

Global 21-cm signal and its implications

Atrideb Chatterjee

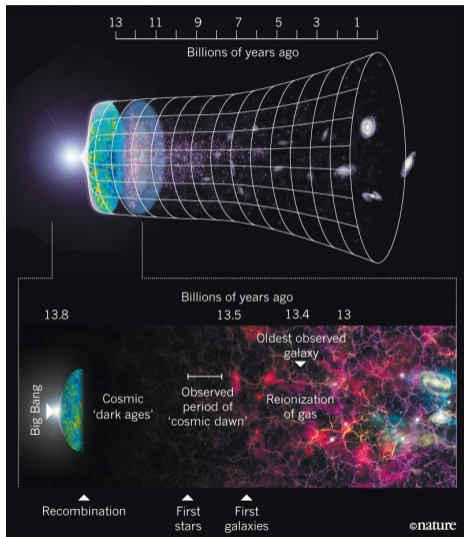
NCRA-TIFR

Collaborator: Tirthankar Roychoudhury, Pratika Dayal, Sourav Mitra, Anne Hutter, Raffaella Schneider

SOTU seminar, TIFR, Dec 2021

- Time evolution of the Universe
- Basics of the Global 21-cm signal
- Redshift Evolution of the signal
- Observation and Existing theory
- Implications of the signal
 - Constraining mass of the Warm Dark Matter
 - Estimating SFRD of PopIII stars
 - CosmoReionMC: A parameter estimation package to constrain cosmological and astrophysical parameters
- Future outlook
- Summary

Time evolution of the Universe



Basics of the Global 21-cm signal

- The cosmic mean 21-cm differential brightness temperature

$$\delta T_b \propto \left(\frac{T_S - T_\gamma}{T_S} \right) x_{HI}$$

(Field 1958)

- The spin temperature T_S can be written as

$$T_S^{-1} \simeq \frac{T_\gamma^{-1} + x_\alpha T_k^{-1}}{1 + x_\alpha}$$

- T_k is the kinetic temperature of the IGM; x_{HI} is the neutral Hydrogen fraction.
- x_α is the Lyman- α coupling coefficient.

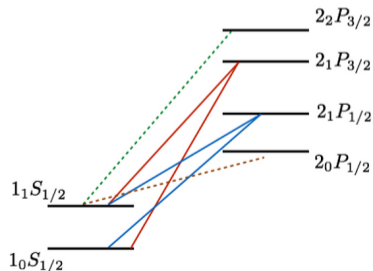
Evolution of kinetic temperature & W-F effect

- The evolution equation of gas kinetic temperature is given by

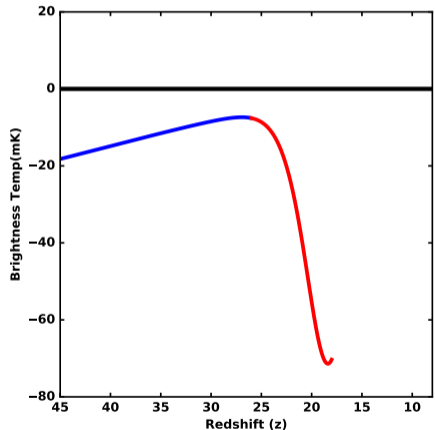
$$\frac{dT_k}{dz} = \frac{2T_k}{1+z} - \frac{2}{3H(z)(1+z)} \frac{\epsilon_X}{k_B n} \quad [\epsilon_X : X - \text{ray emissivity}]$$

(Furlanetto et al. 2006)

- Lyman- α coupling (W-F effect)

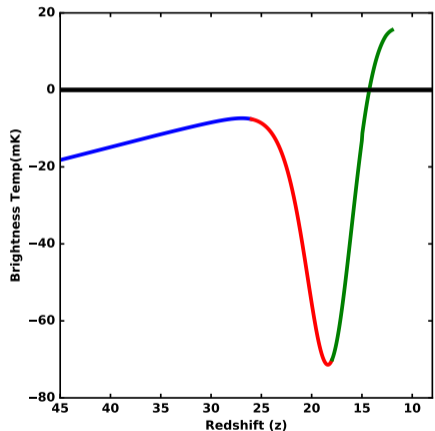


Redshift evolution of the brightness temperature



$$x_{\alpha} \gg 1 \quad T_s \approx T_k \quad \delta T_b = \left(\frac{T_k - T_{\gamma}}{T_k} \right) x_{HI}$$

Redshift evolution of the brightness temperature



$$x_{\alpha} \gg 1$$

$$T_s \approx T_k$$

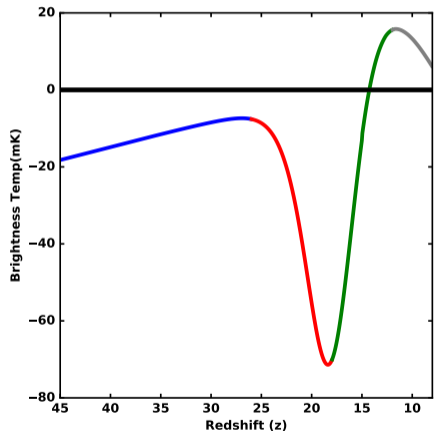
$$\delta T_b = \left(\frac{T_k - T_{\gamma}}{T_k} \right) x_{HI}$$

$$x_{\alpha} \gg 1$$

$$T_s \approx T_k$$

$$\delta T_b = \left(\frac{T_k - T_{\gamma}}{T_k} \right) x_{HI}$$

Redshift evolution of the brightness temperature



$$x_{\alpha} \gg 1$$

$$T_s \approx T_k$$

$$\delta T_b = \left(\frac{T_k - T_{\gamma}}{T_k} \right) x_{HI}$$

$$x_{\alpha} \gg 1$$

$$T_s \approx T_k$$

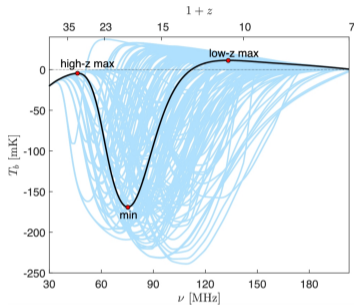
$$\delta T_b = \left(\frac{T_k - T_{\gamma}}{T_k} \right) x_{HI}$$

$$x_{HI} \rightarrow 0.0$$

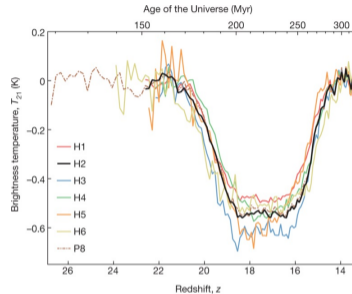
$$T_s \approx T_k$$

$$\delta T_b = \left(\frac{T_k - T_{\gamma}}{T_k} \right) x_{HI}$$

Observation & Existing theory



Model (Cohen et al. 2017)



Observation (Bowman et al. 2018)

Existing theory cannot explain the excess dip in the observed signal.

- Debate over cosmological origin of the signal. (Hills et al. 2018)
- Global 21-cm signal probes the structure formation.

Ruling out 3 keV warm dark matter using 21 cm-EDGES data

In collaboration with Pratika Dayal, Tirthankar Roy Choudhury & Anne Hutter

- Problems with the Λ CDM model
 - Over-predicts the number of satellite and field galaxies compared to the observations.
(Moore et al. 1999)
 - Predicting cuspy halo profiles as opposed to the observationally preferred constant density cores.
(Navarro et al. 1997)
 - Predicts massive concentrated Galactic sub-halos inconsistent with observations.
(Boylan-Kolchin et al. 2012)
 - Limited success with baryonic feedback.
(Boylan-Kolchin et al. 2012)
- Structure formation is delayed in WDM models.
- Global 21-cm signal probes the structure formation.

X-ray heating & Lyman- α coupling

- The energy density injected into the gas per second through X-ray heating is given by

$$\epsilon_X \propto f_{X,h} \times \left(\frac{\dot{\rho}_*}{M_\odot \text{yr}^{-1} \text{Mpc}^{-3}} \right) \text{erg/sec/mpc}^3$$

(Ranalli et al. 1992)

- $\dot{\rho}_*$ is the Star Formation Rate Density (SFRD).
- $f_{X,h} = f_X \times f_h$
- The Lyman- α coupling coefficient can be expressed as

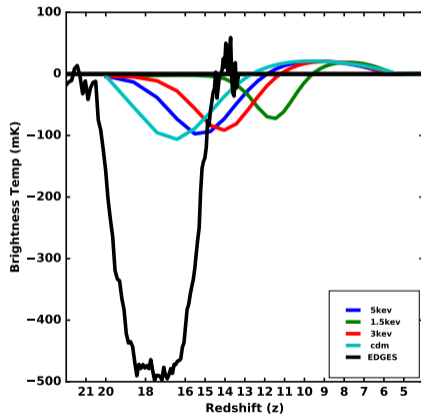
$$x_\alpha = 1.81 \times 10^{11} (1+z)^{-1} S_\alpha J_\alpha$$

J_α is the background Ly α flux.

- The SFRD is taken from a semi-numerical code DELPHI.

(Dayal et al. 2014)

Global 21-cm signal from different DM models



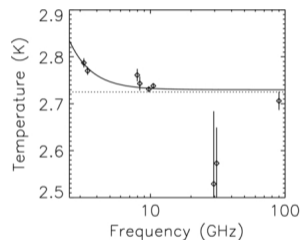
Excess Radio background

- Arcade-2 observed an excess radio background along with the CMB background. (Fixsen et al. 2011)
- Possible origin is faint extragalactic sources although Galactic emission origin can not be excluded. (Seiffert et al. 2011)
- Radio emissivity ϵ_R is modelled as

$$\epsilon_R \propto f_R \dot{\rho}_*$$

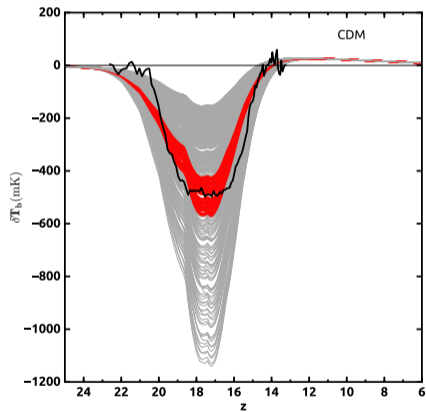
- f_R accounts for any differences between the local and high- z observation.
- Background temperature becomes

$$T_\gamma = T_{CMB} + T_R$$

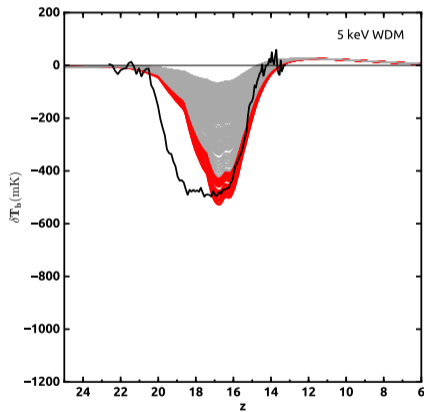


(Fixsen et al. 2011)

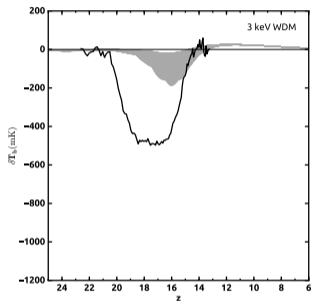
CDM and 5 keV WDM



(a) CDM



(b) 5 keV WDM



3keV WDM

$m_{\chi} \leq 3 \text{ keV}$ is ruled out

(AC, PD, TRC, AH 2019, MNRAS)

Limitation: Only PopII stars no PopIII

A hint on the metal-free star formation rate density from 21-cm-EDGES data

In collaboration with Pratika Dayal, Tirthankar Roy Choudhury & Raffaella Schneider

- The first generation of stars are known as Pop III stars.
- Postulated to explain the metallicity-gap from Big Bang Nucleosynthesis to the metal-rich PopII stars in high redshift galaxies.
- Massive, short-lived and yet to be detected from direct observation.
(Tumlinson 2006)
- 21-cm signal at $z \approx 12 - 21$ probes the epoch when the first stars form.

- The global 21 cm brightness temperature can be computed once SFRD ($\dot{\rho}_*$) is known.
- $\dot{\rho}_*$ is assumed to have contribution from both PopII and PopIII stars

$$\dot{\rho}_* = \dot{\rho}_*^{II} + \dot{\rho}_*^{III}$$

- $\dot{\rho}_*^{III}$ is taken as a fourth degree polynomial

$$\log \left[\frac{\dot{\rho}_*^{III}(z)}{\dot{\rho}_*^{III}(z_{\text{mean}})} \right] = a_1(z - z_{\text{mean}}) + a_2(z - z_{\text{mean}})^2 + a_3(z - z_{\text{mean}})^3 + a_4(z - z_{\text{mean}})^4$$

- We take PopII contribution from DELPHI.

(Dayal et al. 2014)

Contribution from both PopII and PopIII

- The X-ray energy density per unit time, ϵ_X , can be written as

$$\epsilon_X \propto \left(f_{X,h}^{II} \dot{\rho}_*^{II} + f_{X,h}^{III} \dot{\rho}_*^{III} \right)$$

- \dot{n}_ν , the production rate of Ly- α photons per unit frequency per unit comoving volume, get contribution from both PopII and PopIII stars

$$\dot{n}_\nu = \dot{n}_\nu^{II} + \dot{n}_\nu^{III}$$

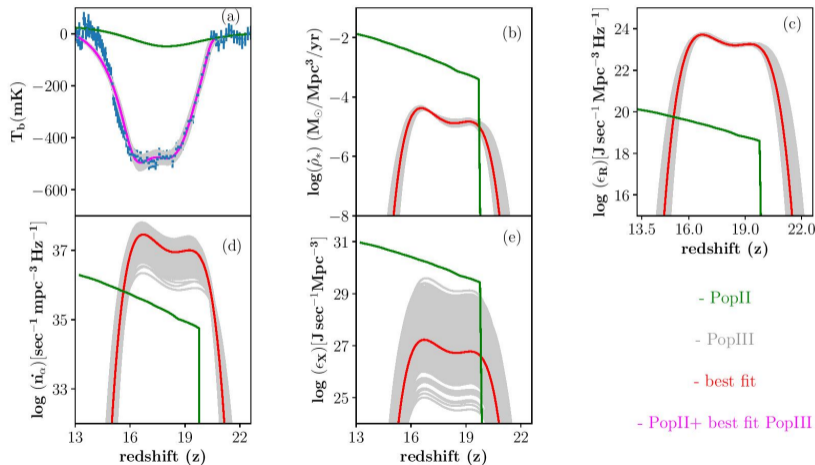
- $\dot{n}_\nu^{III} = f_\alpha^{III} \dot{\rho}_*^{III}(z')$, f_α^{III} being the proportionality factor.

- Radio emissivity ϵ_R can be expressed as

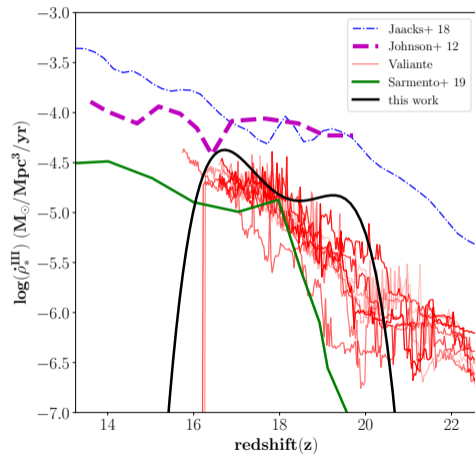
$$\epsilon_R \propto \left(f_R^{II} \dot{\rho}_*^{II} + f_R^{III} \dot{\rho}_*^{III} \right)$$

- PopII parameters are kept fixed.

MCMC analysis



Compare SFRD



CosmoReionMC: A parameter estimation package using Reionization and Cosmic dawn observations

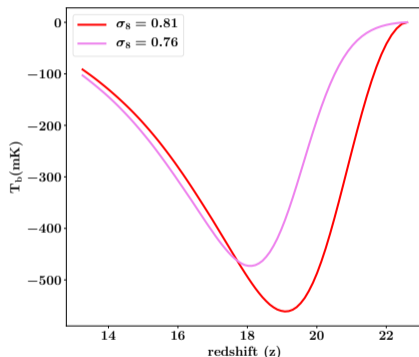
In collaboration with Tirthankar Roy Choudhury & Sourav Mitra

Constraining Reionization model from CMB observations

- Interaction of the CMB photons with the free electrons in the IGM imprinted in the CMB angular power spectrum observations. **(Sunyaev & Zeldovich 1980)**
- This observations along with an assumed reionization history measures the Thomson scattering optical depth τ . **(Planck Collaboration et al. 2016)**
- τ from the observation can then be compared with the prediction from a reionization model to constrain free parameters of the model. **(Pandolfi et al. 2011)**
- Usual approach of comparing observed τ (based on a tanh reionization) with the theoretically estimated τ based on a different reionization model is not a correct approach.
- To overcome this, connect a reionization model with a “CMB anisotropy code” like CAMB or CLASS and constrain cosmological and astrophysical parameters simultaneously. **(Qin et al. 2020)**

Motivation for adding 21-cm signal to constrain parameters

- Modelling of 21-cm signal depends on cosmological parameters.



- Inclusion of observed 21-cm signal (by EDGES) along with CMB and other reionization related observations can jointly put tighter constraint on cosmological parameters.

Reionization Model (Choudhury & Ferrera 2005,2006)

- The over-densities in the low density region of the IGM follows a lognormal distribution and change to a power-law distribution in high density region. (Miralda-Escude 2003)
- The sources of reionization are QSOs, PopII and PopIII stars.
- The collapse fraction for PopII and PopIII stars are assumed to be

$$f_{\text{coll,II}} = \frac{1}{\rho_m} \int_{M_{\text{min}}(z)}^{\infty} dM [1 - f_{\text{III}}(M, z)] M \frac{\partial n(M, z)}{\partial M},$$

$$f_{\text{coll,III}} = \frac{1}{\rho_m} \int_{M_{\text{min}}(z)}^{\infty} dM [f_{\text{III}}(M, z)] M \frac{\partial n(M, z)}{\partial M},$$

- PopIII \rightarrow PopII transition essentially determines the form of $f_{\text{III}}(M, z)$.

$$f_{\text{III}}(M, z) = f_{\text{III}}(z) = \frac{1}{2} \left[1 + \tanh \left(\frac{z - z_{\text{trans}}}{\delta z} \right) \right].$$

Photon production from stars

- The photon production rate from stellar sources is

$$\begin{aligned}\dot{n}_\nu(z) &= \dot{n}_\nu^{\text{II}} + \dot{n}_\nu^{\text{III}} \\ &= \rho_b \left[f_*^{\text{II}} f_{\text{esc}}^{\text{II}} \left(\frac{dN_\nu}{dM} \right)_{\text{II}} \frac{df_{\text{coll,II}}}{dt} + f_*^{\text{III}} f_{\text{esc}}^{\text{III}} \left(\frac{dN_\nu}{dM} \right)_{\text{III}} \frac{df_{\text{coll,III}}}{dt} \right]\end{aligned}$$

- ρ_b : Mean comoving baryon density.
- f_* : Star formation efficiency ; f_{esc} : Escape fraction of the ionizing photons entering IGM
- $\frac{dN_\nu}{dM}$: Number of photons emitted per frequency range per unit mass of stars.
- $\frac{df_{\text{coll}}}{dt}$: Rate of collapse fraction

Observable quantities from the Reionization model

- The reionization model can predict the redshift distribution of the Lyman-limit systems dN_{LL}/dz .

- The hydrogen photoionization rate (Γ_{PI}) can be obtained from

$$\Gamma_{PI}(z) = (1+z)^3 \int_{\nu_H}^{\infty} d\nu \lambda_{mfp}(z) \dot{n}_\nu(z) \sigma_H(\nu),$$

- $\sigma_H(\nu)$: Hydrogen photoionization cross-section.

- The mean free path of the photons (λ_{mfp}) given by

$$\lambda_{mfp} = \frac{\lambda_0}{[1 - F_\nu]^{2/3}}$$

- F_ν : Volume fraction of the ionized region; λ_0 : Free parameter of the model.

- Global 21-cm signal can also be calculated from this model.

MCMC analysis: Parameter space, Likelihood and Constraints

- Free parameters

$$\Theta = \underbrace{\{H_0, \Omega_b h^2, \Omega_c h^2, A_s, n_s\}}_{\Theta_{\text{cosmo}}} \underbrace{\{f_{\text{esc}}^{\text{II}}, \lambda_0\}}_{\Theta_{\text{Re}}} \underbrace{\{f_{\text{Xh}}^{\text{II}}, f_{*,\text{esc}}^{\text{III}}, f_{*,\alpha}^{\text{III}}, f_{*,R}^{\text{III}}, z_{\text{trans}}\}}_{\Theta_{21}}$$

- $f_{*,\text{esc}}^{\text{III}} \equiv f_*^{\text{III}} f_{\text{esc}}^{\text{III}}; \quad f_{*,\alpha}^{\text{III}} \equiv f_*^{\text{III}} f_{\alpha}^{\text{III}}; \quad f_{*,\text{Xh}}^{\text{III}} \equiv f_*^{\text{III}} f_{\text{Xh}}^{\text{III}}; \quad f_{*,R}^{\text{II}} \equiv f_*^{\text{III}} f_R^{\text{III}}$
- Fixed parameters

$$f_{\alpha}^{\text{II}} = f_R^{\text{II}} = 1; \quad f_{\text{Xh}}^{\text{III}} = 0; \quad f_*^{\text{II}} = 0.01$$

- Likelihood: $\mathcal{L} = \mathcal{L}_{\text{P1}} + \mathcal{L}_{\text{Re}} + \mathcal{L}_{21}$

where

$$\mathcal{L}_{21} = \sum_i \left[\frac{\delta T_b^{\text{obs}}(\nu_i) - \delta T_b^{\text{th}}(\nu_i)}{\sigma_i} \right]^2; \quad \mathcal{L}_{\text{Re}} = \frac{1}{2} \sum_{\alpha=1}^{N_{\text{obs}}} \left[\frac{\zeta_{\alpha}^{\text{obs}} - \zeta_{\alpha}^{\text{th}}}{\sigma_{\alpha}} \right]^2.$$

- Constraint

- Reionization must be completed at $z \geq 5.3$

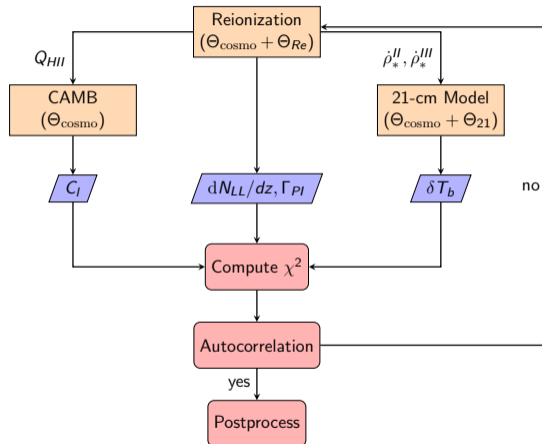
(Kulkarni et al. 2019a, Choudhury et al. 2020)

- Upper limit on the neutral hydrogen fraction coming from the dark fractions in quasar spectra.

(McGreer et al. 2015)

CosmoReionMC

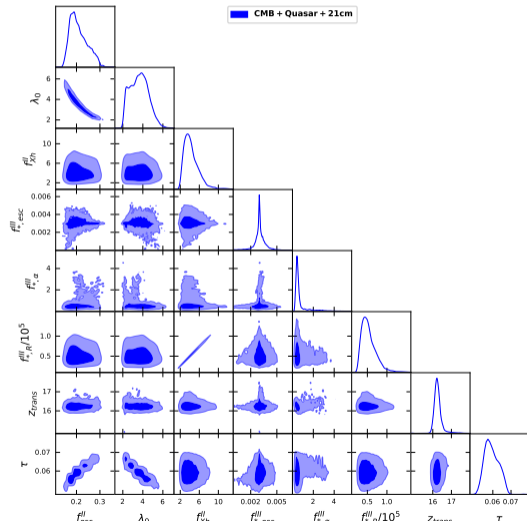
- Written completely in Python.
- Free parameters
 - Cosmological (Θ_{cosmo}) + Reionization (Θ_{Re}) + 21-cm signal (Θ_{21})
- Theoretical model
 - CAMB+ Reionization + 21-cm signal
- Observational data used
 - Planck 2018 + QSO observations + EDGES
- MCMC package used
 - emcee



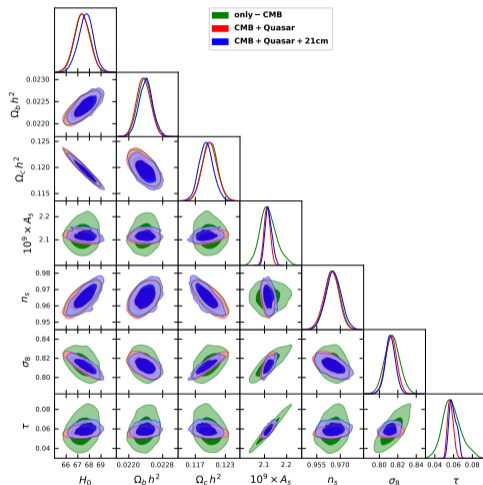
AC, TRC, SM 2021

Constrain on astrophysical parameters

- Strong positive correlation between τ and $f_{\text{esc}}^{\text{II}}$.
- Increasing the value of $f_{\text{esc}}^{\text{II}}$ the reionization will take place earlier so τ will increase.
- Efficiency parameters related to the PopIII stars are required to be non-zero to match the data.
- z_{trans} is well constrained between 15.89 and 16.68.
- Strong correlation between $f_{*,R}^{\text{III}}$ and $f_{\text{Xh}}^{\text{II}}$.



Constraints on cosmological parameters from different analysis



Tighter constraints when reionization observations has been used.

Chatterjee et al. 2021

- Updating CosmoReionMC with more observational data and more accurate reionization model.
- Constraining non-standard cosmology with CosmoReionMC.
- Using Neural Network based machine learning approach to make the MCMC run remarkably efficient.

- Global 21-cm signal is an excellent probe for structure formation in the high redshift Universe.
- EDGES observation of the signal bring renewed interest in this field.
- WDM constraint
 - WDM model with particle mass ≤ 3 keV cannot explain EDGES signal.
- PopIII SFRD
 - We provide for the first time a completely data driven PopIII SFRD using EDGES data.
- CosmoReionMC
 - We introduce an advanced MCMC based parameter estimation package called CosmoReionMC which connects a reionization model (Choudhury & Ferrera 2006) with CAMB and vary both cosmological and model parameters simultaneously to put constraints on them.
 - When reionization observations along with the Planck observation are used, we get tighter constraints on τ and σ_8 .
 - Cosmological parameters do not change significantly when EDGES signal is included.