\bar{q}}$ (= 6 mb best fit)

Recombination revisited

$$\frac{dN_H(p)}{d^3 p} = g_H \int (\prod_i^n d^3 x_i d^3 p_i) \Phi_q(x_i, p_i) W_H(x_{i...n}, p_{i...n}) \delta^3(p - \sum_i^n p_i)$$

Assumptions in $v_2^{(H)}(p_T) \approx n v_2^{(q)}(p_T/n)$

S. Pratt, S.P, PRC 71 ('05)

No spread in relative coordinate and momenta in hadron wave fn.

$$W_H(x, p) = \delta(x) \delta(p)$$

- **Uncorrelated quarks** $\Phi_q = \prod_i^n f_i(x_i, p_i)$
- **sudden (instantaneous) hadronization**

$$f^{(H)}(p_T, \phi, x, t_c) = [f^{(q)}(p_T/n, \phi, x, t_c)]^n$$

→ *Binding potential turned on suddenly (no energy conservation)*

→ *For thermal distribution* $f^{(q)} \propto \exp[-(p_T/n - \mu)/T]$, $\therefore \mu^{(H)} = n\mu^{(q)}$

Phase-space density vs emission volume

Define: $\bar{f} = \int d^3r f_q^2 / \int d^3r f_q$, $\Omega = \left(\int d^3r f_q \right) / \int d^3r f_q^2$

Spectra determined by both average phase-space density and effective emission volume of partons

$$\frac{dN_H(p_T)}{d^3p} \prec \int d^3r [f_q(p_T/n, \phi)]^n \\ \approx \Omega(p_T/n, \phi) [\bar{f}(p_T/n, \phi)]^n$$

If both mean \bar{f} and Ω have small elliptic anisotropies

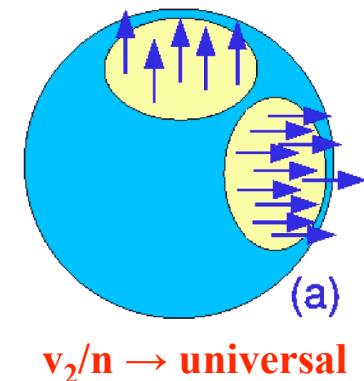
$$\bar{f}(p_T, \phi) = \bar{f}_0(p_T) [1 + 2 v_2^{(\bar{f})} \cos(2\phi)]$$

$$\Omega(p_T, \phi) = \Omega_0(p_T) [1 + 2 v_2^{(\Omega)} \cos(2\phi)]$$

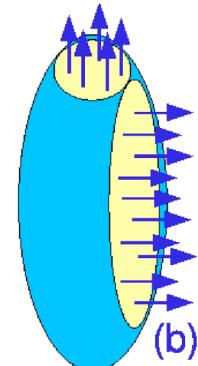
The anisotropy for the spectra obtained from $\langle \cos(2\phi) \rangle$:

$$v_2^{(H)}(p_T) = v_2^{(\Omega)}(p_T/n) + n v_2^{(\bar{f})}(p_T/n)$$

Quark number scaling violated if
effective emission volume depends on φ



$v_2/n \rightarrow \text{universal}$



quark-number scaling violated

Blast wave parametrization

- Hydro parametrization for the phase-space density at kinetic freeze-out
- Boost invariance with transverse anisotropic collective flow field

Emission probability for n quarks in blast wave with azimuthally symmetric

$$\frac{dN}{d^3 p} \underset{\text{shell}}{\int} d\phi_u d\eta \cosh\eta \rho_s(\phi_u) \exp\left[\{-\cosh\eta m_T \sqrt{1+u_T(\phi_u)} + u_T(\phi_u) p_T \cos(\phi - \phi_u) + n\mu(\phi_u)\}/T(\phi_u)\right]$$

Anisotropies is added via: ρ_s , μ , T , u_T

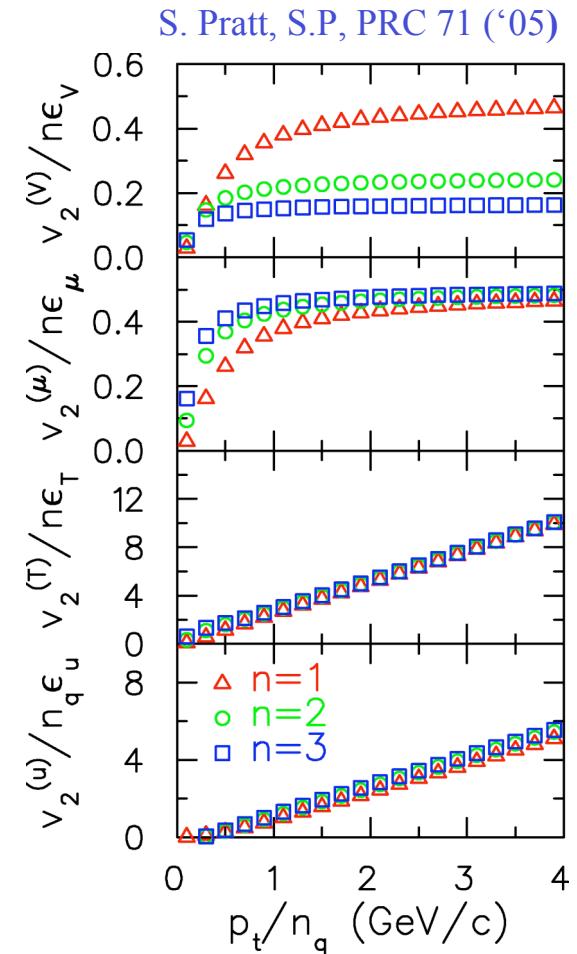
$$\rho_s(\phi_u) = (1 + \varepsilon_v \cos 2\phi_u) \rho_0$$

$$\exp[\mu(\phi_u)/T(\phi_u)] = (1 + \varepsilon_\mu \cos 2\phi_u) \exp(\mu/T)$$

$$T(\phi_u) = (1 + \varepsilon_T \cos 2\phi_u) T$$

$$u_T(\phi_u) = (1 + \varepsilon_u \cos 2\phi_u) u_T$$

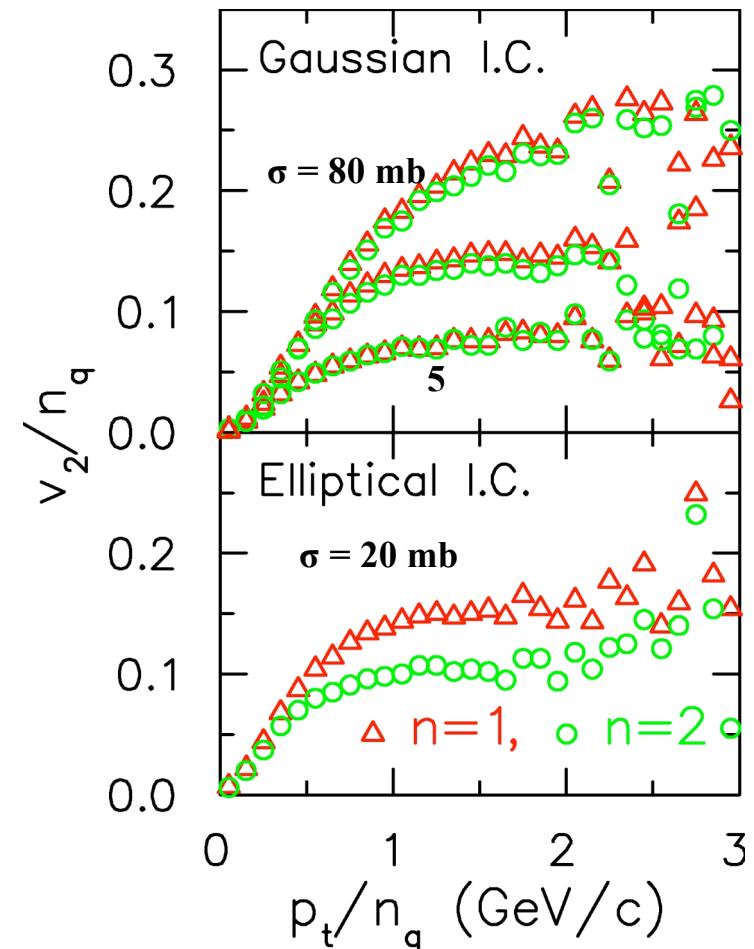
Schnedermann et al, PRC 48 ('93)
Huovinen et al, PLB 503 ('01)



Microscopic simulation

GROMIT: Cheng, Pratt, Csizmadia, PRC 65 ('02) 024901

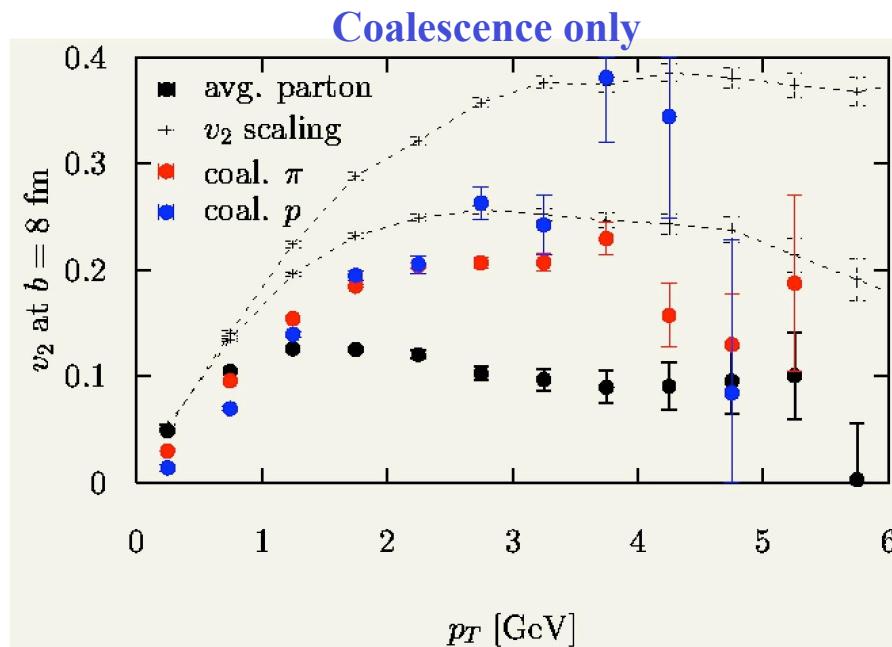
- Boltzmann description provide complicated final phase-space density
- Constituent quarks initially thermalized
- Quarks collide with simple s-wave cross section of $\sigma = 5, 20, 80$
- Initial transverse density distributions:
gaussian and elliptic
Quark number scaling violated
for elliptic quark density profile



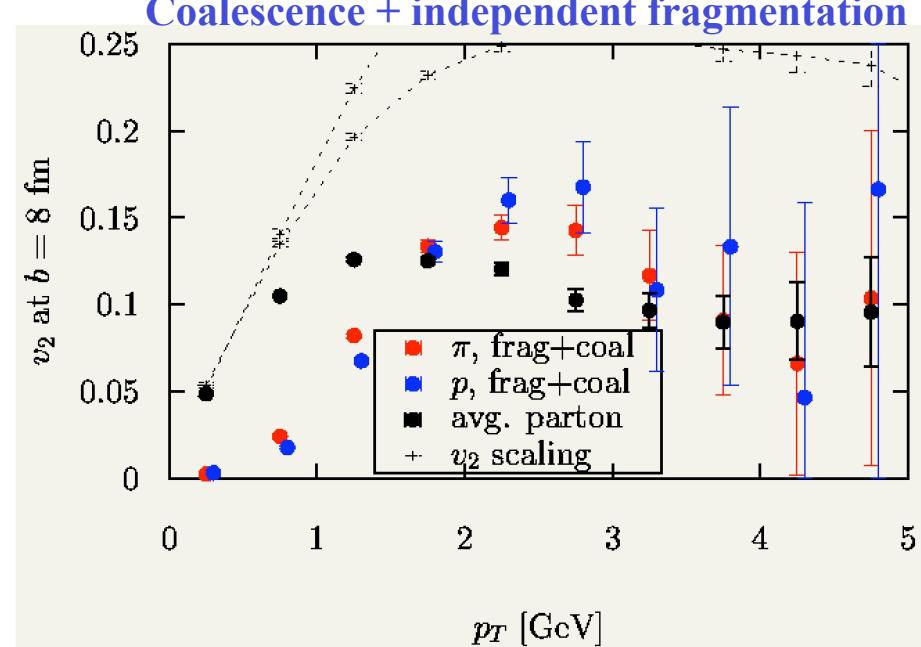
Quark scaling violation in MPC

Au+Au @ 200A GeV, $b = 8$ fm, $\sigma_{gg} = 10$ mb

Molnar, J.Phys G30 ('04)



v_2 for $\pi(p)$ reduced by 20(40)%
relative to v_2 scaling law



Fragmentation quench
spectra and smears out
anisotropy

→ v_2 scaling further violated

➤ parton $v_2(p_T)$ extracted from hadron v_2 will be underpredicted

Summary

- ❖ Hadronization of dense phase of partons can be described by recombination
 - ❖ For (semi)central Au+Au collision at RHIC, hadron production is dominated via recombination at intermediate $1.5 < p_T < 4\text{-}6 \text{ GeV}$
 - ❖ Simple recombination formalism explains, large baryon/meson ratio, elliptic flow systematic for various hadron species & elliptic flow scaling with quark number
 - ❖ Quark number scaling can be significantly affected by anisotropy in emission volume, emission profile, space-time dynamics (x-p correlations)
- Recombination challenges:**
- Without jet quenching recombination may not be seen
 - Simple formula for rare process, current formalism not suited at low p_T
 - No energy conservation, decreasing entropy ?
 - Effects of hadron scattering, jet-like correlations...
 - Space-time evolution not well understood in heavy ion collisions
→ requires study in a consistent dynamical calculation ☺

Charm elliptic flow

$T_q=170 \text{ MeV}$, $v_T = 0.5-0.65$

