E and B modes of the CMB y-type distortions: Polarised kinetic Sunyaev-Zeldovich effect

State of the Universe Seminar



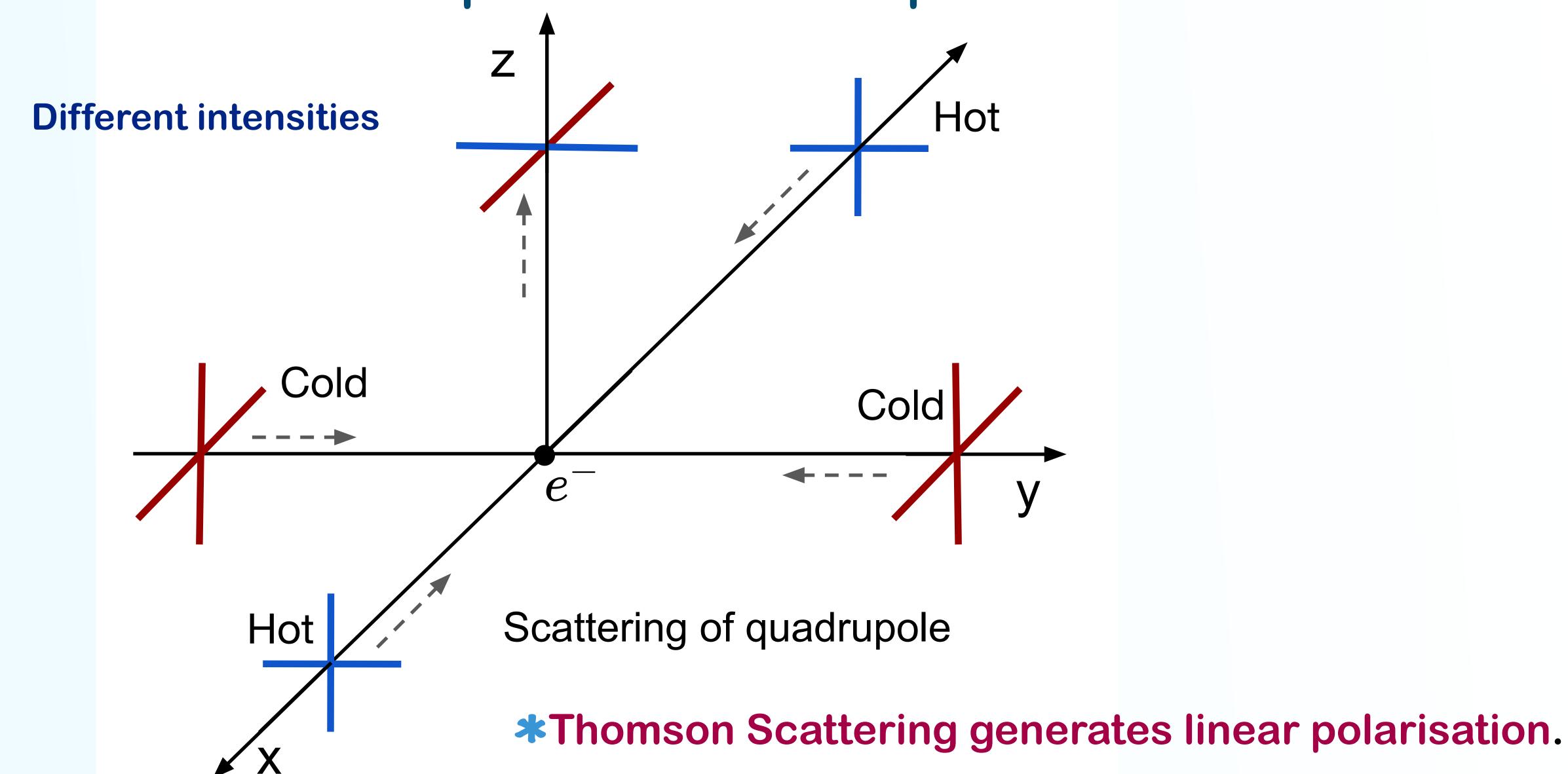
arXiv:2208.02270

September 2022



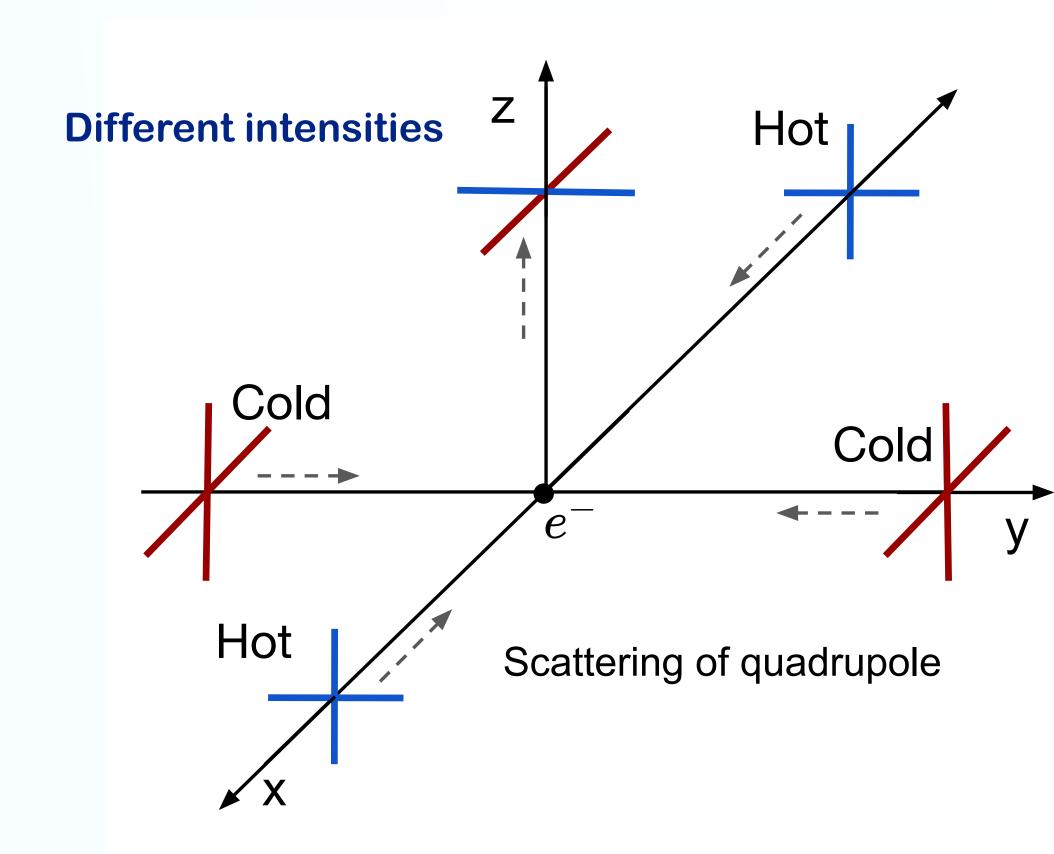


Electron peculiar velocities at second order generate E and B mode polarisation: The pkSZ effect



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- * Free electrons produced during reionisation, have peculiar velocities (\overrightarrow{v})
- * In the electron rest frame, the CMB is not isotropic. Has a quadrupolar anisotropy.



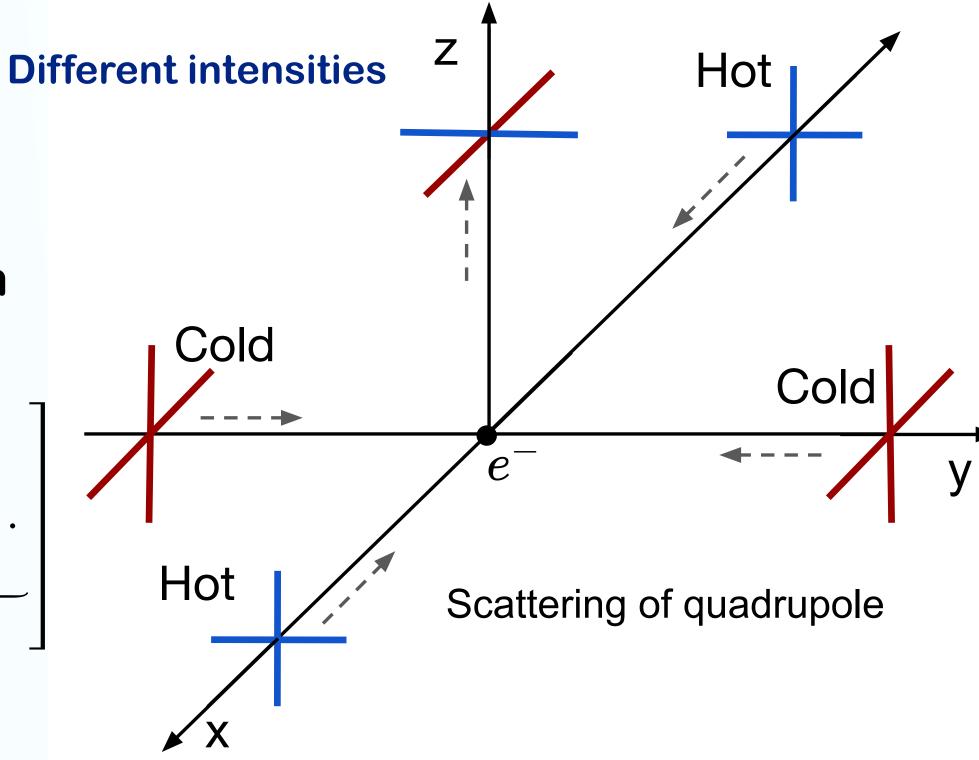
* Predicted by Sunyaev and Zeldovich in 1980. (MNRAS, 190:413-420)

Electron peculiar velocities at second order generate E and B mode polarisation: The pkSZ effect

- * Free electrons produced during reionisation, have peculiar velocities (\overrightarrow{v})
- * In the electron rest frame, the CMB is not isotropic. Has a quadrupolar anisotropy.
 - * Non-linear nature of Relativistic Doppler shift.
 - * A non-linear relation between temperature and intensity in the Planck spectrum

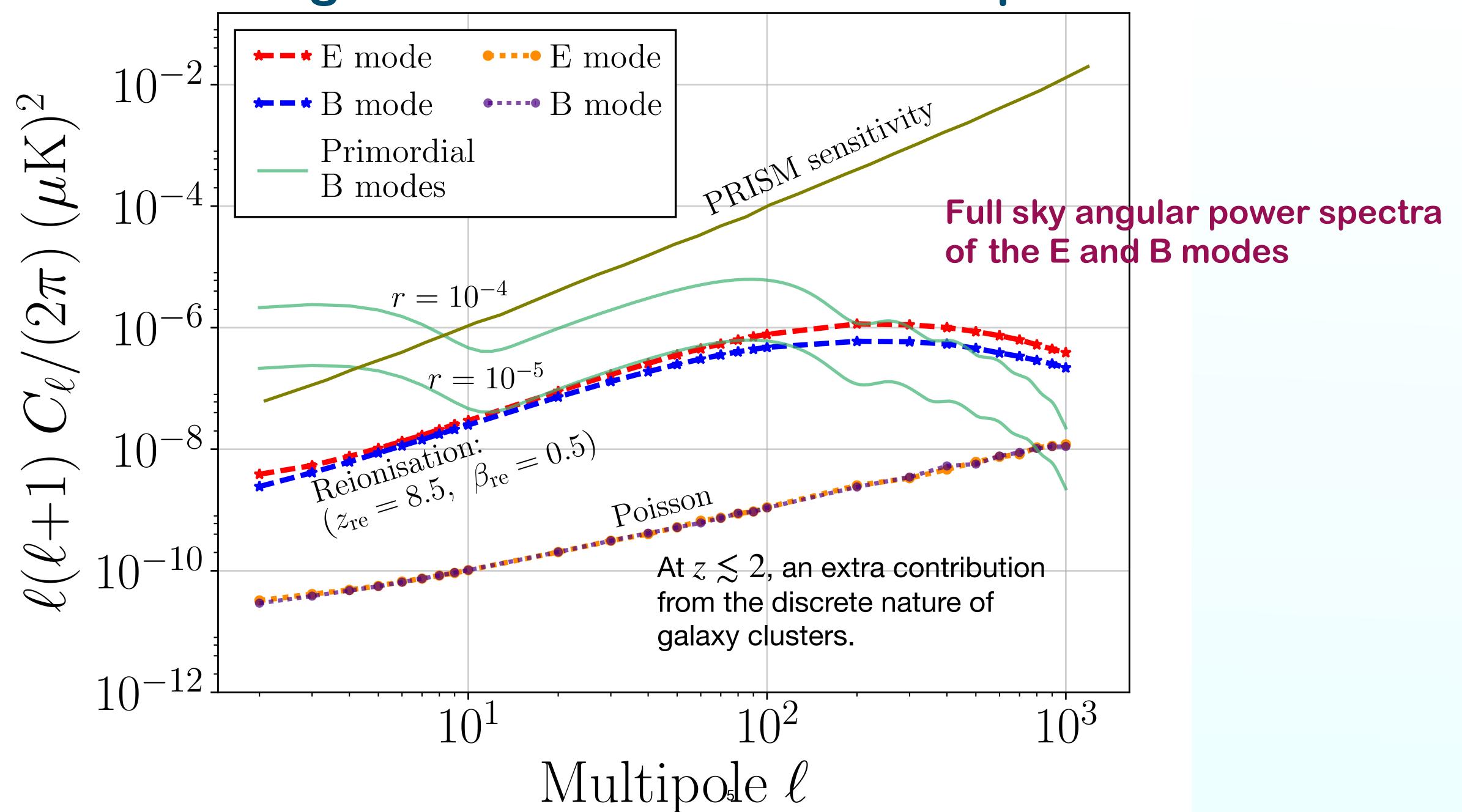
$$T\left(\mathbf{r},\hat{\mathbf{n}}',\eta\right) = \frac{T_0(\eta)}{\gamma\left(1+\mathbf{v}(\mathbf{r},\eta)\cdot\hat{\mathbf{n}}'\right)} = T_0(\eta)\left[1+\frac{1}{2}v^2-\mathbf{v}\cdot\hat{\mathbf{n}}'+\left(\mathbf{v}\cdot\hat{\mathbf{n}}'\right)^2+\mathcal{O}\left(\left(\mathbf{v}\cdot\hat{\mathbf{n}}'\right)^3\right)+\cdots\right]$$

Planck Spectrum: $n_{pl}(x) = \frac{1}{(e^x - 1)}$ $x = \frac{h\nu}{k_B T_0}$



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Beating the cosmic variance with pkSZ effect



Beating the cosmic variance with pkSZ effect

- * The spectrum consists of the y-type distortions part and a blackbody part.
- Differentiates it from primary CMB signals with blackbody spectrum and other SZ-type signals, which are unpolarised.

* Free from the cosmic variance of the primary CMB polarisation signal and lensing B modes.

Sensitive to reionisation central redshift, width, and the matter velocity power spectrum.

The scattered spectrum has a y-type distortion

- *Photons from different blackbody spectra with different temperatures mix.
- *Scattered spectrum not only has a differential blackbody but also a y-type distortion also.

$$\left(\frac{\delta I}{I}\right)\Big|_{\text{(quadrupolar)}} = 2\left(\mathbf{v}\cdot\hat{\mathbf{n}}'\right)^2 \mathbf{g}(\mathbf{x}) + \frac{1}{2}\mathbf{v}(\mathbf{x})\left(\mathbf{v}\cdot\hat{\mathbf{n}}'\right)^2$$

$$\delta n_{\nu} = \frac{1}{2h\nu^{3}} \delta I_{\nu} = \left(\theta + \theta^{2}\right) \left(T \frac{\partial n_{pl}}{\partial T}\right) \bigg|_{T_{0}} + \frac{\theta^{2}}{2} \left(T^{4} \frac{\partial}{\partial T} \left(\frac{1}{T^{2}} \frac{\partial n_{pl}}{\partial T}\right)\right) \bigg|_{T_{0}} + \mathcal{O}(\theta^{3}) \cdots$$

$$n_{pl}(x) = \frac{1}{(e^x - 1)}$$
 $x = \frac{h\nu}{k_B T_0}$

$$g(x) = \frac{xe^x}{(e^x - 1)}$$

$$y(x) = \frac{xe^x}{(e^x - 1)} \left(x \frac{e^x + 1}{e^x - 1} - 4 \right)$$

The scattered spectrum has a y-type distortion

- Distinguishable from the primary polarisation signals which only have a blackbody spectrum.
- *Differentiable from other y-type signals, such as the thermal SZ effect which are unpolarised.
- *The blackbody part will act as a foreground for primordial B modes for $r \lesssim 3 \times 10^{-5}$ for $\ell \gtrsim 100$.

Polarisation depends on square of transverse velocity

* The polarisation field is a spin-2 field.

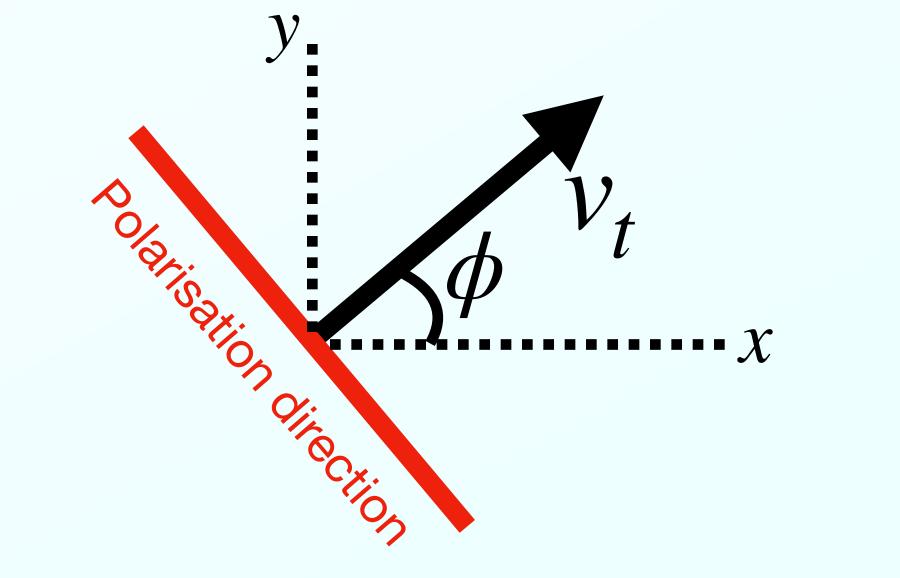
$$\left(\mathcal{Q} \pm i\mathcal{U}\right)\left(\hat{\mathbf{n}}\right) \equiv P_{\pm}\left(\hat{\mathbf{n}}\right)$$

- 1. Electron number density only a function of time.
- 2. Square of transverse velocity.

$$P_{\pm}\left(\hat{\mathbf{n}}\right) = -\frac{\sqrt{6}\sigma_{\mathrm{T}}}{10} \int_{0}^{\chi} d\chi \, a(\chi) \, e^{-\tau(\chi)} \mathbf{n}_{\mathrm{e}}(\chi) \sum_{\lambda=-2}^{2} \, \pm_{2} Y_{2\lambda}\left(\hat{\mathbf{n}}\right) \int d^{2}\hat{\mathbf{n}}' \, Y_{2\lambda}^{*}\left(\hat{\mathbf{n}}'\right) \frac{\left(\mathbf{v}(\mathbf{r},\chi) \cdot \hat{\mathbf{n}}'\right)^{2}}{\left(\mathbf{v}(\mathbf{r},\chi) \cdot \hat{\mathbf{n}}'\right)^{2}}.$$

Quadrupole

$$P_{+}\left(\hat{\mathbf{z}}\right) = -\frac{\sigma_{\mathrm{T}}}{10} \int_{0}^{\chi_{i}} d\chi \ e^{-\tau(\chi)} \ n_{\mathrm{e}}(\chi) \ a(\chi) \frac{v_{t}^{2}}{v_{t}} e^{-2i\phi}$$



Polarisation field and angular power spectra

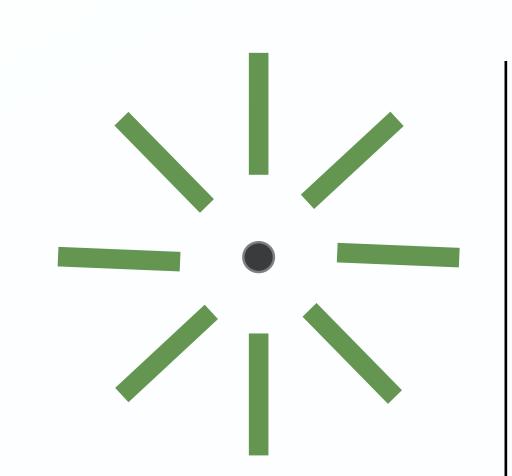
* The polarisation field is a spin-2 field.

$$(\mathcal{Q} \pm i\mathcal{U})(\hat{\mathbf{n}}) \equiv P_{\pm}(\hat{\mathbf{n}}) \qquad a_{\ell m} = \int P_{+}(\hat{\mathbf{n}}) \ _{2}Y_{\ell m}^{*}(\hat{\mathbf{n}}) d^{2}\hat{\mathbf{n}}$$

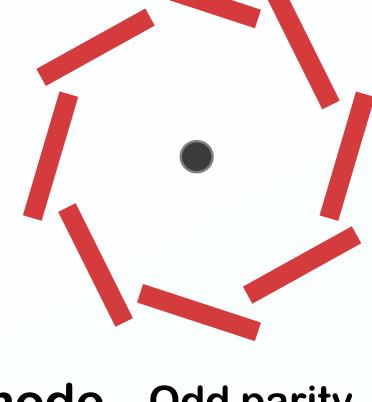
$$P_{\pm}(\hat{\mathbf{n}}) = \sum_{\ell,m} (e_{\ell,m} + ib_{\ell,m}) \ _{\pm 2}Y_{\ell,m}$$

E mode

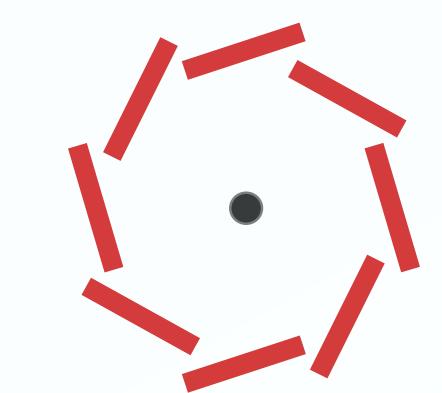
B mode



E mode - Even parity







*Construct spin-0 fields related to polarisation field.

$$e_{\ell m} = \frac{1}{2} \left(a_{\ell m} + (-1)^m a_{\ell - m}^* \right) \qquad b_{\ell m} = \frac{-i}{2} \left(a_{\ell m} - (-1)^m a_{\ell - m}^* \right)$$

* The E and B mode power spectra:

$$\langle e_{\ell m} e_{\ell' m'}^* \rangle = C_{\ell}^{EE} \, \delta_{\ell,\ell'} \, \delta_{m,m'} \qquad \langle b_{\ell m} b_{\ell' m'}^* \rangle = C_{\ell}^{BB} \, \delta_{\ell,\ell'} \, \delta_{m,m'}$$

The power spectra at second order is a high dimensional integral.

$$\begin{split} C_{\ell}^{BB} &= \frac{T_{CMB}^2}{2} \left[(4\pi) \left(\frac{4\pi}{3} \right)^2 \sqrt{\frac{3}{2\pi}} \frac{\sqrt{6}\sigma_{\mathrm{T}}}{10} \right]^2 \sum_{\lambda,\lambda' = -2}^2 (-1)^{(\lambda + \lambda')} \int_0^\chi d\chi \ e^{-\tau(\chi)} \ a(\chi) \int_0^\chi d\chi' \ e^{-\tau(\chi')} \times \\ & a(\chi') \mathbf{n}_{\mathrm{e}}(\chi) \mathbf{n}_{\mathrm{e}}(\chi') \sum_{\substack{L,M \\ \ell_1,M' \neq_1, \nu_2}} \sum_{p_1, p_2} i^{(L-L')} \left(\frac{1}{p_1} \frac{1}{p_2} \frac{2}{-\lambda} \right) \left(\frac{1}{p_1'} \frac{1}{p_2'} \frac{2}{-\lambda'} \right) \int \int \frac{k_1^2 dk_1 \, k_2^2 dk_2}{(2\pi)^6} \times \\ & P_{uu}(k_1) P_{uu}(k_2) \, j_L(k\chi) \, j_L(k'\chi') \int d\Omega_{\mathbf{k}_1} \int d\Omega_{\mathbf{k}_2} \, Y_{LM}^*(\hat{\mathbf{k}}) Y_{L'M}(\hat{\mathbf{k}}) Y_{1p_1}^*(\hat{\mathbf{k}}_1) Y_{1p_2}^*(\hat{\mathbf{k}}_2) \times \\ & Y_{1p_1'}(\hat{\mathbf{k}}_1') Y_{1p_2'}(\hat{\mathbf{k}}_2') \, A_{\ell m}^{\lambda LM} \, A_{\ell m}^{\lambda' L'M'} \left(1 - (-1)^{(L+\ell)} \right) \left(1 - (-1)^{(L'+\ell)} \right) \end{split}$$

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Matter velocity power spectrum

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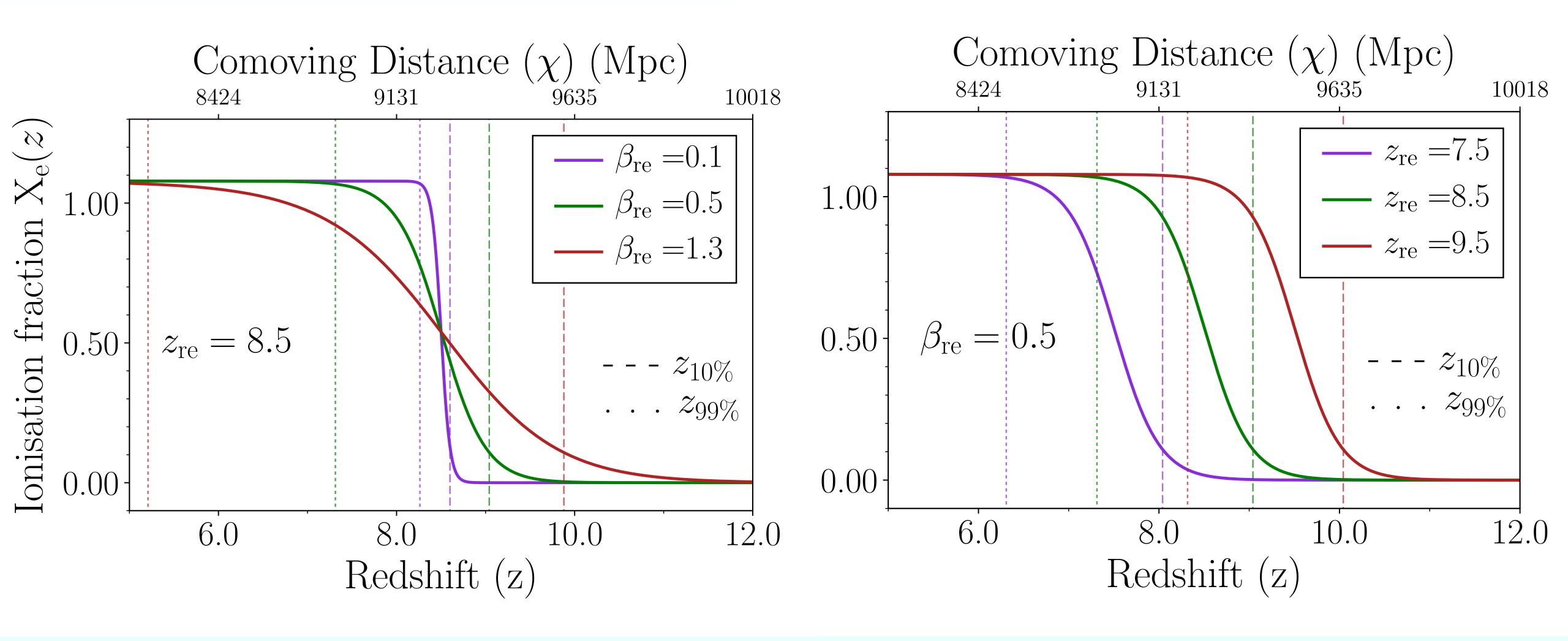
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Matter velocity power spectrum

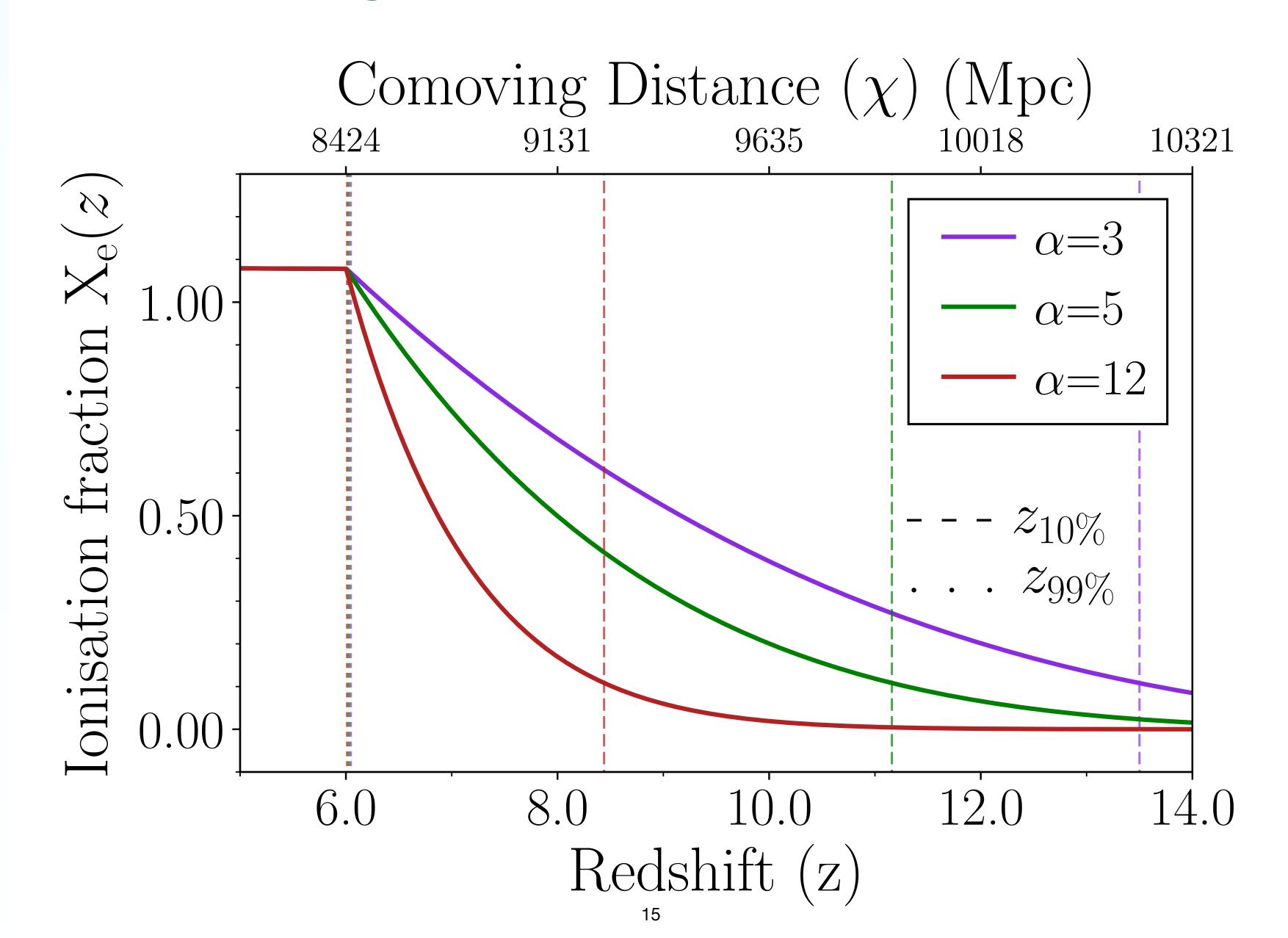
Electron number density—(Reionisation history)

$$A_{\ell m}^{\lambda LM} = \sqrt{\frac{5(2L+1)(2\ell+1)}{4\pi}} (-1)^{(m)} \begin{pmatrix} L & 2 & \ell \\ 0 & -2 & 2 \end{pmatrix} \begin{pmatrix} L & 2 & \ell \\ M & \lambda & -m \end{pmatrix}$$

Symmetric Reionisation

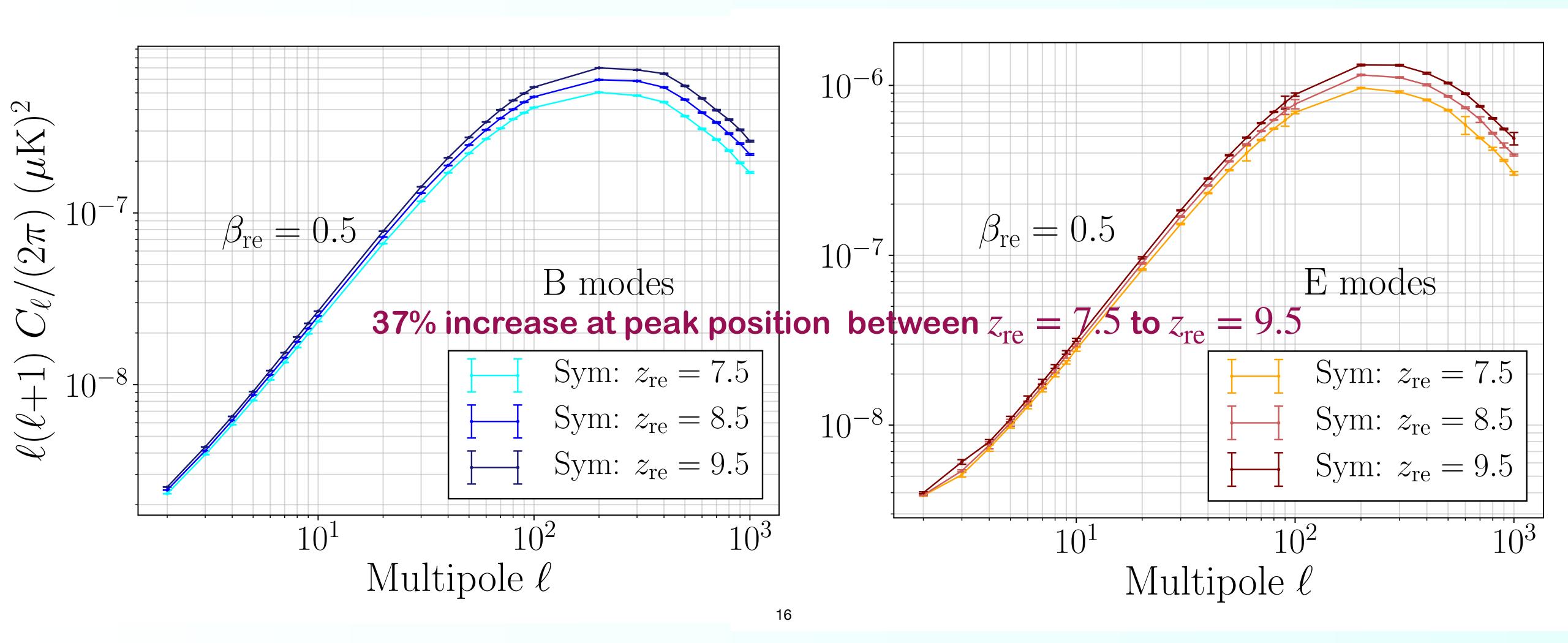


Asymmetric Reionisation



pkSZ effect is sensitive to the redshift of central reionization

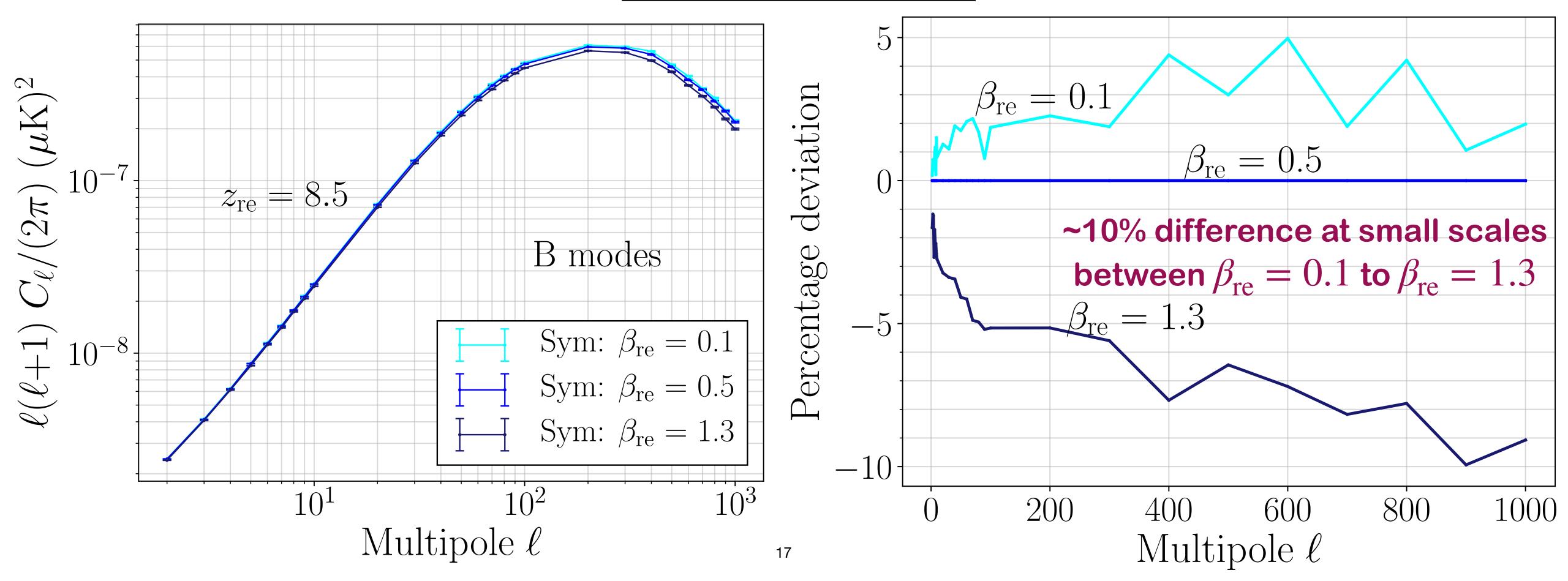
- * The power spectra increase with the increase in the central redshift of reionisation.
- Increasing the central redshift increases the total Thomson optical depth.



pkSZ effect is sensitive to the reionisation width

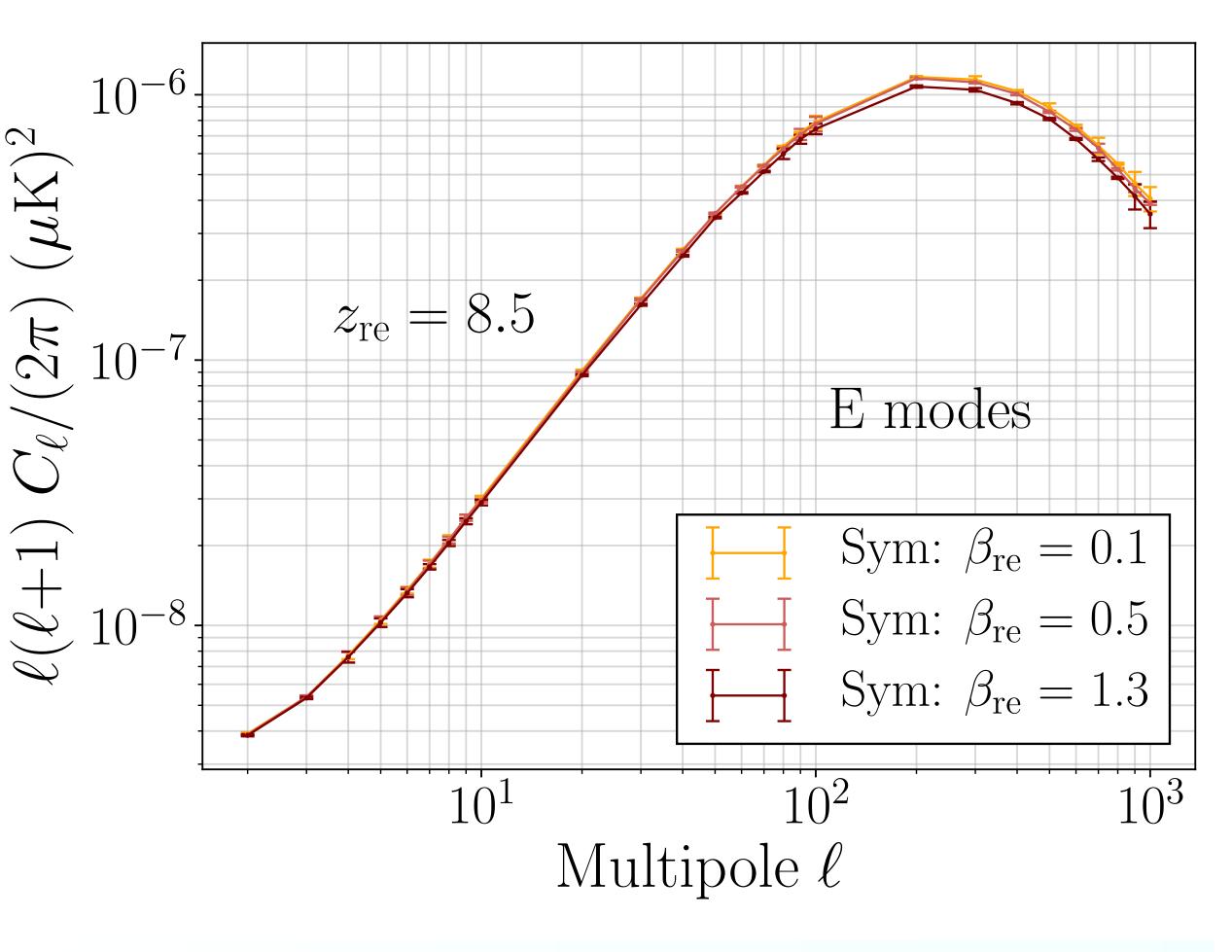
- *Changing the width at a fixed central redshift has a negligible effect on the optical depth
- *The power spectra still decrease with the increase in the duration of reionisation.

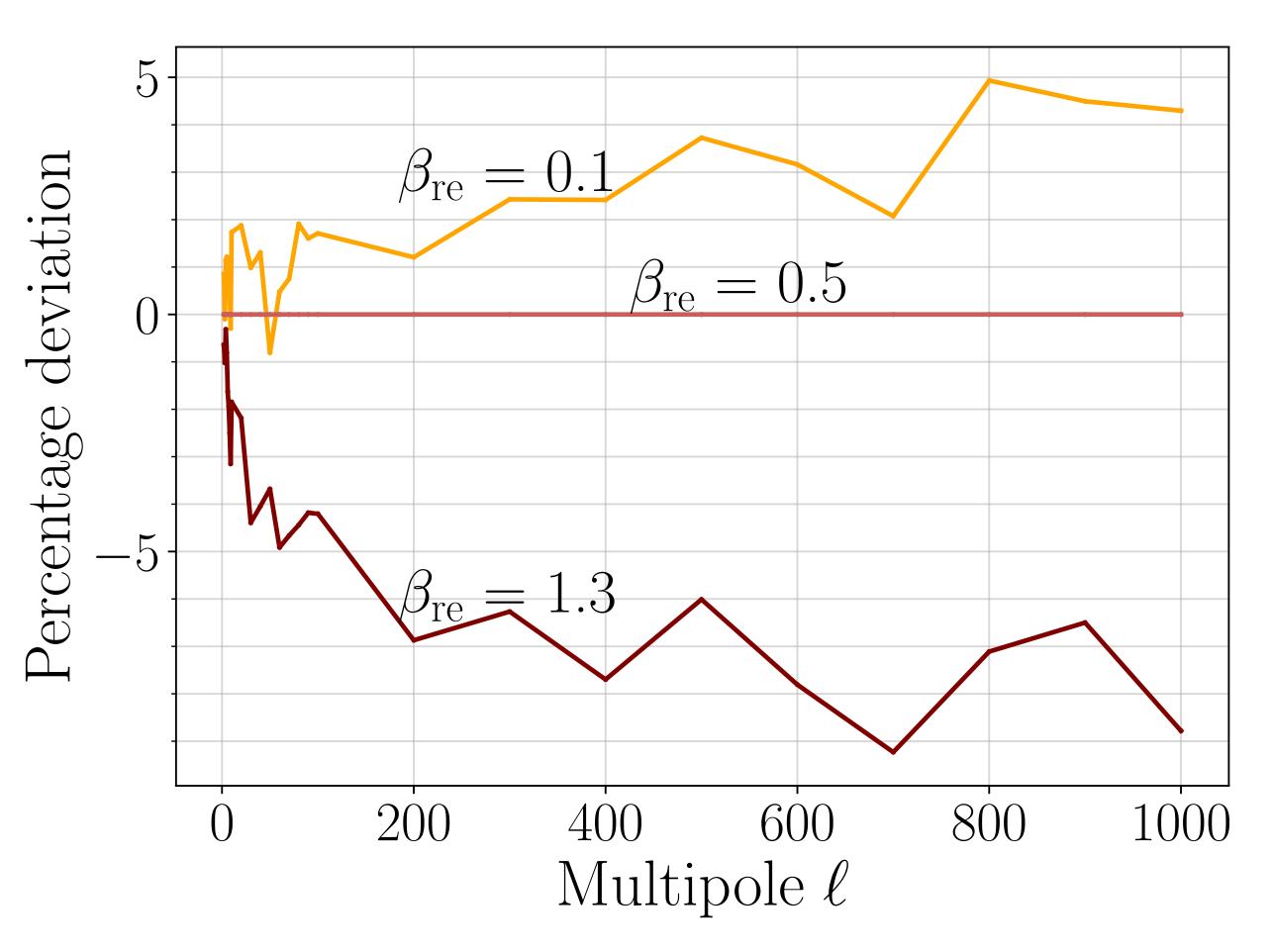
Width =
$$z_{99\%} - z_{10\%}$$



pkSZ effect is sensitive to the reionisation width

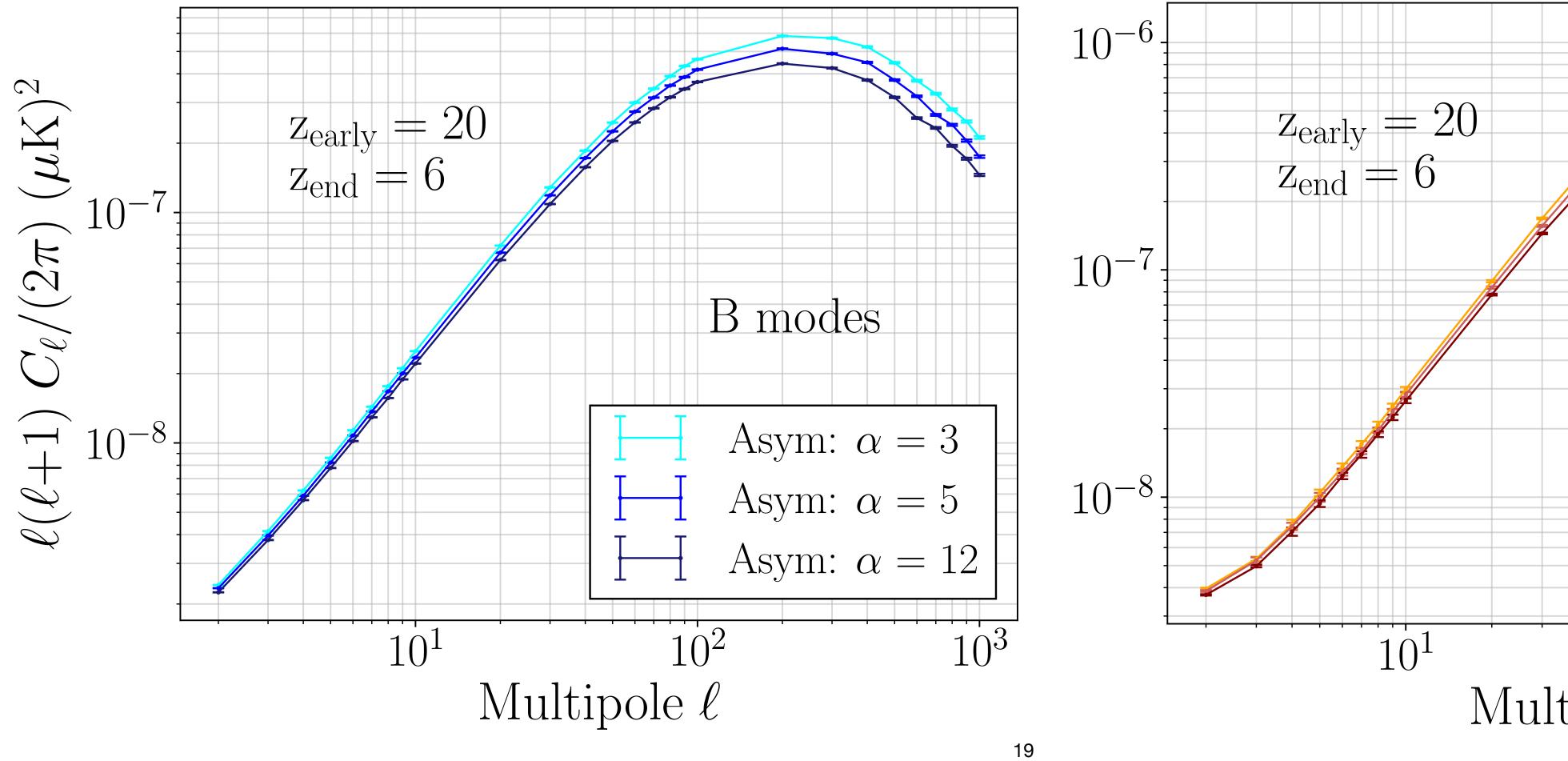
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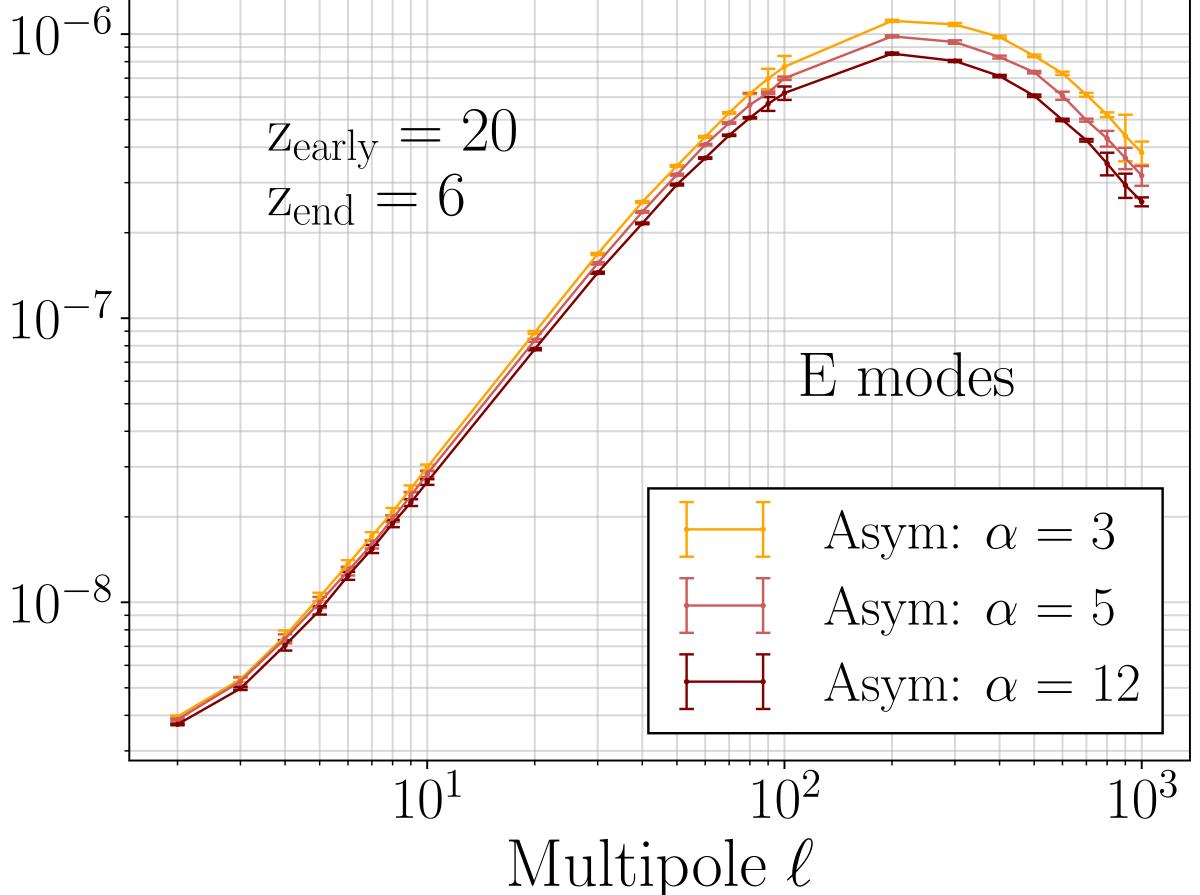




pkSZ effect is sensitive to the rapidity of reionisation

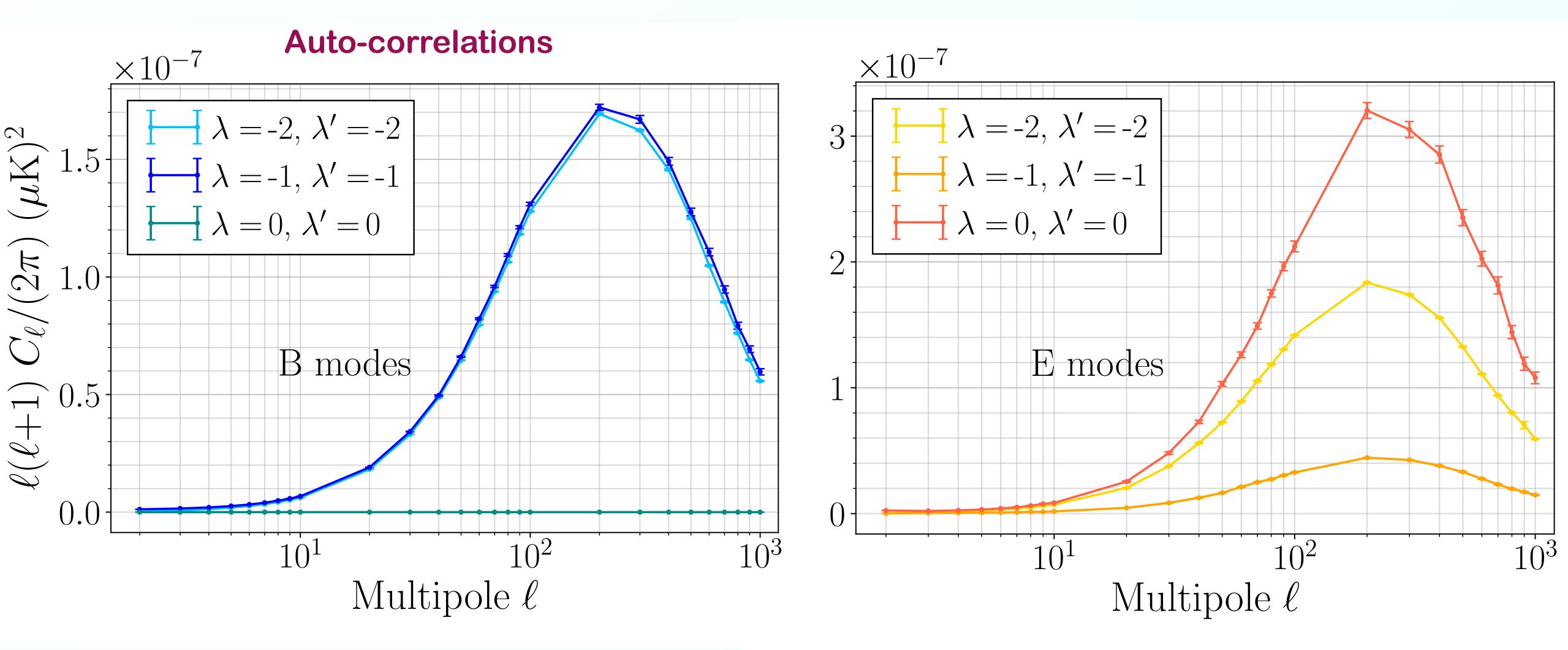
* In the case of asymmetric reionisation, the power spectra are sensitive to how quickly reionisation occurs.





E modes are greater than the B modes

Scalar ($\lambda = 0$), Vector ($\lambda = 1$) and Tensor ($\lambda = 2$) Decomposition



Previous Works

* Renaux-Petel et al. arXiv:1312.4448: They claimed to observe only a 2% effect on the width of the Reionisation but we observed a much larger effect.

* Kamionkowski et al. arXiv: 2203.12503: They concentrated on a non-Gaussianity and birefringence and not on reionisation

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Primary CMB anisotropies are sensitive to only the total reionisation optical depth but the pkSZ effect is sensitive to the Reionisation history.

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Thank You!!