

Self-interacting neutrinos in the light of cosmological and IceCube observations

Priyank Parashari

**THEPH Division,
Physical Research Laboratory, Ahmedabad**

Work in Collaboration with : Arindam Mazumdar, Subhendra Mohanty

January 29, 2021



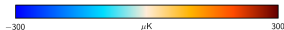
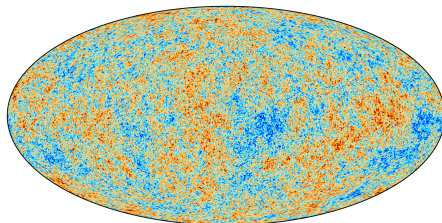


- ① H_0 Tension?
- ② Self interacting neutrinos
- ③ IceCube and self-interacting neutrinos
- ④ Self-interacting sterile neutrinos



- CMB \Rightarrow “Cosmic Microwave Background”

Snapshot of photons in our universe leftover after it's decoupling.



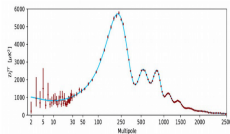
Λ CDM model

- CDM, baryons, dark energy, radiation
- Six standard parameters (ω_b , ω_c , A_s , n_s , τ_{reio} , θ)
- Two derived parameters (H_0 and σ_8)

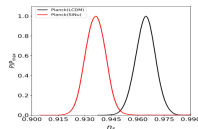
- Angular power spectrum

Model
(e.g. Λ CDM)

Einstein-Boltzmann
Solver



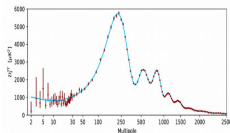
Data & MCMC
analysis



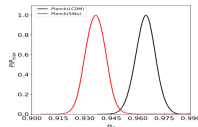


Model
(e.g. Λ CDM)

Einstein-Boltzmann
Solver



Data & MCMC
analysis



Λ CDM model

- CDM, baryons, dark energy, radiation
- Six standard parameters (ω_b , ω_c , A_s , n_s , τ_{reio} , θ)
- Two derived parameters (H_0 and σ_8)

Most of the parameters
determined at (sub-) percent level.

Parameter	Λ CDM
ω_b	$2.212 \pm 0.022 \times 10^{-2}$
ω_c	0.1202 ± 0.0021
A_s	$2.091 \pm 0.034 \times 10^{-9}$
n_s	0.9643 ± 0.0057
τ_{reio}	0.0522 ± 0.0079
100θ	1.04185 ± 0.00045
H_0	67.08 ± 0.92
σ_8	0.8101 ± 0.0092

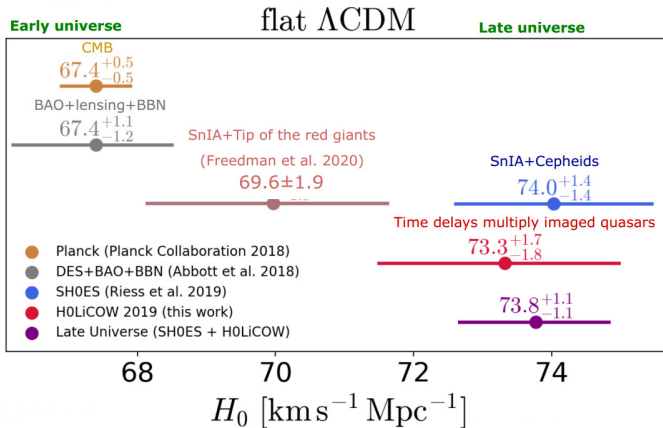
Λ CDM is a good fit to the Planck data.



- ❑ Overall Λ CDM is remarkably consistent with CMB, large scale structure and supernova data. However, there are few discrepancies within Λ CDM framework.
- ❑ Within Λ CDM cosmology, the derived value of σ_8 from Planck CMB observation is in tension with the value of σ_8 obtained from low redshift observations (cluster count, weak lensing, etc.) at $2\text{-}\sigma$ level.
- ❑ Another discrepancy within Λ CDM framework is the mismatch between the H_0 value inferred from Planck CMB observation and those obtained from direct local measurements.

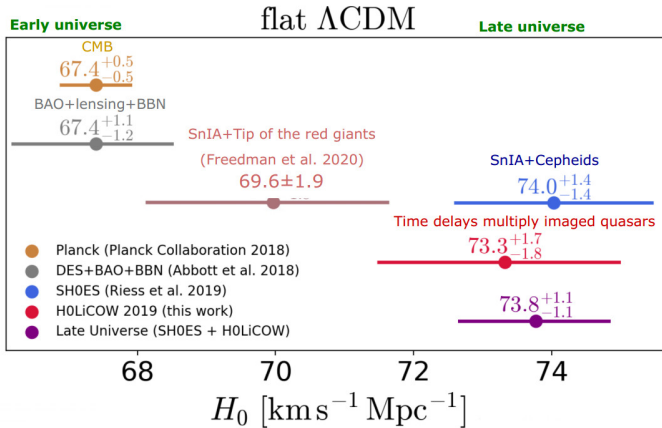


- The value of the Hubble constant (H_0) inferred from Planck CMB data is lower than the value obtained by the direct local measurements by around 4σ .





- The value of the Hubble constant (H_0) inferred from Planck CMB data is lower than the value obtained by the direct local measurements by around 4σ .

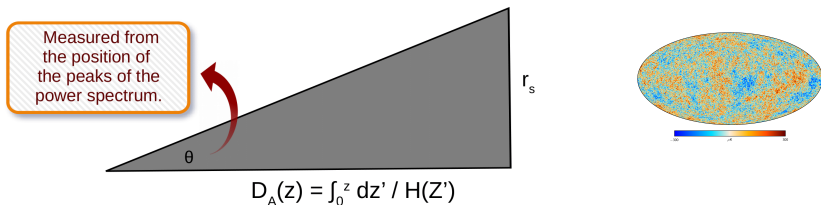


Systematic or
Hint of new
physics??



- Calculate the sound horizon at recombination : assumes model for sound speed and expansion of the universe before recombination

$$r_s = \int_{z_s}^0 \frac{c_s(z)}{H(z)} dz$$



- Infer angular diameter distance to last scattering surface \Leftrightarrow depends on H_0
- Expansion rate after recombination

$$H(z)^2 = H_0^2 [\Omega_m * (1+z)^3 + \Omega_r * (1+z)^4 + \Omega_\Lambda]$$

This is model dependent \Rightarrow New physics can change H_0 significantly.



- Peak position in the angular power spectrum is fixed from the observations $\Rightarrow \theta_s$ is fixed.

- **Late Universe solution:**

(e.g modified DE equation of state, dynamical dark energy (Planck 2015, 2018), Interacting dark matter dark energy (Di Valentino+2019), decaying DM (Poulin+2018, Pettis+2019), Modified gravity (Raveri+ 2019)etc.)

- Change in late-time dynamics of dark matter and/or dark energy

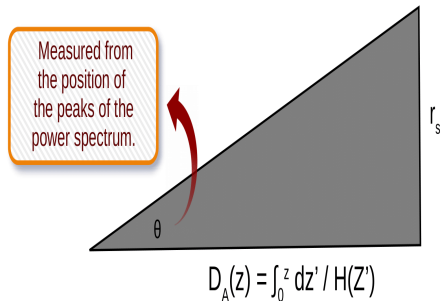


r_s is fixed and $H(z)$ at small z changes.



Increase H_0 to keep $D_A(z)$ fixed.

- Highly constrained by BAO, Supernovae and other observations.





Early Universe solution:

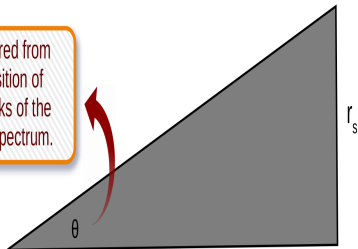
- Change the sound horizon

$$r_s = \int_{z_s}^0 \frac{c_s(z)}{H(z)} dz$$

- Change the expansion history $H(z)$ of the Universe at early times.
(e.g. Early dark energy (Polin+ 2018, Agarwal+2019), dark radiation, etc.)
- Change the Sound speed $c_s(z)$
(e.g. DM-baryon interaction (Boddy+ 2018))
- Change the time of decoupling.
(e.g. changing recombination (Hart&chluba 2019, Ivanov+ 2020))

- Problem with many other observations

Measured from the position of the peaks of the power spectrum.



$$D_A(z) = \int_0^z dz' / H(z')$$

- Change the inference of θ_s from CMB.
(e.g. Neutrino interaction (Kreisch et.al. 2019), Dark neutrino interaction (Ghosh et.al. 2019))



- ❑ Neutrinos free stream through the photon-baryon plasma at early times.
- ❑ Free-streaming neutrinos \Rightarrow Phase shift in the CMB power spectra towards larger scales and slight suppression in the amplitude.



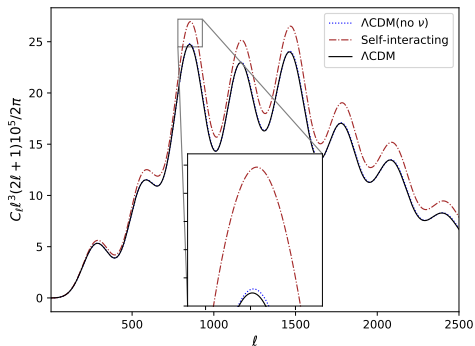
- ❑ Neutrinos free stream through the photon-baryon plasma at early times.
- ❑ Free-streaming neutrinos \Rightarrow Phase shift in the CMB power spectra towards larger scales and slight suppression in the amplitude.
- ❑ Neutrino self-interactions can be mediated by scalars and gauge bosons which are motivated from different particle physics models.

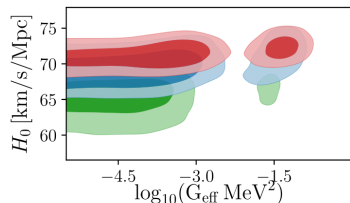
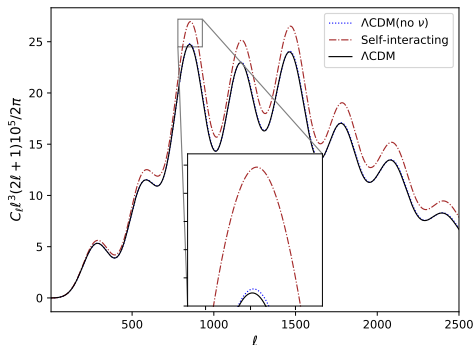
$$\mathcal{L} = g_\phi \sum_{\alpha\beta} g_{\alpha\beta} \phi \bar{\nu}_\alpha \nu_\beta, \quad \mathcal{L} = g_X \sum_{\alpha\beta} \bar{\nu}_\alpha g_{\alpha\beta} \gamma^\mu P_L \nu_\beta X_\mu,$$

- ❑ Neutrino self interaction delays the time at which neutrinos begin to free-stream.

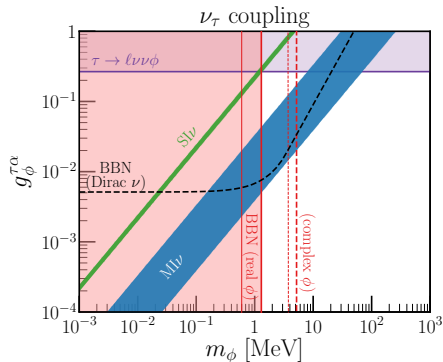
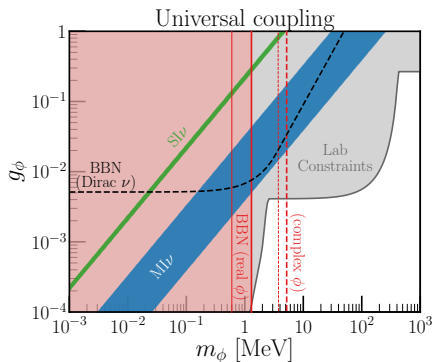


Modes entering the horizon while neutrinos are still tightly-coupled will not receive the associated phase shift and amplitude reduction.





- Bimodal posterior (Moderate self-interaction (MI) and strong self-interaction (SI) region)
- Flavour universal self interaction is strongly constrained by BBN and Lab bounds



- Flavor-universal G_{eff} is excluded by lab and BBN constraints .
- Only some part of the cosmological allowed MI region of ν_τ interaction is allowed by lab and BBN constraints.
- ν_e and ν_μ self-interactions is also disfavored by the lab and BBN constraints.
- Cosmological bounds assume flavour universal interaction.



- Flavour specific self interaction

$$\mathcal{L} = g_\phi \sum_{\alpha,\beta} g_{\alpha\beta} \phi \bar{\nu}_\alpha \nu_\beta, \quad \mathcal{L} = g_X \sum_{\alpha,\beta} \bar{\nu}_\alpha g_{\alpha\beta} \gamma^\mu P_L \nu_\beta X_\mu,$$

- Four cases : ν_e, ν_μ, ν_τ and the flavour universal self-interactions.

The matrix $g_{\alpha\beta}$ defines the flavour dependence of the interactions. $g_{\alpha\beta}$ will be $\delta_{\alpha\beta}$ for universal interaction and for interaction in a particular flavour it will be a diagonal matrix with only one among $g_{ee}, g_{\mu\mu}$ or $g_{\tau\tau}$ set to be one.



- Flavour specific self interaction

$$\mathcal{L} = g_\phi \sum_{\alpha\beta} g_{\alpha\beta} \phi \bar{\nu}_\alpha \nu_\beta, \quad \mathcal{L} = g_X \sum_{\alpha\beta} \bar{\nu}_\alpha g_{\alpha\beta} \gamma^\mu P_L \nu_\beta X_\mu,$$

- Four cases : ν_e, ν_μ, ν_τ and the flavour universal self-interactions.

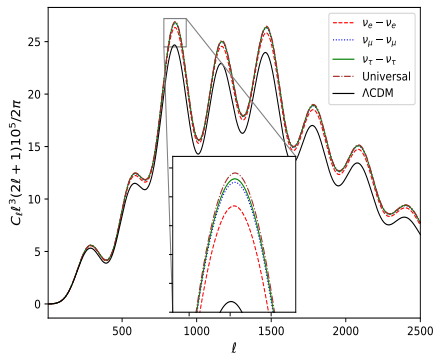
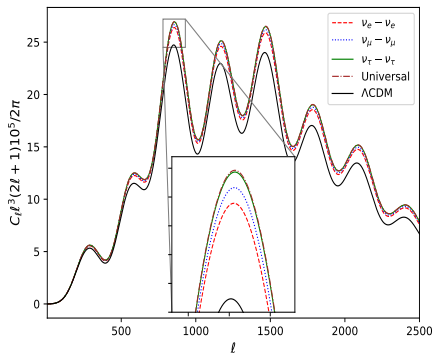
The matrix $g_{\alpha\beta}$ defines the flavour dependence of the interactions. $g_{\alpha\beta}$ will be $\delta_{\alpha\beta}$ for universal interaction and for interaction in a particular flavour it will be a diagonal matrix with only one among $g_{ee}, g_{\mu\mu}$ or $g_{\tau\tau}$ set to be one.

- The perturbed Boltzmann equation

$$\frac{\partial \Psi_i}{\partial \tau} + i \frac{q(\vec{k} \cdot \hat{n})}{\epsilon} \Psi_i + \frac{d \ln f_0}{d \ln q} \left[\dot{\eta} - \frac{\dot{h} + 6\dot{\eta}}{2} (\hat{k} \cdot \hat{n})^2 \right] = -\Gamma_{ij} \Psi_j,$$

where $\Gamma_{ij} = U^\dagger g_{\alpha\beta} U \tau_\nu^{-1}$. and

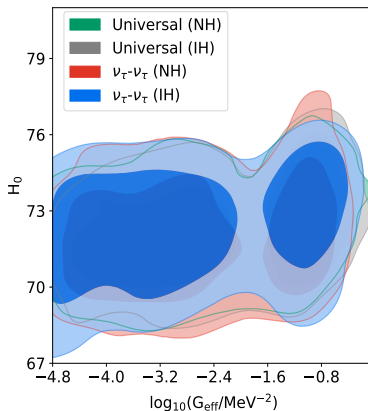
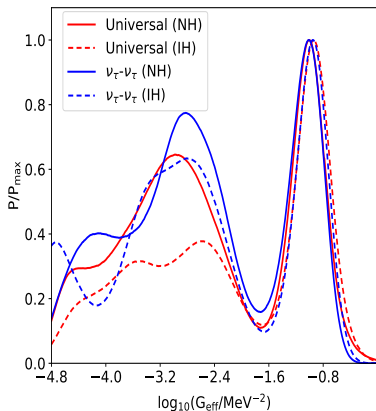
$$\tau_\nu^{-1} = a n_\nu \langle \sigma \nu \rangle = \frac{3}{2} \frac{\zeta(3)}{\pi^2} a G_{\text{eff}}^2 T_\nu^5,$$



Effect of self-interaction $\propto \sum_{i,j} \Gamma_{ij}$.

Hierarchy	Universal Int.	ν_τ Int.	ν_μ Int.	ν_e Int.
NH	$3\tau_\nu^{-1}$	$2.965 \tau_\nu^{-1}$	$2.643 \tau_\nu^{-1}$	$2.308 \tau_\nu^{-1}$
IH	$3\tau_\nu^{-1}$	$2.88116 \tau_\nu^{-1}$	$2.809 \tau_\nu^{-1}$	$2.309 \tau_\nu^{-1}$

Parameters in flavour specific self interaction



- Similar to the Universal case, G_{eff} has a bimodal posterior. We call the two modes strong interaction (MI) and moderate interaction (MI) region.
- In case of $\nu_{\tau}-\nu_{\tau}$ interaction, $H_0 = 73.3 \pm 1.8$ and $H_0 = 73.3^{+2.0}_{-1.4}$ for NH and IH respectively.

□ Parameters for the Moderate self Interaction.

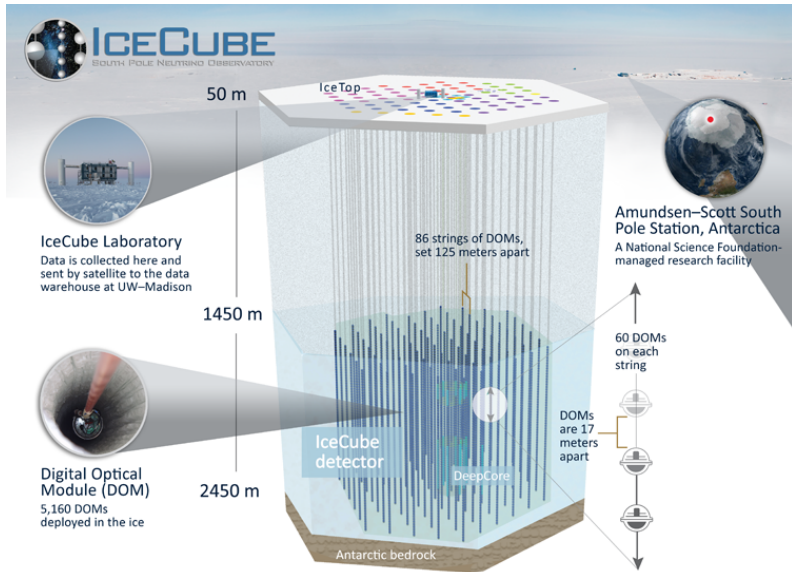
Parameters	Universal Int.	Universal Int.	$\nu_\tau - \nu_\tau$ Int.	$\nu_\tau - \nu_\tau$ Int.
	(NH)	(IH)	(NH)	(IH)
$100\theta_s$	$1.0412^{+0.0009}_{-0.0011}$	$1.0411^{+0.0010}_{-0.0012}$	$1.0413^{+0.0010}_{-0.0011}$	1.0411 ± 0.0010
n_s	0.986 ± 0.014	$0.988^{+0.015}_{-0.013}$	$0.984^{+0.016}_{-0.013}$	$0.986^{+0.015}_{-0.013}$
m_0	$0.057^{+0.022}_{-0.050}$	$0.058^{+0.025}_{-0.055}$	$0.052^{+0.017}_{-0.049}$	$0.051^{+0.018}_{-0.042}$
$\log_{10} G_{\text{eff}}$	$-3.25^{+0.94}_{-0.65}$	-3.23 ± 0.75	$-3.20^{+1.3}_{-1.5}$	$-3.33^{+1.1}_{-0.75}$
N_{eff}	$3.75^{+0.31}_{-0.38}$	$3.77^{+0.29}_{-0.32}$	$3.69^{+0.28}_{-0.33}$	$3.74^{+0.33}_{-0.37}$
H_0	$72.0^{+1.7}_{-1.9}$	71.8 ± 1.7	71.6 ± 1.8	71.9 ± 1.9

□ Parameters for the strong self Interaction.

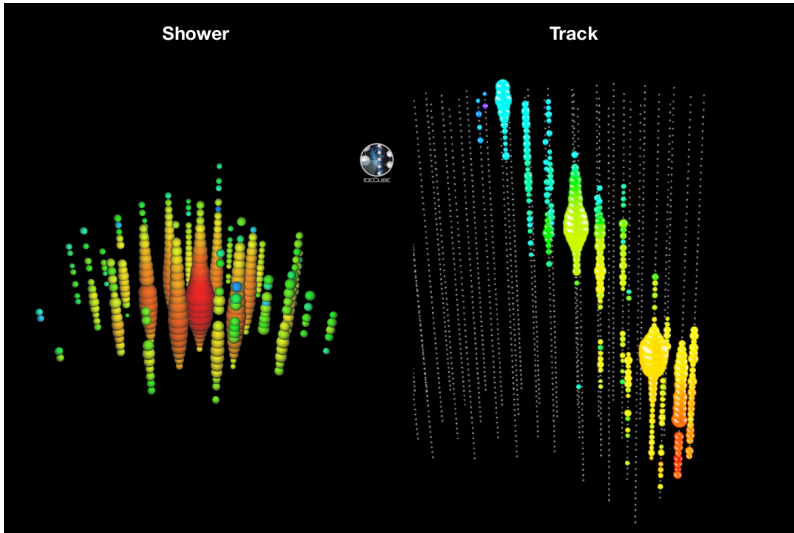
$100\theta_s$	1.0456 ± 0.0012	$1.0454^{+0.0012}_{-0.0014}$	1.0453 ± 0.0011	1.0451 ± 0.0010
n_s	0.959 ± 0.014	0.959 ± 0.013	$0.962^{+0.015}_{-0.011}$	0.964 ± 0.012
m_0	$0.081^{+0.038}_{-0.069}$	$0.082^{+0.045}_{-0.054}$	$0.084^{+0.044}_{-0.060}$	0.092 ± 0.062
$\log_{10} G_{\text{eff}}$	$-1.01^{+0.21}_{-0.18}$	$-0.95^{+0.22}_{-0.18}$	$-1.05^{+0.26}_{-0.17}$	$-0.97^{+0.19}_{-0.12}$
N_{eff}	3.93 ± 0.35	3.97 ± 0.36	$4.07^{+0.39}_{-0.35}$	$4.09^{+0.42}_{-0.21}$
H_0	72.8 ± 1.9	72.9 ± 1.7	73.3 ± 1.8	$73.3^{+2.0}_{-1.4}$

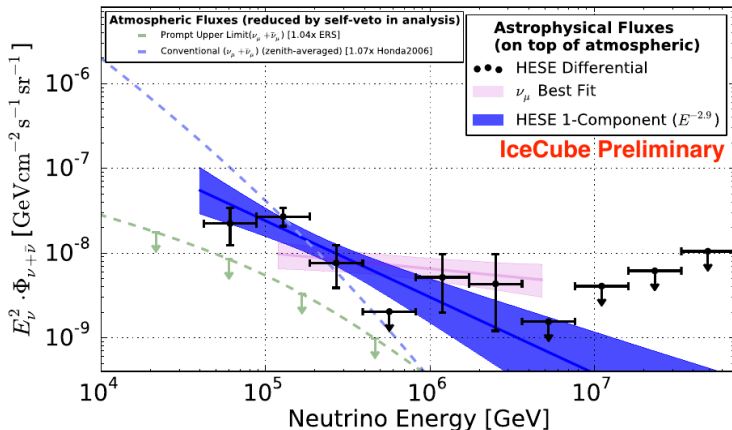
IceCube and self-interacting neutrinos

IceCube Neutrino Observatory



¹From: <https://icecube.wisc.edu/science/icecube/detector>



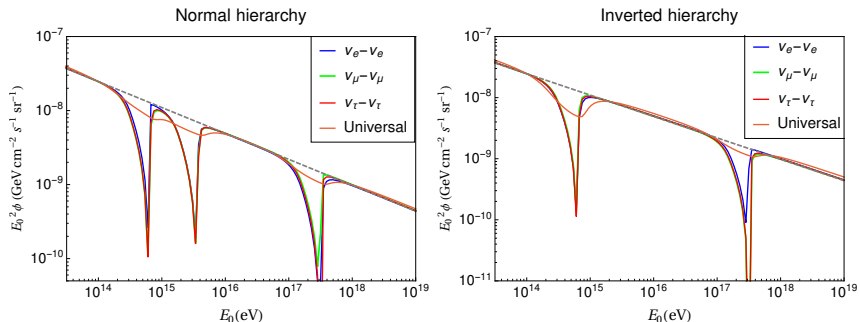


The best fit for single power law flux is

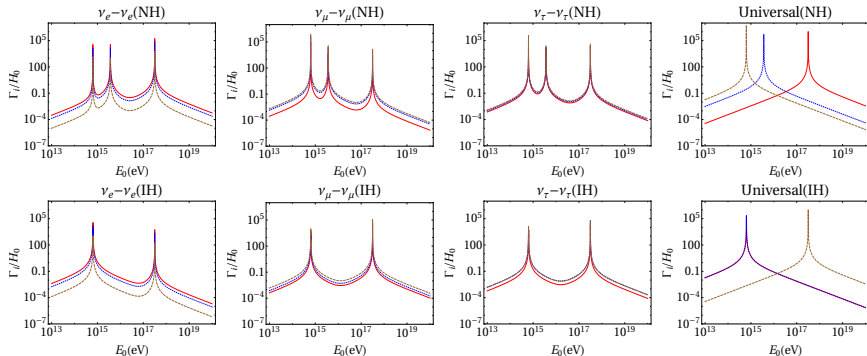
$$E_0^2 \phi = (2.46 \pm 0.8) \times 10^{-8} \left(\frac{E_0}{100 \text{ TeV}} \right)^{-0.92} \text{ GeV cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$



- The astrophysical neutrinos can interact with cosmic neutrinos and produce an on shell mediator $\nu\nu \rightarrow \phi$, when the neutrino energy $E_\nu = m_\phi^2/(2m_\nu)$.
- The resonant absorption of astrophysical neutrinos will show up as dips in the IceCube flux.



$$m_\phi = 10^{6.9} \text{eV}, g_\phi = 10^{-1.5}, \gamma = 2.35 \text{ and } m_0 = 10^{-4} \text{eV}$$

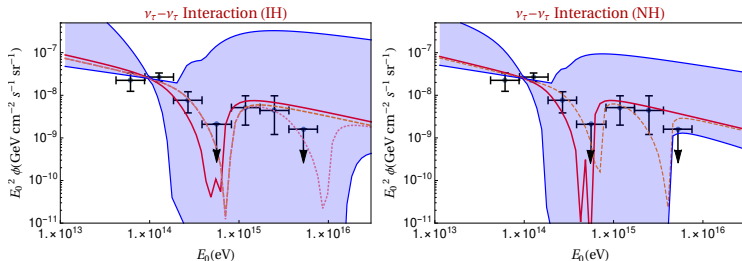


$$\sigma_{ijkl} = \sigma(\bar{\nu}_i \nu_j \rightarrow \bar{\nu}_k \nu_l) = \frac{1}{4\pi} |g_{kl}|^2 |g_{ij}|^2 \frac{g_\phi^4 s_j}{(s_j - m_\phi^2)^2 + m_\phi^2 \Gamma_\phi^2},$$

where $s \approx 2E_i m_i$ and $g_{ij} = g_{\alpha\beta} U_{\alpha i}^* U_{\beta j}$ $\Gamma_i = \sum_j \tilde{n}_j \sigma_{ij}$



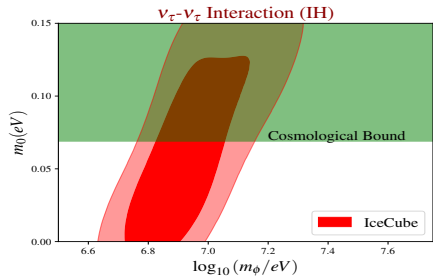
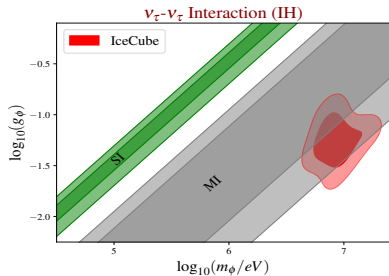
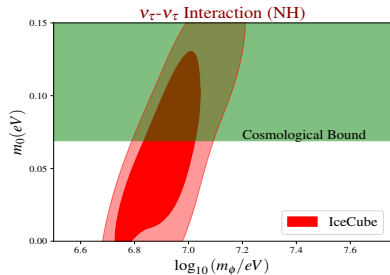
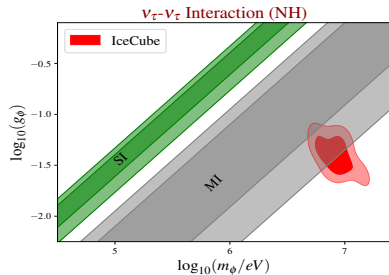
- Four parameters (m_ϕ , g_ϕ , m_0 and γ) have been varied to do the MCMC analysis with IceCube data.



Parameters	NH	IH
m_0 (eV)	$0.062^{+0.042}_{-0.046}$	$0.062^{+0.038}_{-0.058}$
$\log_{10} m_\phi / \text{eV}$	6.92 ± 0.11	$6.936^{+0.089}_{-0.14}$
$\log_{10} g_\phi$	-1.40 ± 0.13	-1.25 ± 0.17
γ	2.50 ± 0.17	$2.46^{+0.18}_{-0.16}$

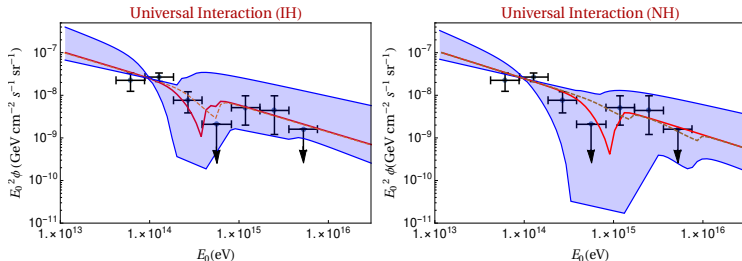
- Both one dip and two dip solution are possible.

Parameter estimation from flux at IceCube





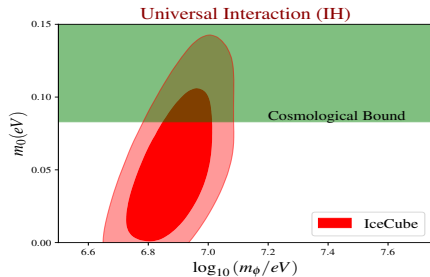
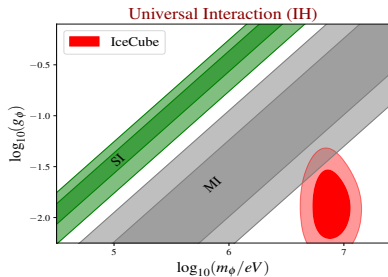
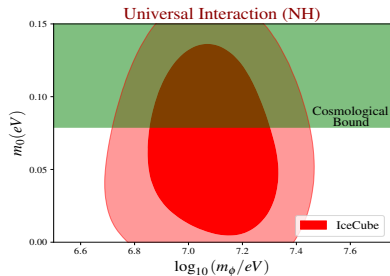
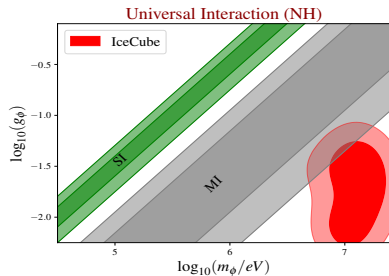
- Four parameters (m_ϕ , g_ϕ , m_0 and γ) have been varied to do the MCMC analysis with IceCube data.



Parameters	NH	IH
$m_0(\text{eV})$	$0.067^{+0.038}_{-0.046}$	$0.052^{+0.034}_{-0.040}$
$\log_{10} m_\phi/\text{eV}$	7.09 ± 0.15	6.883 ± 0.091
$\log_{10} g_\phi$	-1.75 ± 0.31	-1.87 ± 0.20
γ	$2.66^{+0.21}_{-0.18}$	2.63 ± 0.17

Only single dip solution possible.

Parameter estimation from flux at IceCube



Self-interacting sterile neutrinos



- ❑ Short baseline neutrino experiments, like LSND and MiniBooNE experiments, pointed towards the existence of eV mass scale sterile neutrinos. ([PhysRevD.64.112007](#), [PhysRevLett.12.221801](#))
- ❑ There is a long-standing problem of accommodating eV mass sterile neutrinos in the cosmological context.
- ❑ Non zero mixing between active and sterile neutrinos \Rightarrow Thermal equilibrium ($T_s = T_\nu$)



Contribute to N_{eff}

Current bound

$$N_{\text{eff}} = 2.96^{+0.34}_{-0.33} \text{ (CMB + lensing + BAO)}$$

$$N_{\text{eff}} = 2.93 \pm 0.23 \text{ (CMB + lensing + BAO + BBN)}$$



- ❑ Short baseline neutrino experiments, like LSND and MiniBooNE experiments, pointed towards the existence of eV mass scale sterile neutrinos. ([PhysRevD.64.112007](#), [PhysRevLett.12.221801](#))
- ❑ There is a long-standing problem of accommodating eV mass sterile neutrinos in the cosmological context.
- ❑ Non zero mixing between active and sterile neutrinos \Rightarrow Thermal equilibrium ($T_s = T_\nu$)



Contribute to N_{eff}

Current bound

$$N_{\text{eff}} = 2.96^{+0.34}_{-0.33} \text{ (CMB + lensing + BAO)}$$

$$N_{\text{eff}} = 2.93 \pm 0.23 \text{ (CMB + lensing + BAO + BBN)}$$

- ❑ In simple 3+1 neutrinos scenario with eV sterile neutrinos and non-zero mixing angle
 $N_{\text{eff}} \approx 4$

Conflict with the existing bounds \Rightarrow Hints for new physics??



- Self-interaction between sterile neutrinos has been proposed to alleviate this tension. ¹

$$\mathcal{L}_{\text{int}} = g_X \bar{\nu}_s \gamma_\mu \frac{1}{2} (1 - \gamma_5) \nu_s X^\mu$$

- The self-interaction generates a effective potential that suppresses the effective mixing of the sterile and active states in the early universe.
- Mixing angle is large at low temperature, sterile and active neutrino recouple at low temperature.

³Hannestad et al. Phys. Rev. Lett. 112(3):031802

Dasgupta et al. Phys. Rev. Lett. 112(3):031803



- Self-interaction between sterile neutrinos has been proposed to alleviate this tension. ¹

$$\mathcal{L}_{\text{int}} = g_X \bar{\nu}_s \gamma_\mu \frac{1}{2} (1 - \gamma_5) \nu_s X^\mu$$

- The self-interaction generates a effective potential that suppresses the effective mixing of the sterile and active states in the early universe.
- Mixing angle is large at low temperature, sterile and active neutrino recouple at low temperature.
- Quantitatively the effect of self interaction can be studied using QKEs ([Phys. Rev. D86 \(2012\) 053009](#))

$$i \frac{d\varrho}{dt} = [\Omega, \varrho] + C[\varrho]$$

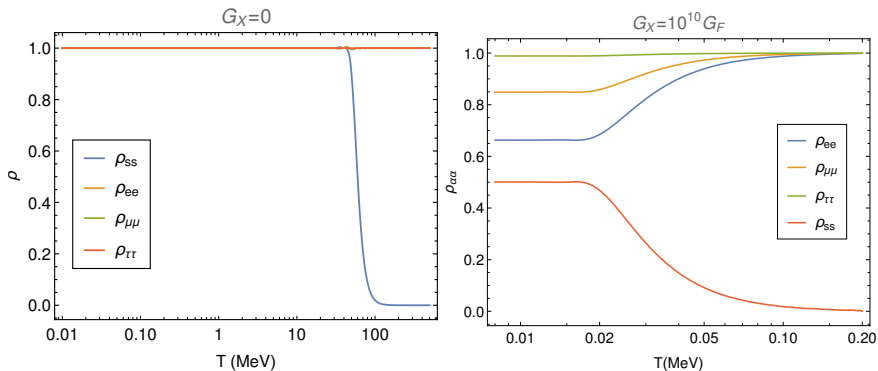
where $C[\varrho]$ is the collision term and

$$\Omega = \frac{1}{2p} U^\dagger M^2 U - \sqrt{2} \frac{8p G_F}{3} \left[\left(\frac{\mathcal{E}_l}{M_W^2} + \frac{\mathcal{E}_\nu}{M_Z^2} \right) \right] - \sqrt{2} G_X \left[\frac{8p}{3} \frac{\mathcal{E}_s}{M_X^2} \right]$$

³Hannestad et al. Phys. Rev. Lett. 112(3):031802

Dasgupta et al. Phys. Rev. Lett. 112(3):031803

Self interacting sterile neutrinos



- 1 Self interaction delays the mixing between active and sterile neutrinos till sub MeV temperatures.
- 2 N_{eff} is always within the allowed range.



- The self-interaction between the sterile neutrinos induces scattering among them and the scattering rate is given as

$$\Gamma_X = n_{\nu_s} \langle \sigma v \rangle \simeq G_X^2 T_{\nu_s}^5$$

- Free-streaming neutrinos suppress the growth of perturbations.
- Self interaction delays the free-stream regime .
- Sufficiently large $G_x \Rightarrow$ No free stream regime till neutrino turn non-relativistic.

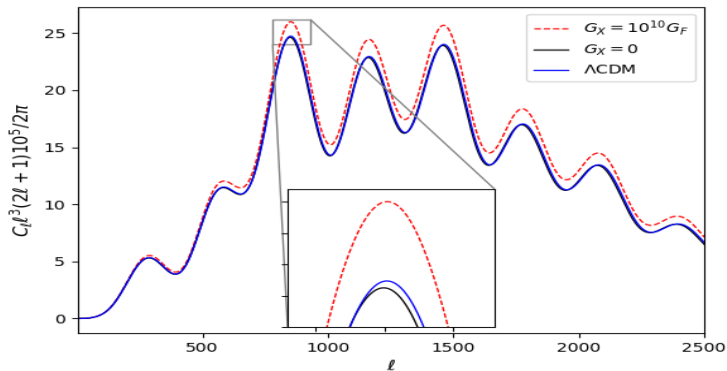
$$G_X^2 T_{\nu_s}^5 \approx H(T_\gamma). \quad \Rightarrow G_X \sim 10^{10} G_F$$

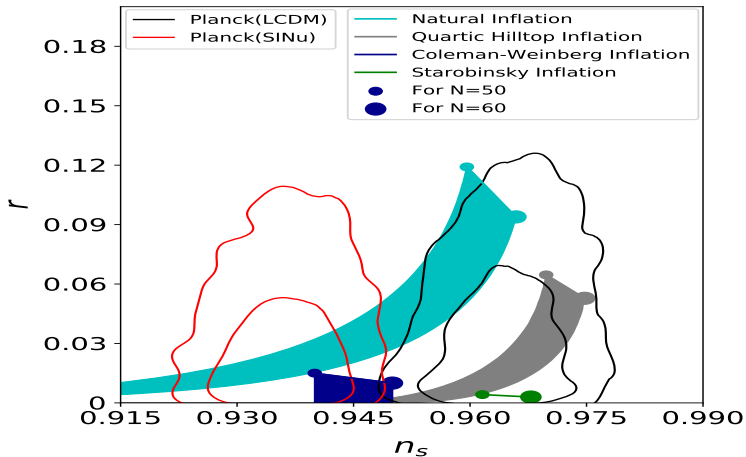
⁵Mirizzi et al. PhysRevD.91.025091

Xiaoyong Chu et al. JCAP10(2105)011



Effect of self interacting sterile neutrinos on the CMB temperature anisotropy power spectrum.







Parameter	Λ CDM	$SI\nu$ model
ω_b	$2.212 \pm 0.022 \times 10^{-2}$	$2.177 \pm 0.022 \times 10^{-2}$
ω_{cdm}	0.1202 ± 0.0021	0.1220 ± 0.0021
$100\theta_s$	1.04185 ± 0.00045	1.04481 ± 0.00045
n_s	0.9643 ± 0.0057	0.9361 ± 0.0055
τ_{reio}	0.0522 ± 0.0079	0.0476 ± 0.0074
r	< 0.0460	< 0.0383
H_0	67.08 ± 0.92	66.00 ± 0.91
A_s	$2.091 \pm 0.034 \times 10^{-9}$	$1.988 \pm 0.030 \times 10^{-9}$
σ_8	0.8101 ± 0.0092	0.7769 ± 0.0086

Table: The best fit values with $1\text{-}\sigma$ error for all the parameters obtained from the MCMC analyses for both the models considered are listed here.



- 1 Flavour specific self interaction (ν_τ - ν_τ) interaction can resolve the H_0 tension. Posterior of G_{eff} is bimodal in nature.
- 2 Flavour specific self interaction can also explain the dip in neutrino flux at IceCube.
- 3 The values of self-interaction coupling constant and mediator mass required for explaining the IceCube dips are inconsistent with the strong neutrino self-interactions preferred by the cosmological data.
- 4 The required amount of self-interaction between tau neutrinos (ν_τ) in inverted hierarchy for explaining IceCube dips is consistent with the moderate self-interaction region of cosmological bounds at 1 - σ level.
- 5 For the case of other interactions and hierarchies, the IceCube preferred amount of self-interaction is consistent with moderate self-interaction region of cosmological bounds at 2 - σ level only.
- 6 Self interaction between sterile neutrinos had been proposed to make eV mass sterile neutrinos viable with cosmological observations .
- 7 The value of spectral index (n_s) shifts to a lower value in $SI\nu$ model, which makes a significant impact on the validity of inflation models.



- 1 Flavour specific self interaction (ν_τ - ν_τ) interaction can resolve the H_0 tension. Posterior of G_{eff} is bimodal in nature.
- 2 Flavour specific self interaction can also explain the dip in neutrino flux at IceCube.
- 3 The values of self-interaction coupling constant and mediator mass required for explaining the IceCube dips are inconsistent with the strong neutrino self-interactions preferred by the cosmological data.
- 4 The required amount of self-interaction between tau neutrinos (ν_τ) in inverted hierarchy for explaining IceCube dips is consistent with the moderate self-interaction region of cosmological bounds at 1 - σ level.
- 5 For the case of other interactions and hierarchies, the IceCube preferred amount of self-interaction is consistent with moderate self-interaction region of cosmological bounds at 2 - σ level only.
- 6 Self interaction between sterile neutrinos had been proposed to make eV mass sterile neutrinos viable with cosmological observations .
- 7 The value of spectral index (n_s) shifts to a lower value in $SI\nu$ model, which makes a significant impact on the validity of inflation models.

Thanks You!!