Heavy-Quark Diffusion in the Quark-Gluon Plasma

Hendrik van Hees

Texas A&M University

February 13, 2008

with M. Mannarelli, V. Greco and R. Rapp





1 Heavy Quark interactions in the sQGP

- 2 Heavy-Quark Diffusion in the Quark-Gluon Plasma
- 3 Non-photonic electrons at RHIC
- 4 Summary and Outlook

Heavy Quarks in Heavy-Ion collisions



hard production of HQs described by PDF's + pQCD (PYTHIA)

c, b quark



HQ rescattering in QGP: Langevin simulation drag and diffusion coefficients from non-perturbative many-body T matrix (sQGP)



Hadronization to D,B mesons via quark coalescence + fragmentation V. Greco, C. M. Ko, R. Rapp, PL B **595**, 202 (2004)



semileptonic decay ⇒ "non-photonic" <mark>electron observables</mark>

Static potentials from lattice QCD



- color-singlet free energy from lattice
- use internal energy

$$U_1(r,T) = F_1(r,T) - T \frac{\partial F_1(r,T)}{\partial T},$$

$$V_1(r,T) = U_1(r,T) - U_1(r \to \infty,T)$$

• Casimir scaling for other color channels [Nakamura et al 05; Döring et al 07]

$$V_{\overline{3}} = \frac{1}{2}V_1, \quad V_6 = -\frac{1}{4}V_1, \quad V_8 = -\frac{1}{8}V_1$$

T-matrix

• Brueckner many-body approach for elastic Qq, $Q\bar{q}$ scattering



- reduction scheme: 4D Bethe-Salpeter \rightarrow 3D Lipmann-Schwinger
- S- and P waves
- same scheme for light quarks (self consistent!)
- Relation to invariant matrix elements

$$\sum |\mathcal{M}(s)|^2 \propto \sum_q d_a \left(|T_{a,l=0}(s)|^2 + 3|T_{a,l=1}(s)|^2 \cos \theta_{\rm cm} \right)$$

T-matrix



- resonance formation at lower temperatures $T \simeq T_c$
- melting of resonances at higher $T! \Rightarrow sQGP$
- P wave smaller
- resonances near T_c : natural connection to quark coalescence [Ravagli, Rapp 07]
- model-independent assessment of elastic Qq, $Q\bar{q}$ scattering
- problems: uncertainties in extracting potential from IQCD in-medium potential V vs. F?

Elastic pQCD processes

• Lowest-order matrix elements [Combridge 79]



• In-medium Debye-screening mass for *t*-channel gluon exchange: $\mu_g = gT$, $\alpha_s = 0.4$

Heavy-Quark diffusion

• Fokker Planck Equation

$$\frac{\partial f(t,\vec{p})}{\partial t} = \frac{\partial}{\partial p_i} \left[p_i A(t,p) + \frac{\partial}{\partial p_j} B_{ij}(t,\vec{p}) \right] f(t,\vec{p})$$

• drag (friction) and diffusion coefficients

$$p_i A(t, \vec{p}) = \langle p_i - p'_i \rangle$$

$$B_{ij}(t, \vec{p}) = \frac{1}{2} \langle (p_i - p'_i)(p_j - p'_j) \rangle$$

$$= B_0(t, p) \left(\delta_{ij} - \frac{p_i p_j}{p^2} \right) + B_1(t, p) \frac{p_i p_j}{p^2}$$

• transport coefficients defined via ${\cal M}$

$$\langle X(\vec{p}') \rangle = \frac{1}{\gamma_c} \frac{1}{2E_p} \int \frac{\mathrm{d}^3 \vec{q}}{(2\pi)^3 2E_q} \int \frac{\mathrm{d}^3 \vec{q}'}{(2\pi)^3 2E_{q'}} \int \frac{\mathrm{d}^3 \vec{p}'}{(2\pi)^3 2E_{p'}} \\ \sum |\mathcal{M}|^2 (2\pi)^4 \delta^{(4)}(p+q-p'-q') \hat{f}(\vec{q}) X(\vec{p}')$$

• correct equil. lim. \Rightarrow Einstein relation: $B_1(t,p) = T(t)E_pA(t,p)$

Transport coefficients



• from non-pert. interactions reach $A_{non-pert} \simeq 1/(7 \text{ fm}/c \simeq 4A_{pQCD})$

- A decreases with higher temperature
- higher density (over)compensated by melting of resonances!
- spatial diffusion coefficient

$$D_s = \frac{T}{mA}$$

increases with temperature

Non-photonic electrons at RHIC

- same model for bottom
- quark coalescence+fragmentation $\rightarrow D/B \rightarrow e + X$



Summary and Outlook

• Summary

- Heavy quarks in the sQGP
- non-perturbative interactions via IQCD potentials parameter free
- resonance formation at $T > T_c \Rightarrow$ strong coupling
- res. melt at higher temperatures \Leftrightarrow consistency betw. R_{AA} and $v_2!$
- also provides "natural" mechanism for quark coalescence
- uncertainties
 - $\bullet~$ extraction of V from lattice data
 - potential approach at finite T: F, V or combination?
- Outlook
 - include inelastic heavy-quark processes (gluon-radiation processes)
 - other heavy-quark observables like charmonium suppression/regeneration