

# Interpreting the high-redshift 21-cm signal observations

---

**Sambit Giri**

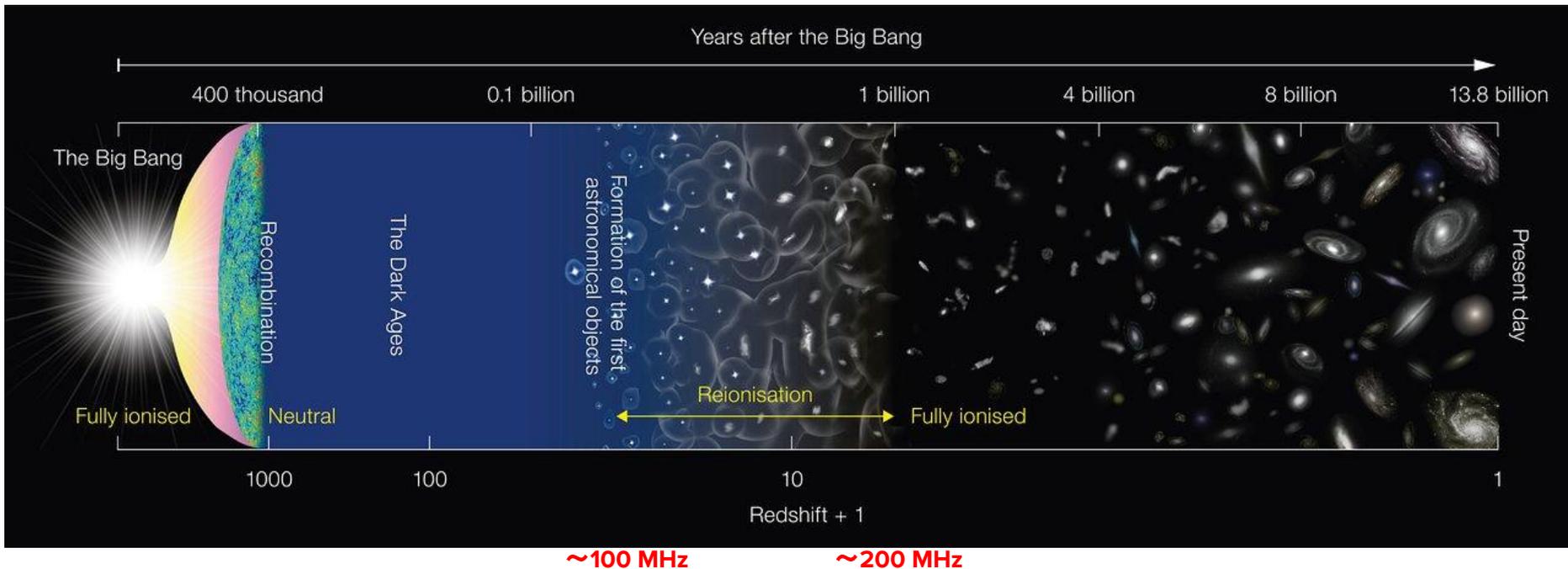


**Universität  
Zürich<sup>UZH</sup>**

# Outline

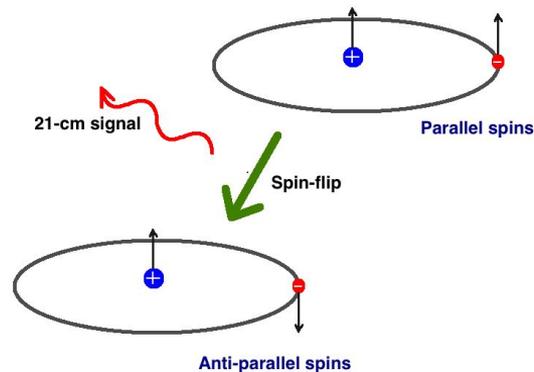
- Basics of 21-cm signal
- Constraints from current 21-cm observations
- Forecast for upcoming observations
  - Halo model based framework

# Timeline of the Universe



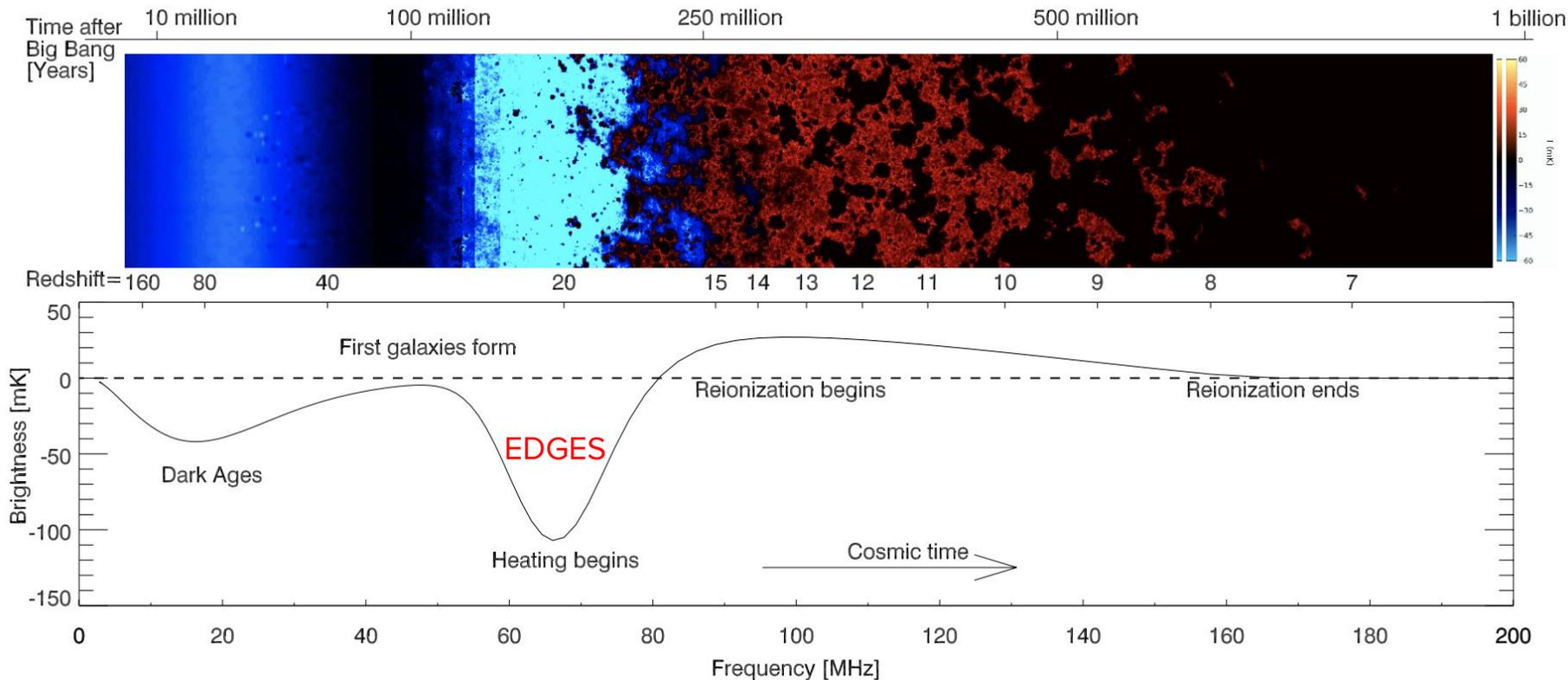
Credit: NAOJ

# Observing the intergalactic medium

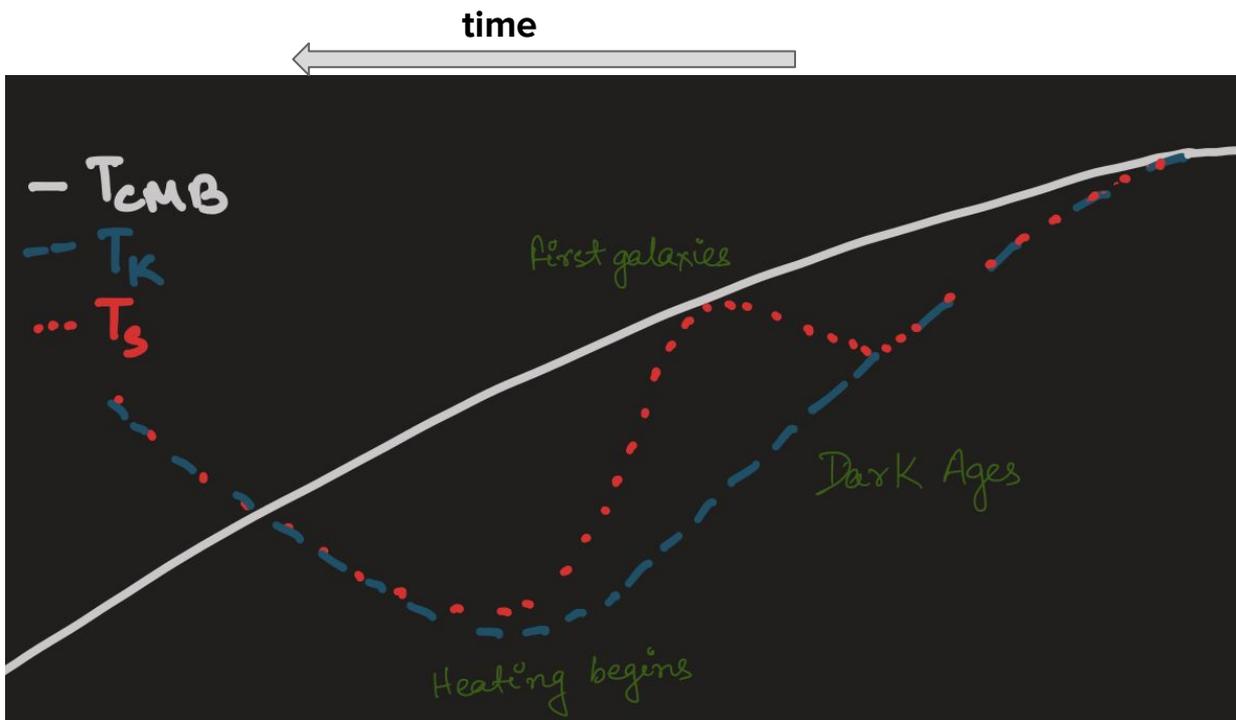


$$\delta T_b \propto x_{\text{HI}}(1 + \delta_b) \left( 1 - \frac{T_{\text{CMB}}}{T_{\text{S}}} \right)$$

# Observing the intergalactic medium

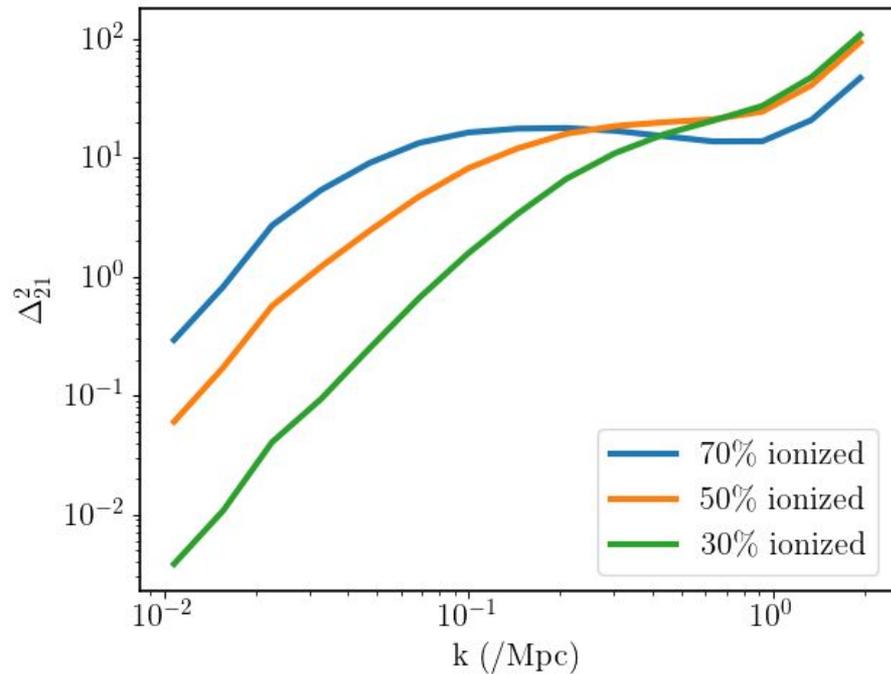


# Spin temperature



$$\delta T_b \propto \left(1 - \frac{T_{\text{CMB}}}{T_{\text{S}}}\right)$$

# Power spectrum

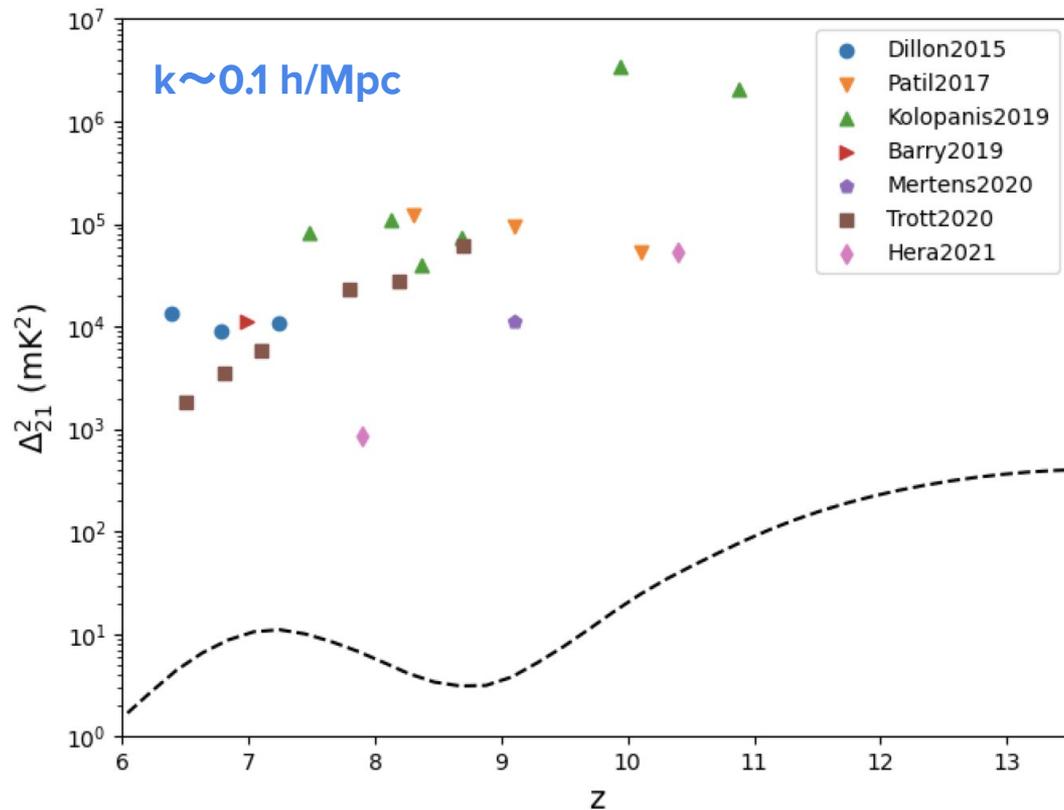


Observed by  
**LOFAR, MWA,**  
**HERA, SKA, etc.**

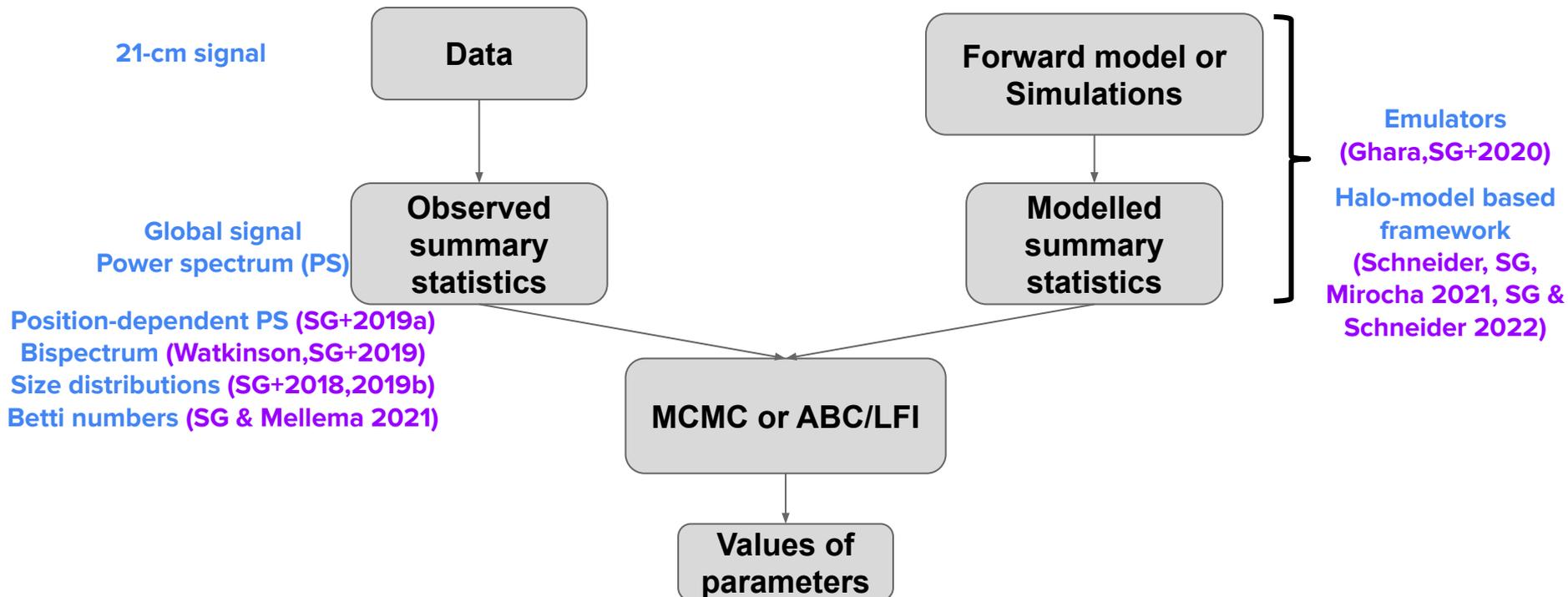
$$\delta T_b \propto x_{\text{HI}}(1 + \delta_b) \left( 1 - \frac{T_{\text{CMB}}}{T_S} \right)$$

$$\Delta_{21}^2(k) = \frac{k^3}{2\pi^2} \langle \delta T_b(k)^* \delta T_b(k) \rangle$$

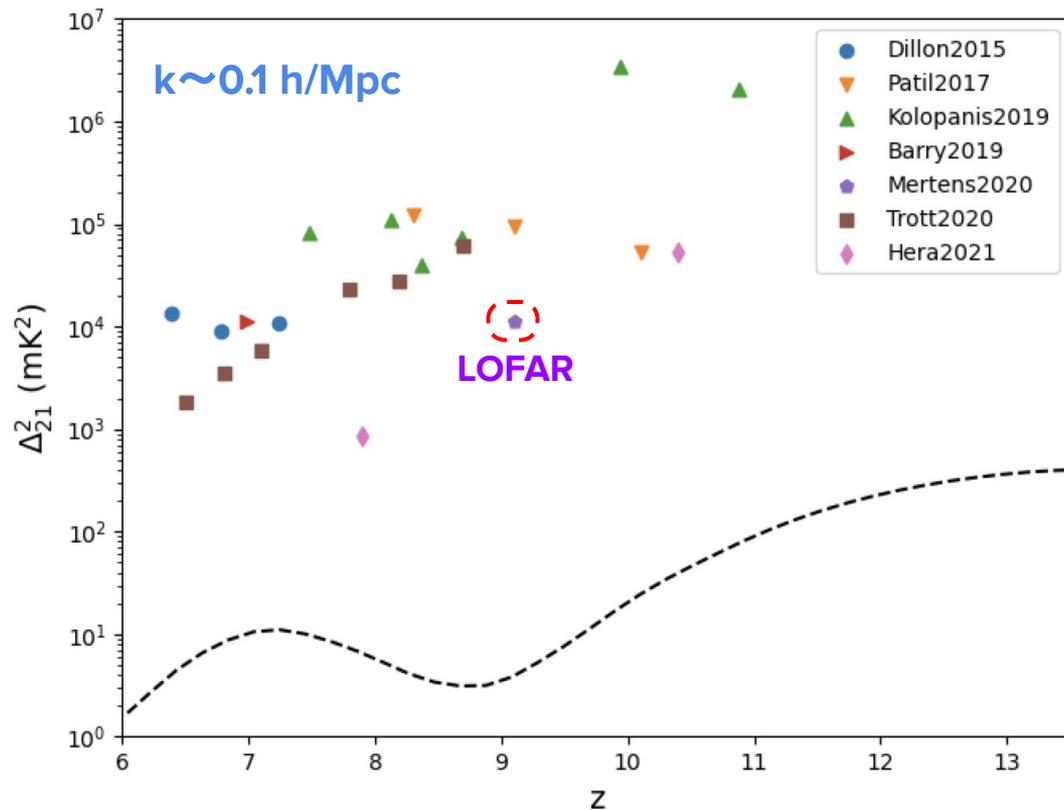
# Current 21-cm signal observations



# Inference from 21-cm observations



# Constraints from upper limits



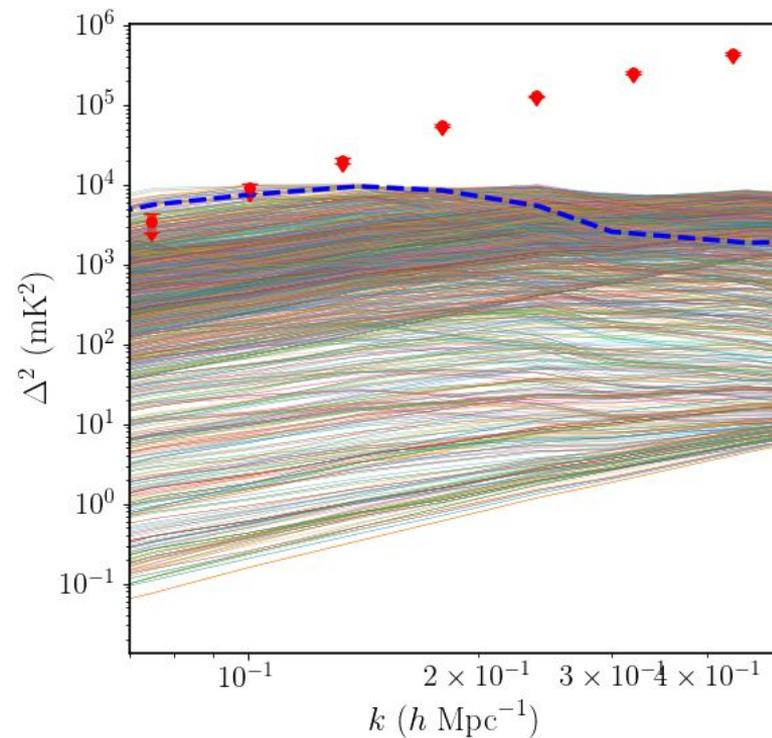
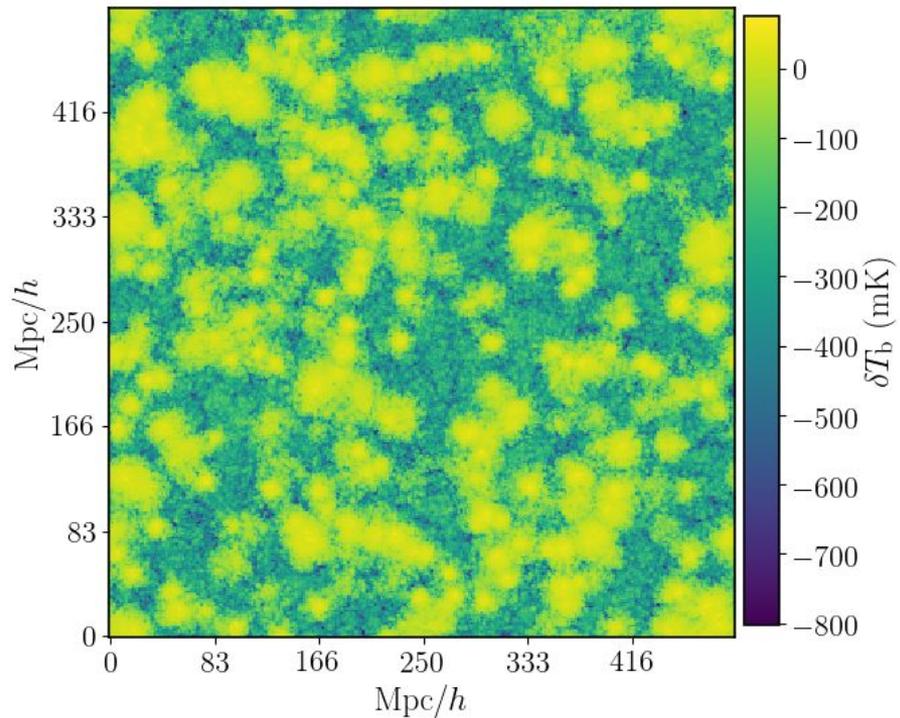
# We need models which boosts large scale fluctuations in 21-cm signal

$$\delta T_b \propto x_{\text{HI}}(1 + \delta) \left( 1 - \frac{T_{\text{CMB}}}{T_S} \right)$$

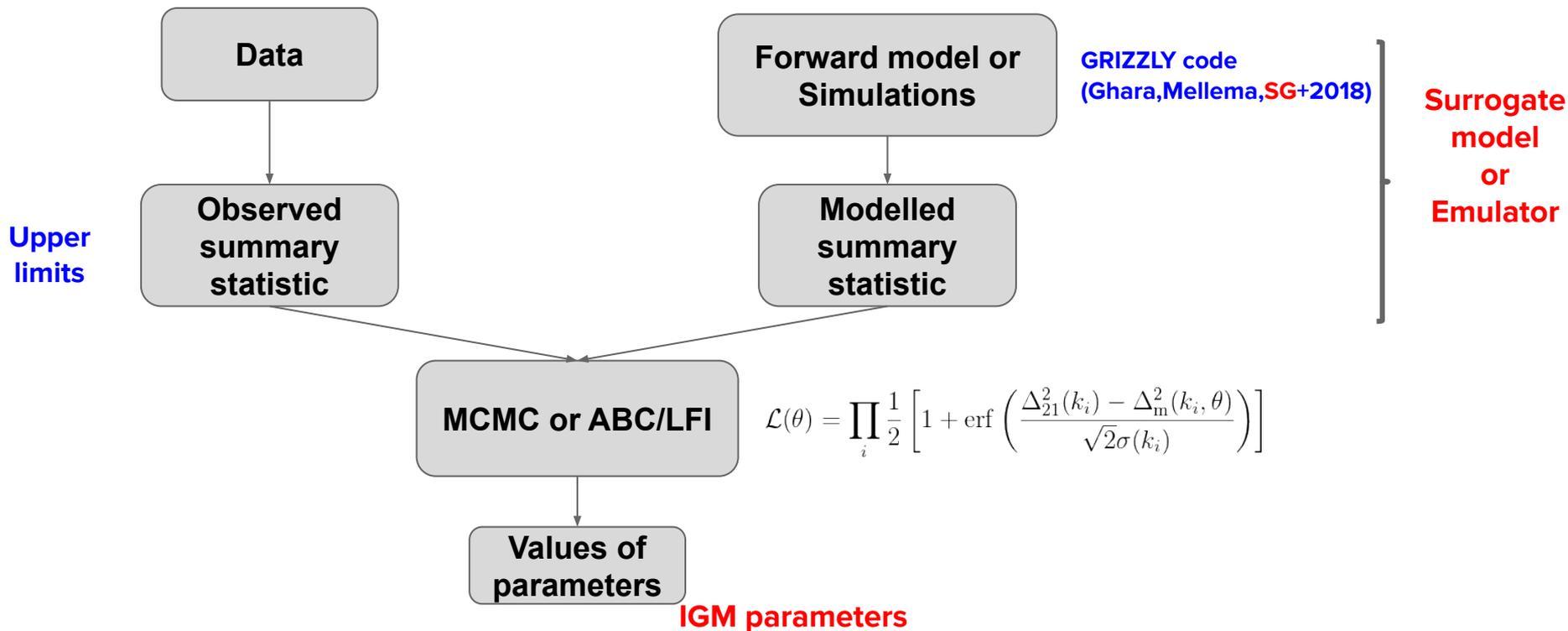
To zeroth order, these fluctuations are scales of  $\sim 50$  Mpc.  
To achieve such fluctuations, we need extreme source models.

# Simulated models at $z=9.1$

(Ghara, SG+2020)



# Inference from 21-cm observations

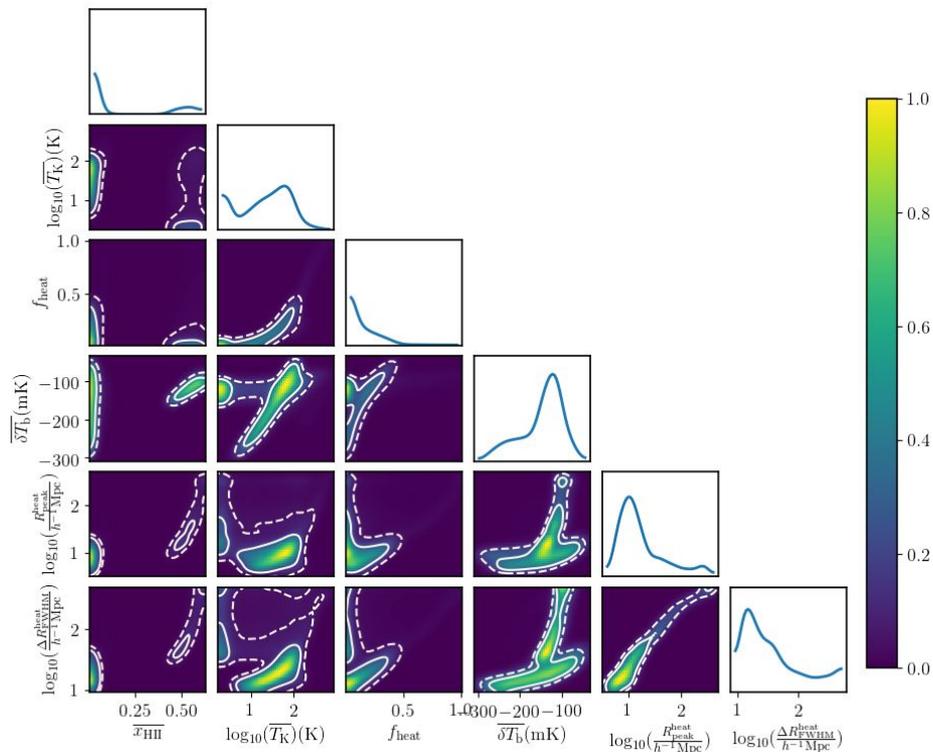


# IGM parameters

Average ionization fraction	$\overline{x_{\text{HIII}}}$
Average gas temperature	$\overline{T_{\text{K}}} \text{ (K)}$
Fraction of heated region	$f_{\text{heat}}$
Global 21-cm signal	$\overline{\delta T_{\text{b}}} \text{ (mK)}$
Peak of the heated bubble PDF	$R_{\text{peak}}^{\text{heat}} \text{ (} h^{-1} \text{Mpc)}$
Width of the heated bubble PDF	$\Delta R_{\text{FWHM}}^{\text{heat}} \text{ (} h^{-1} \text{Mpc)}$

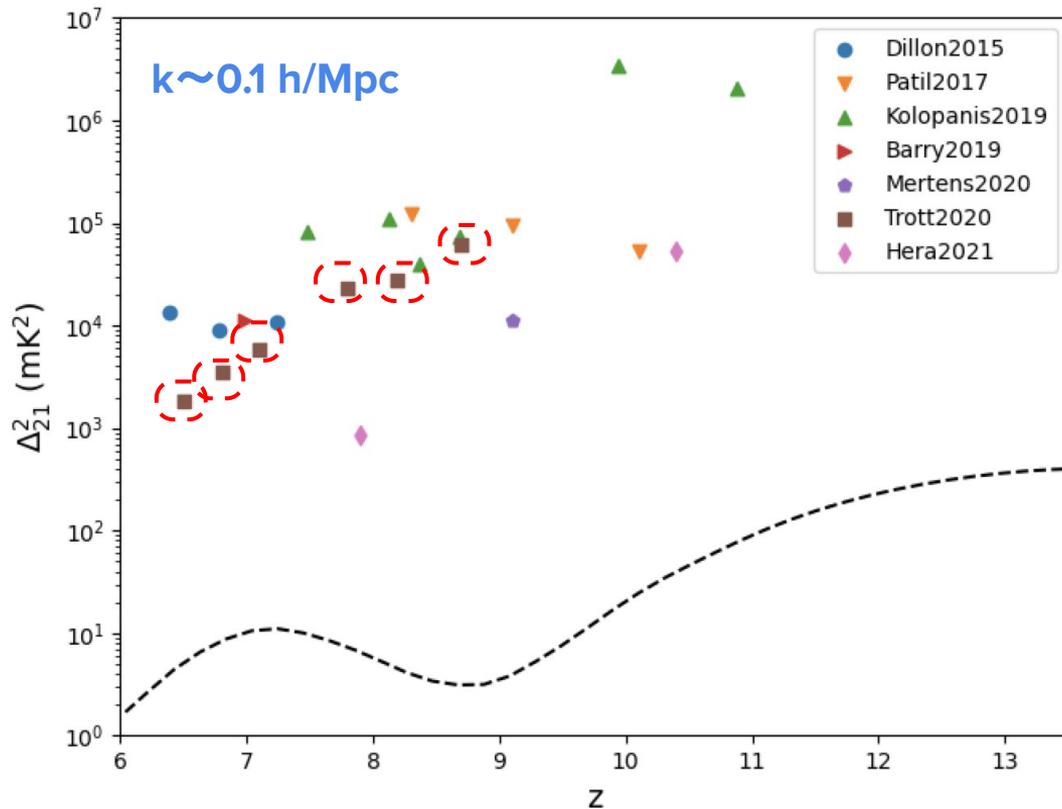
# Excluded IGM parameters

(Ghara, SG+2020)



IGM Parameters of non-uniform $T_S$ scenario	95% credible interval of the excluded models
$\overline{x}_{\text{HII}}$	[0, 0.08], [0.45, 0.62]
$\overline{T_K}$ (K)	[7.41, 158.48], [2.10, 3.55]
$f_{\text{heat}}$	[0, 0.34]
$\overline{\delta T_b}$ (mK)	[-234.15, -65.53]
$R_{\text{peak}}^{\text{heat}}$ ( $h^{-1}\text{Mpc}$ )	[3.50, 69.82]
$\Delta R_{\text{FWHM}}^{\text{heat}}$ ( $h^{-1}\text{Mpc}$ )	[0, 113.76]

# MWA 21-cm signal observations

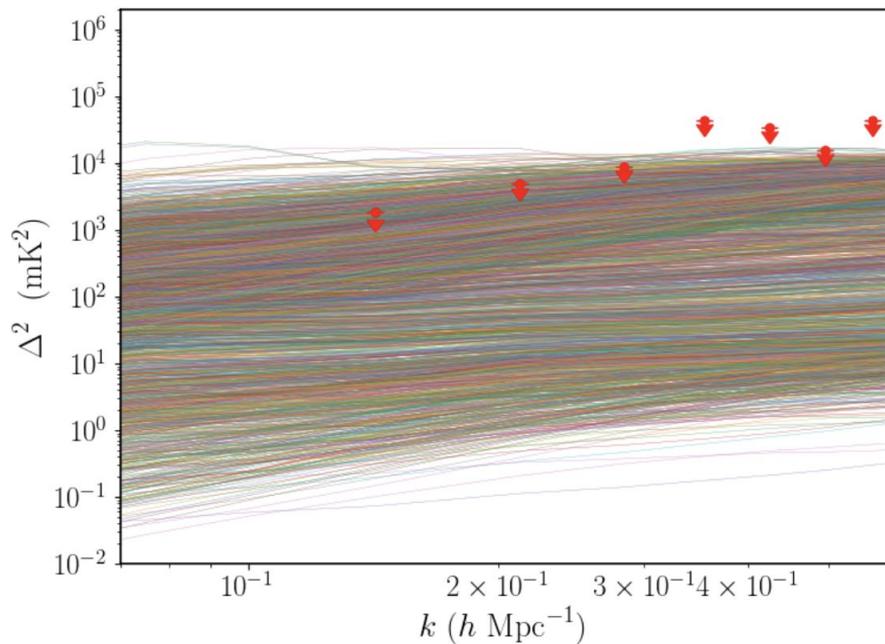


$$\delta T_b \propto x_{\text{HI}}(1 + \delta_b) \left( 1 - \frac{T_{\text{CMB}}}{T_S} \right)$$
$$\Delta_{21}^2(k) = \frac{k^3}{2\pi^2} \langle \delta T_b(k)^* \delta T_b(k) \rangle$$

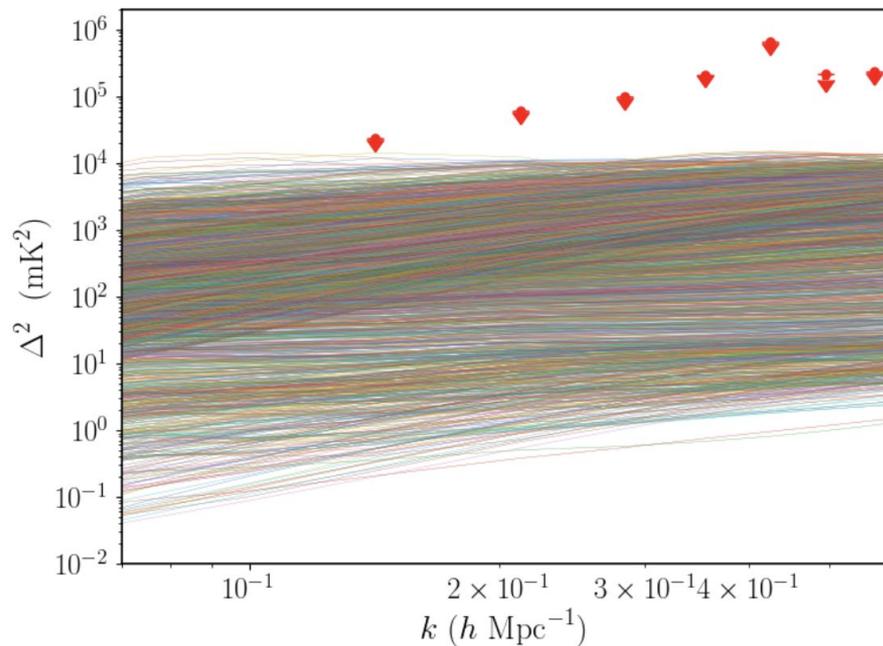
# Simulated models

(Ghara, SG+2021)

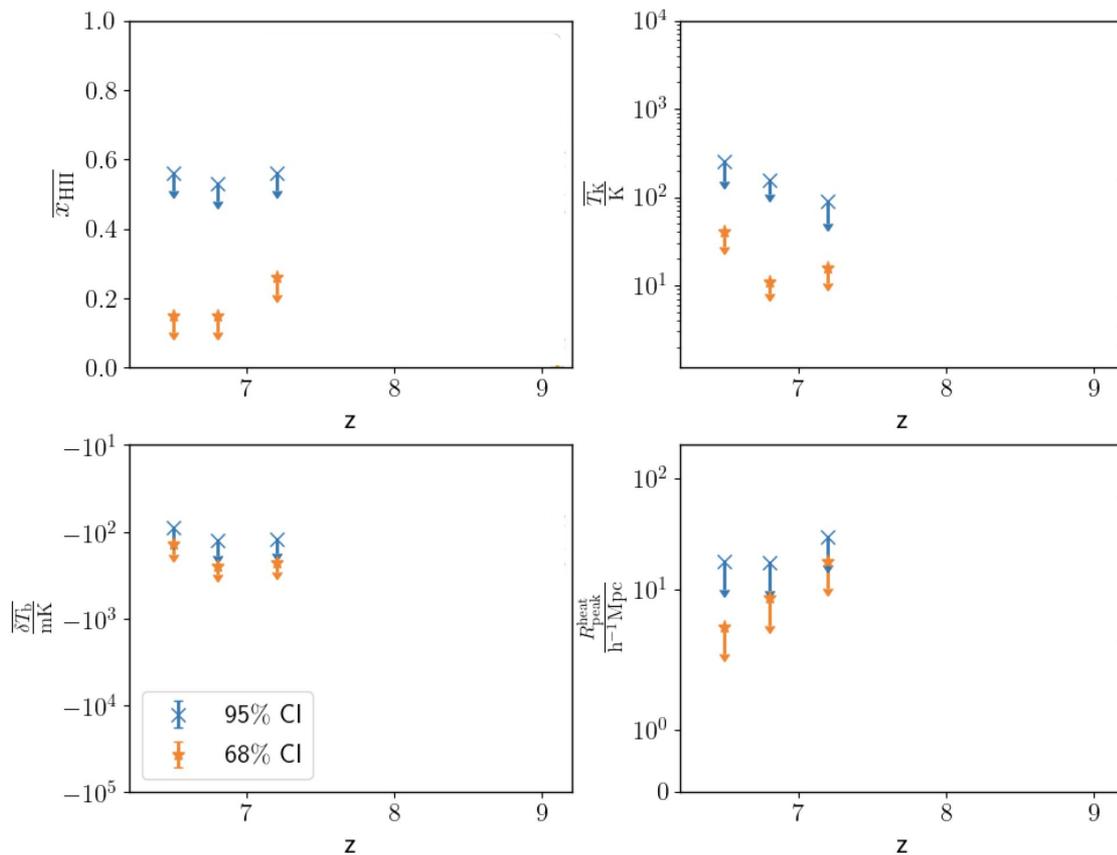
$z = 6.5$



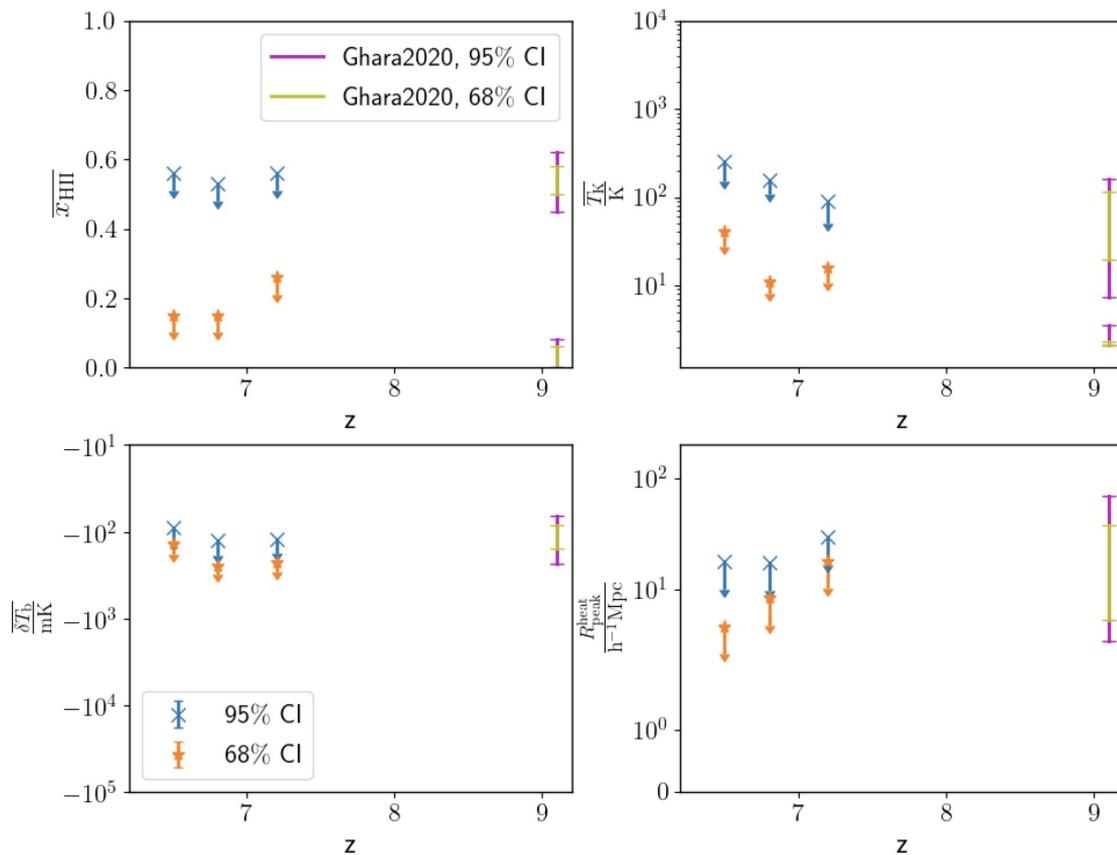
$z = 7.8$



# Excluded IGM parameters

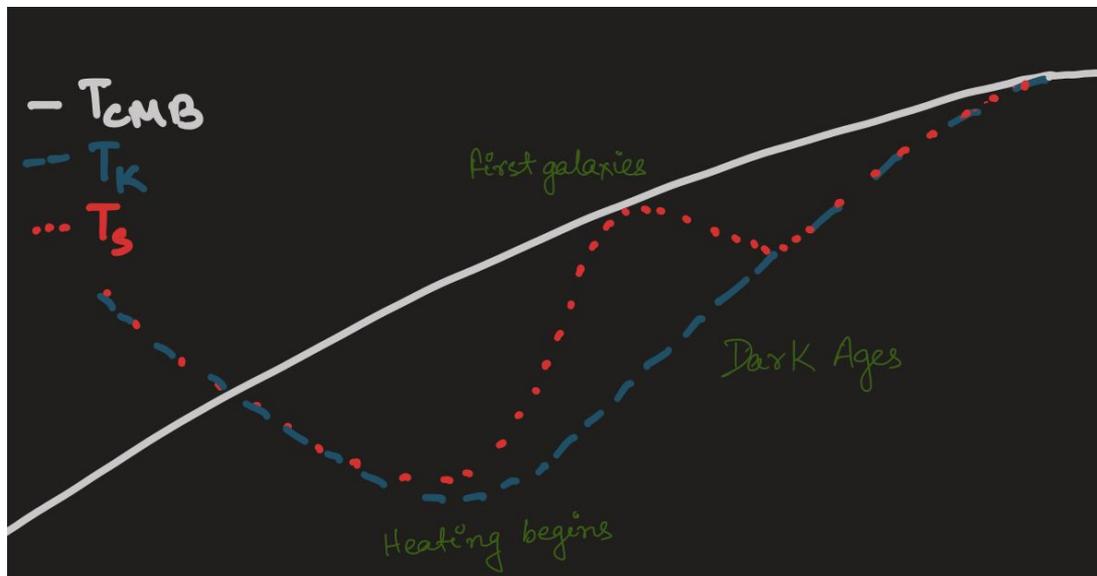


# Excluded IGM parameters



# Excess radio background

$$\delta T_b \propto x_{\text{HI}}(1 + \delta) \left( 1 - \frac{T_{\text{rad}}}{T_S} \right)$$



# Excess radio background

$$\delta T_b \propto x_{\text{HI}}(1 + \delta) \left( 1 - \frac{T_{\text{rad}}}{T_S} \right)$$

Model parameter

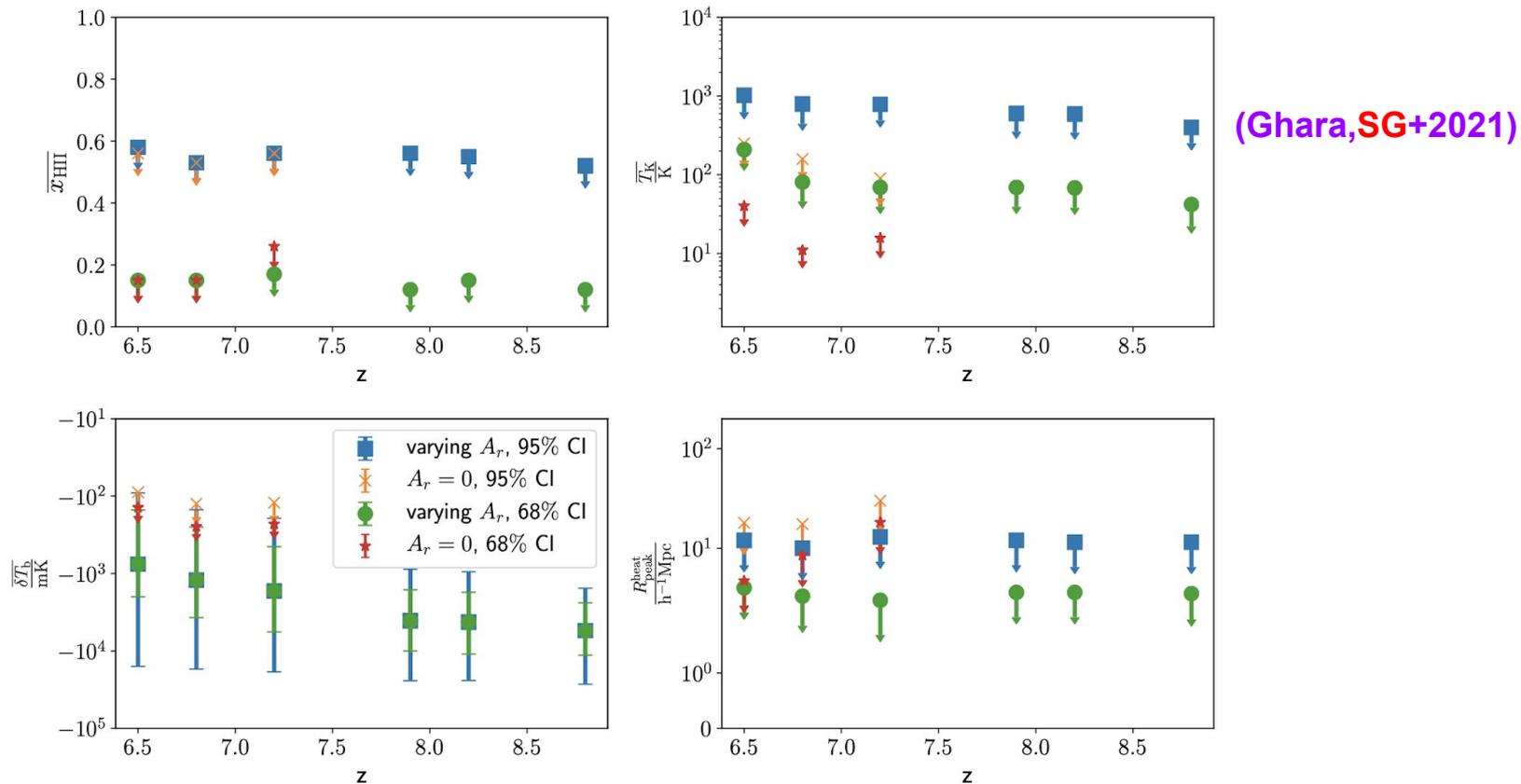
$$T_{\text{rad}} = T_{\text{CMB}}(1 + z) \left[ 1 + A_r \left( \frac{\nu_{\text{obs}}}{78 \text{ MHz}} \right)^\beta \right]$$

Fialkov & Barkana 2019  
Mondal...SG+2020

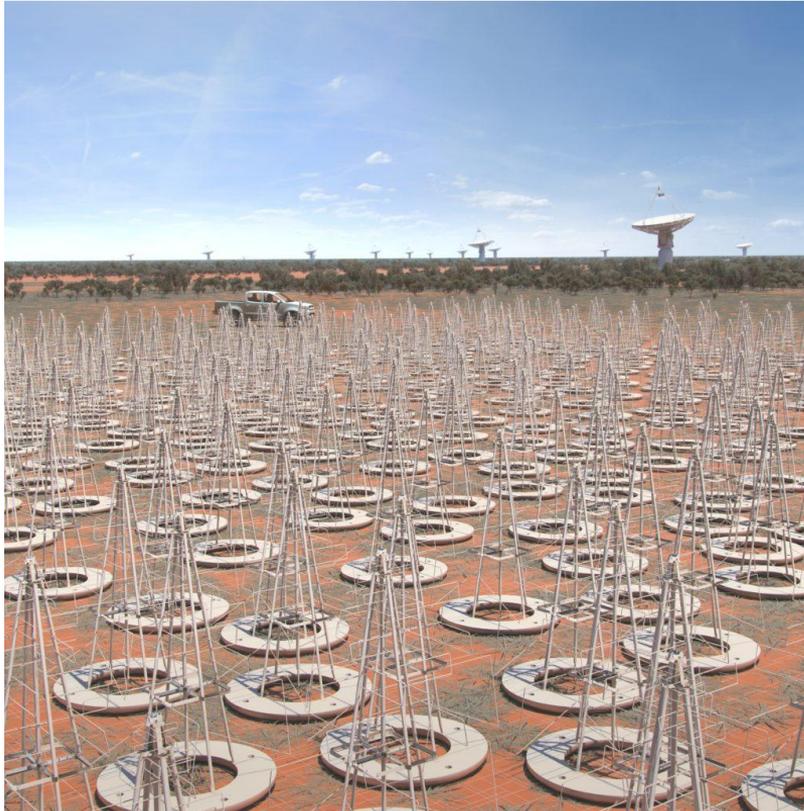
$$\beta = -2.6$$

Value that agrees with LWA1  
observations (Dowell & Taylor 2018)

# Excluded IGM states with excess radio background

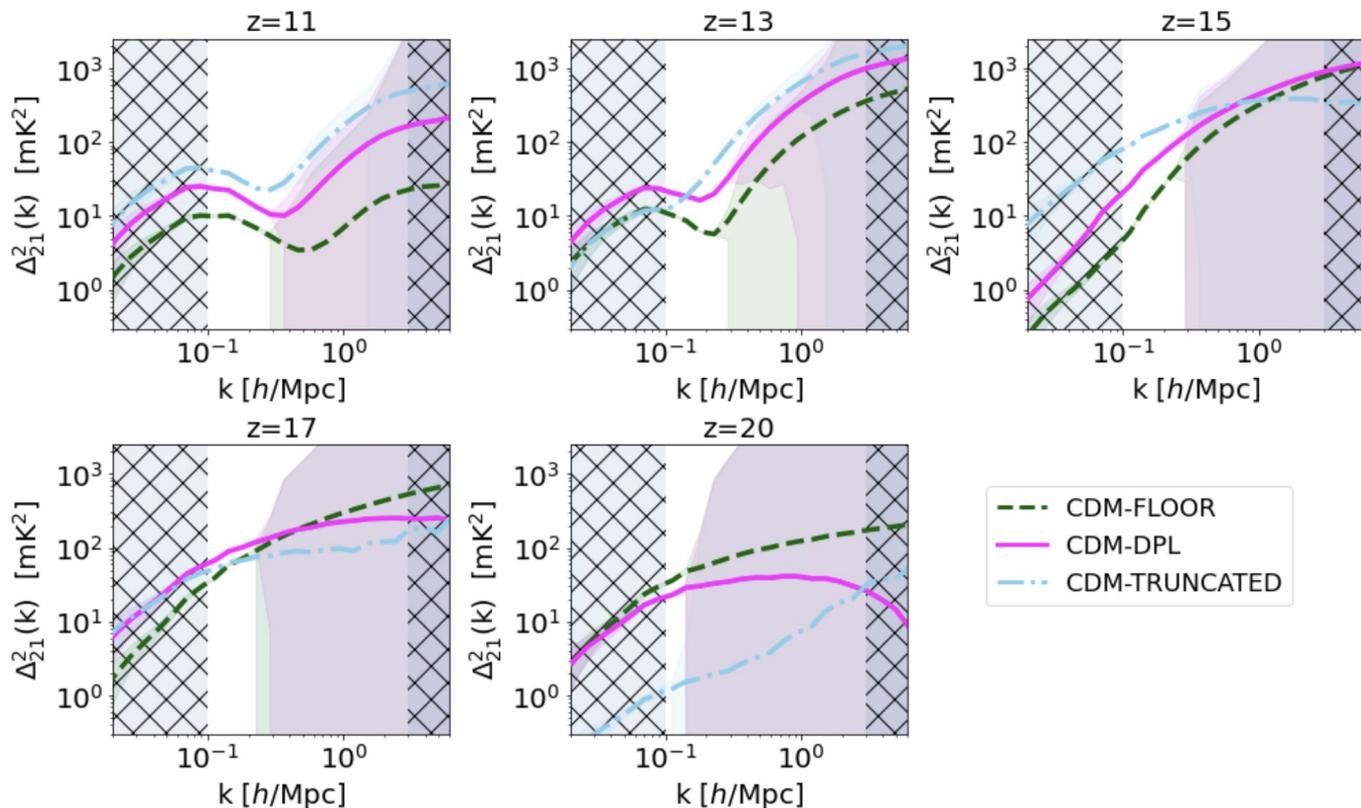


# Forecast study: SKA-Low



*Credit: SKA Organisation*

# Mock observation at cosmic dawn

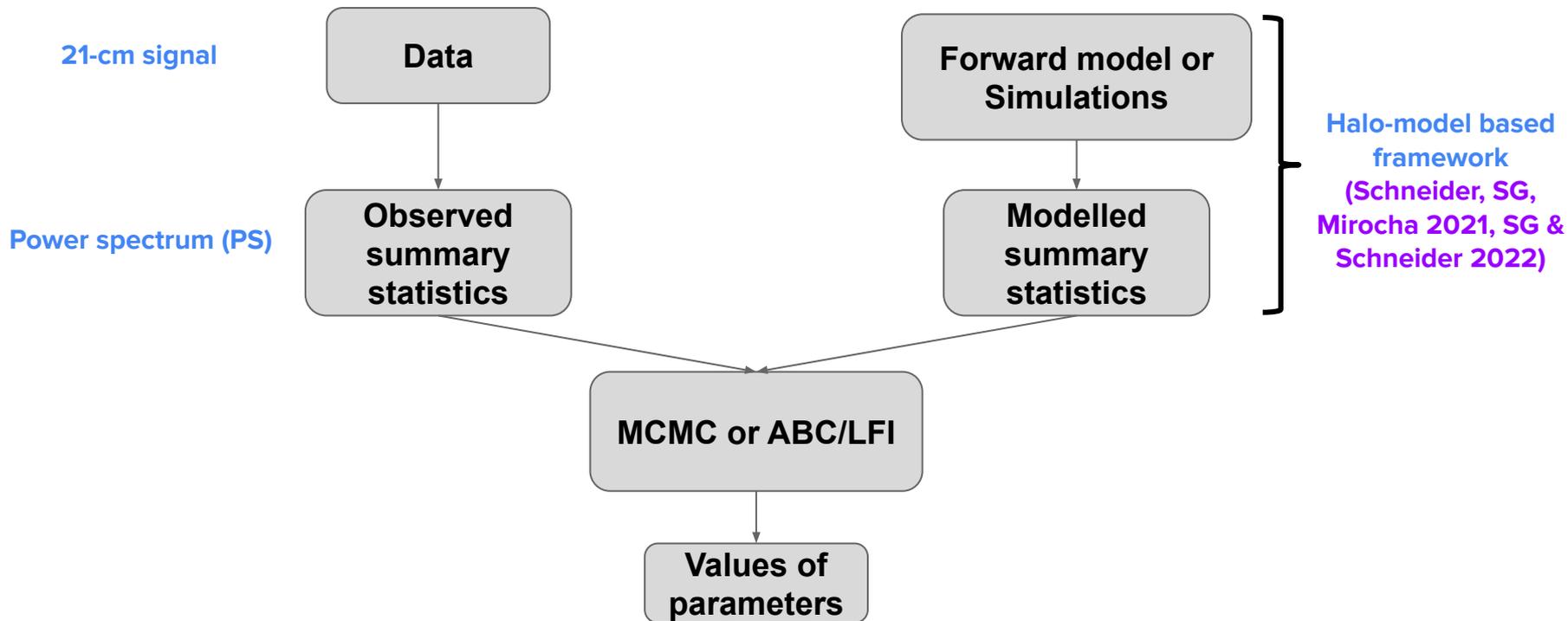


1000 hour observations

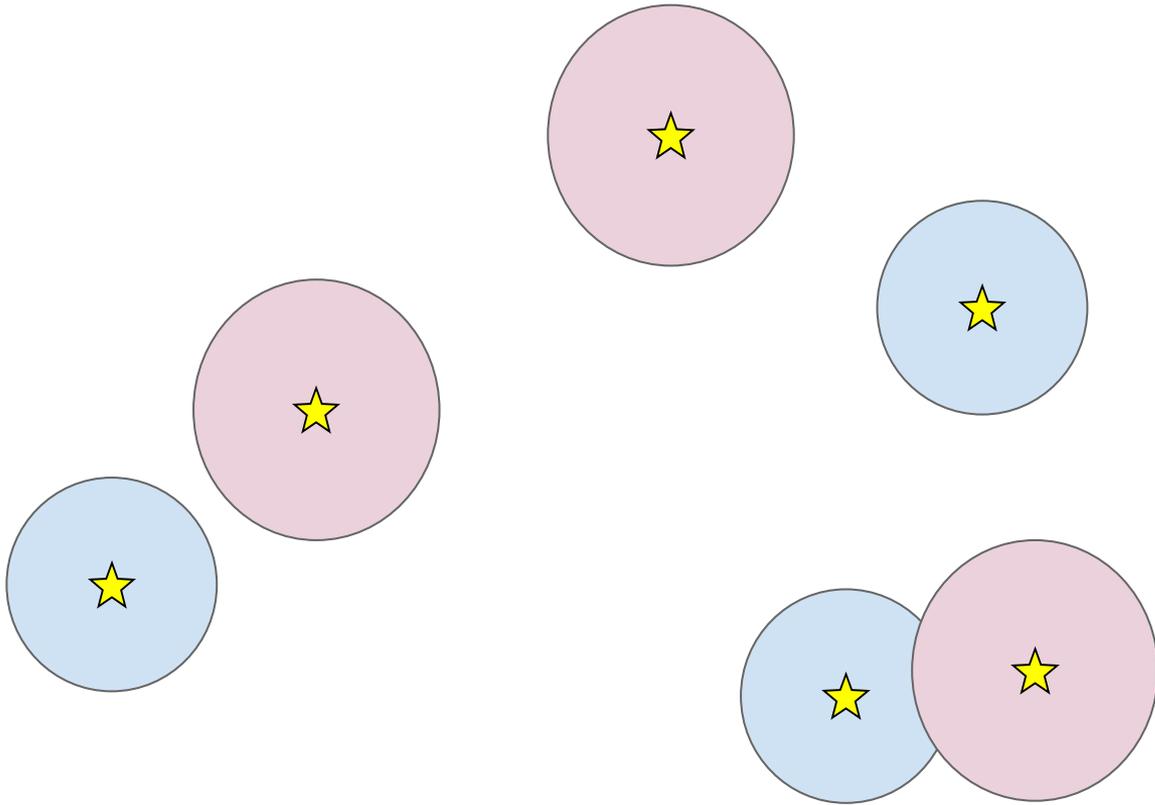
Instrumental effects are calculated using Tools21cm (SG+2020)

(SG & Schneider 2022)

# Inference from 21-cm observations



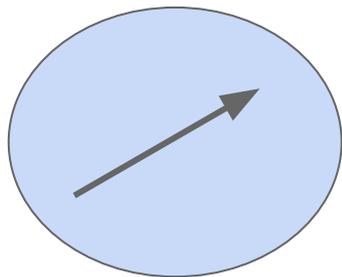
# Halo model



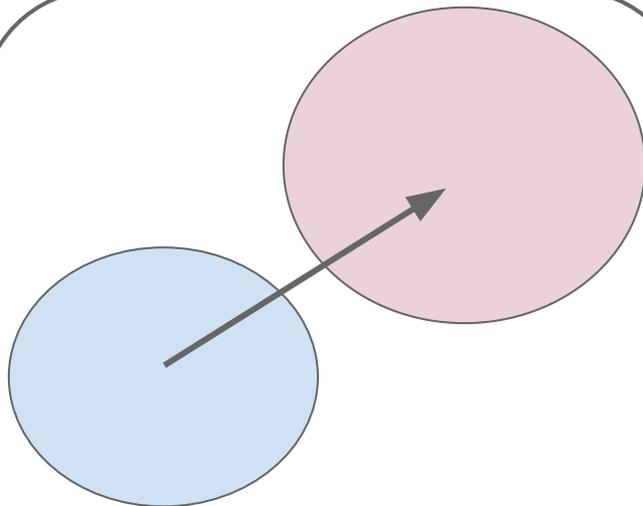
Seljak (2000)  
Cooray & Sheth (2002)

...

# Halo model: power spectrum



$$P_{XY}^{1,h}(k, z) = \frac{\beta_X \beta_Y}{(\bar{\rho} f_{\text{coll}})^2} \int dM \frac{dn}{dM} \tilde{f}_*^2 M^2 |u_X| |u_Y|,$$



$$P_{XY}^{2,h}(k, z) = \frac{\beta_X}{(\bar{\rho} f_{\text{coll}})} \int dM \frac{dn}{dM} \tilde{f}_* M |u_X| b_X \\ \times \frac{\beta_Y}{(\bar{\rho} f_{\text{coll}})} \int dM \frac{dn}{dM} \tilde{f}_* M |u_Y| b_Y \times P_{\text{lin}}$$

$$P_{XY}(k, z) = P_{XY}^{1,h}(k, z) + P_{XY}^{2,h}(k, z),$$

# Ingredients for the halo model

Linear power spectrum

Halo mass function

Mass accretion

Halo bias

Stellar to halo mass relation

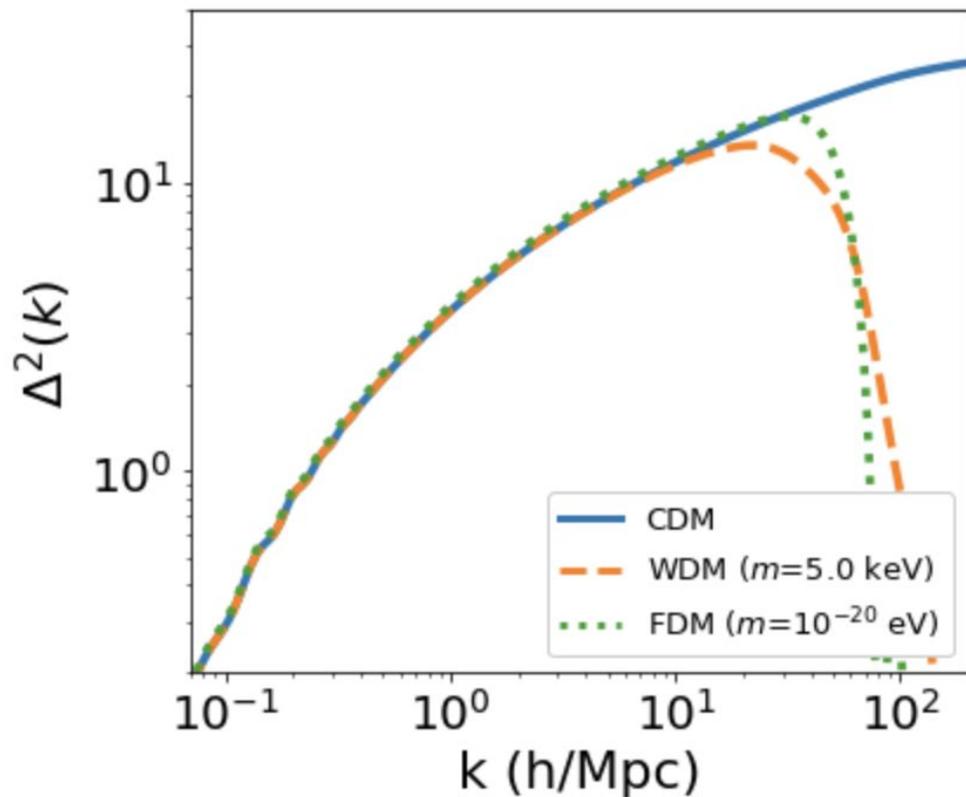
Flux profiles

$$P_{XY}^{1h}(k, z) = \frac{\beta_X \beta_Y}{(\bar{\rho} f_{\text{coll}})^2} \int dM \frac{dn}{dM} \tilde{f}_*^2 M^2 |u_X| |u_Y|,$$

$$P_{XY}^{2h}(k, z) = \frac{\beta_X}{(\bar{\rho} f_{\text{coll}})} \int dM \frac{dn}{dM} \tilde{f}_* M |u_X| b_X \\ \times \frac{\beta_Y}{(\bar{\rho} f_{\text{coll}})} \int dM \frac{dn}{dM} \tilde{f}_* M |u_Y| b_Y \times P_{\text{lin}},$$

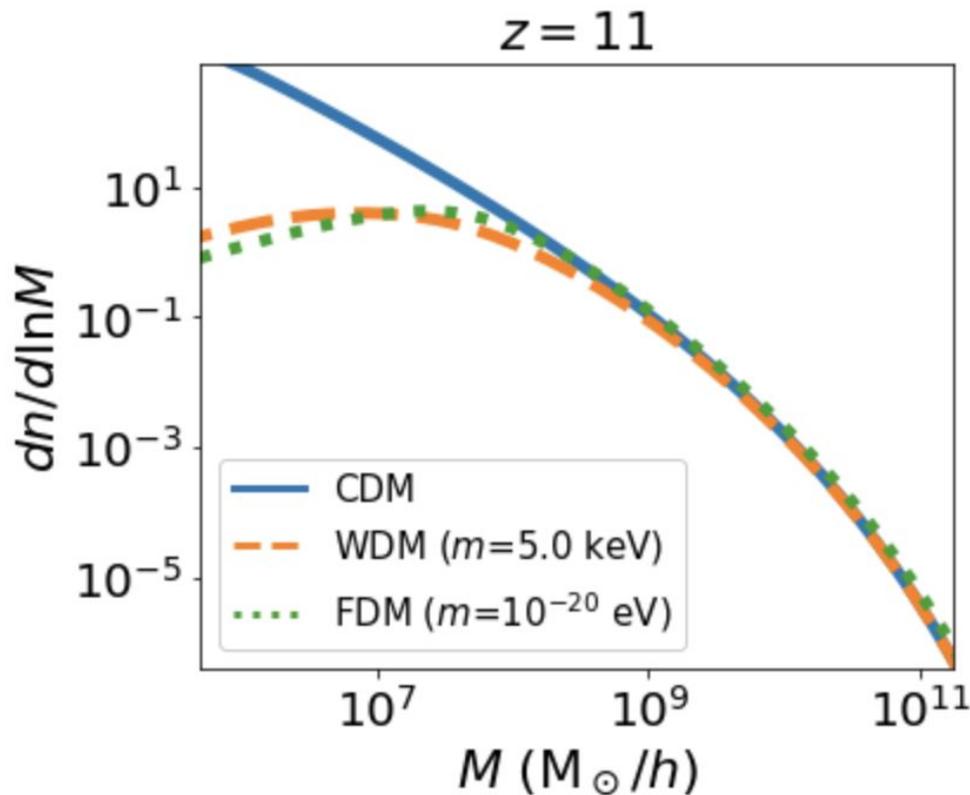
$$P_{XY}(k, z) = P_{XY}^{1h}(k, z) + P_{XY}^{2h}(k, z),$$

# Linear power spectra



$z = 0$

# Halo mass function



$$\frac{dn}{d\ln M} = -\frac{\bar{\rho}}{M} \nu f(\nu) \frac{d\ln \sigma}{d\ln M},$$

$$M = \frac{4\pi}{3} \bar{\rho} (cR)^3$$

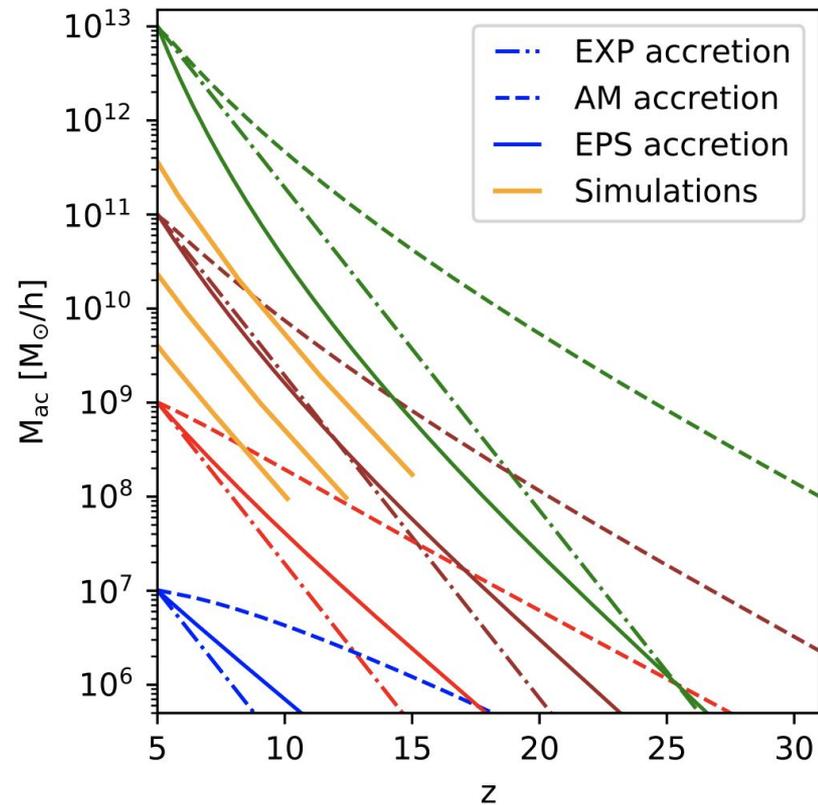
$$f(\nu) = A \sqrt{\frac{2q\nu}{\pi}} (1 + \nu^{-p}) e^{-q\nu/2}$$

$$\sigma^2(R, z) = \int \frac{dk^3}{(2\pi)^3} P_{\text{lin}}(k) \mathcal{W}(k|R)$$

$$\mathcal{W}(k|R) = \frac{1}{1 + (kR)^\beta}.$$

Leo+2018

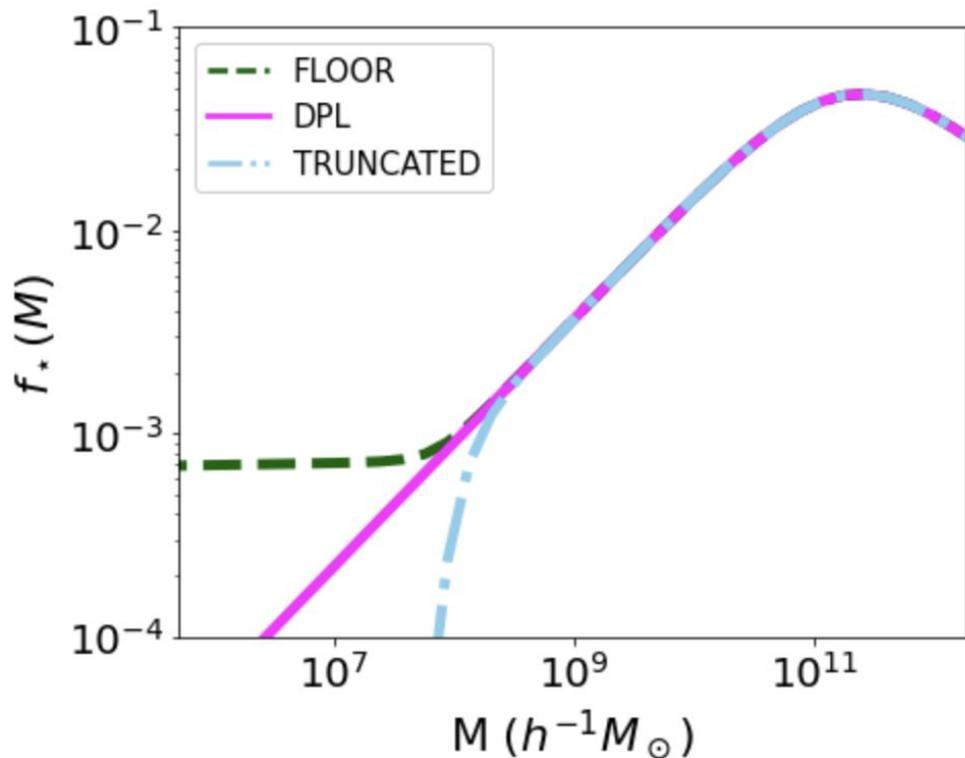
# Mass accretion rate



# Halo bias

$$b(M) = 1 + \frac{q\nu - 1}{\delta_c(z)} + \frac{2p}{\delta_c(z)[1 + (q\nu)^p]}.$$

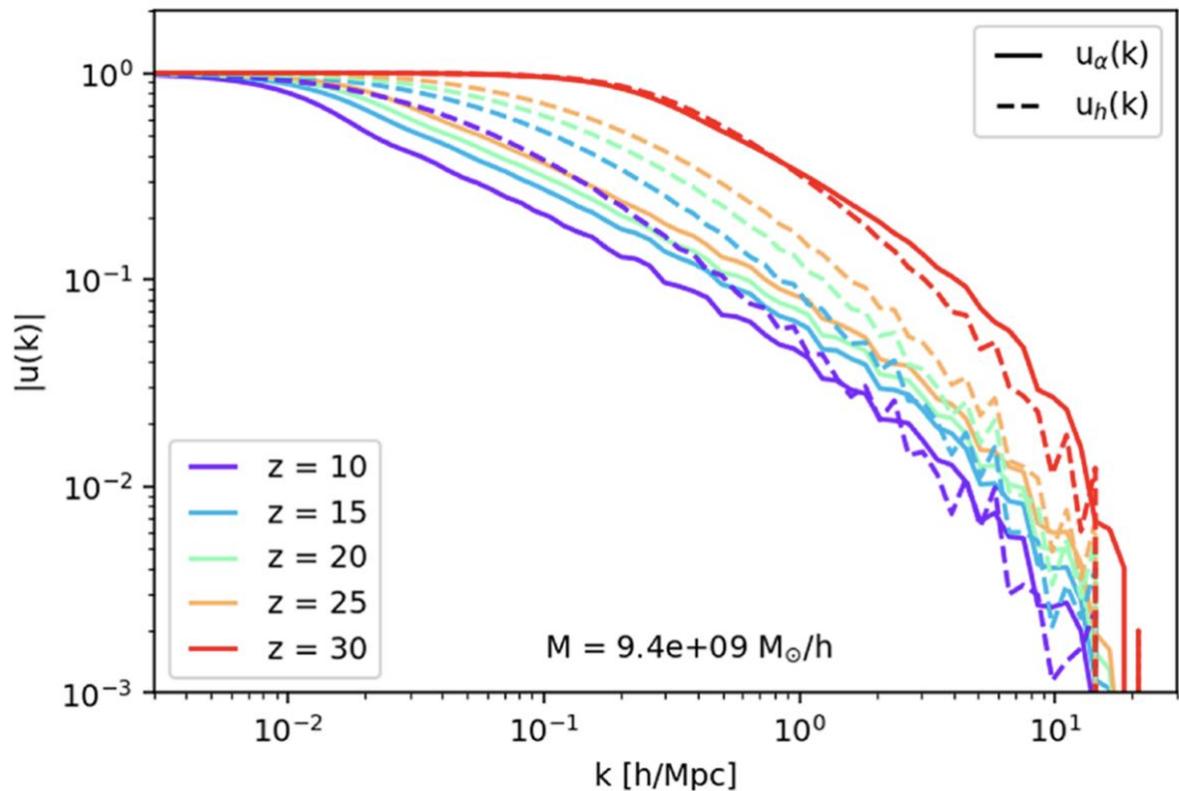
# Stellar to halo mass relation



$$f_*(M) = \frac{2(\Omega_b/\Omega_m)f_{*,0}}{(M/M_p)^{\gamma_1} + (M/M_p)^{\gamma_2}} \times S(M)$$

$$S(M) = [1 + (M_t/M)^{\gamma_3}]^{\gamma_4},$$

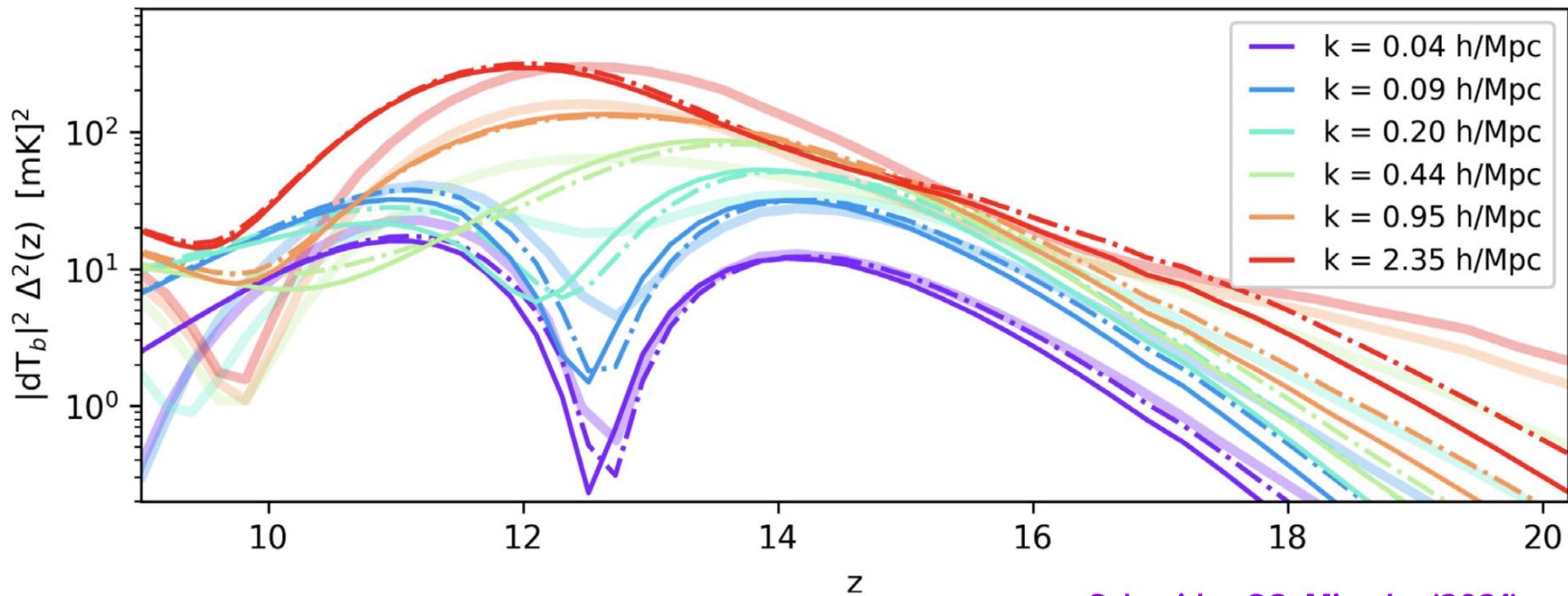
# Flux profiles



$$P_{21} = P_{\alpha\alpha} + P_{hh} + P_{pp} + P_{bb} \\ + 2(P_{\alpha h} + P_{\alpha p} + P_{\alpha b} + P_{hp} + P_{hb} + P_{pb}) \\ + \frac{2}{3}(P_{\alpha m} + P_{hm} + P_{pm} + P_{bm}) + \frac{1}{5}P_{mm}.$$

Schneider, SG, Mirocha (2021)

# Validity of the approach



Schneider, SG, Mirocha (2021)

# Constraining mixed dark matter models

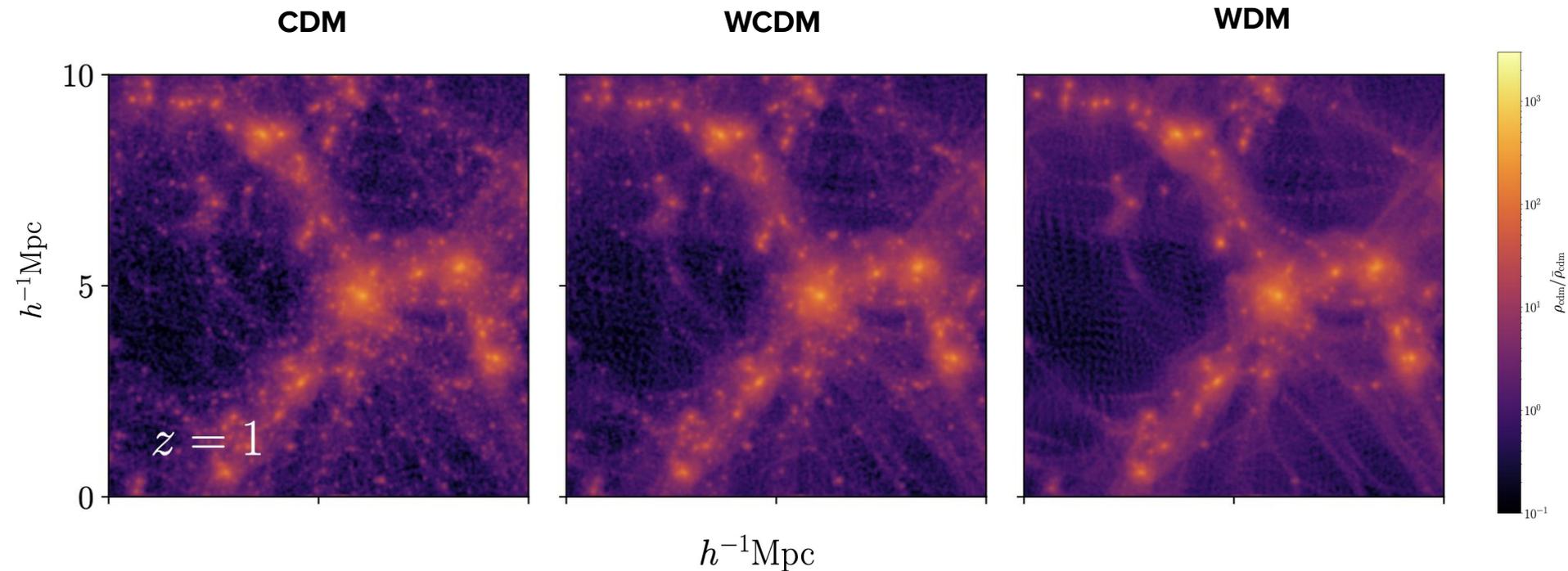
Contains a mixtures of two components

- Cold DM
- Non-cold: WDM/FDM

$$f_{n\text{DM}} = \frac{\Omega_{n\text{DM}}}{\Omega_{\text{DM}} + \Omega_{n\text{DM}}}$$

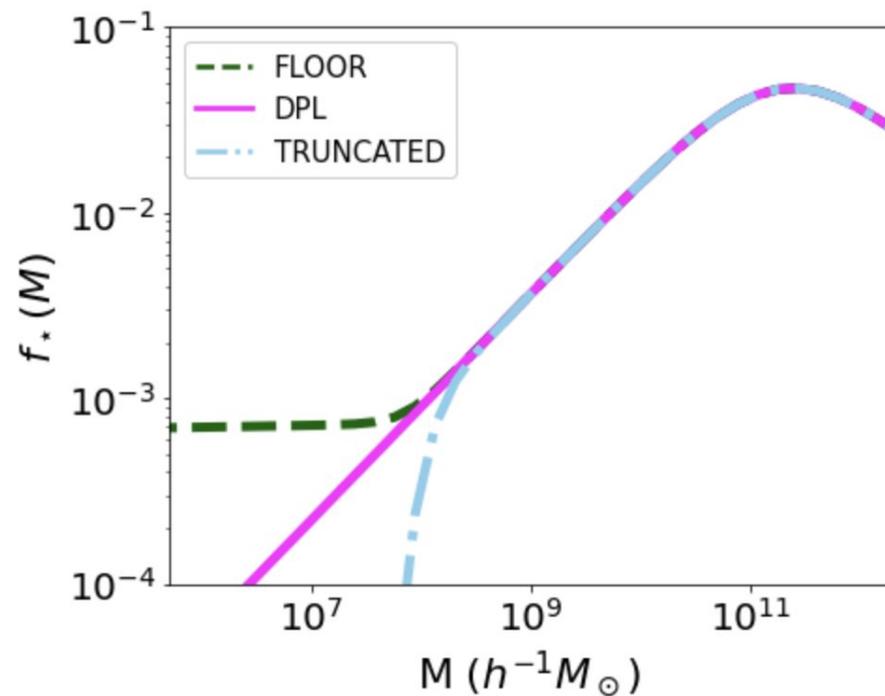
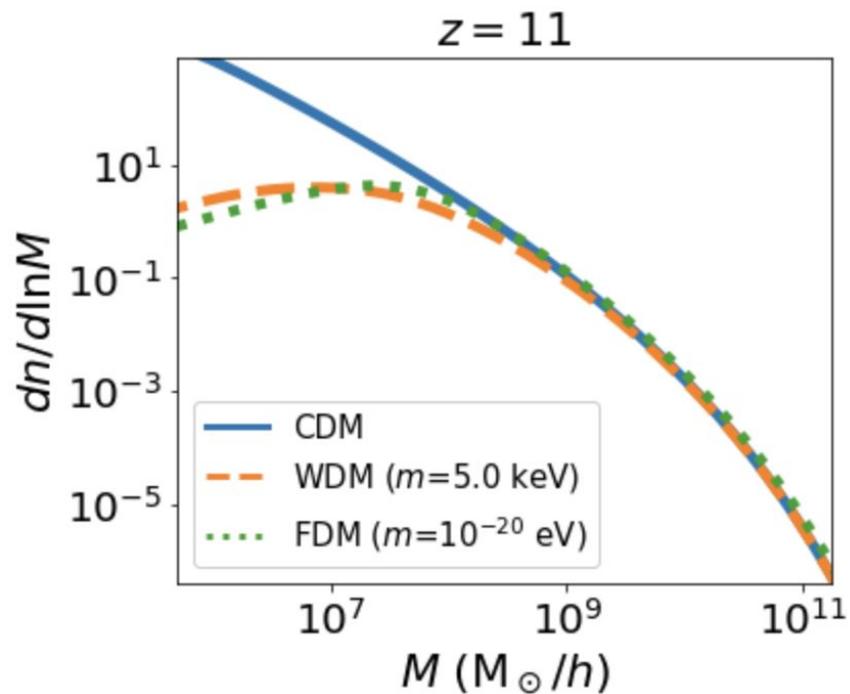
$$f_{n\text{DM}} < 20\% \text{ (Boyarsky+2009)}$$

# Snapshots of Cold-Warm dark matter



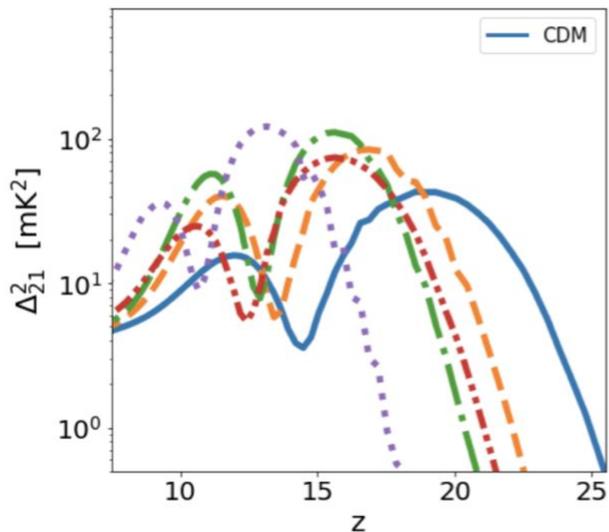
Paribelli, Scelfo, SG+2021

# Constraining dark matter models

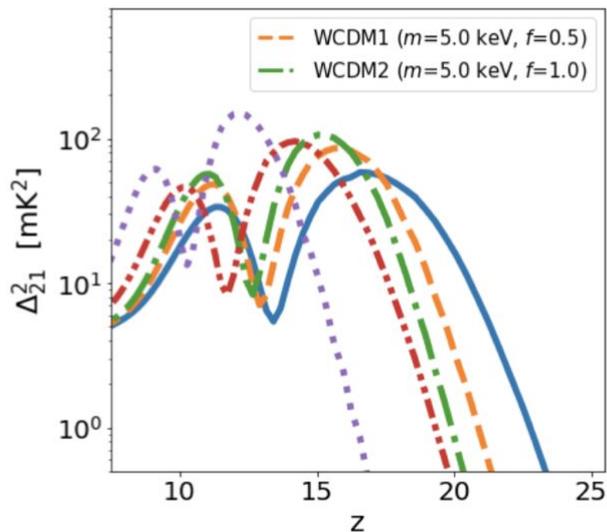


# Evolution of power spectra

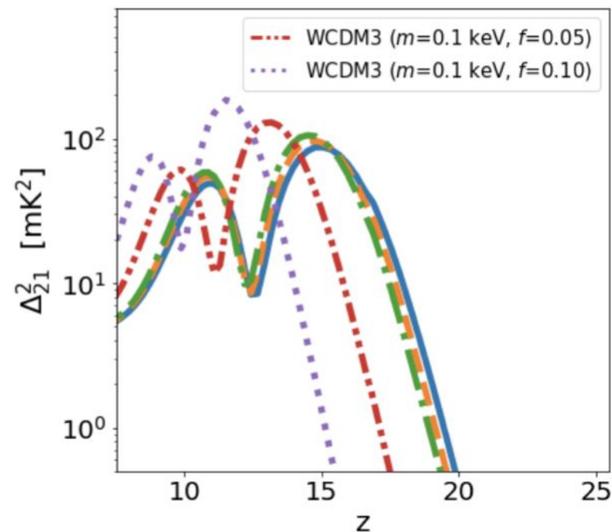
FLOOR



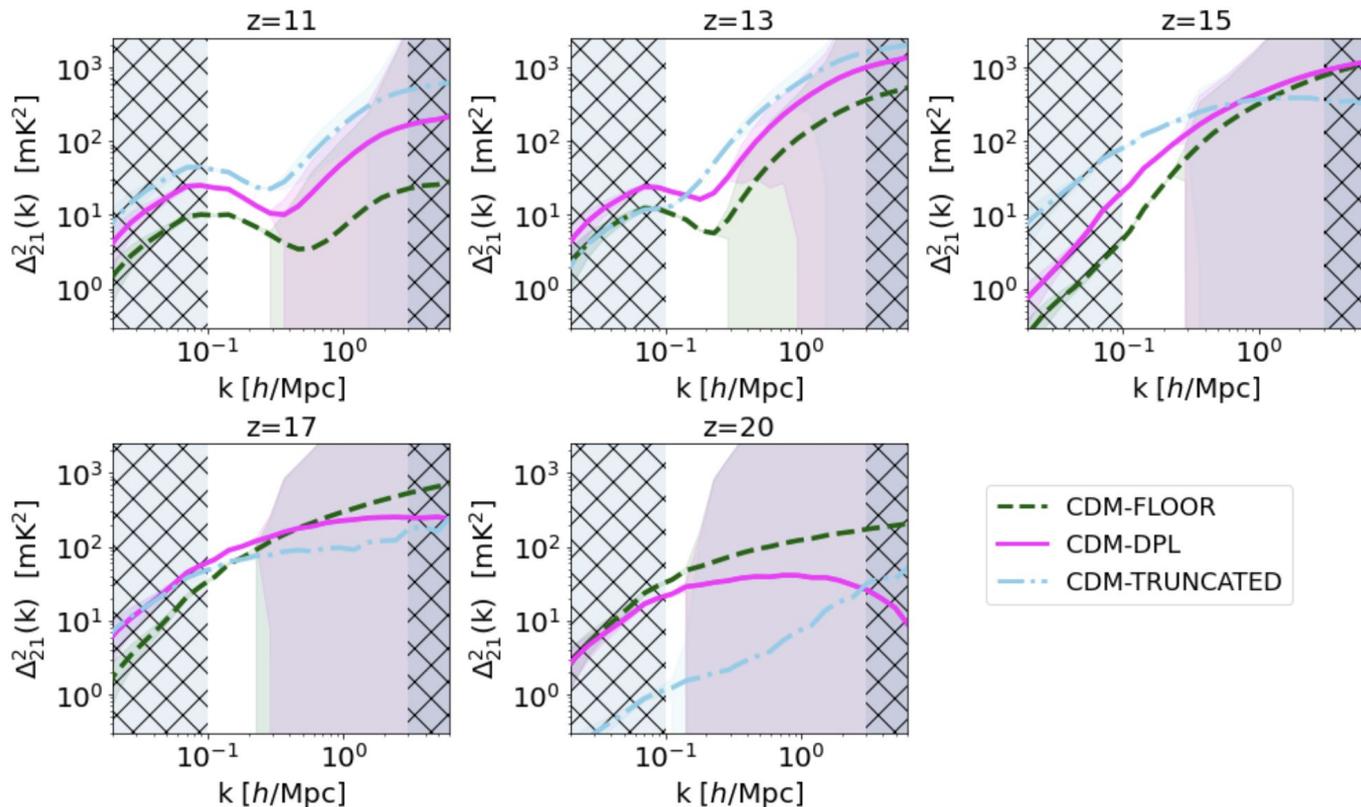
DPL



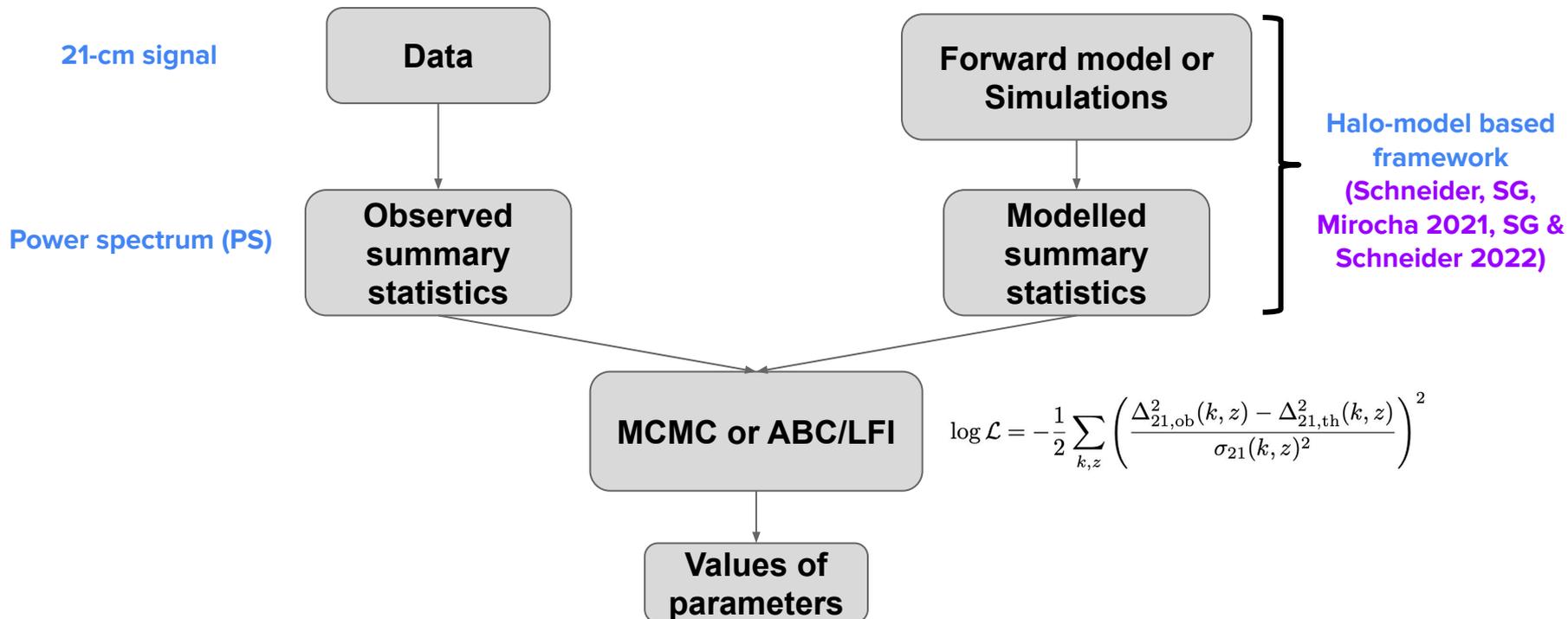
TRUNCATED



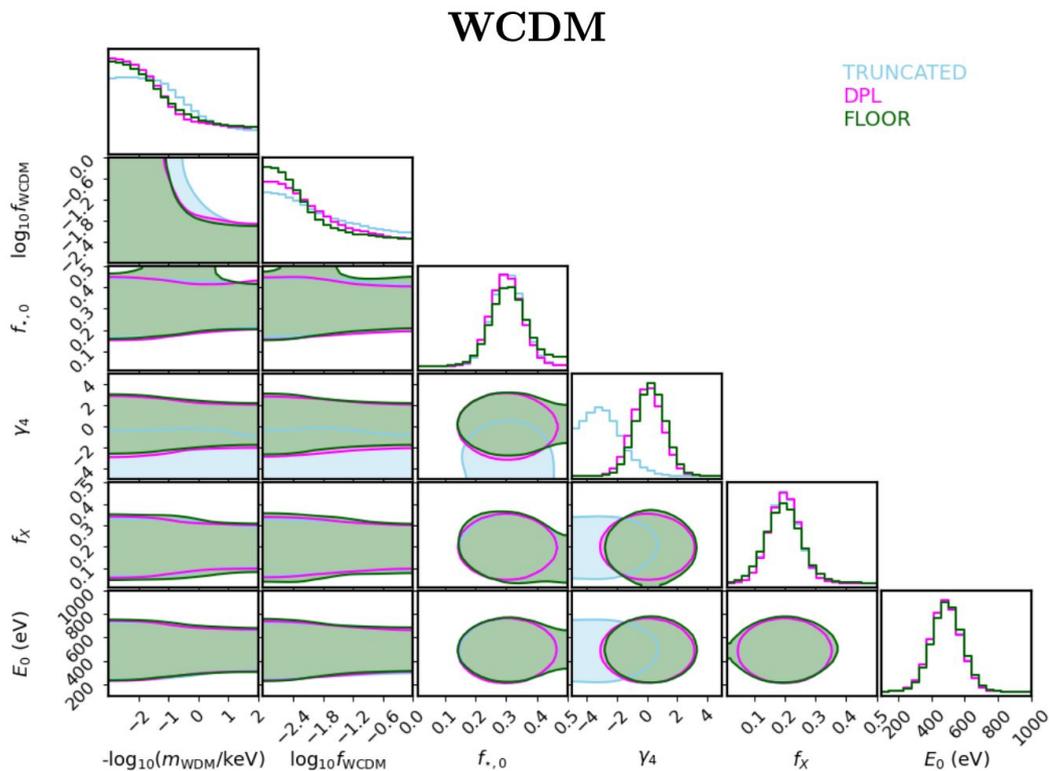
# Mock observation at cosmic dawn



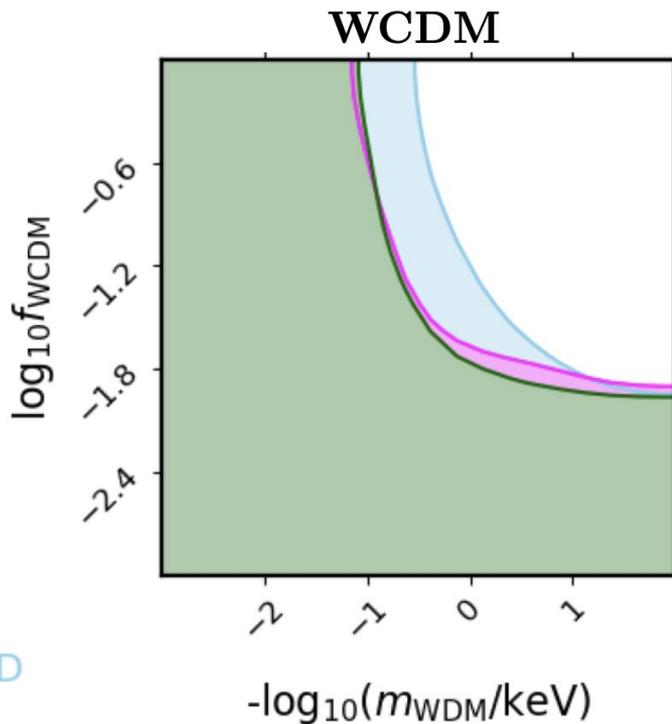
# Inference from 21-cm observations



# Corner showing the posterior distribution



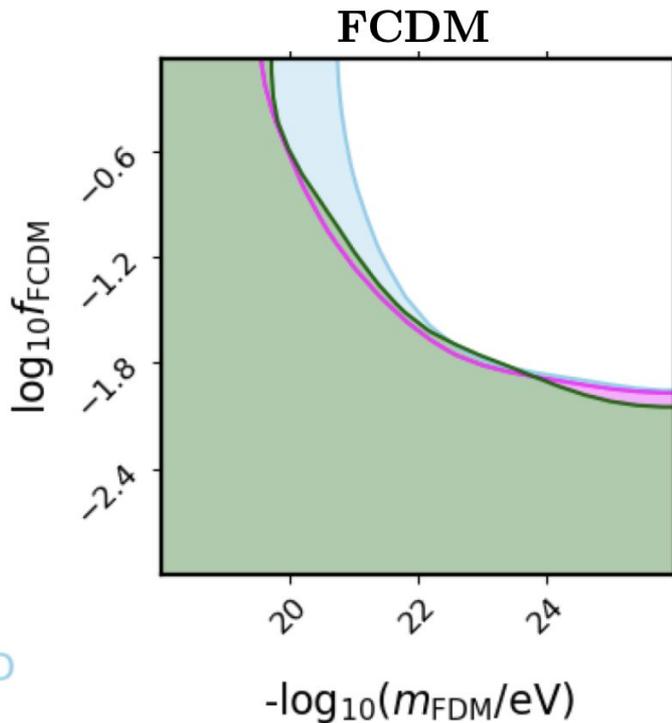
# Constraints on cold + warm DM



TRUNCATED  
DPL  
FLOOR

$f \sim 1 : m_{\text{WDM}} \gtrsim 15 \text{ keV}$  (FLOOR, DPL),  
 $\gtrsim 4 \text{ keV}$  (TRUNCATED)  
CDM + hot relic :  $f \lesssim 1\%$  (FLOOR, DPL, TRUNCATED)

# Constraints on cold + fuzzy DM



TRUNCATED  
DPL  
FLOOR

$f \sim 1 : m_{\text{FDM}} \gtrsim 2 \times 10^{-20}$  eV (FLOOR, DPL),  
 $\gtrsim 2 \times 10^{-21}$  eV (TRUNCATED)  
CDM + hot relic :  $f \lesssim 1\%$  (FLOOR, DPL, TRUNCATED)

# Summary

- With the current upper limits, we are able to exclude extreme states of the IGM
- Models with excess radio background can achieve very high amplitude and therefore they will be easier to constrain with these upper limits
- Halo-model based approach gives a fast and flexible way of exploring many cosmological and astrophysical models
- We can put constraints on non-cold dark matter models using SKA observations of the cosmic dawn