

# Cosmological signatures of hidden sector containing dark radiation

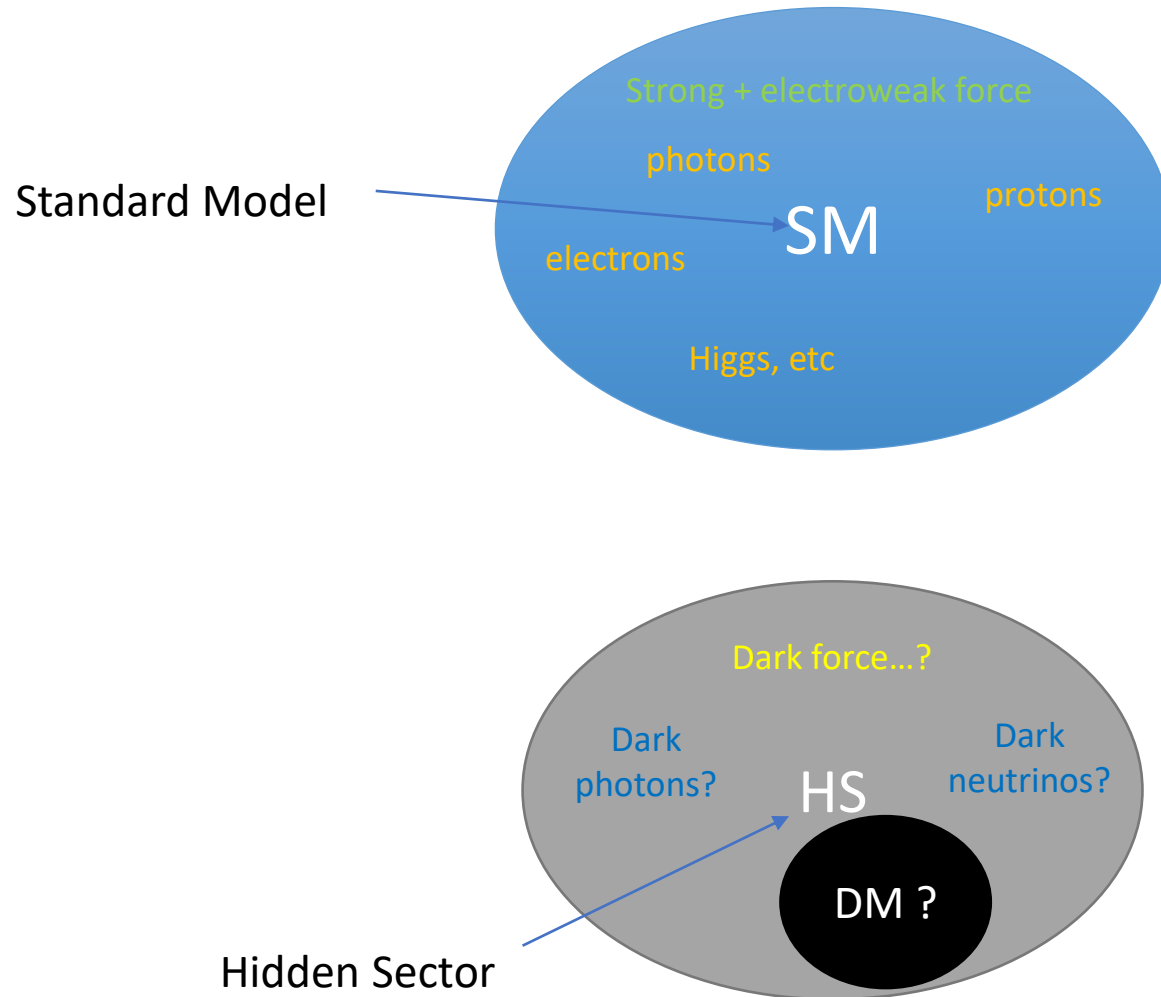
arxiv:2206.13530, 2006.01165

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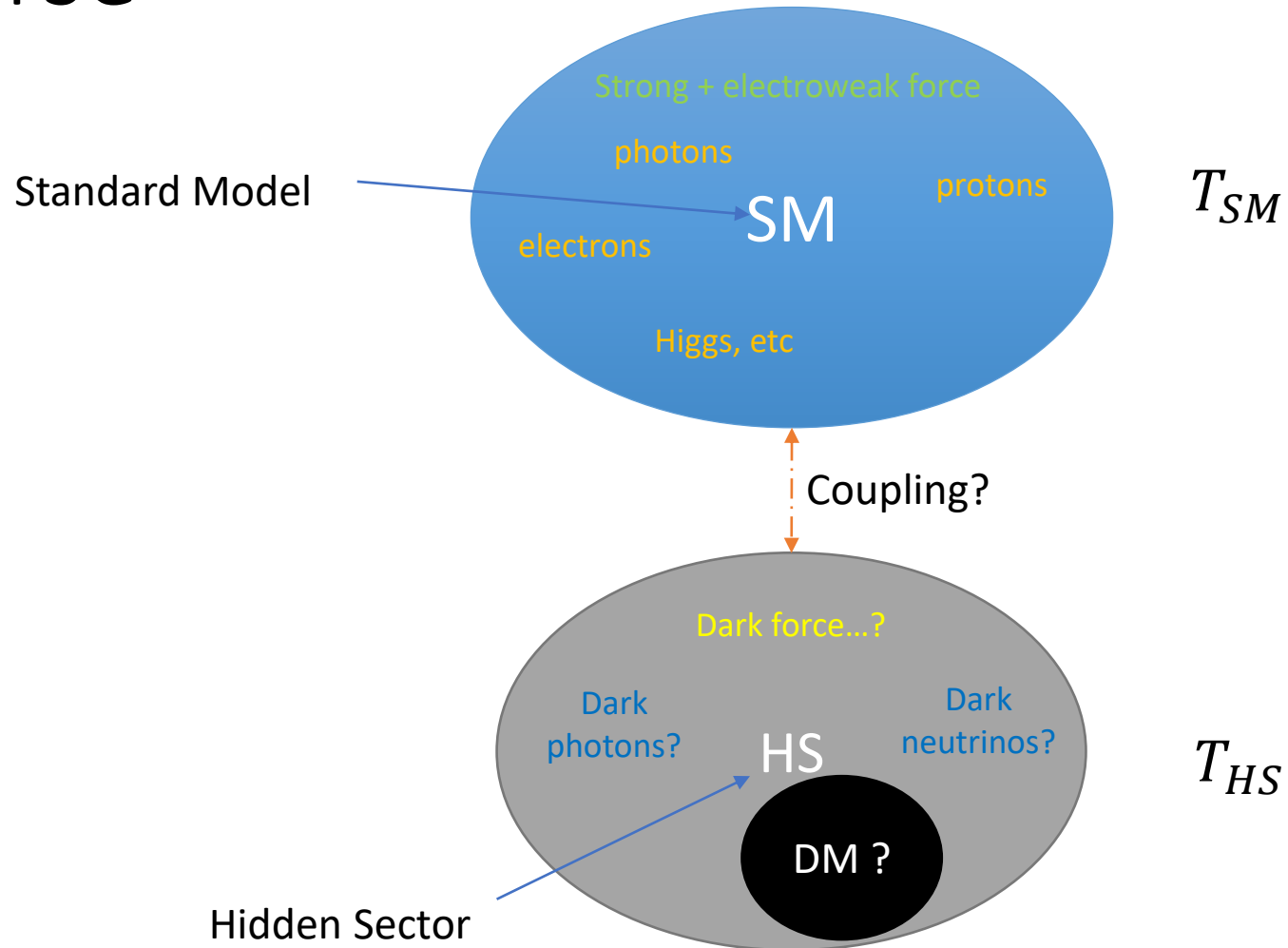
-Pranjal Ralegankar  
University of Illinois at Urbana-Champaign

Collaborators: Peter Adshead and  
Jessie Shelton

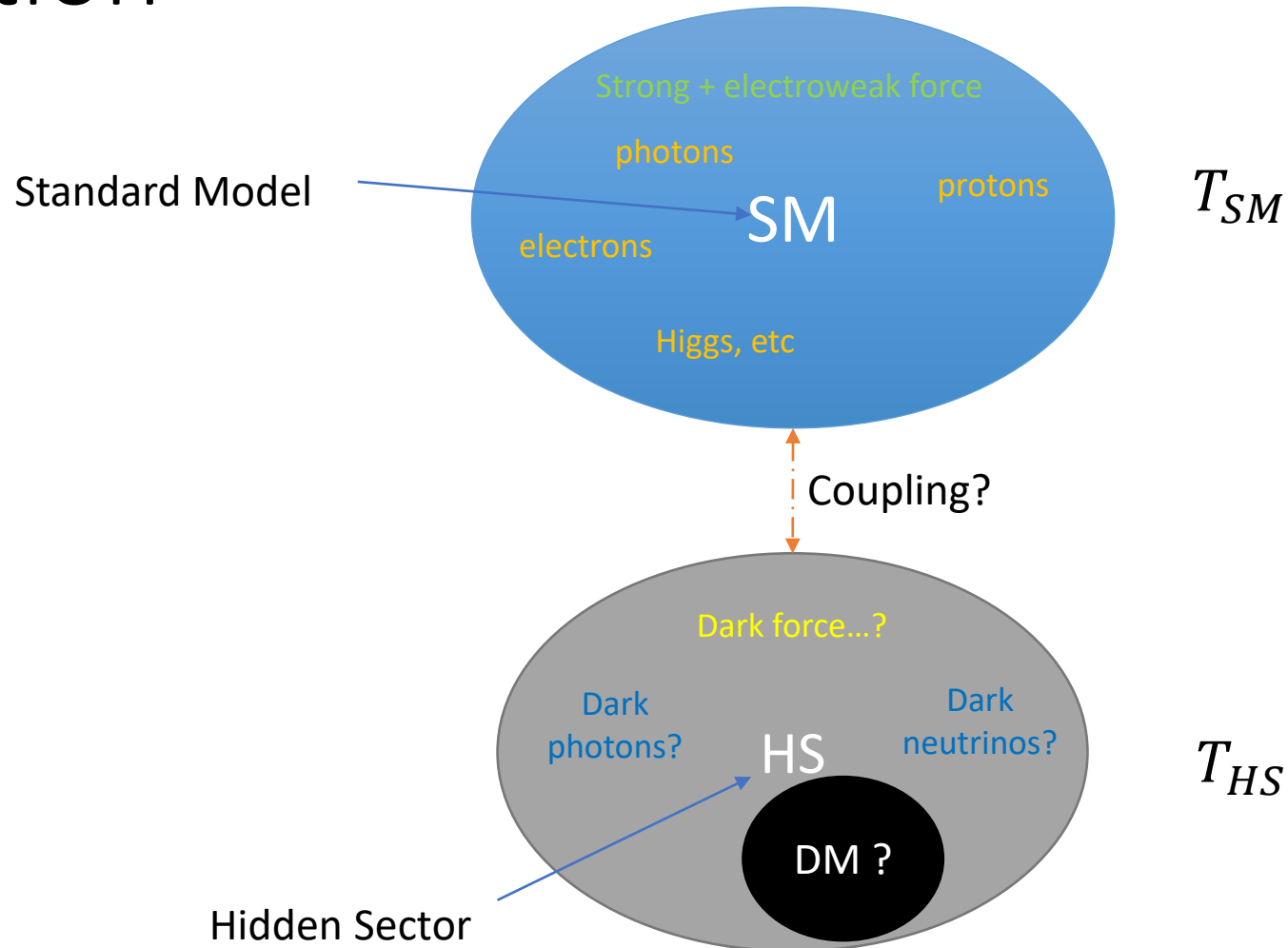
# Hidden sector 101



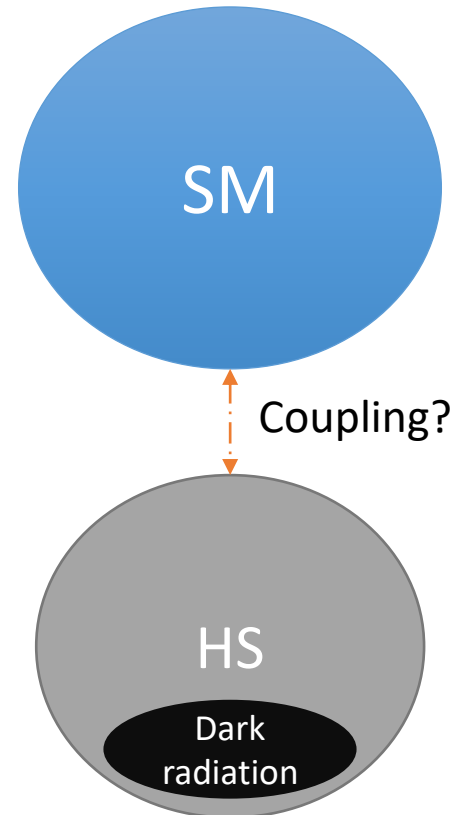
# Hidden sector 101: Two baths in the early universe



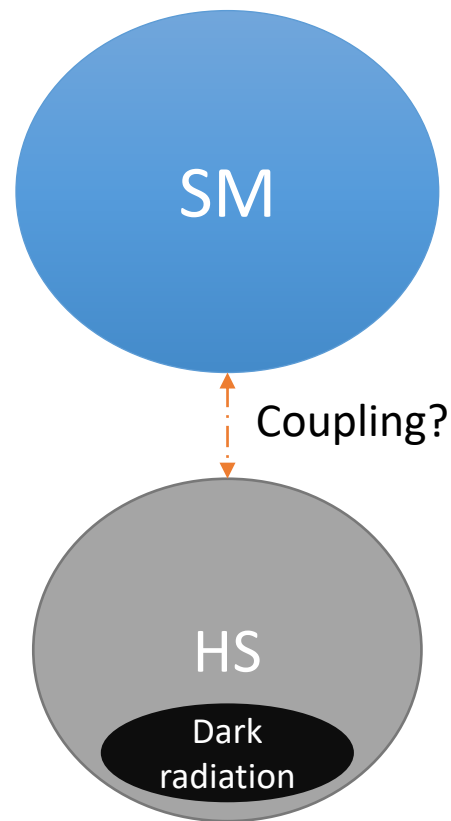
# Hidden sector 101: generically have dark radiation



Goal: Probe models containing dark radiation using cosmological observations

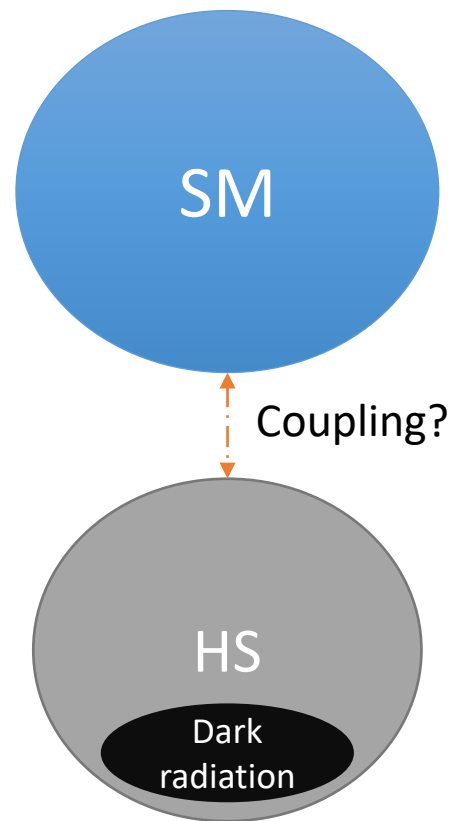


# Goal: Probe models containing dark radiation using cosmological observations



- If no dark radiation detected in future: Constrain portal interactions with hidden sectors  
(Adshead, Ralegankar, Shelton 2022)
- If dark radiation detected: Search for interaction by checking adiabaticity of perturbations during BBN  
(Adshead, Holder, Ralegankar 2020)

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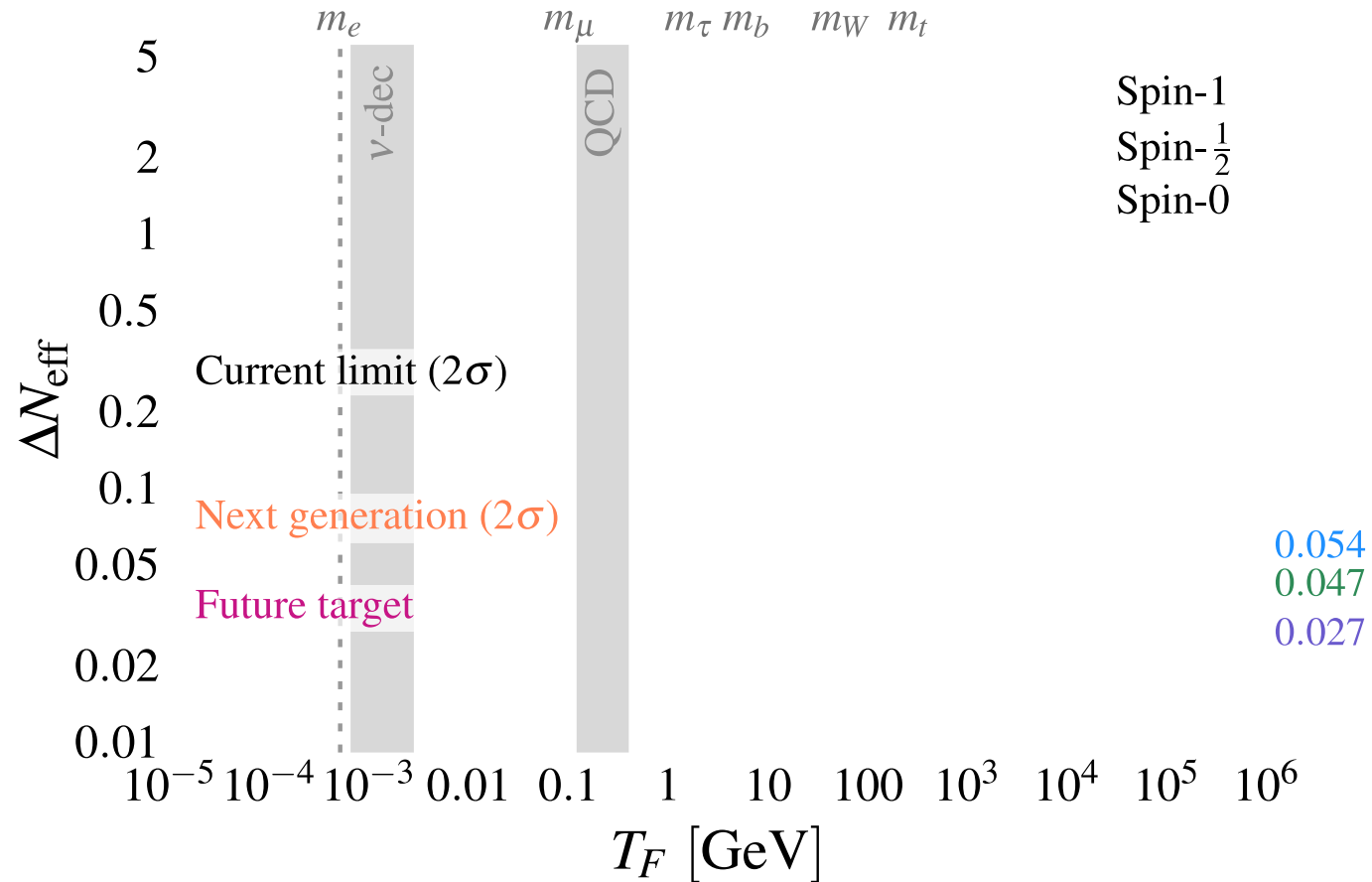
# $N_{\text{eff}}$ : A parametrization for dark radiation abundance

- $N_{\text{eff}}$  : effective number of neutrino species
- CMB sensitive to  $\rho_\nu$ :  $N_{\text{eff}}^{\text{SM}} = \frac{8}{7} \left(\frac{11}{4}\right)^{4/3} \frac{\rho_\nu}{\rho_\gamma} = 3.044$
- CMB also sensitive to  $\rho_{\text{dr}}$ :  $\Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4}\right)^{4/3} \frac{\rho_{\text{dr}}}{\rho_\gamma}$



# $\Delta N_{\text{eff}}$ : Typically discussed as constraint on decoupling temperature

$$\Delta N_{\text{eff}} = \frac{8}{7} \left( \frac{11}{4} \right)^{4/3} \frac{\rho_{\text{dr}}}{\rho_{\gamma}}$$

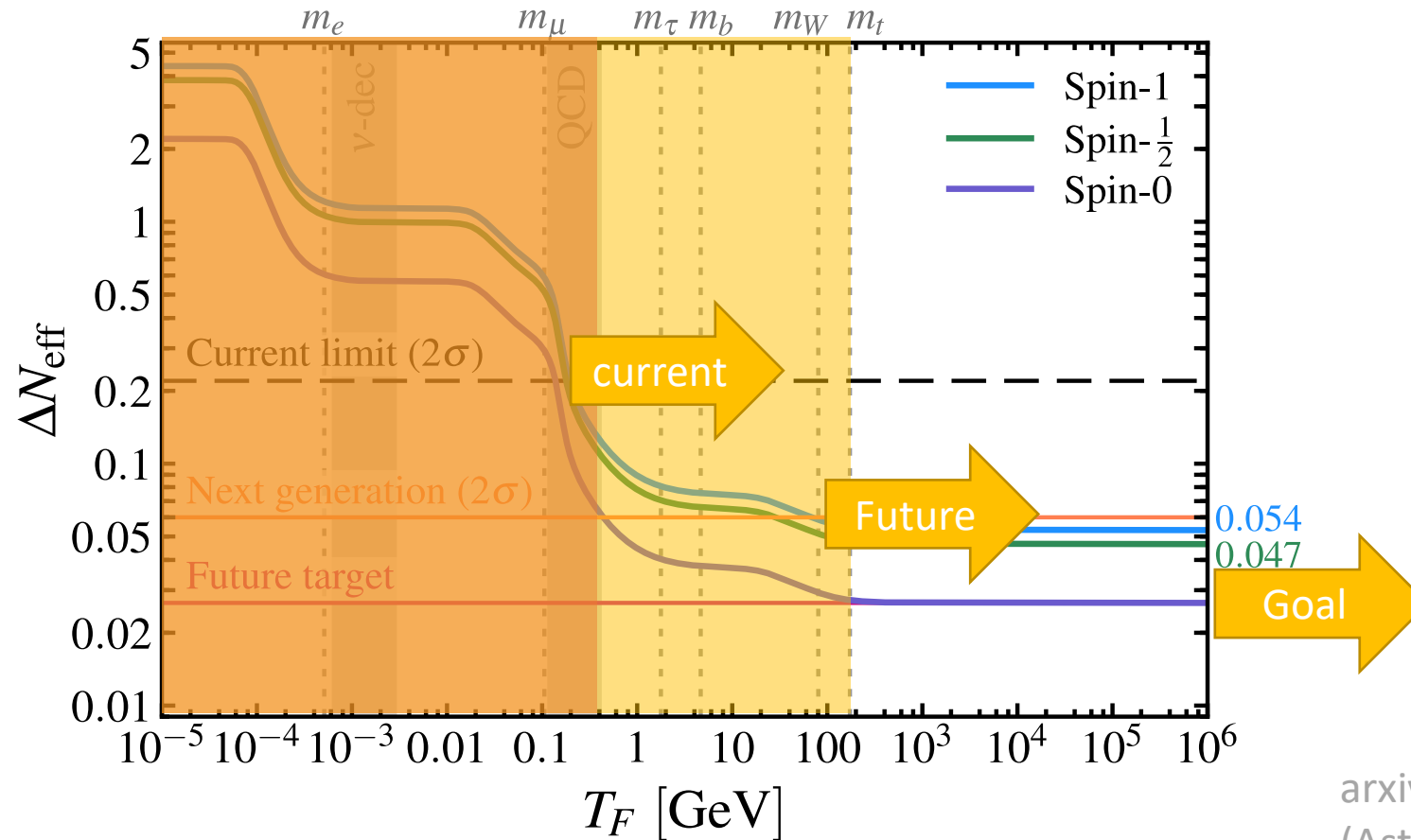


Decoupling temperature of new light particle

arxiv:1903:04763  
(Astro 2020 white paper)

# Reinterpreting $\Delta N_{\text{eff}}$ : Constraint on interactions with out-of-equilibrium sectors

$$\Delta N_{\text{eff}} = \frac{8}{7} \left( \frac{11}{4} \right)^{4/3} \frac{\rho_{\text{dr}}}{\rho_{\gamma}}$$



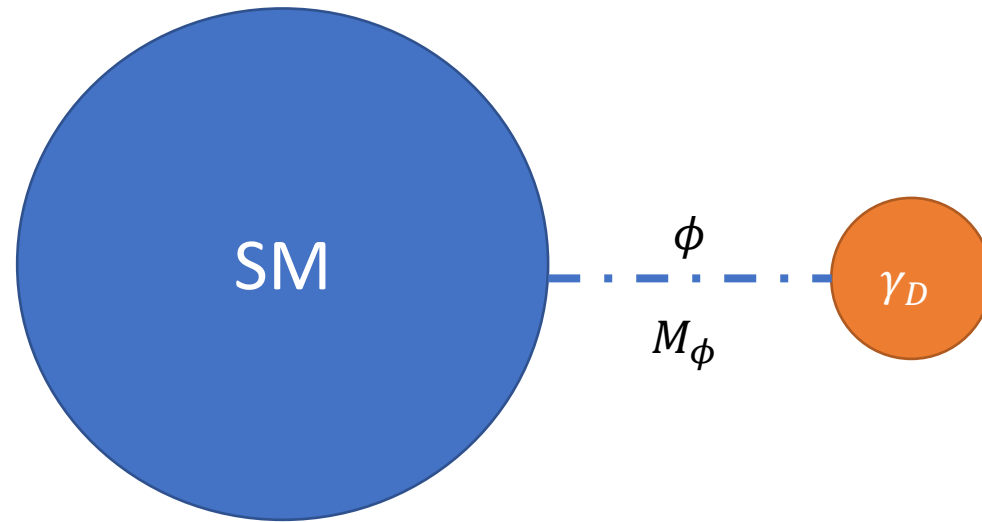
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Unified treatment for calculating  $N_{\text{eff}}$   
constraints on beyond SM interactions

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Strong Implications for model building with  
HS containing dark radiation

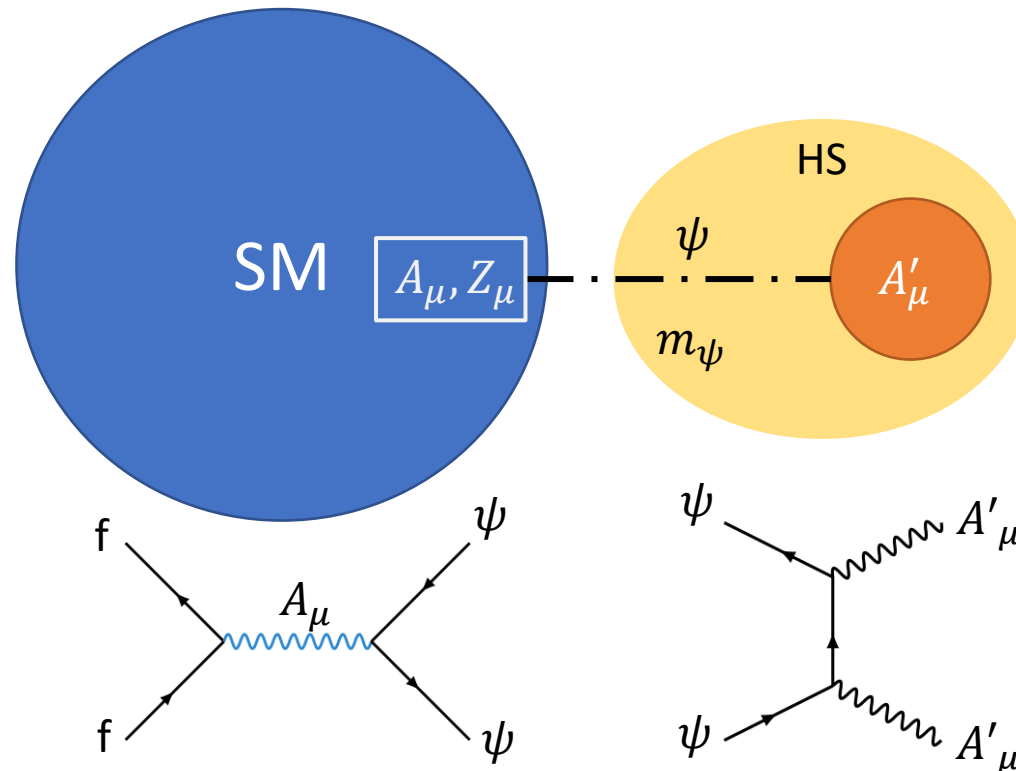
# Application of $N_{\text{eff}}$ constraint : Relevant types of interaction



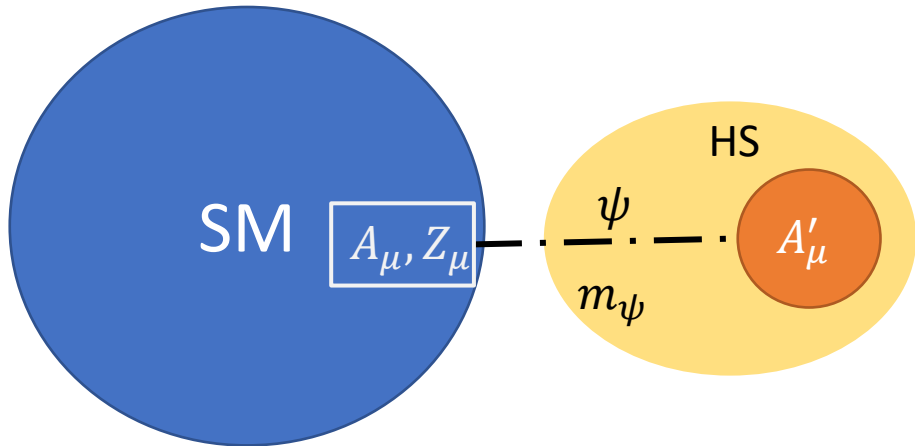
# Application of $N_{\text{eff}}$ constraint: Millicharged particle example

$$L_{\text{int}} \supset -\frac{\epsilon}{2} B_{\mu\nu} F^{\mu\nu'} + e' A'_\mu \bar{\psi} \gamma^\mu \psi - m_\psi \bar{\psi} \psi$$

$$L_{\text{int}} \supset -eQA_\mu \bar{\psi} \gamma^\mu \psi + e' A'_\mu \bar{\psi} \gamma^\mu \psi + eQZ_\mu \tan \theta_W \bar{\psi} \gamma^\mu \psi - m_\psi \bar{\psi} \psi$$



# Physics behind dark radiation production: Boltzmann equations



Boltzmann equations:

$$\frac{d\rho_{SM}}{dt} + 3H(1 + w_{SM})\rho_{SM} = -C$$

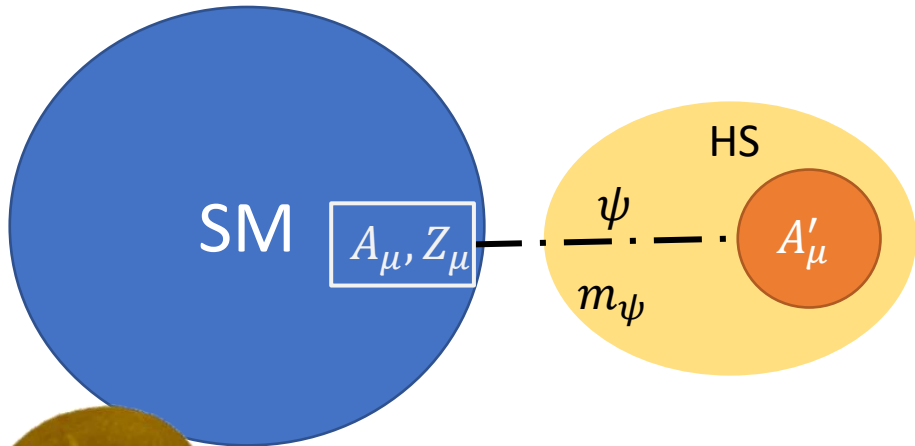
$$\frac{d\rho_{HS}}{dt} + 3H(1 + w_{HS})\rho_{HS} = C$$

$$H = \frac{\sqrt{\rho_{SM} + \rho_{HS}}}{\sqrt{3}M_{\text{Pl}}}$$

$$C = \frac{1}{32\pi^4} \sum_f \int ds (s - 4m_f^2) s \sigma_{ff \rightarrow \psi\psi} [T_{SM} G(\sqrt{s}/T_{SM}) - T_{HS} G(\sqrt{s}/T_{HS})] + \dots$$

↑  
Energy transfer collision term

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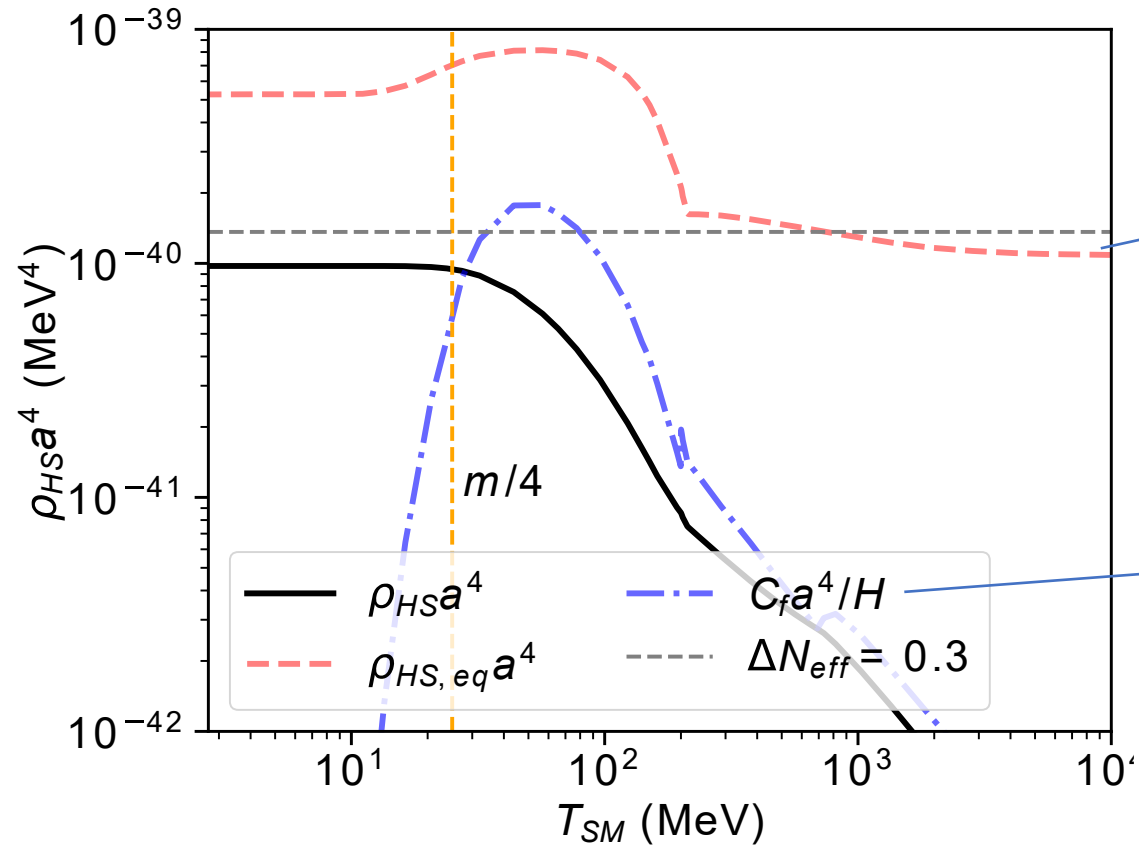
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Energy transfer collision term





# Physics behind dark radiation production: Plots!

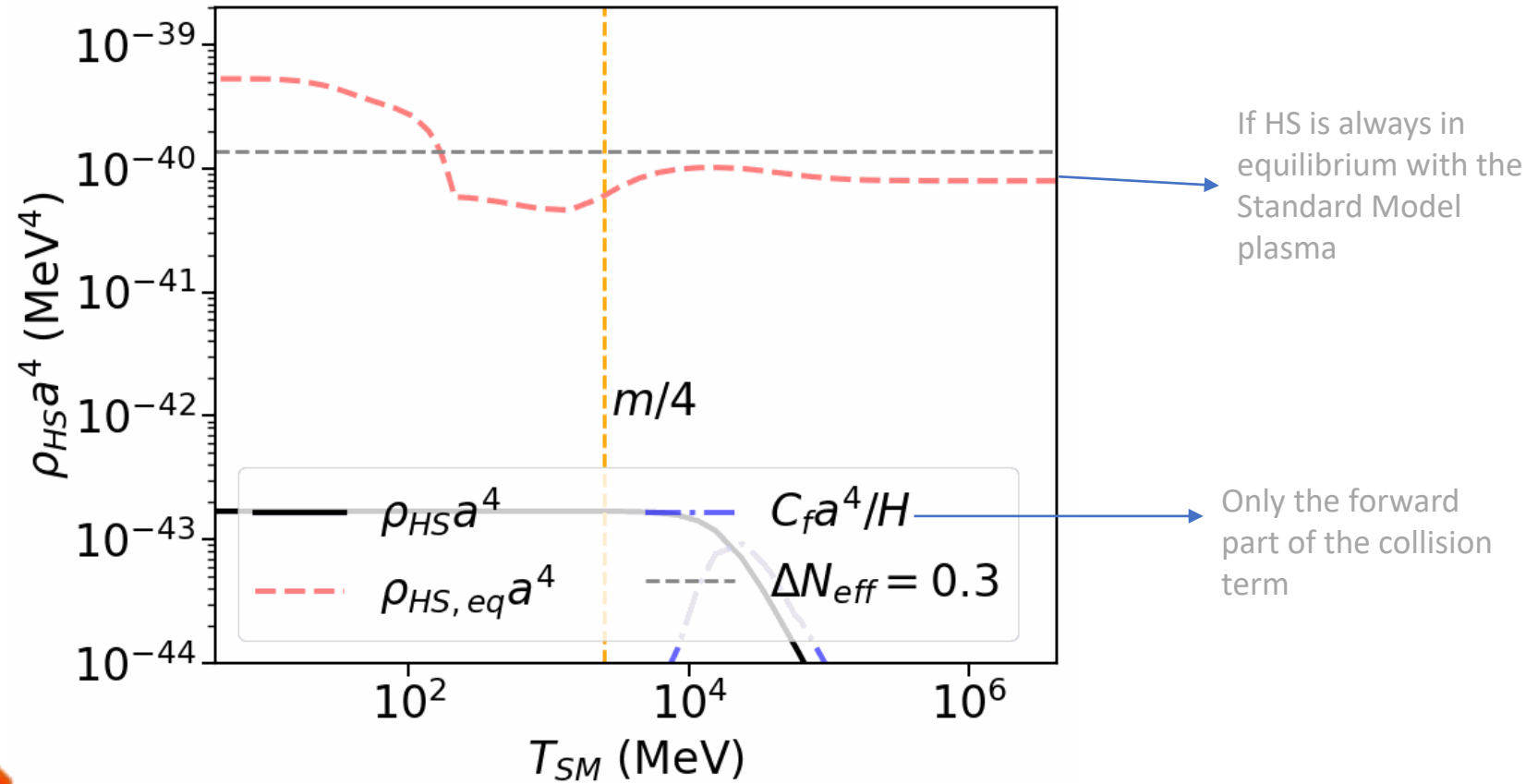


If HS is always in equilibrium with the Standard Model plasma

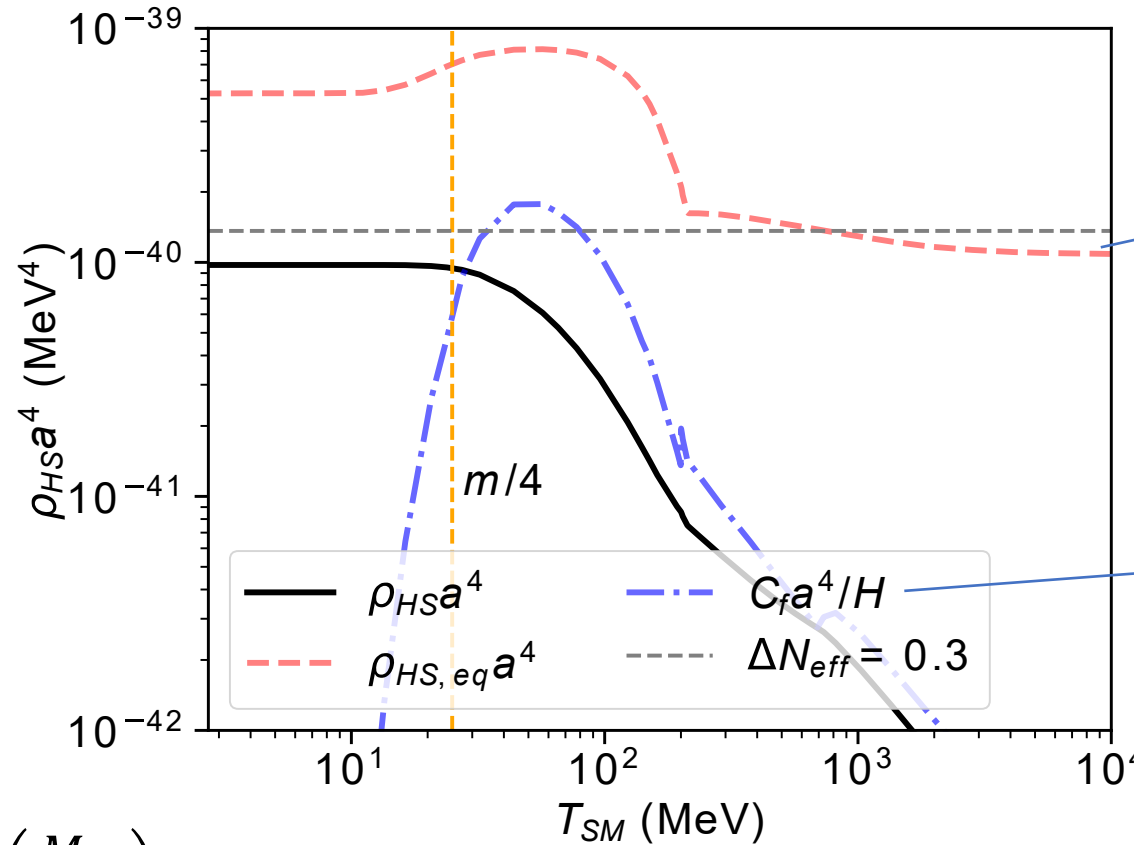
Only the forward part of the collision term



# Physics behind dark radiation production: GIF!



# Physics behind dark radiation production: Out-of-equilibrium $\rho_{\text{DR}}$ proportional to portal coupling



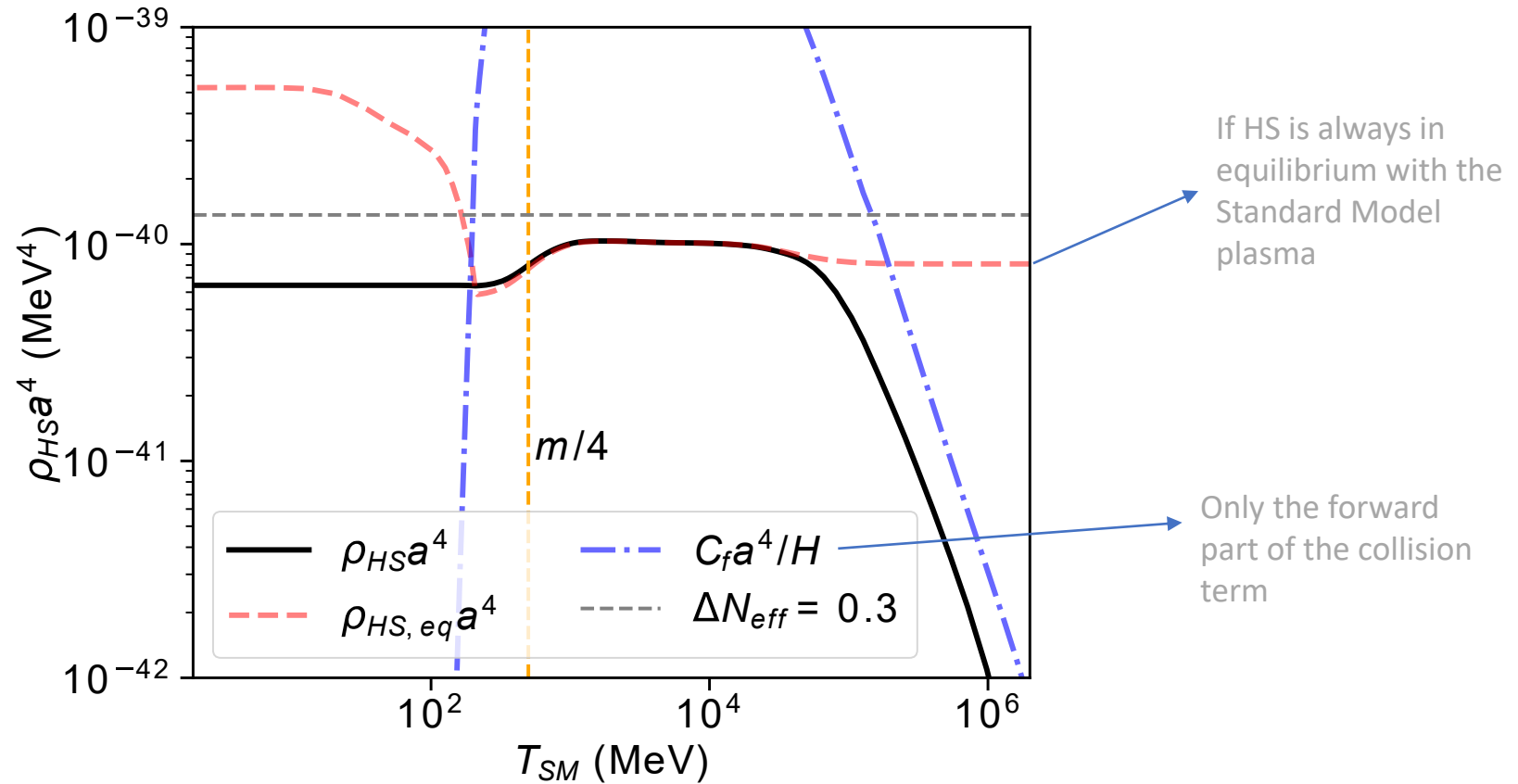
If HS is always in equilibrium with the Standard Model plasma

Only the forward part of the collision term

$$\Gamma_E = \frac{C_f}{\rho_{HS,eq}}$$

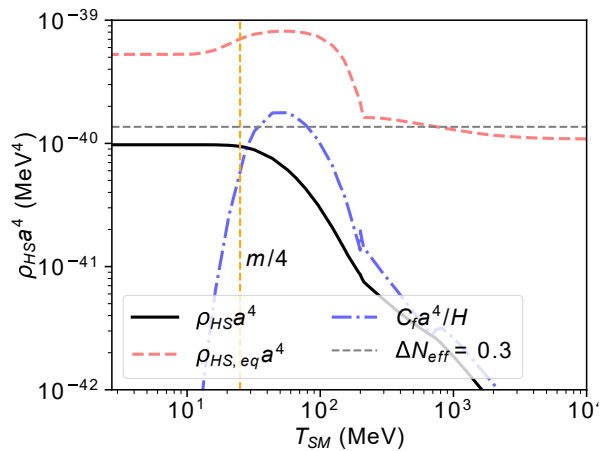
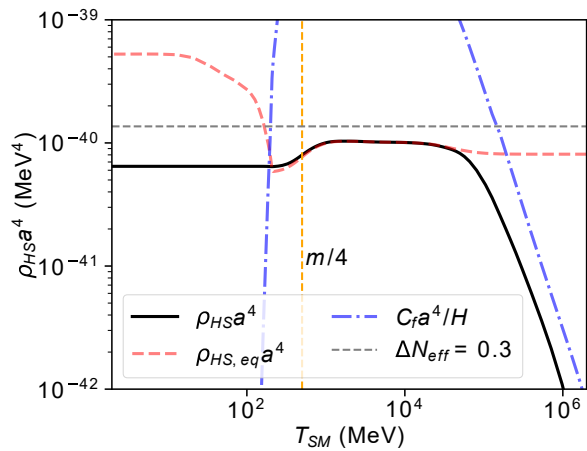
$$\frac{\rho_{HS}}{\rho_{SM}} \sim \left( \frac{\Gamma_E}{H} \right)_{T_{SM} \sim \frac{m}{4}} \propto Q^2 \left( \frac{M_{\text{Pl}}}{m/4} \right)$$

# Physics behind dark radiation production: Thermalized $\rho_{\text{DR}}$ insensitive to portal coupling

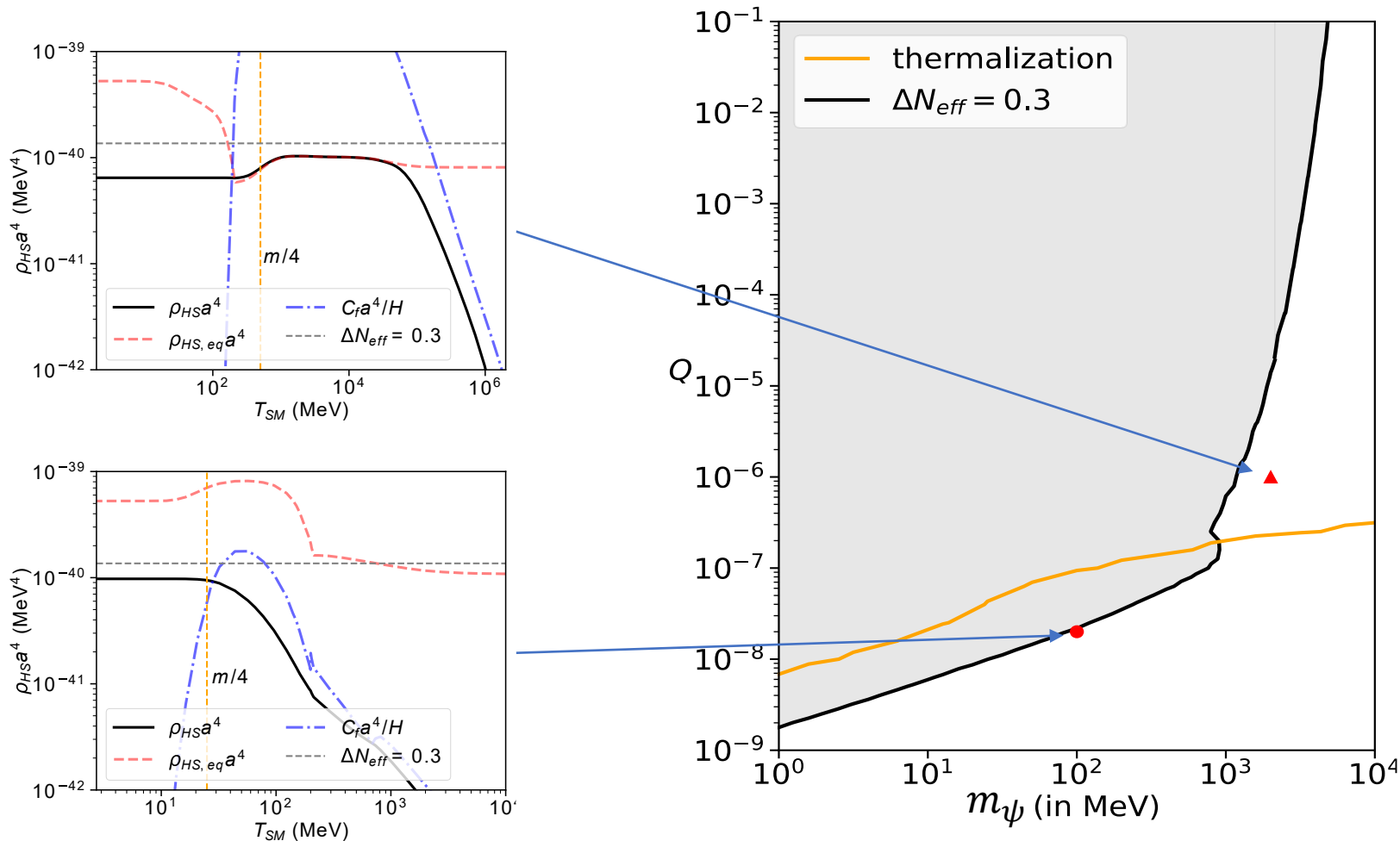


$$\frac{\rho_{HS}}{\rho_{SM}} \propto \frac{g_{HS}}{g_{*SM}(T_d)} \quad \begin{matrix} \rightarrow \\ \rightarrow \end{matrix} \text{Degrees of freedom}$$

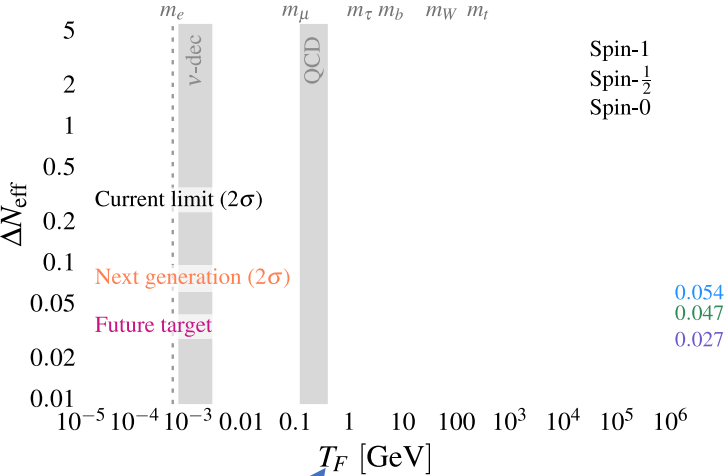
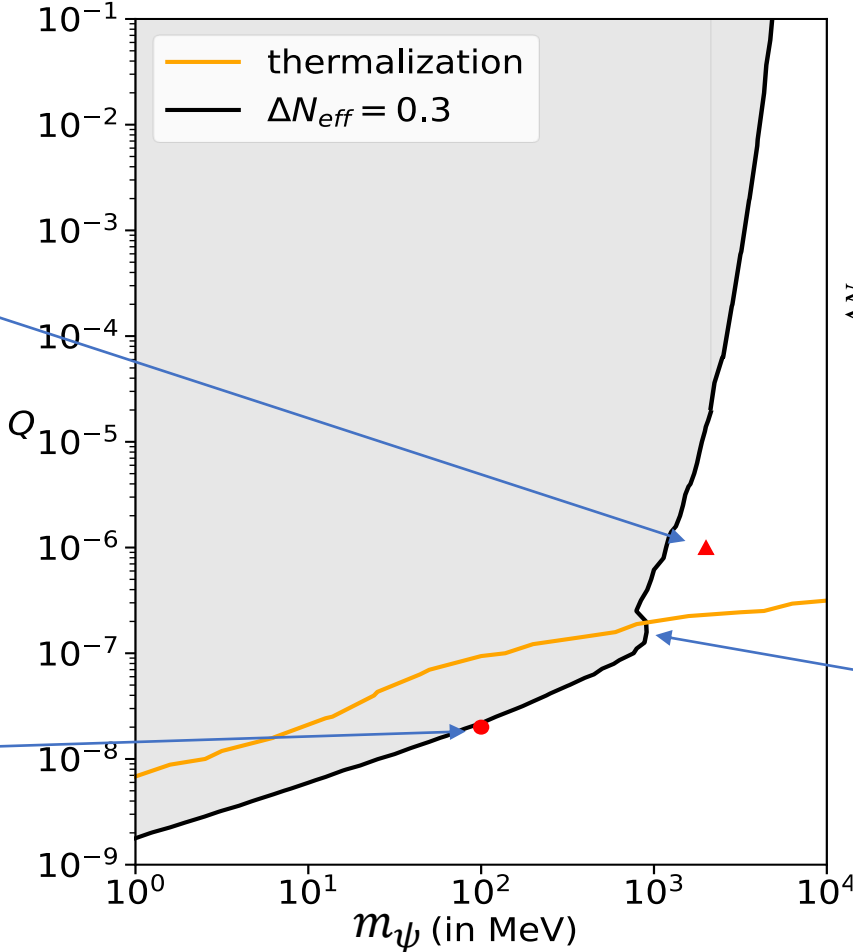
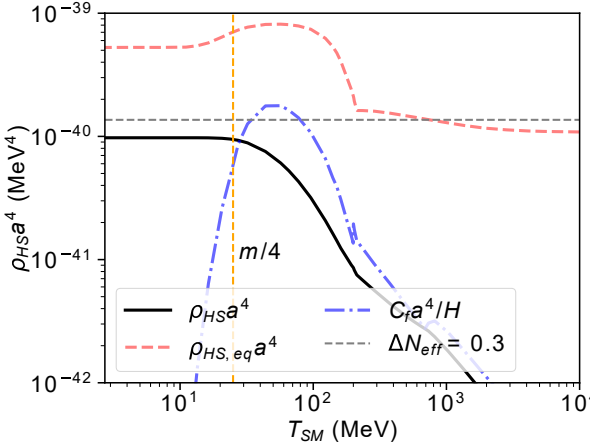
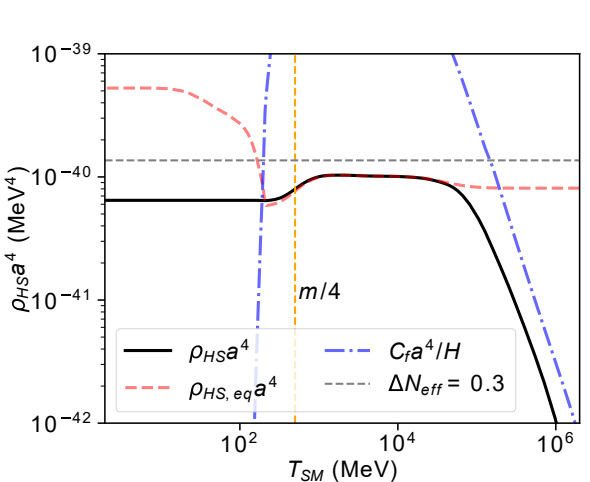
# Physics behind dark radiation production: Translating to constraints



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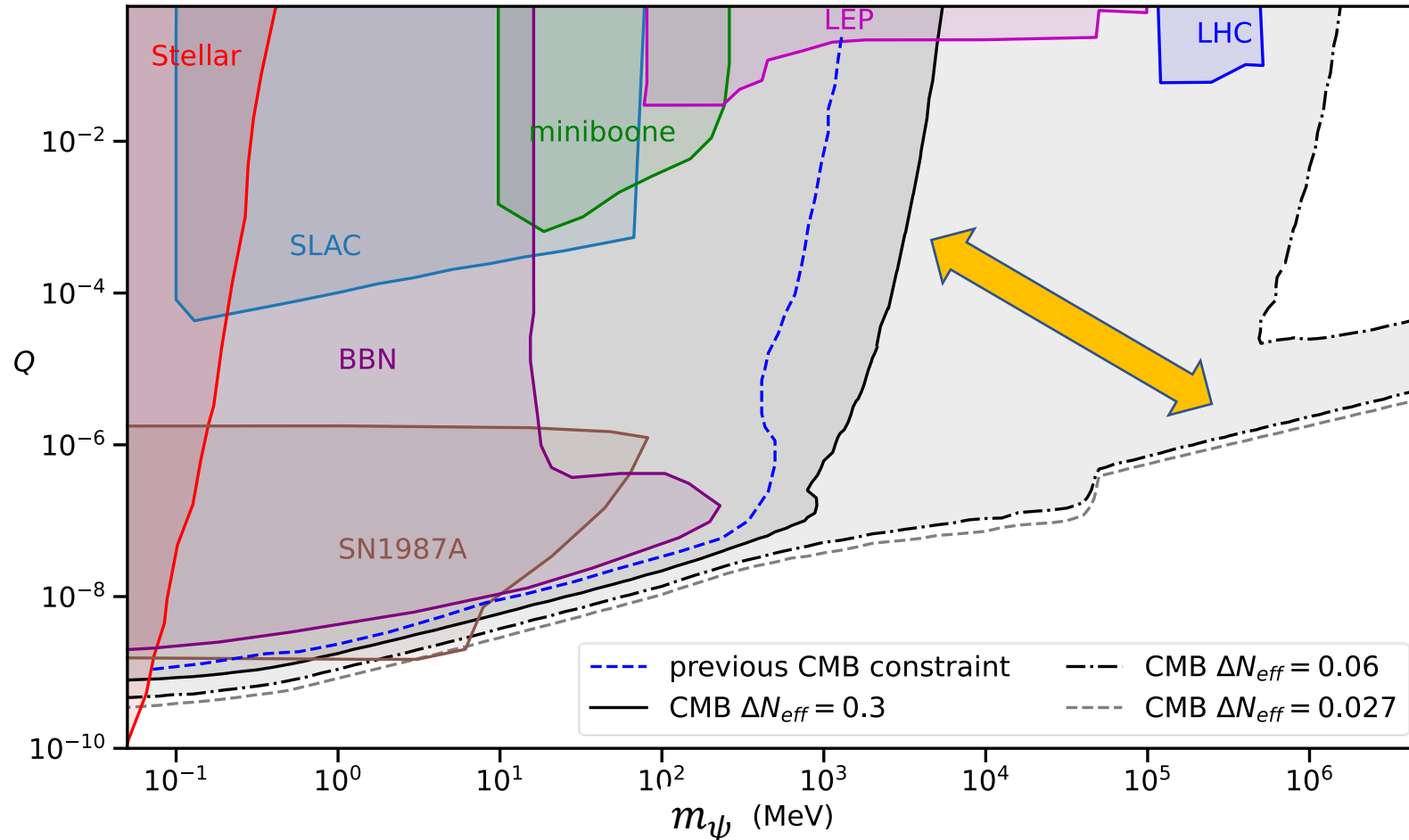


# Physics behind dark radiation production : Most relevant when thermally decoupled



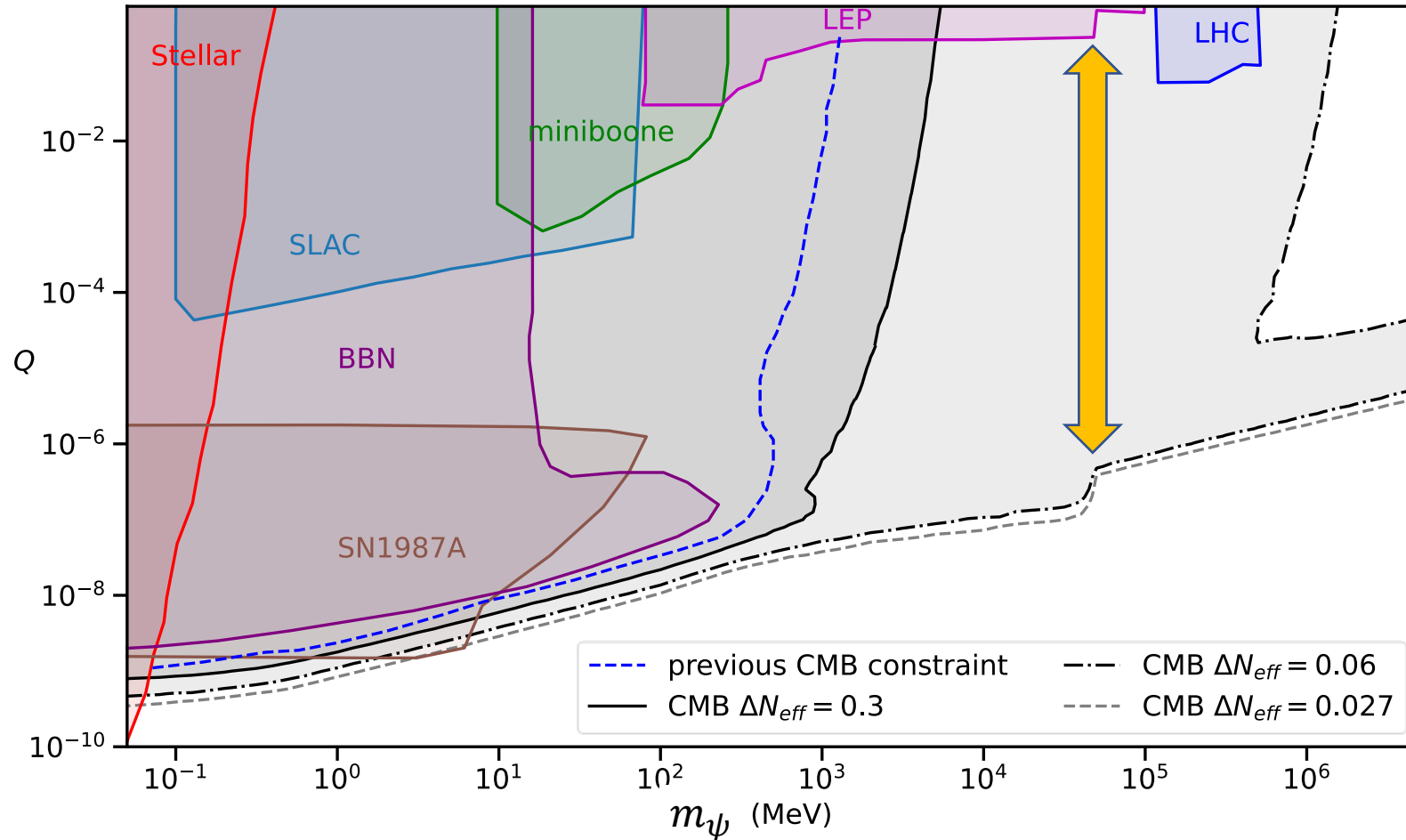
$m_{th} = 4T_{d,s}$

# Comparing $N_{eff}$ constraints: Future constraint will extend to much larger parameter space

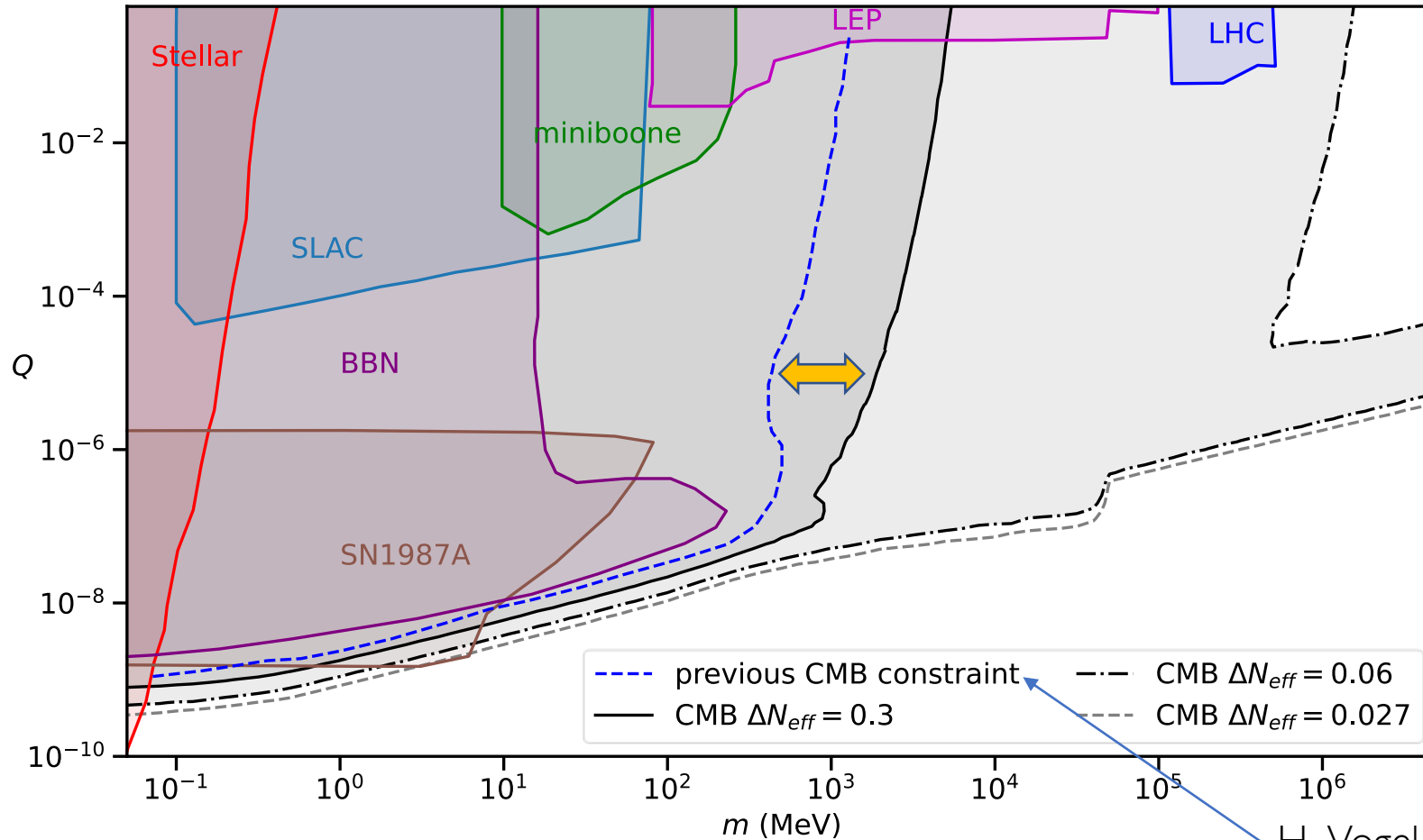




# Comparing Neff constraints: Dominant for $M_\psi > 0.1$ MeV

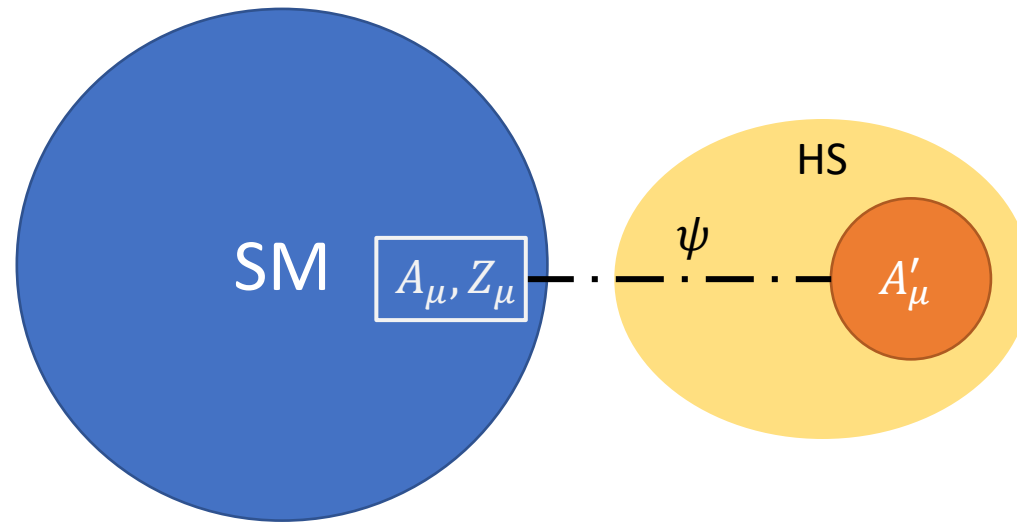


# Comparing Neff constraints: Updating previous constraint

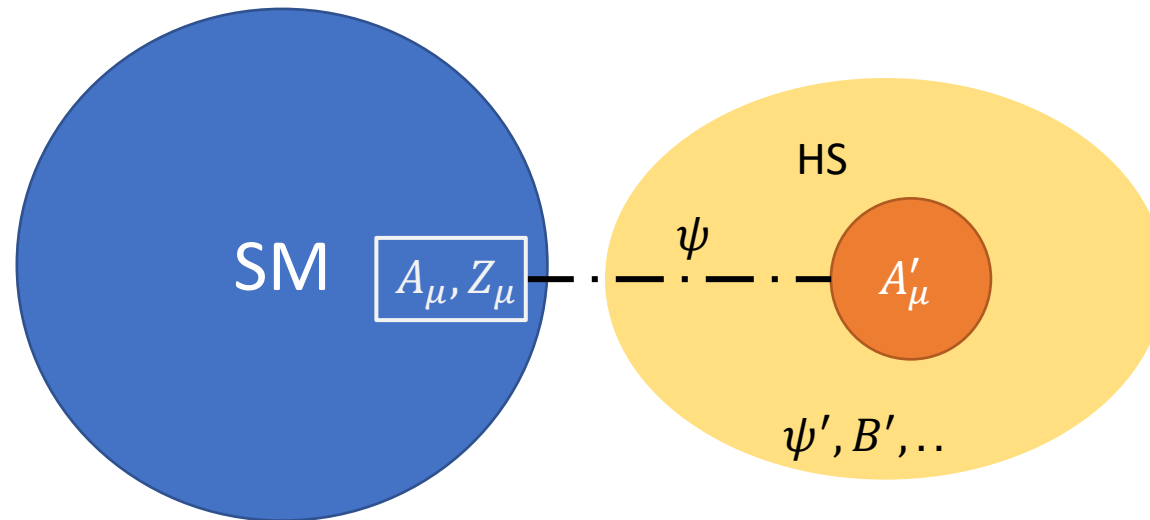


H. Vogel and J. Redondo, JCAP 02 (2014) 029.

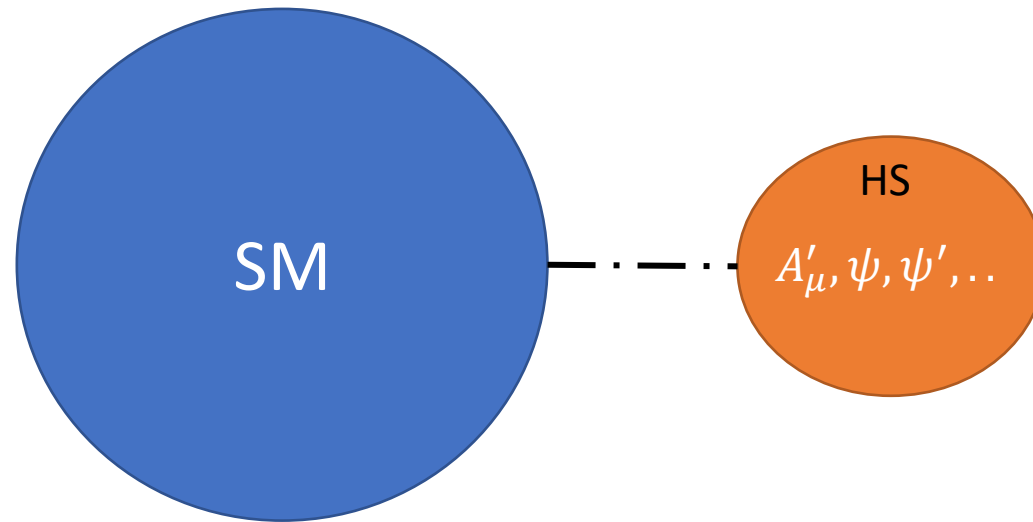
# Extending to general hidden sector



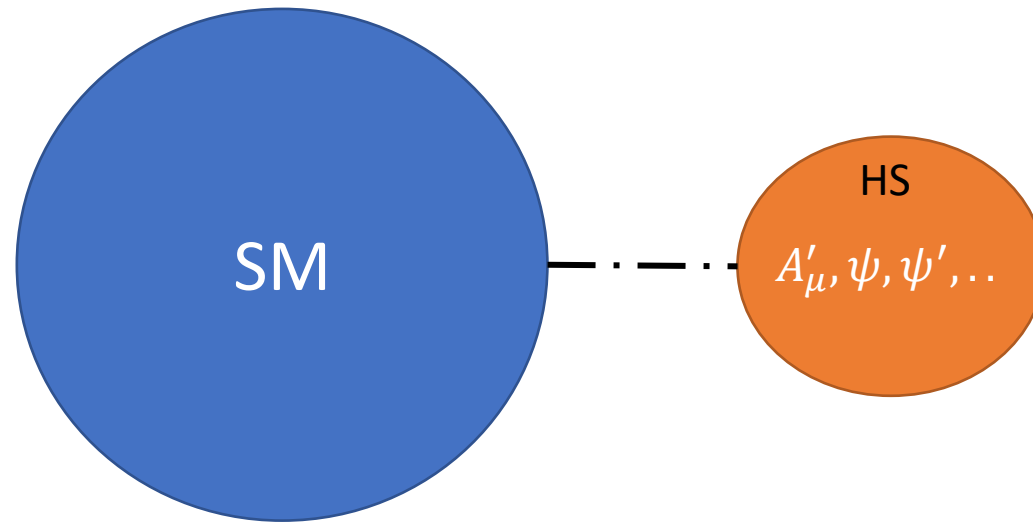
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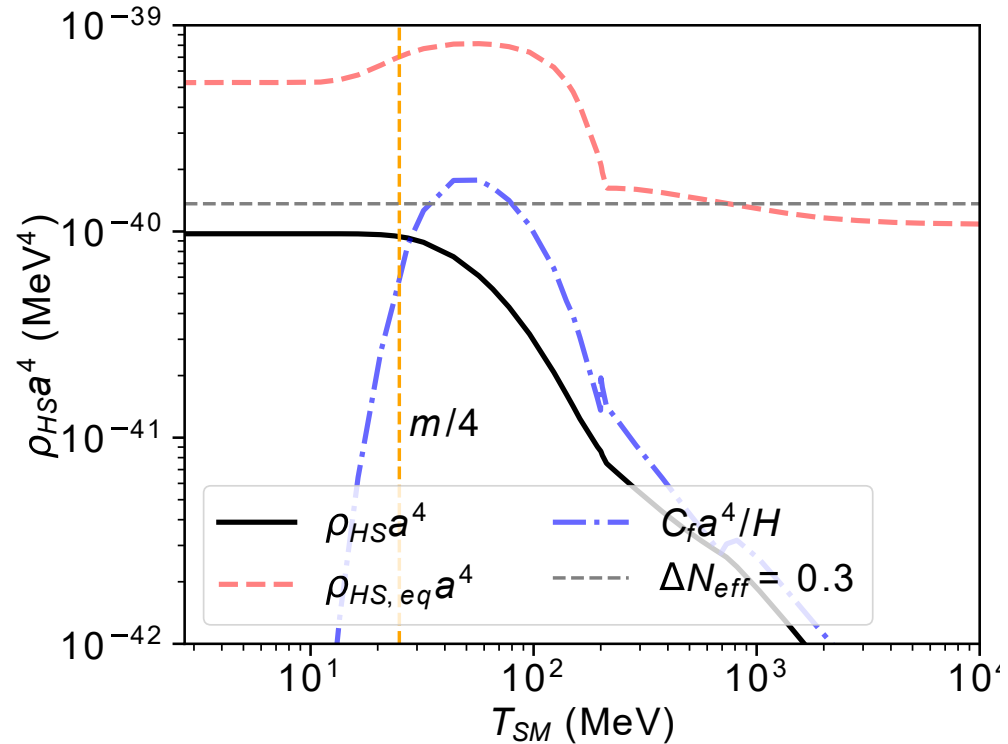
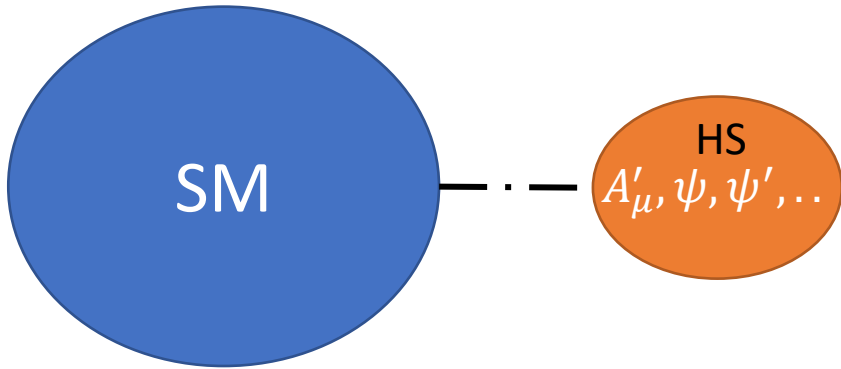


# Extending to general hidden sector

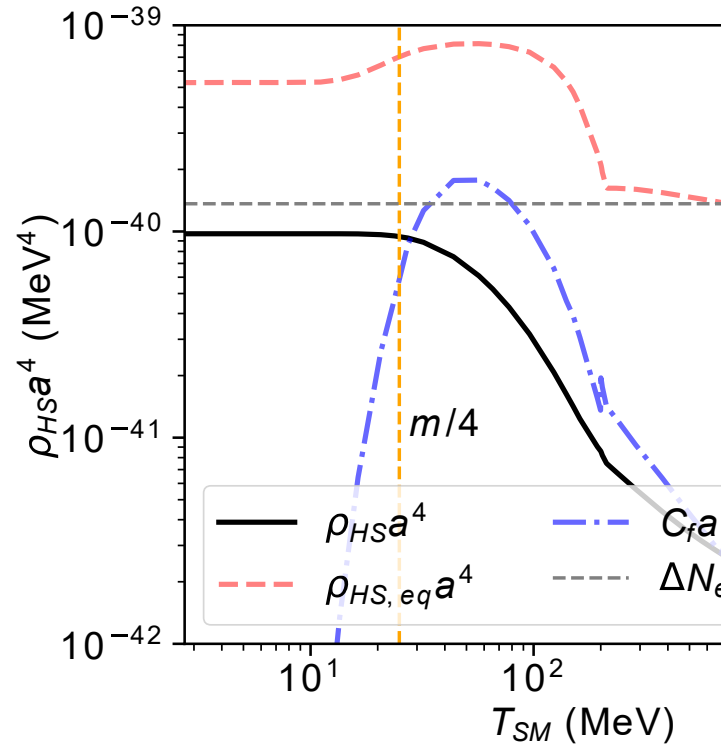
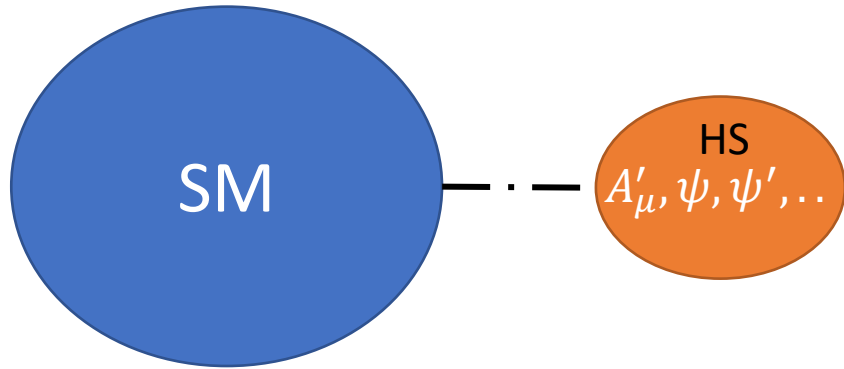


One can calculate a conservative  $N_{\text{eff}}$  constraint on the millicharge interaction that is independent of details of hidden sector.

# Minimum leaked energy independent of details within HS

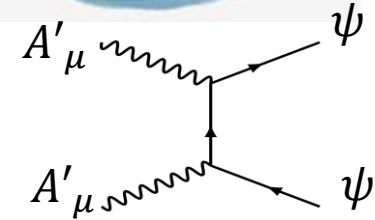
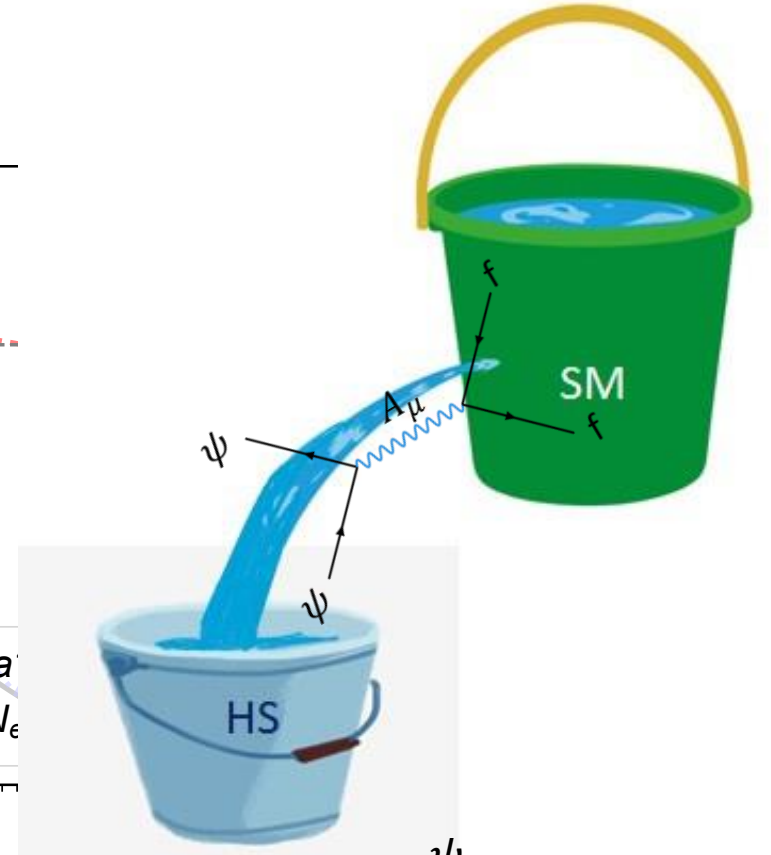
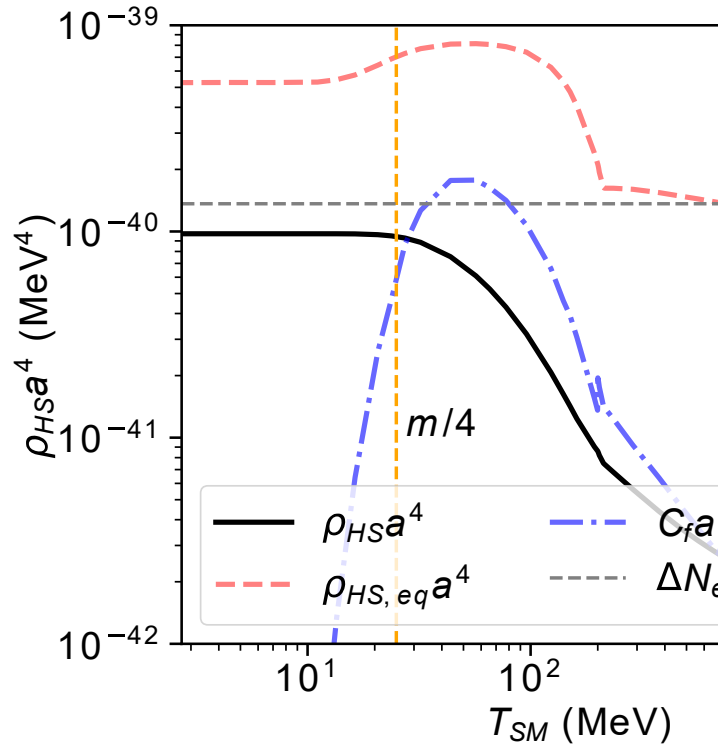
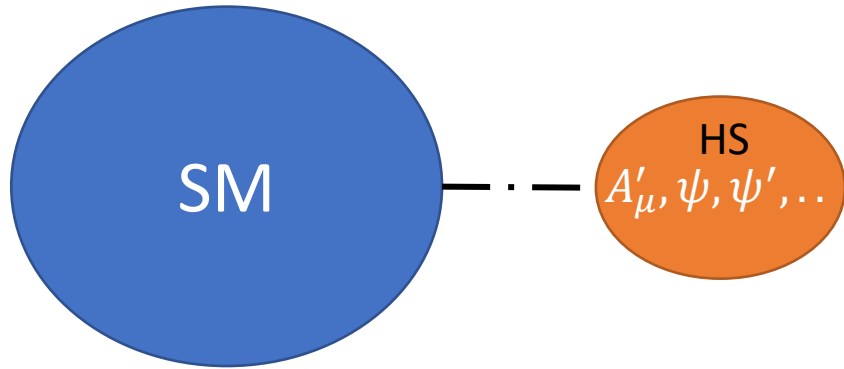


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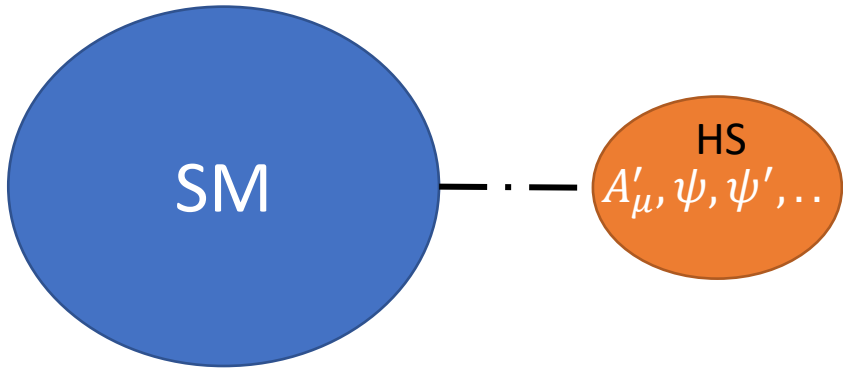




# Minimum leaked energy independent of details within HS

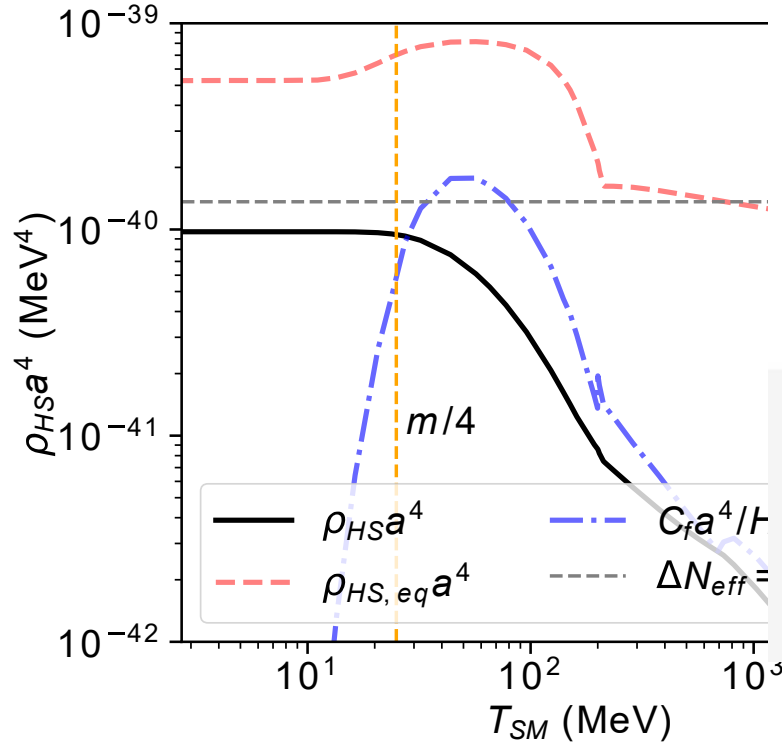


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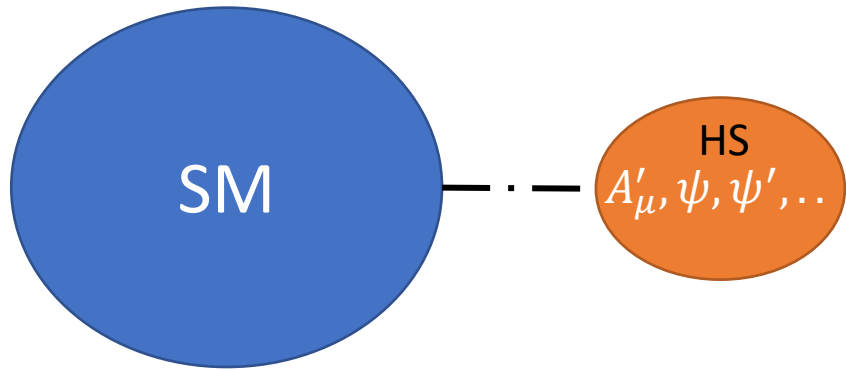


$$\left(\frac{\rho_{HS}}{\rho_{SM}}\right)_{\min} \propto \frac{M_{\text{Pl}}}{m_\psi} L$$

Leak factor  $L = m_\psi \int ds \frac{s - 4m_f^2}{s\sqrt{s}} \sigma_{ff \rightarrow \psi\psi} \propto Q^2$



# Minimum leaked energy independent of details within HS

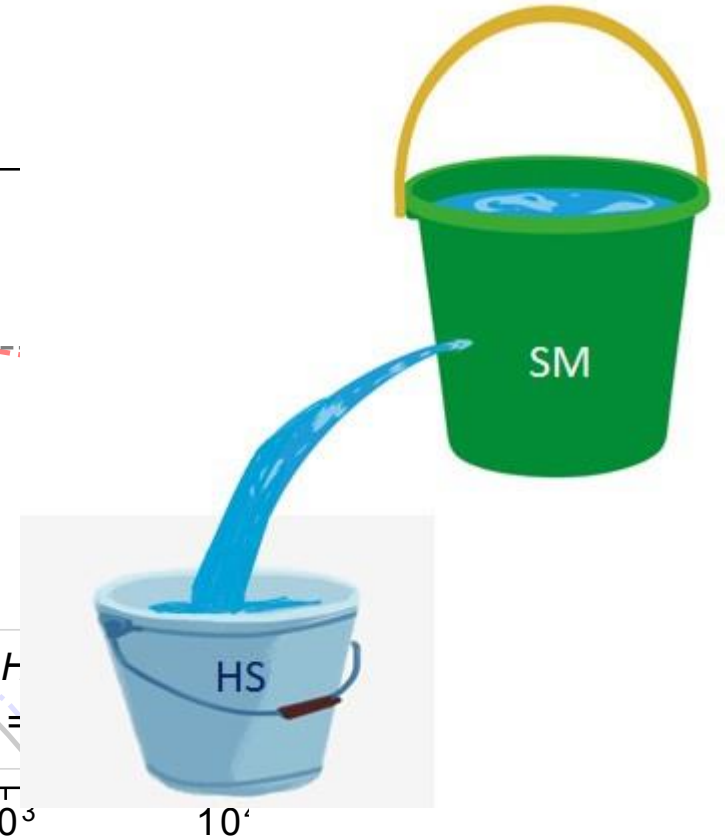
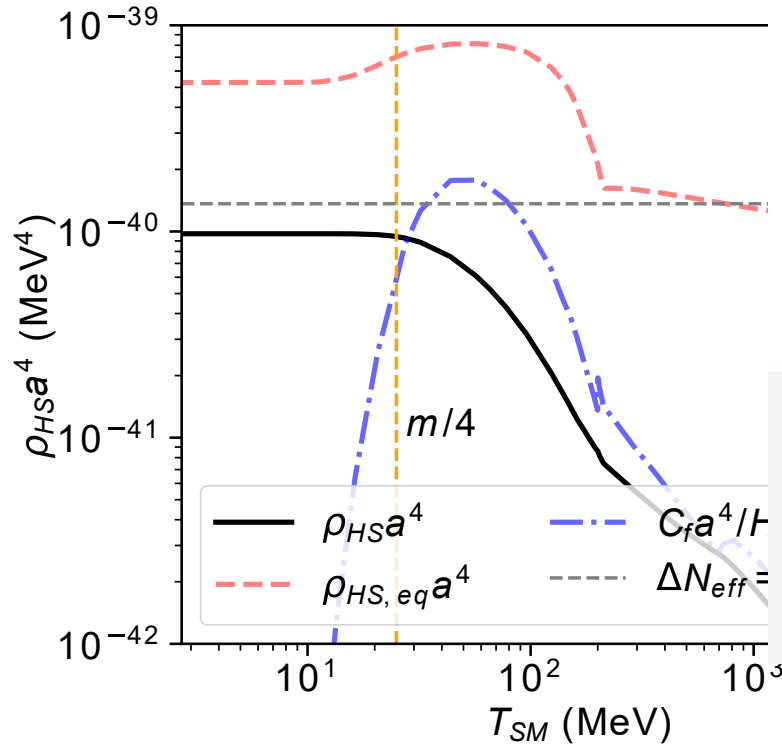


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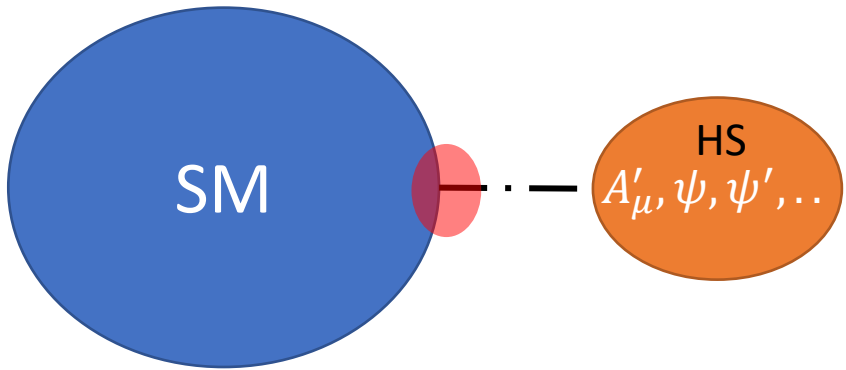
Leak factor  $\uparrow$

$$L = m_\psi \int ds \frac{s - 4m_f^2}{s\sqrt{s}} \sigma_{ff \rightarrow \psi\psi} \propto Q^2$$

Independent of  $g_{HS}$  or  $T_{HS}$  or value of dark charge.



# Minimum leaked energy independent of details within HS: Depends only on one BSM coupling

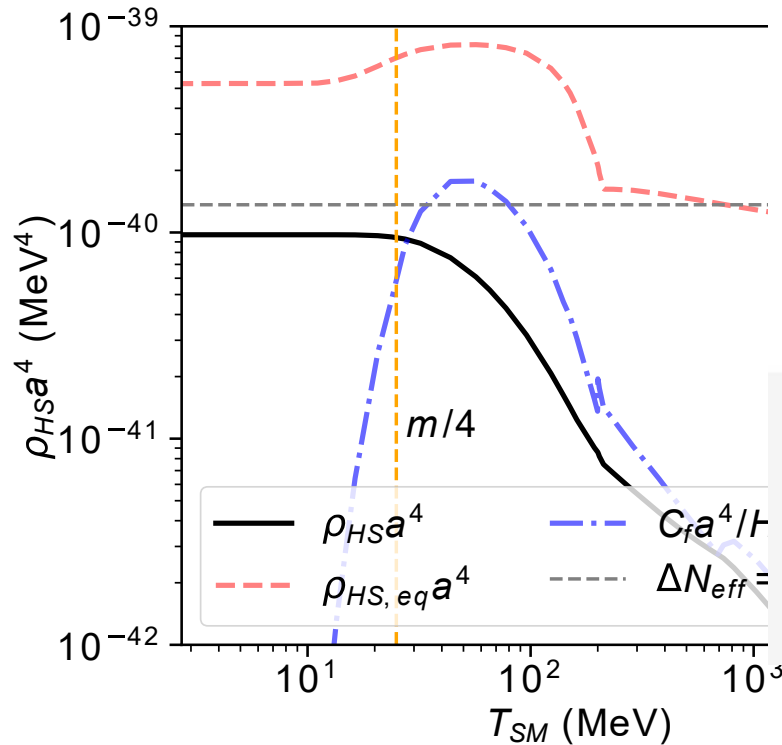


$$\left(\frac{\rho_{HS}}{\rho_{SM}}\right)_{\min} \propto \frac{M_{\text{Pl}}}{m_\psi} L$$

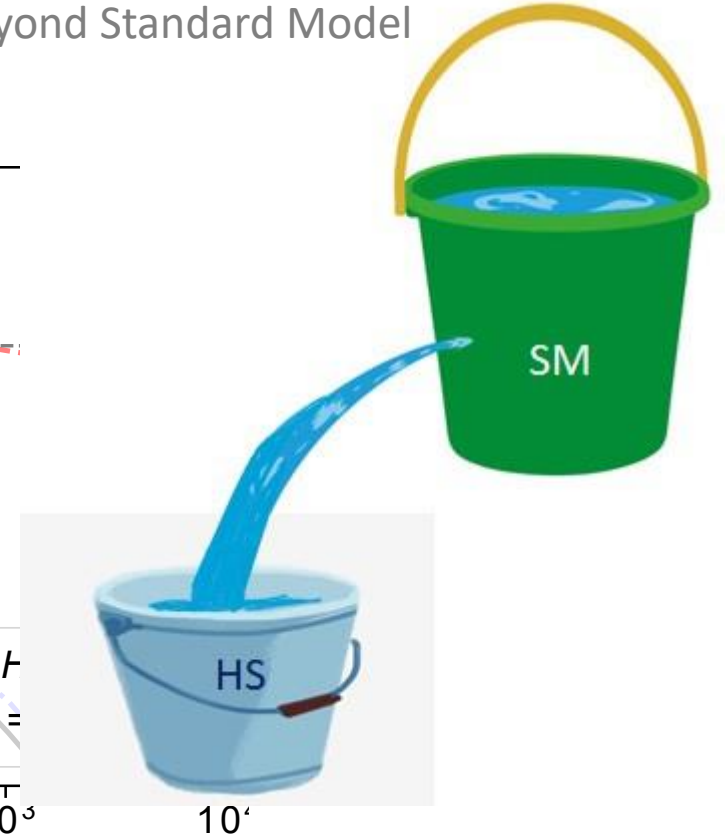
$$L = m_\psi \int ds \frac{s - 4m_f^2}{s\sqrt{s}} \sigma_{ff \rightarrow \psi\psi} \propto Q^2$$

Leak factor

Independent of  $g_{HS}$  or  $T_{HS}$  or value of dark charge.

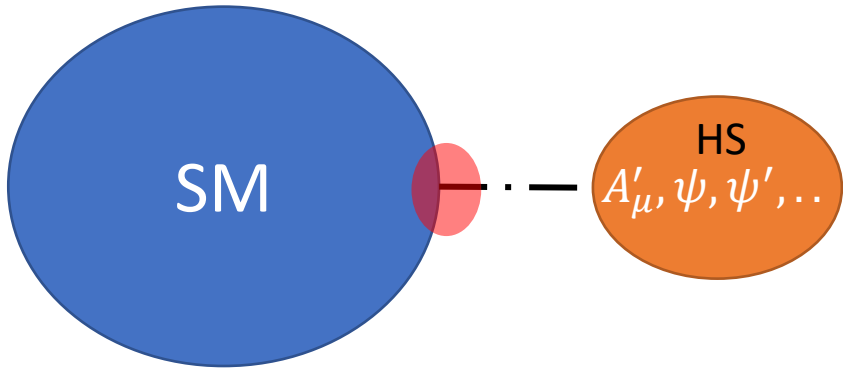


Beyond Standard Model



# Minimum leaked energy independent of details within HS: Conservative constraint on BSM coupling

Beyond Standard Model

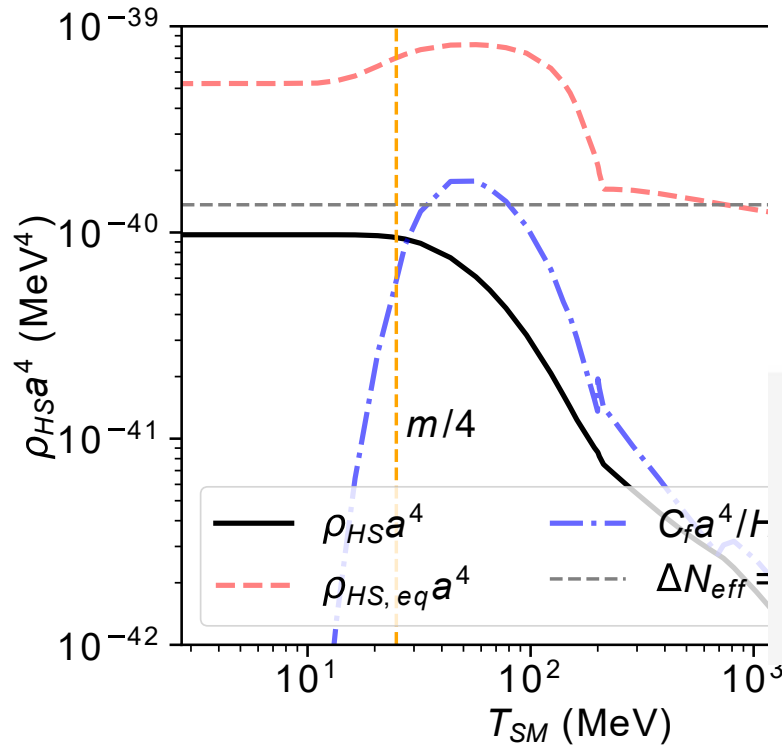


$$\left(\frac{\rho_{HS}}{\rho_{SM}}\right)_{\min} \propto \frac{M_{\text{Pl}}}{m_\psi} L$$

$$L = m_\psi \int ds \frac{s - 4m_f^2}{s\sqrt{s}} \sigma_{ff \rightarrow \psi\psi} \propto Q^2$$

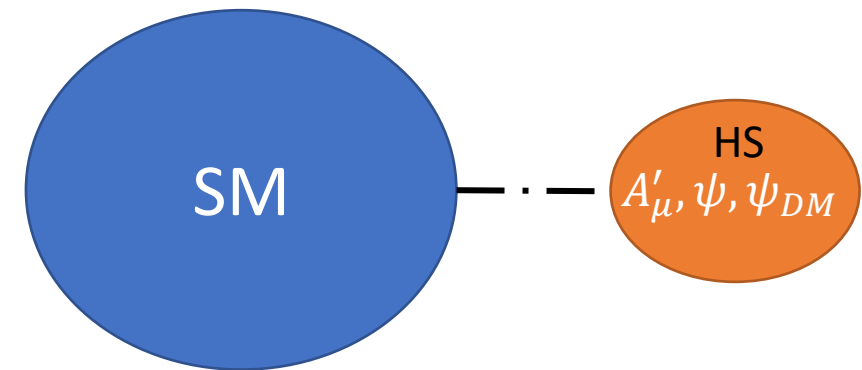
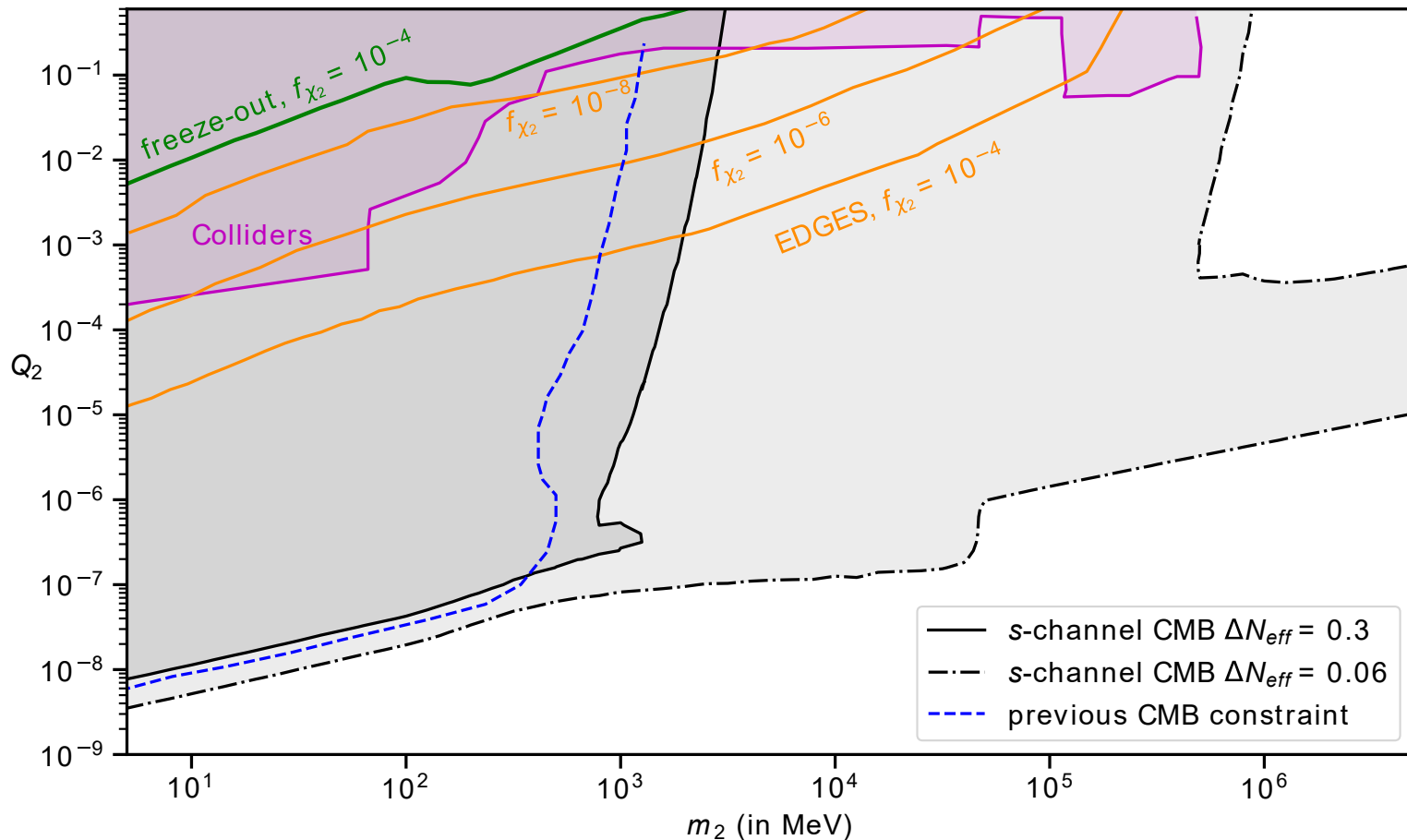
Leak factor

Independent of  $g_{HS}$  or  $T_{HS}$  or value of dark charge.



$$L < g_*^{3/4}(m) g_*^{1/3} \left(\frac{m}{4}\right) \frac{m/4}{M_{\text{Pl}}} (\Delta N_{\text{eff}})_{\text{max}}$$

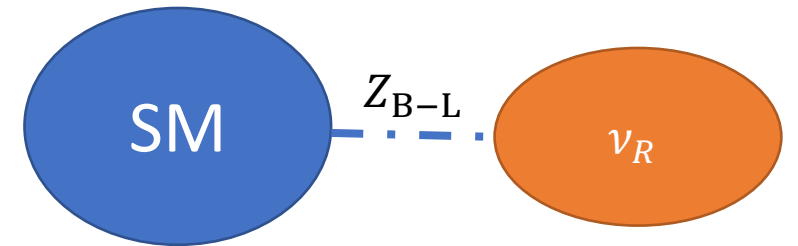
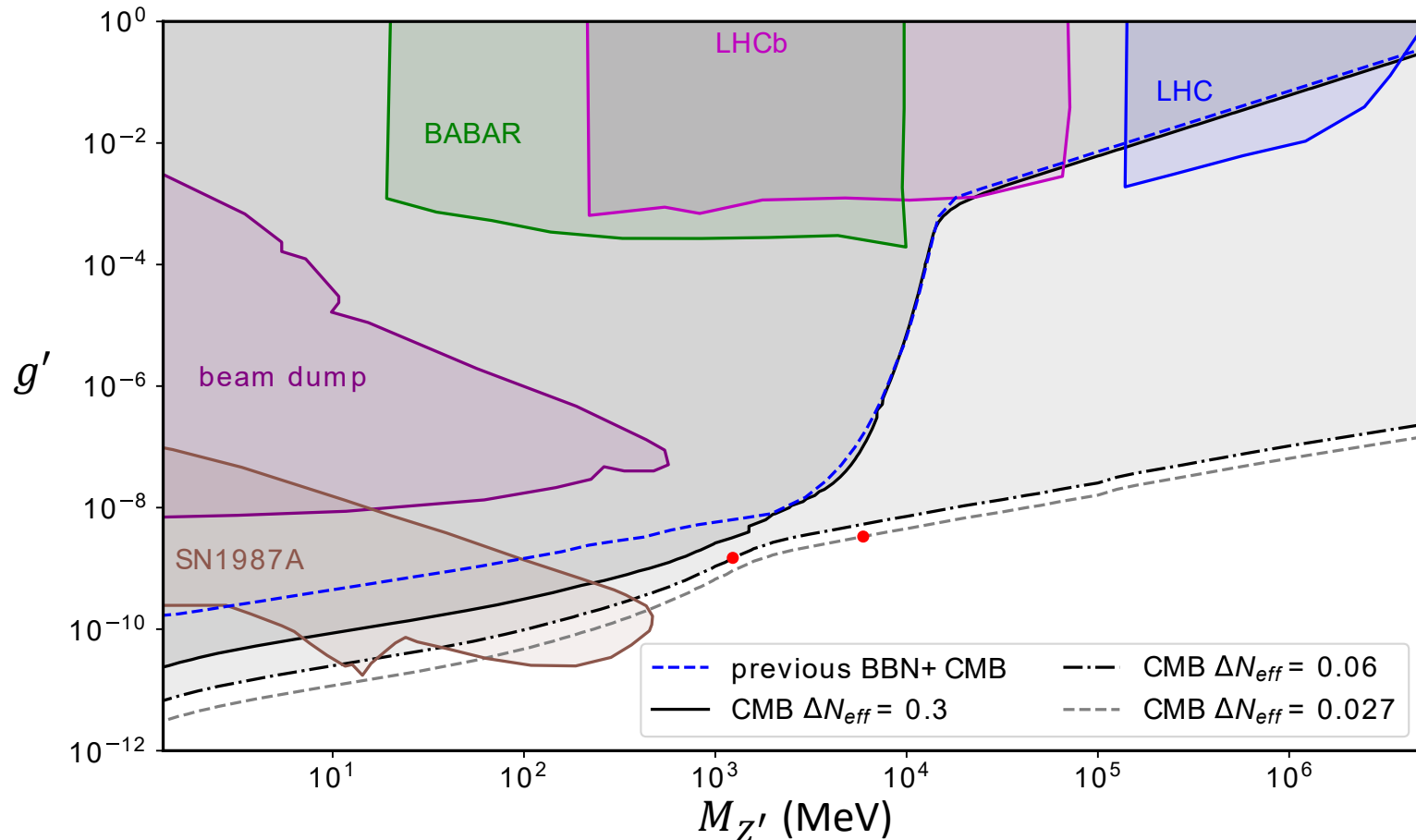
# Neff constraints applicable for wide class of hidden sectors: Application to EDGES



H. Liu, N. J. Outmezguine, D. Redigolo, and T. Volansky, Phys. Rev. D 100 no. 12, (2019) 123011.

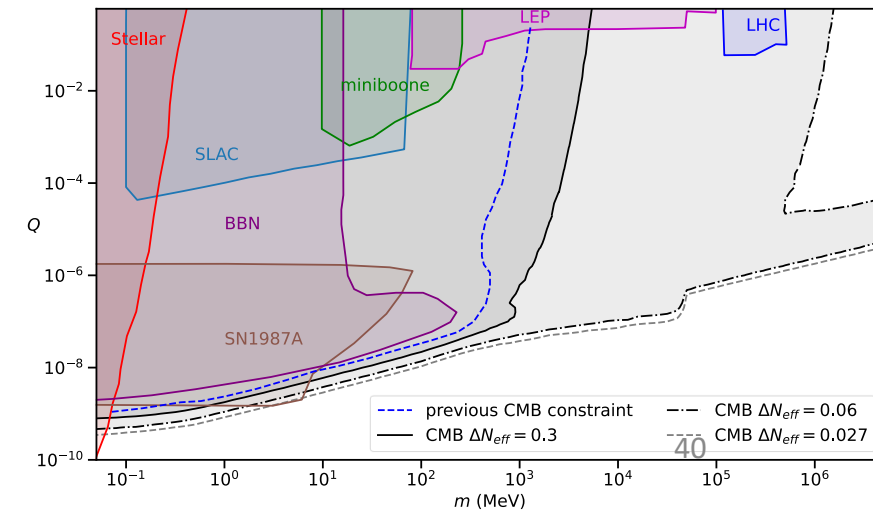
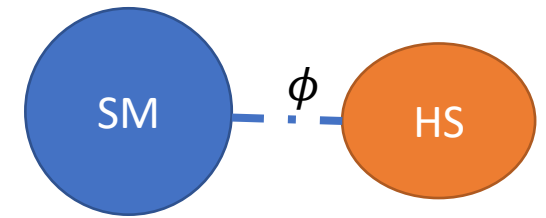
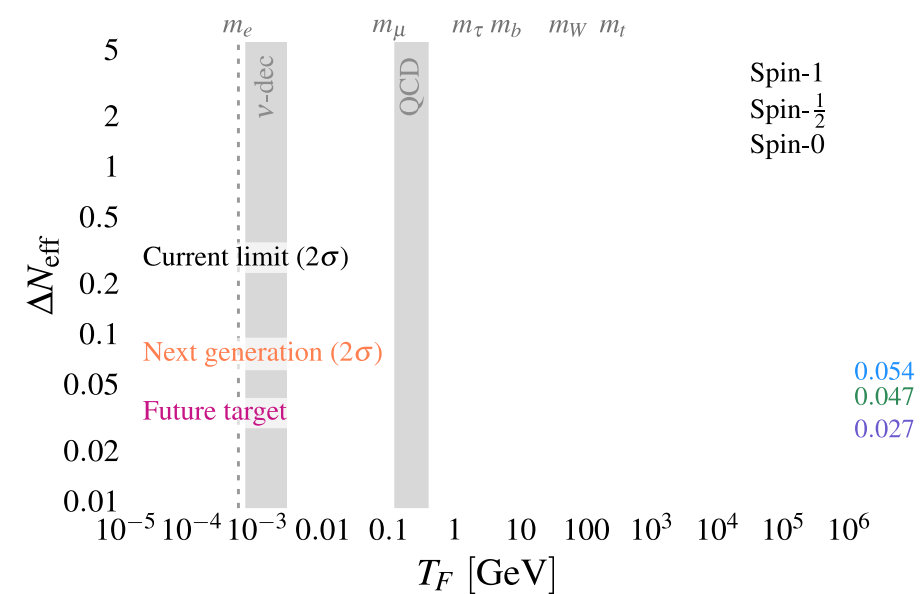
# Neff constraints applicable for wide class of hidden sectors: B-L model

$$L_{int} \supset -\frac{1}{4}F'_{\mu\nu}F^{\mu\nu'} + g'Z'_\mu J_{B-L,SM}^\mu - g'Z'_\mu \sum_i \bar{\nu}_{R,i}\gamma^\mu\nu_{R,i} + \frac{1}{2}M_{Z'}^2 Z'^\mu Z'_\mu$$



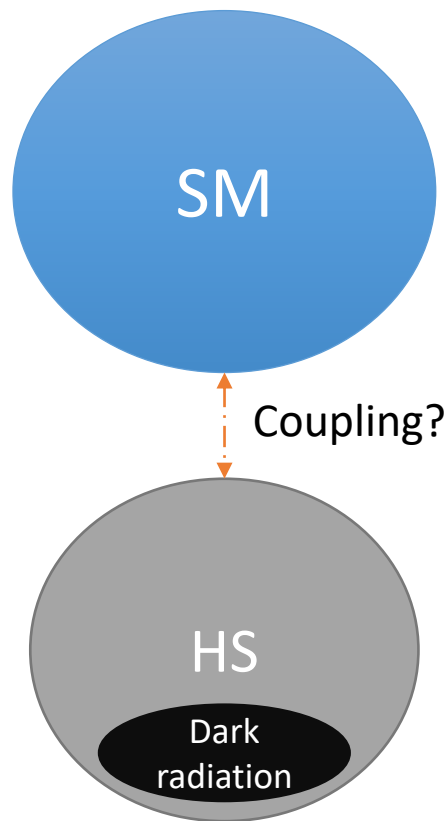
# Summary

- With improving  $N_{\text{eff}}$  measurements, we should interpret them as constraints on portal interactions with out-of-equilibrium sectors
- $N_{\text{eff}}$  constraints on out-of-equilibrium particles are:
  - Most relevant for portal interactions mediated by a particle heavier than 0.1 MeV
  - Orders of magnitude stronger than collider experiments
  - Constraints largely independent of internal hidden sector model
- Simple way to calculate:  $L < g_*^{3/4} (4\Lambda) g_*^{1/3} (\Lambda) \frac{\Lambda}{M_{\text{Pl}}} (\Delta N_{\text{eff}})_{\text{max}}$



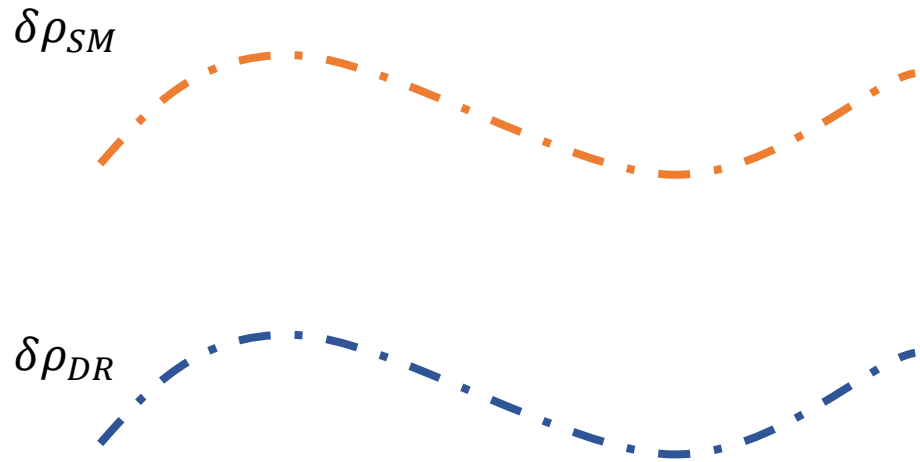


# Goal: Probe models containing dark radiation using cosmological observations



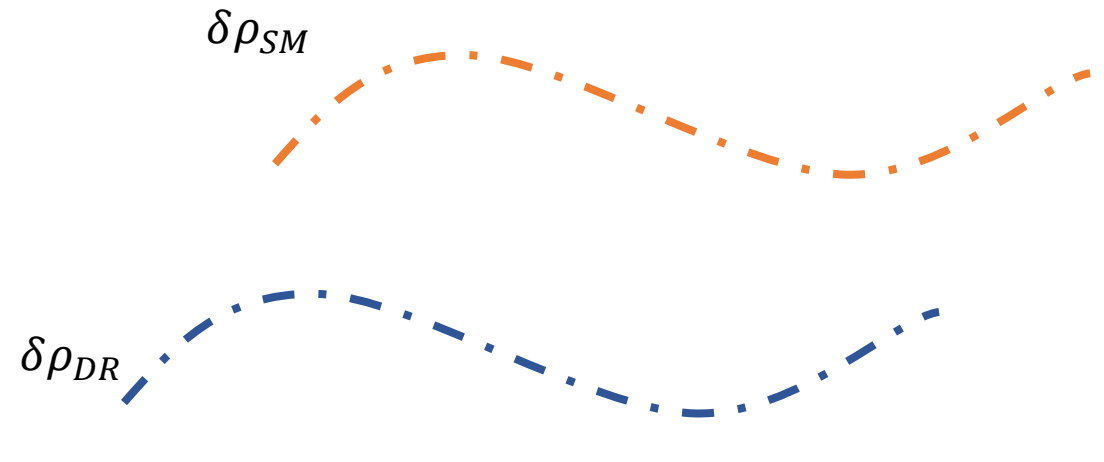
- If no dark radiation detected in future:  
Constrain portal interactions with hidden sectors  
(Adshead, Ralegankar, Shelton 2022)
- If dark radiation detected: Search for interaction by checking adiabaticity of perturbations during BBN  
(Adshead, Holder, Ralegankar 2020)

# Adiabaticity: 'sync' between perturbations



Adiabatic perturbations

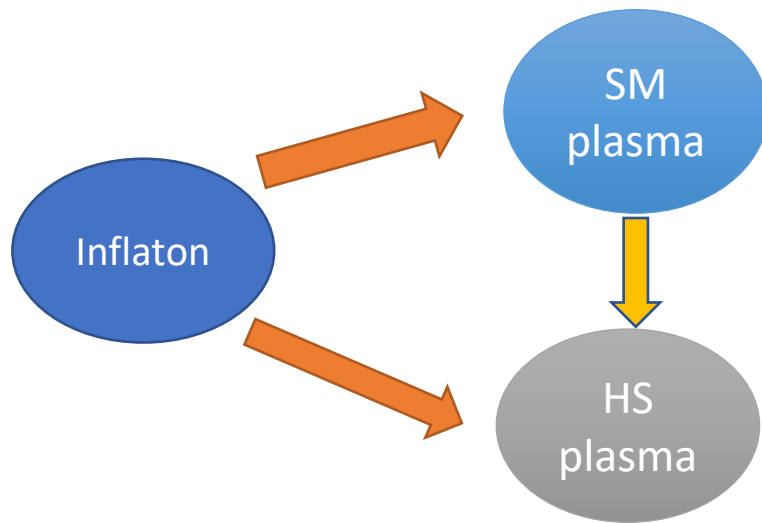
$$\frac{\delta\rho_{SM}}{\rho_{SM}} = \frac{\delta\rho_{DR}}{\rho_{DR}}$$



Isocurvature component in perturbations

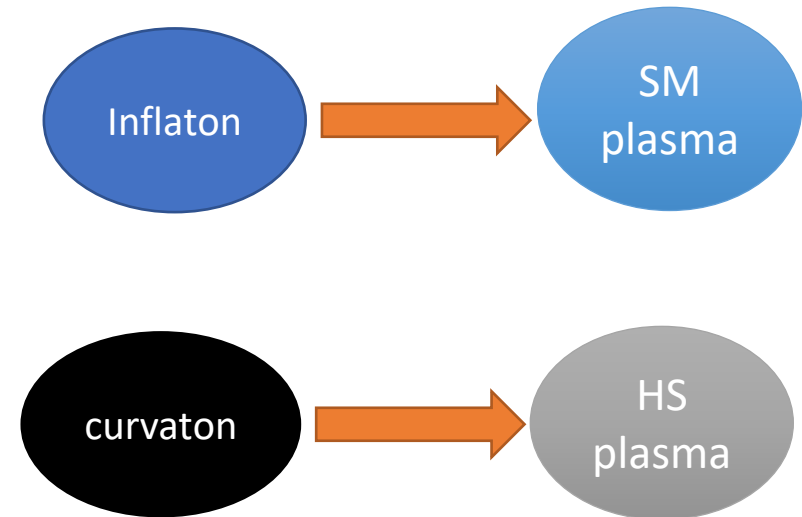
$$S_{DR} = \frac{4}{3} \left( \frac{\delta\rho_{DR}}{\rho_{DR}} - \frac{\delta\rho_{SM}}{\rho_{SM}} \right)$$

# Adiabaticity: Related to early universe production



Adiabatic perturbations

$$\frac{\delta\rho_{SM}}{\rho_{SM}} = \frac{\delta\rho_{DR}}{\rho_{DR}}$$

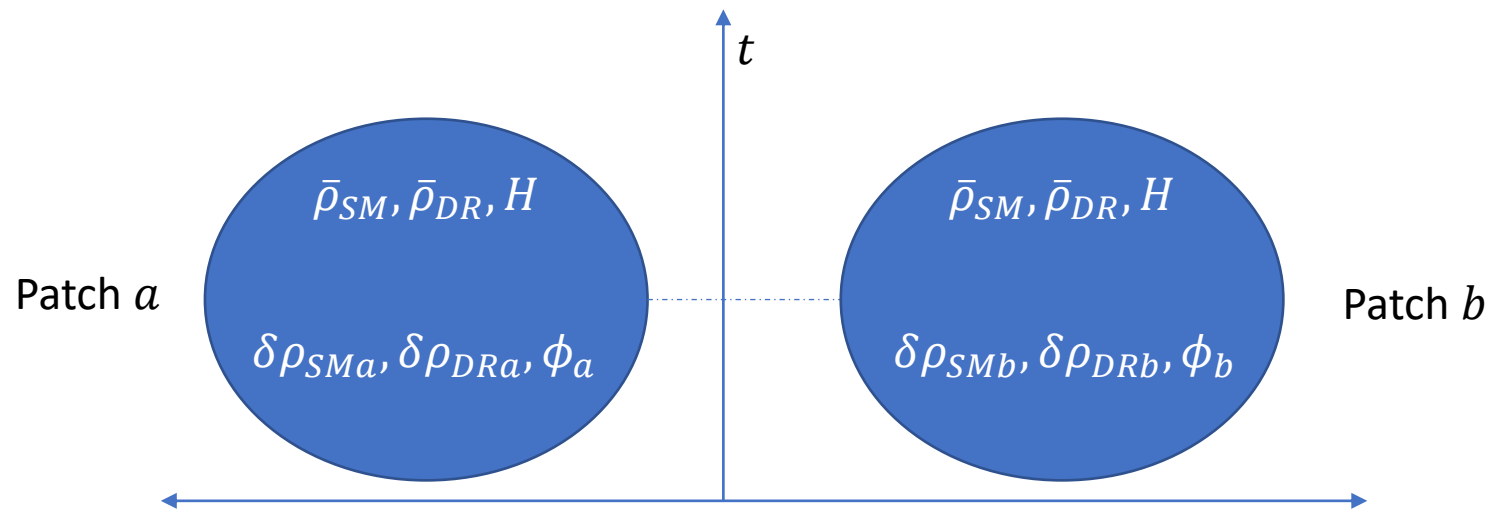
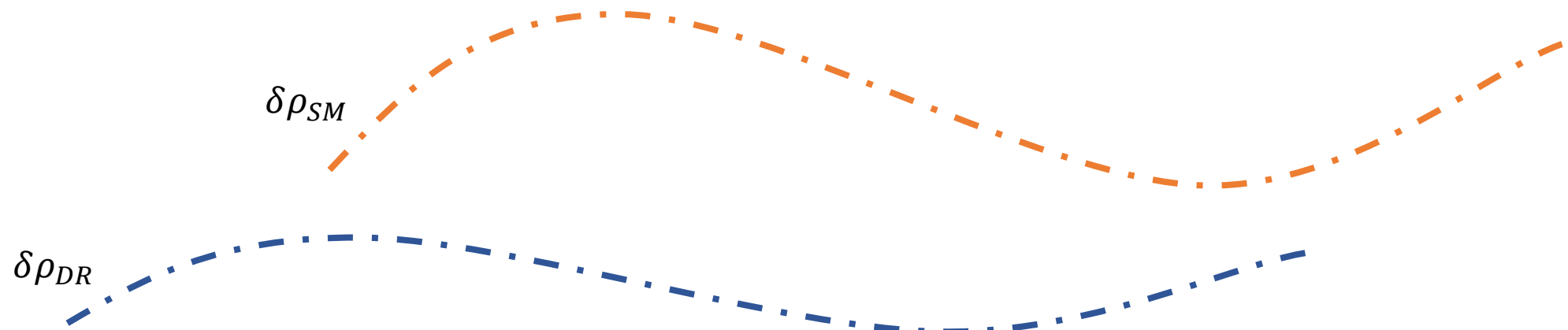


Isocurvature component in perturbations

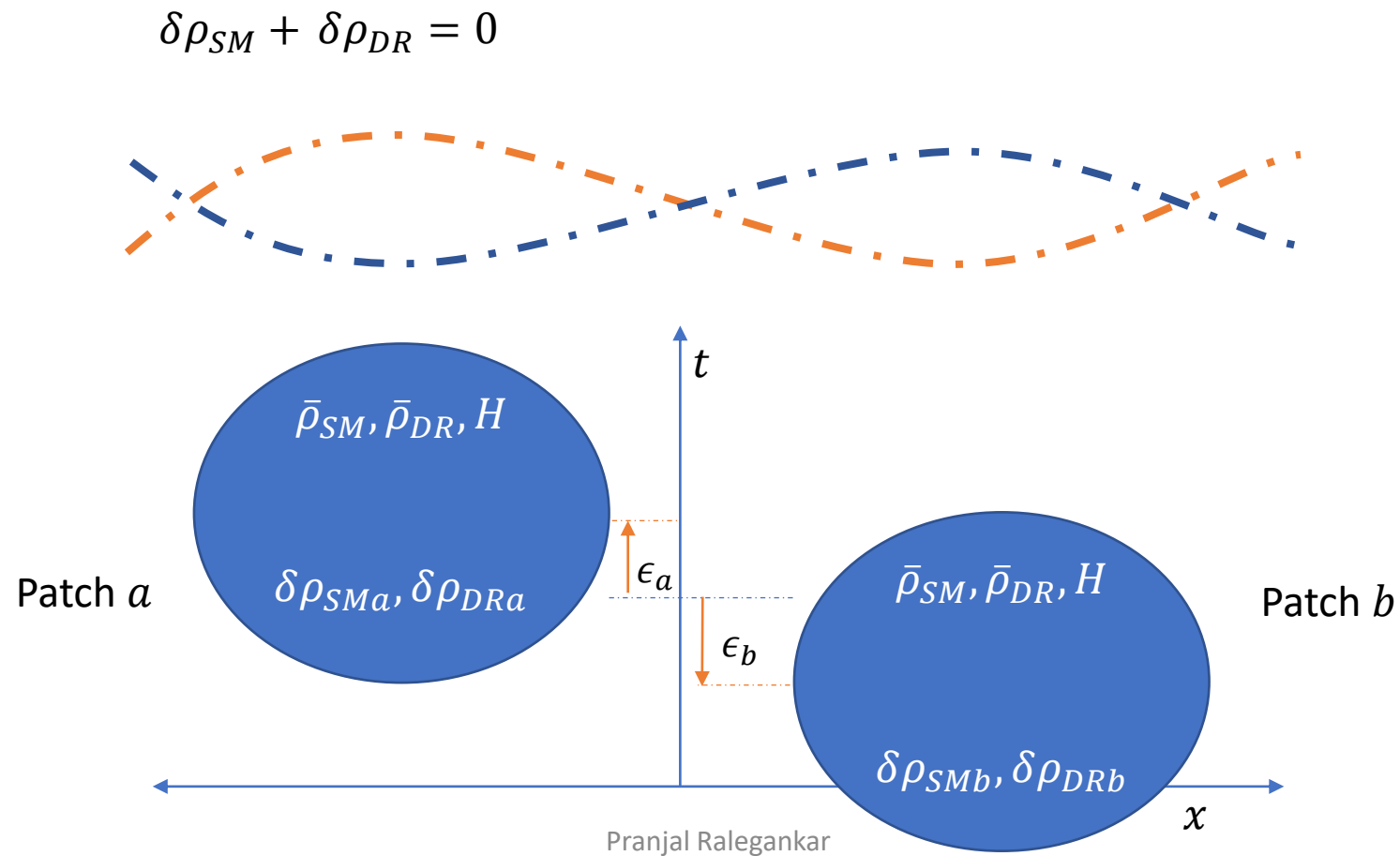
$$S_{DR} = \frac{4}{3} \left( \frac{\delta\rho_{DR}}{\rho_{DR}} - \frac{\delta\rho_{SM}}{\rho_{SM}} \right)$$

Goal: Use BBN observations to find dark radiation isocurvature

# Dark radiation isocurvature => Variation in $N_{\text{eff}}$

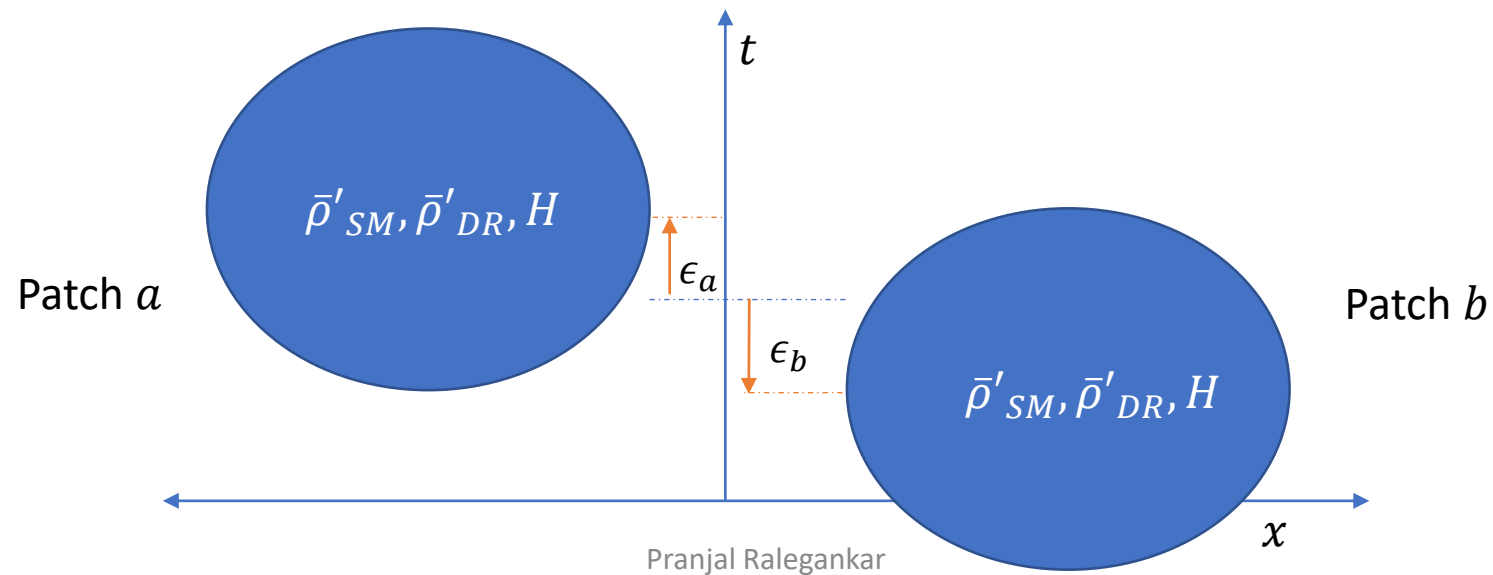


# Dark radiation isocurvature => Variation in $N_{\text{eff}}$



# Dark radiation isocurvature => Variation in $N_{\text{eff}}$

$$\bar{\rho}' = \bar{\rho} + \delta\rho$$

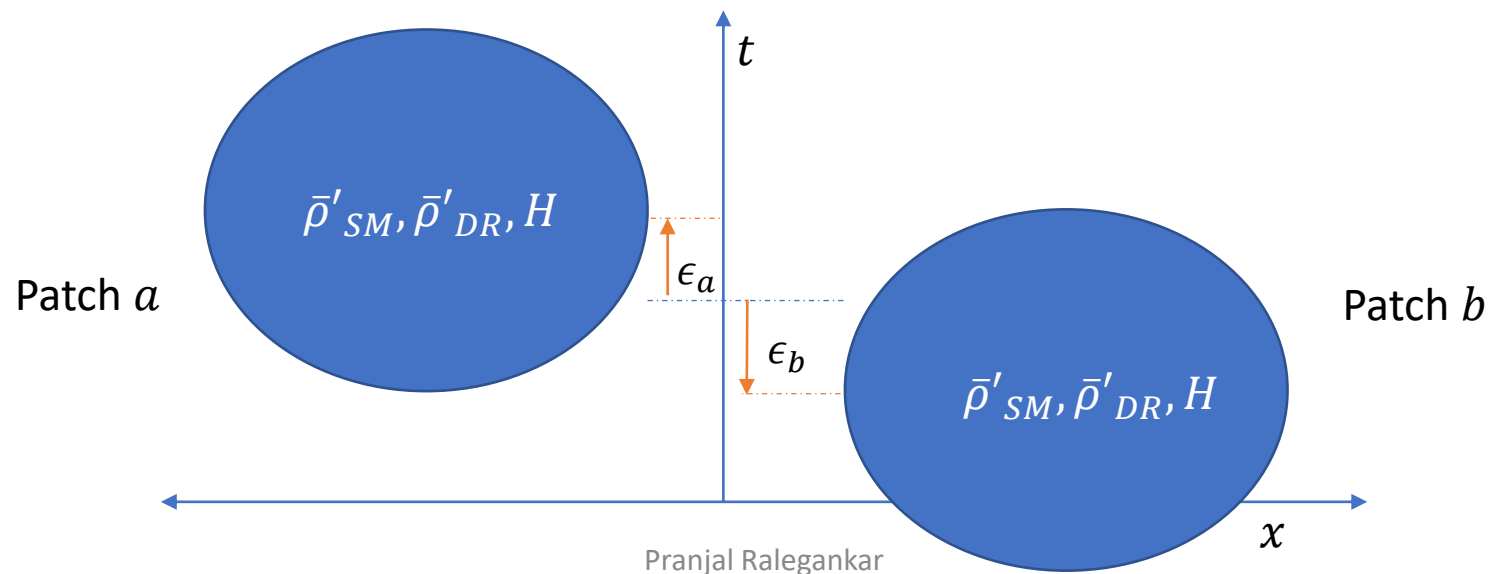


# Dark radiation isocurvature => Variation in $N_{\text{eff}}$

$$\bar{\rho}' = \bar{\rho} + \delta\rho$$

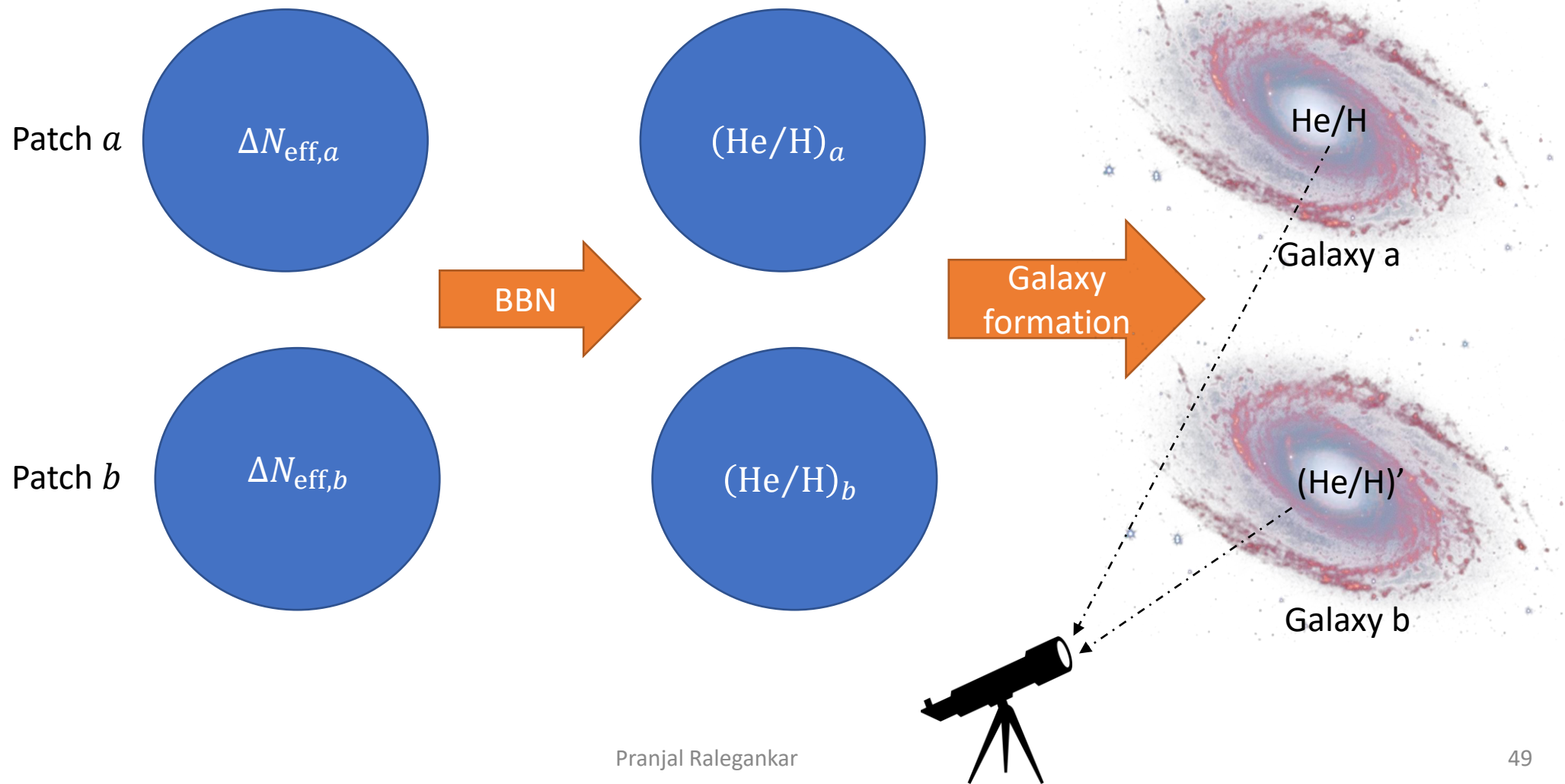
$$N_{\text{eff}} \propto \frac{\bar{\rho}'_{SM}}{\bar{\rho}'_{DR}} \quad \longrightarrow \quad \Delta N_{\text{eff}}(\vec{x}) = \Delta \bar{N}_{\text{eff}} \left( 1 + \frac{4}{3} S_{DR}(\vec{x}) \right)$$

$$S_{DR} = \frac{4}{3} \left( \frac{\delta\rho_{DR}}{\rho_{DR}} - \frac{\delta\rho_{SM}}{\rho_{SM}} \right)$$

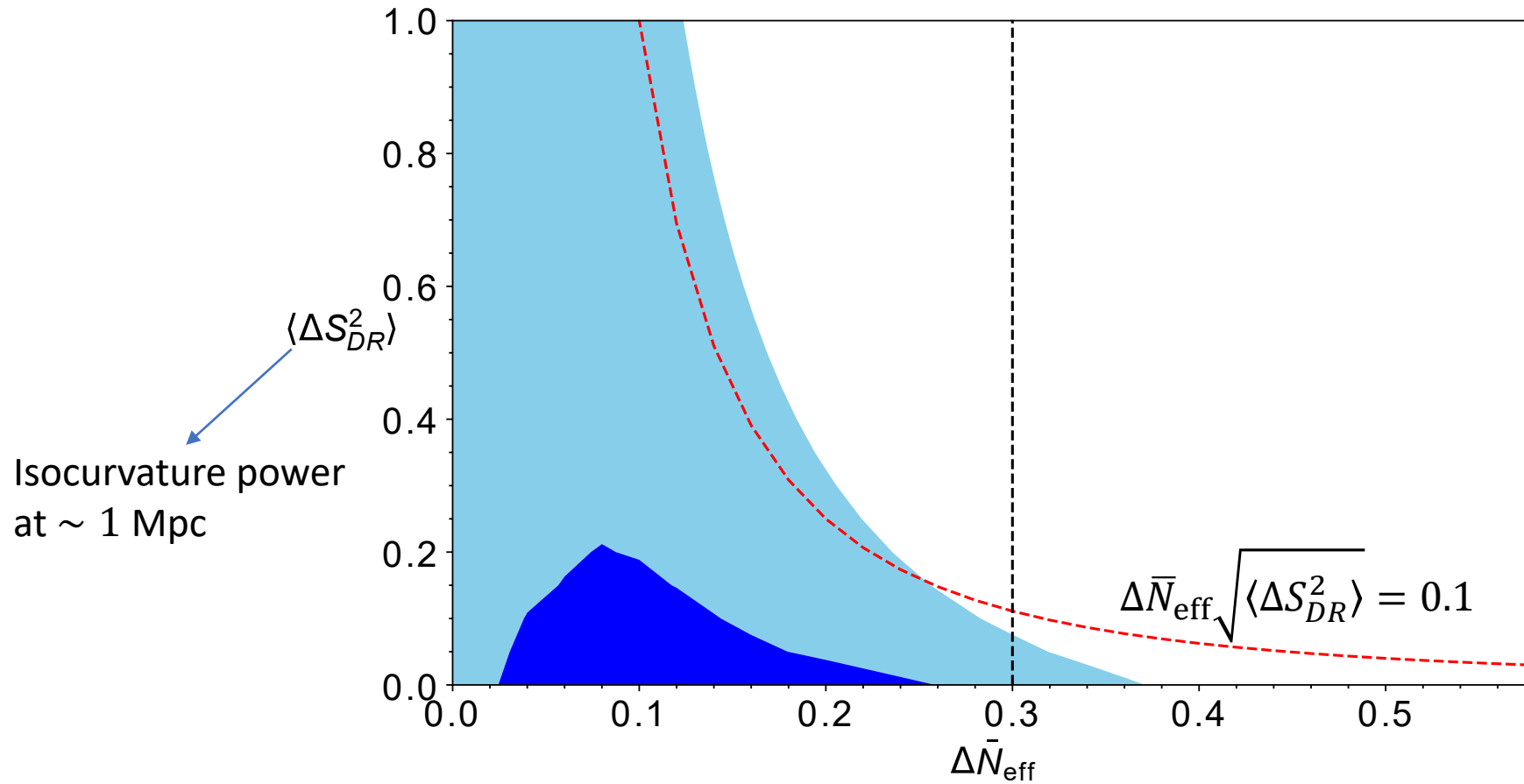




# Dark radiation isocurvature => Variation in He/H

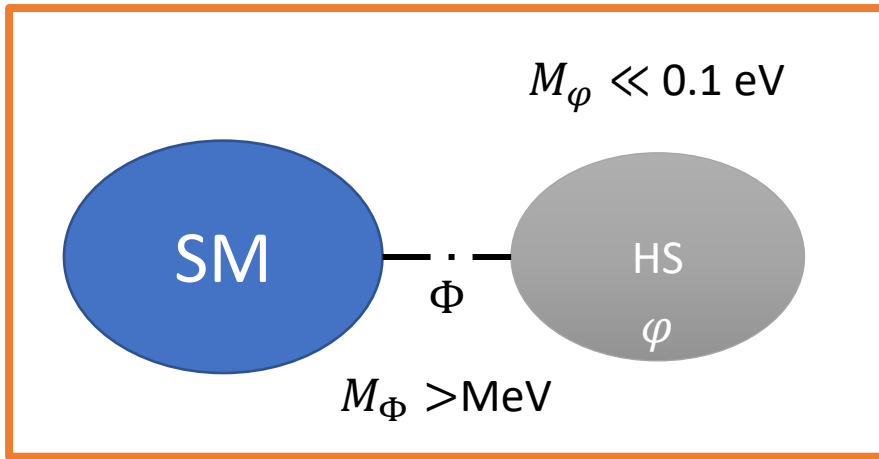


# Constraint on dark radiation isocurvature

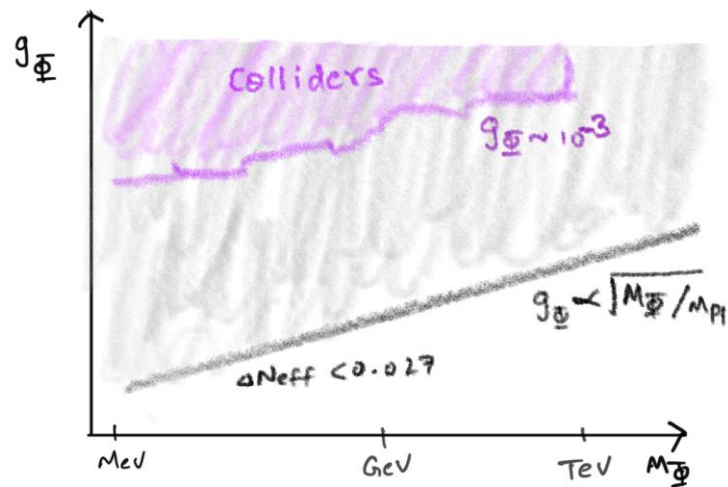


# Combined implication of projects

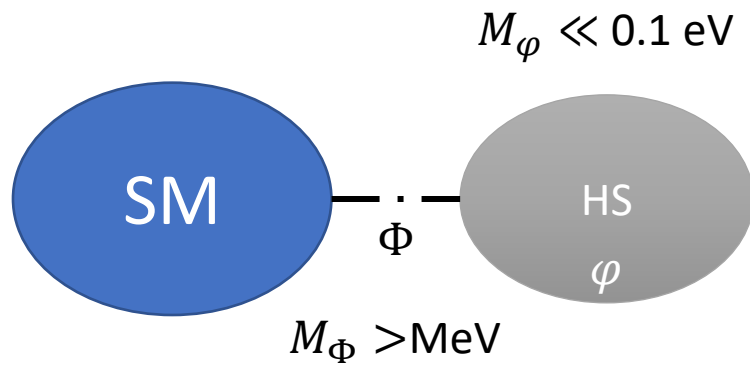
# Massless lightest particle + portal interaction => Impact CMB $N_{\text{eff}}$



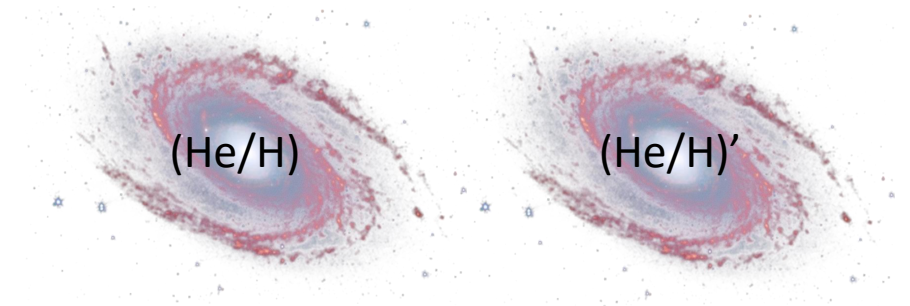
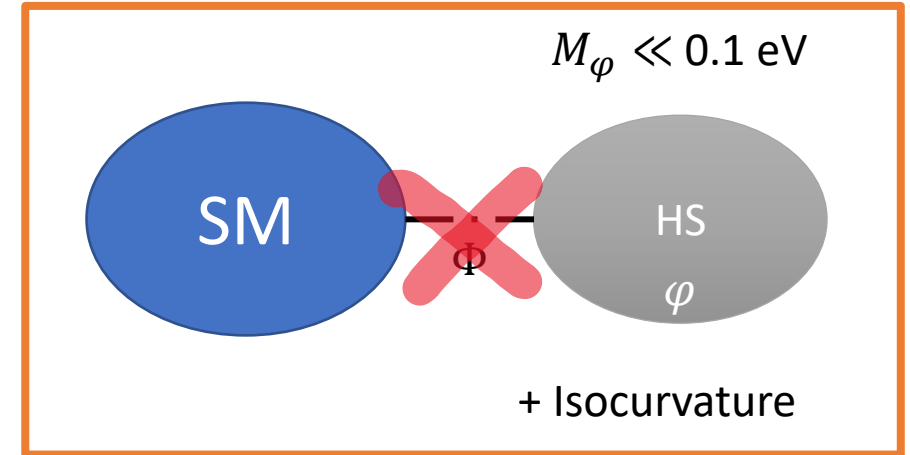
$\phi$  = dark radiation



# Massless lightest particle + no interaction => inhomogeneous BBN

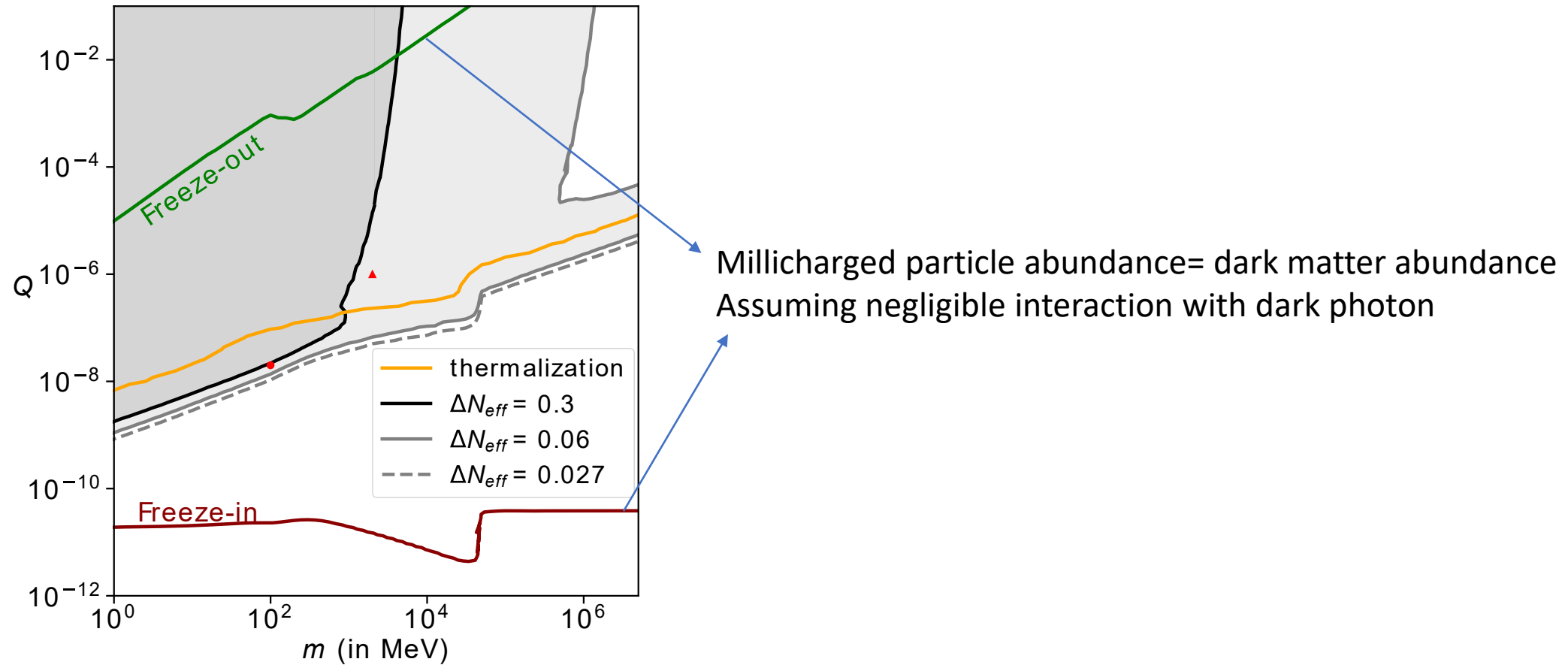


$\varphi = \text{dark radiation}$



Backup slides

# Millicharged particles must dominantly annihilate into dark photons



# BBN constraints on dark radiation isocurvature



# Likelihood analysis: Deuterium

D/H value from  $i^{th}$  gas cloud

$$P(D_i | \{\bar{D}, \sigma_d\}) = \frac{1}{\sqrt{2\pi(\sigma_{n,i}^2 + \sigma_d^2)}} \exp\left(-\frac{(D_i - \bar{D})^2}{2(\sigma_{n,i}^2 + \sigma_d^2)}\right)$$

Theory prediction for average and variance of D/H

Noise in measurement

$$\mathcal{L}_0(\Delta N_{\text{eff}}, \sigma_{\Delta N_{\text{eff}}}) = \int_0^\infty \left[ \prod_i P(D_i | \{\bar{D}, \sigma_d\}) \right]_{\Delta N_{\text{eff}}, \Omega_b h^2, \sigma_{N_{\text{eff}}}} \frac{\exp\left(-\frac{(\Omega_b h^2 - \Omega_b h^2)^2}{2\sigma_{\Omega_b}^2}\right)}{\sqrt{2\pi\sigma_{\Omega_b}^2}} d(\Omega_b h^2).$$

Variance in  $\Delta N_{\text{eff}}$  due to isocurvature

Marginalizing over baryon-to-photon ratio using Planck data

# Error in Deuterium nuclear rates

- Still uncertainty in nuclear reaction of  $d(p, \gamma)^3\text{He}$ :  
$$d + p \rightarrow \gamma + {}^3\text{He}$$
- In my BBN code I used ab-initio calculated rate of above reaction (from Marcucci et al. 2016) in Parthenope.

