

# Hadronic Screening Lengths : A window to **Quark-Gluon Plasma**

*Rajiv V. Gavai \**  
*T. I. F. R., Mumbai*

*\*In collaboration with Sourendu Gupta, TIFR, Mumbai and Robert Lacaze, SPhT, Saclay.*

# Hadronic Screening Lengths : A window to Quark-Gluon Plasma

*Rajiv V. Gavai \**  
*T. I. F. R., Mumbai*

Introduction & Motivation

Hadronic Screening Lengths

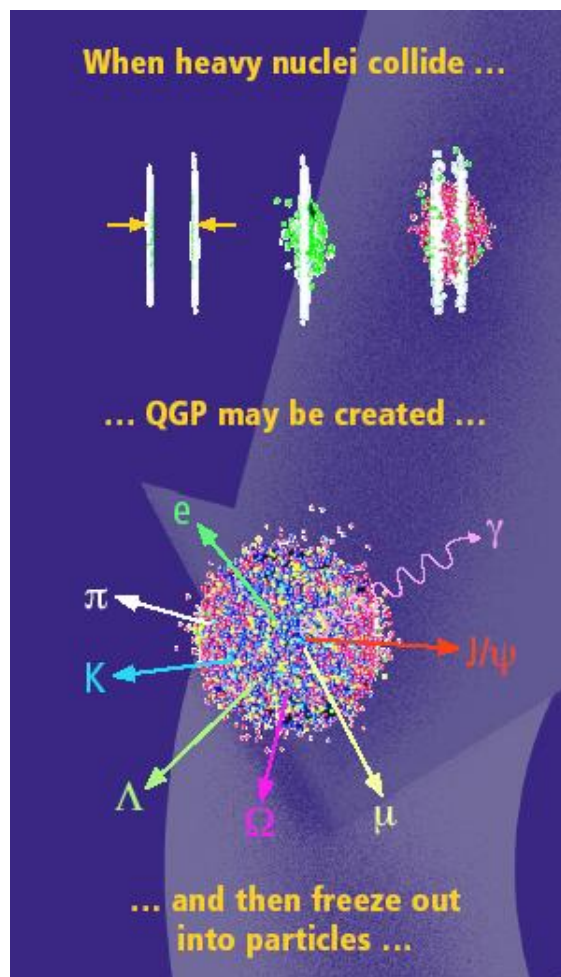
Our Results

Summary

*\*In collaboration with Sourendu Gupta, TIFR, Mumbai and Robert Lacaze, SPhT, Saclay.*

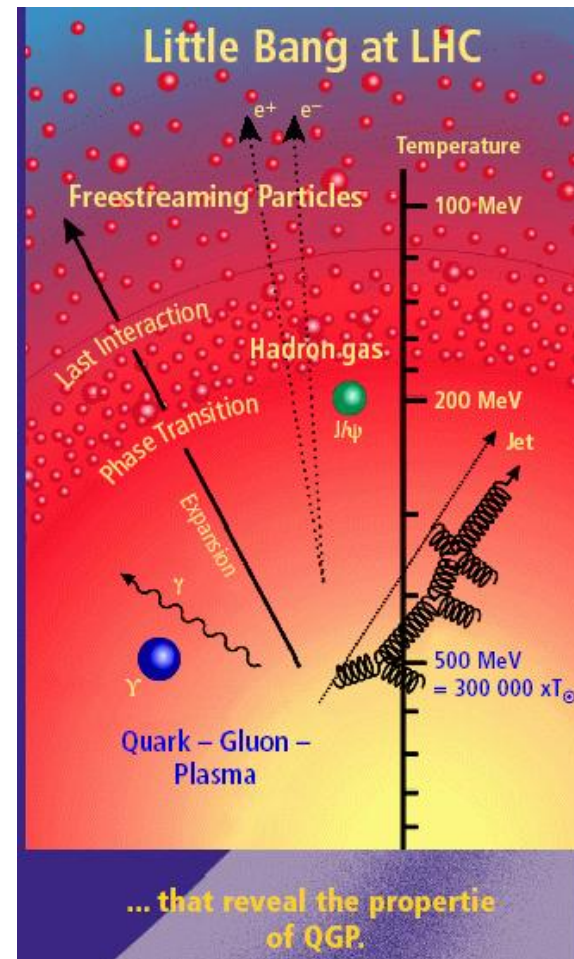
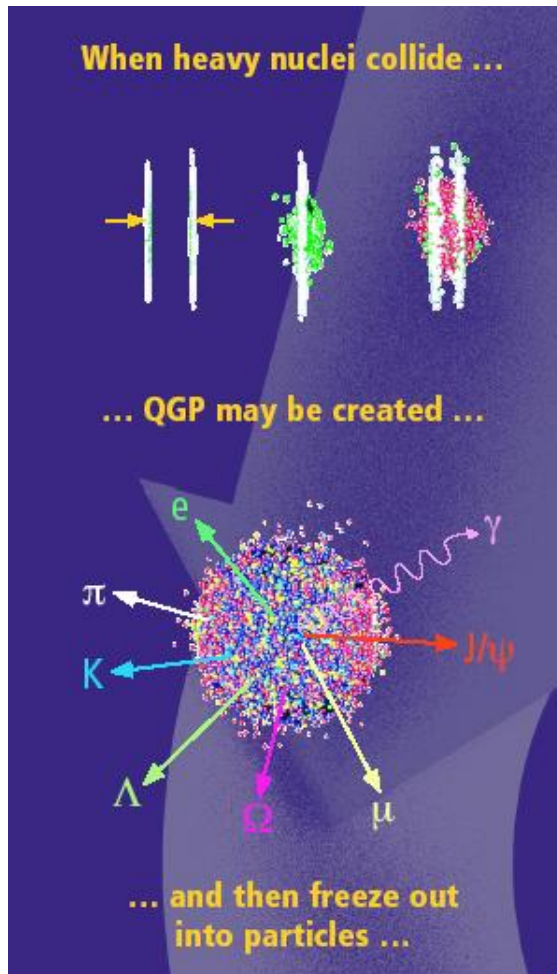
# Introduction

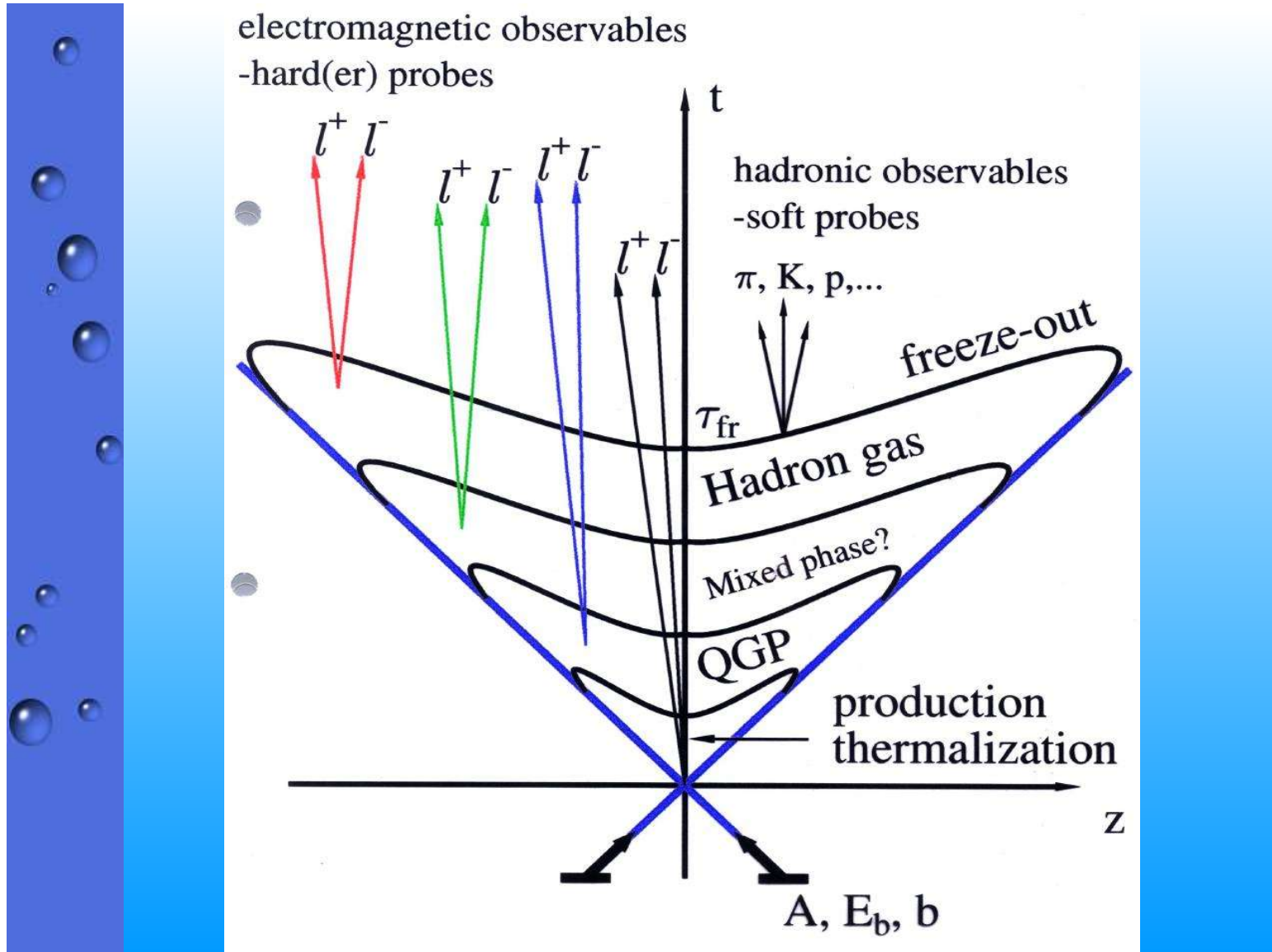
♠ Quest for Quark-Gluon Plasma : Heavy Ion Collisions at SPS, RHIC and LHC.



# Introduction

♠ Quest for Quark-Gluon Plasma : Heavy Ion Collisions at SPS, RHIC and LHC.





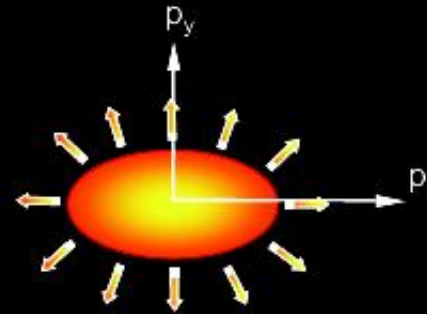
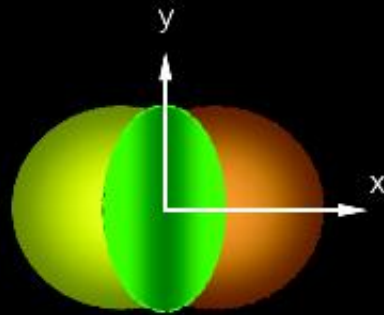


# Anisotropy Parameter $v_2$

coordinate-space-anisotropy



momentum-space-anisotropy



$$\varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

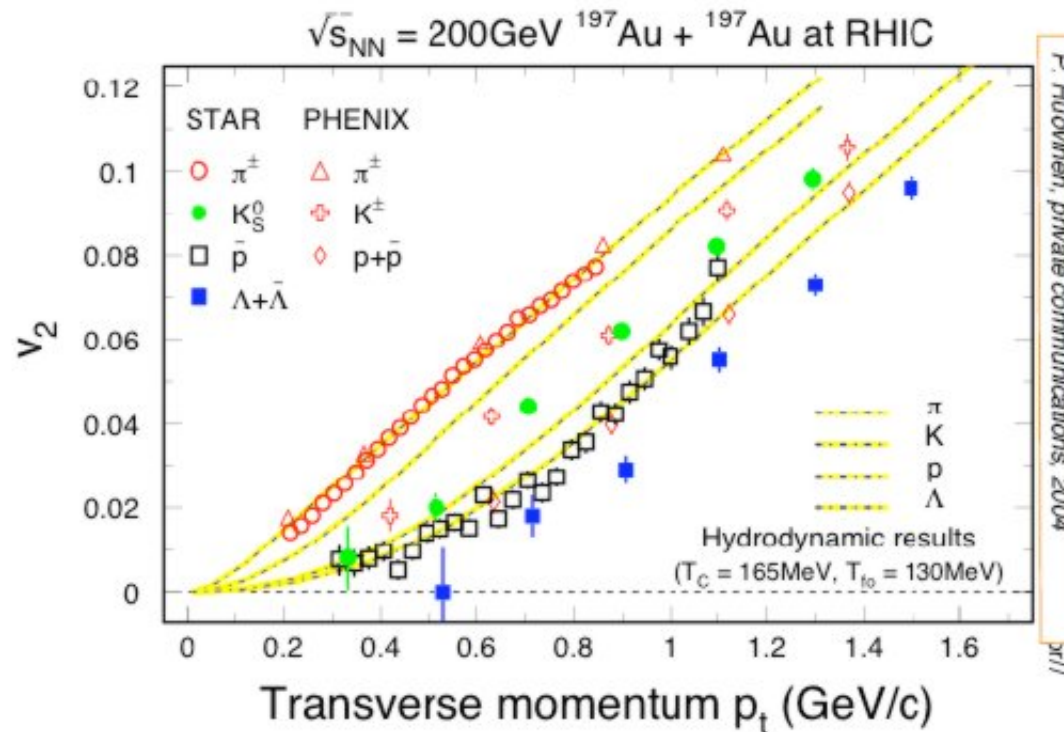
$$v_2 = \langle \cos 2\varphi \rangle, \quad \varphi = \tan^{-1}\left(\frac{p_y}{p_x}\right)$$

**Initial/final conditions, EoS, degrees of freedom**

"ICHEP 2006" Moscow, Russia, July 26 - August 2, 2006

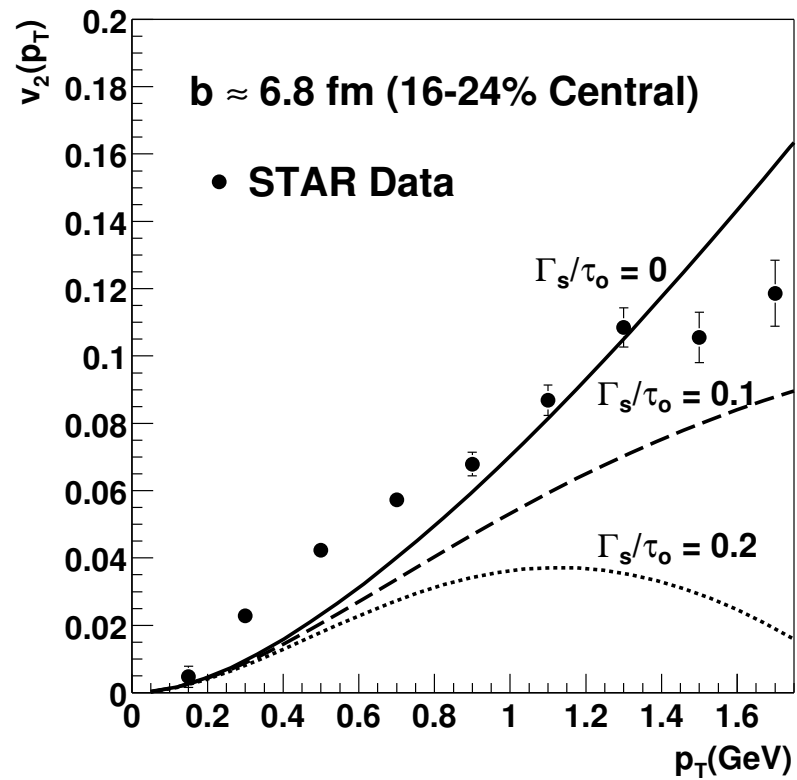


# $v_2$ at Low $p_T$ Region



- Minimum bias data! At low  $p_T$ , model result fits mass hierarchy well!
- Details do not work, need more flow in the model!

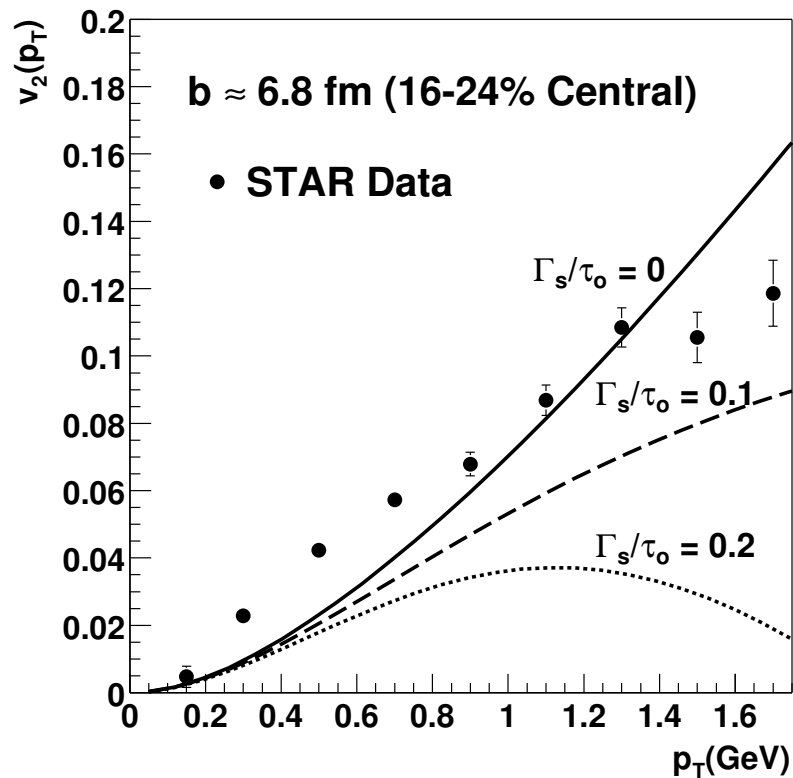
# QGP - (Almost) Perfect Liquid





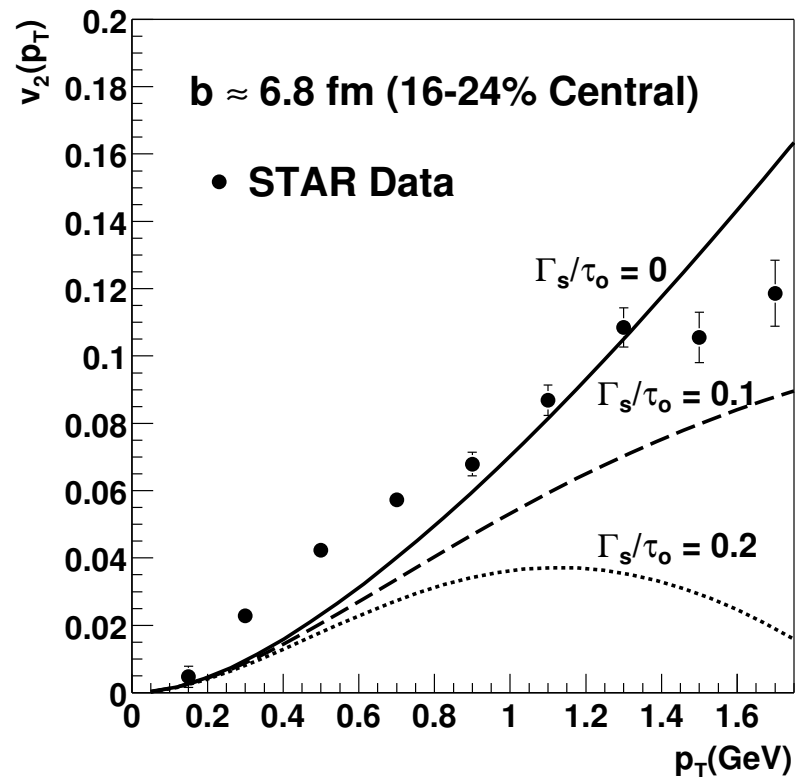
# QGP - (Almost) Perfect Liquid

Elliptic flow consistent with Ideal Hydrodynamics  $\rightsquigarrow$  bound on Shear viscosity. (D. Teaney, nucl-th/03010099; PRC 2003)



# QGP - (Almost) Perfect Liquid

Elliptic flow consistent with Ideal Hydrodynamics  $\rightsquigarrow$  bound on Shear viscosity. (D. Teaney, nucl-th/03010099; PRC 2003)

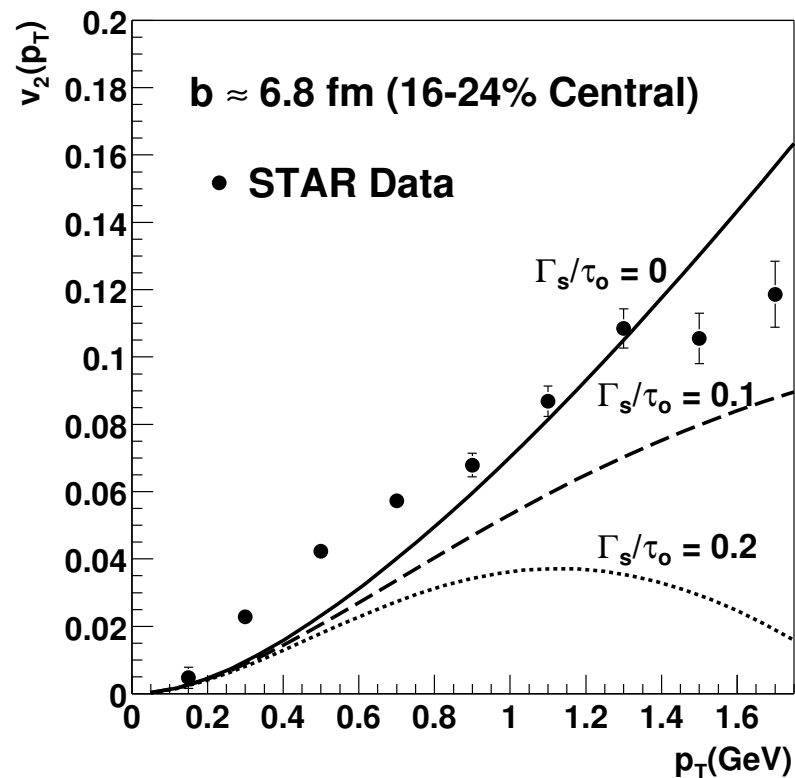


$$\Gamma_s = \frac{4}{3} \frac{\eta}{sT}, \quad (1)$$

where  $\eta$  is Shear Viscosity and  $s$  is entropy density;  $\tau = \sqrt{t^2 - z^2}$  is the time scale of expansion.

# QGP - (Almost) Perfect Liquid

Elliptic flow consistent with Ideal Hydrodynamics  $\rightsquigarrow$  bound on Shear viscosity. (D. Teaney, nucl-th/03010099; PRC 2003)



$$\Gamma_s = \frac{4}{3} \frac{\eta}{sT}, \quad (1)$$

where  $\eta$  is Shear Viscosity and  $s$  is entropy density;  $\tau = \sqrt{t^2 - z^2}$  is the time scale of expansion.

Perturbation theory  $\Rightarrow$  Large  $\eta/s$   
Small  $\eta/s \rightarrow$  Strongly Coupled Liquid.

# Lattice QCD : What it can do

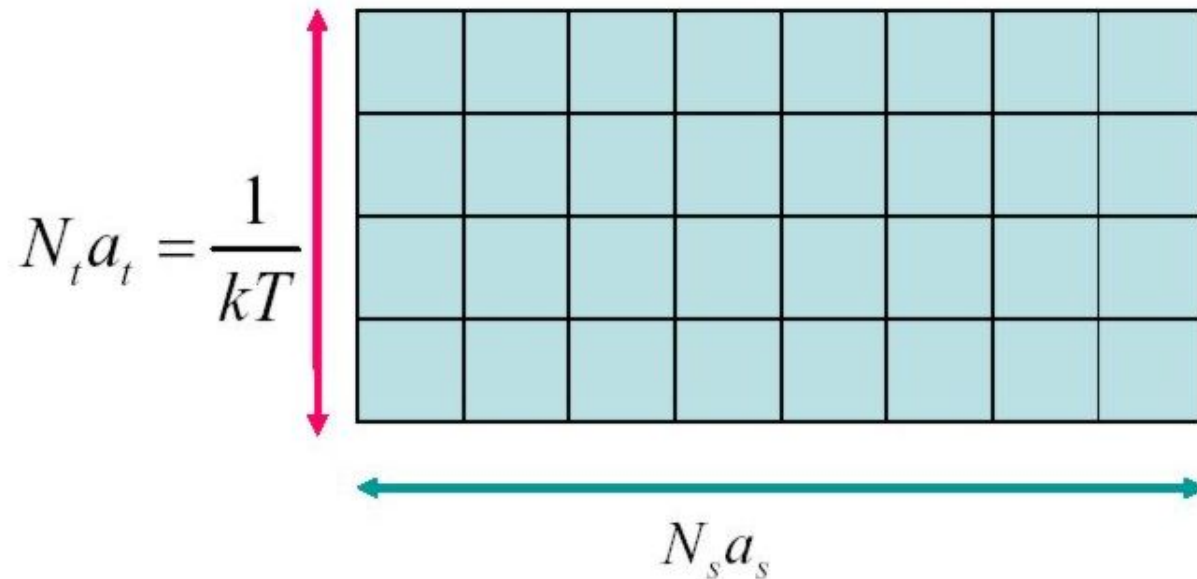
- Transition temperature, Critical energy density, Order of Phase Transition, Properties of QGP (EoS, Excitation types, Screening..)

# Lattice QCD : What it can do

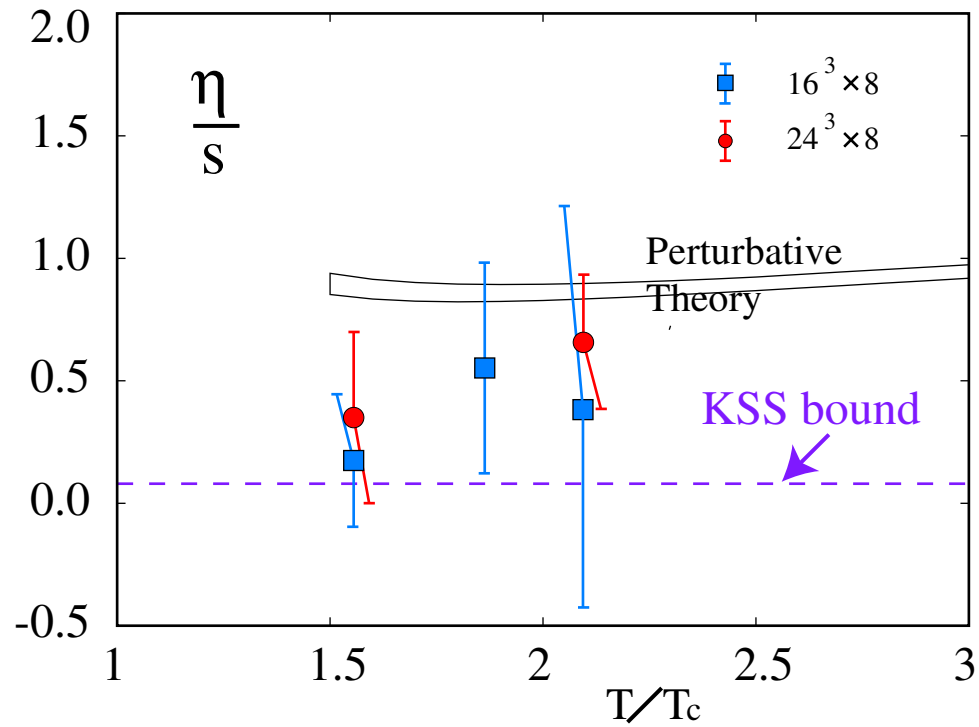
- Transition temperature, Critical energy density, Order of Phase Transition, Properties of QGP (EoS, Excitation types, Screening..)
- Completely parameter-free :  $\Lambda_{QCD}$  and quark masses from hadron spectrum.

# Lattice QCD : What it can do

- Transition temperature, Critical energy density, Order of Phase Transition, Properties of QGP (EoS, Excitation types, Screening..)
- Completely parameter-free :  $\Lambda_{QCD}$  and quark masses from hadron spectrum.

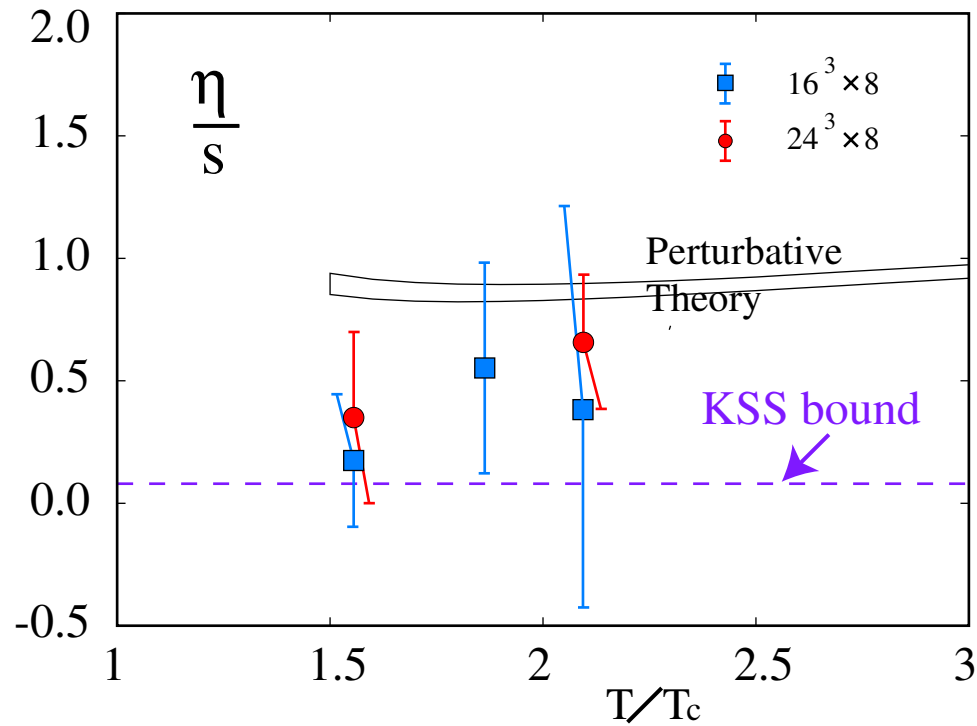


Need  $N_s \gg N_t$  for thermodynamic limit and large  $N_t$  for continuum limit.



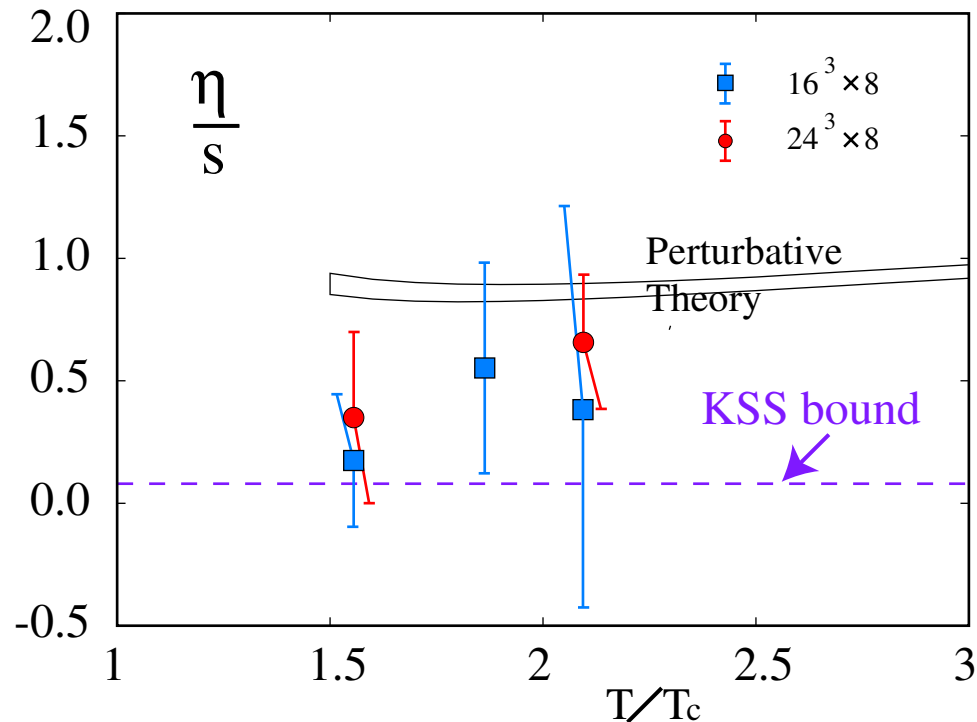
Nakamura and Sakai, PRL 94 (2005).

- Kubo's Linear Response Theory : Transport Coefficients in terms of equilibrium correlation functions.



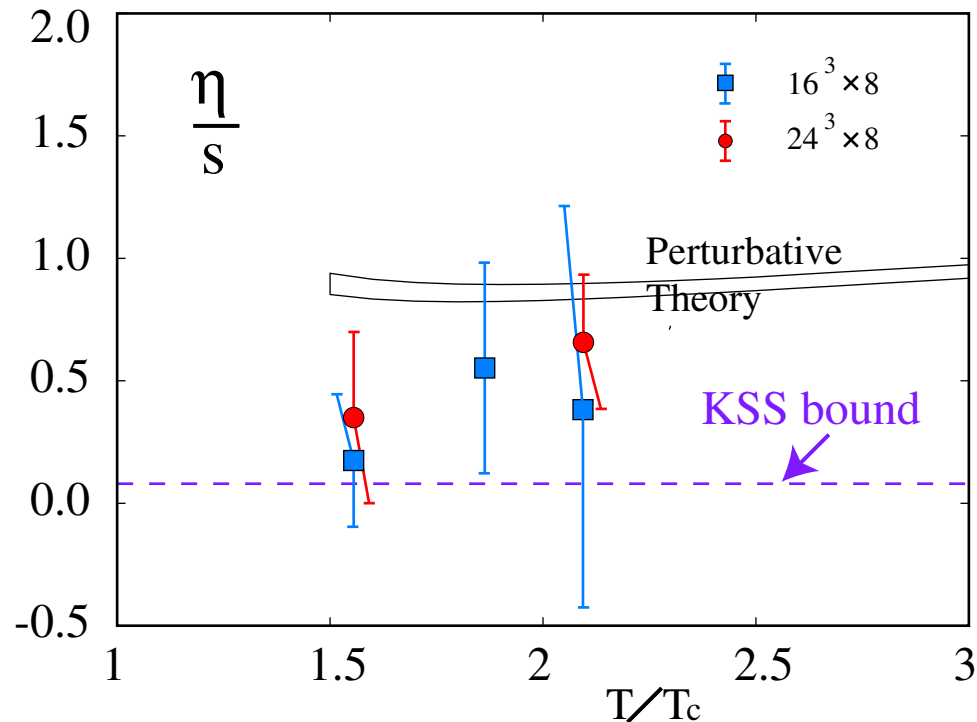
Nakamura and Sakai, PRL 94 (2005).





Nakamura and Sakai, PRL 94 (2005).

- Kubo's Linear Response Theory : Transport Coefficients in terms of equilibrium correlation functions.
- Obtain Energy-Momentum Correlation functions on Lattice (at discrete Matsubara frequencies).
- Continue them to get Retarded ones  $\rightsquigarrow$  Shear, Bulk Viscosities.

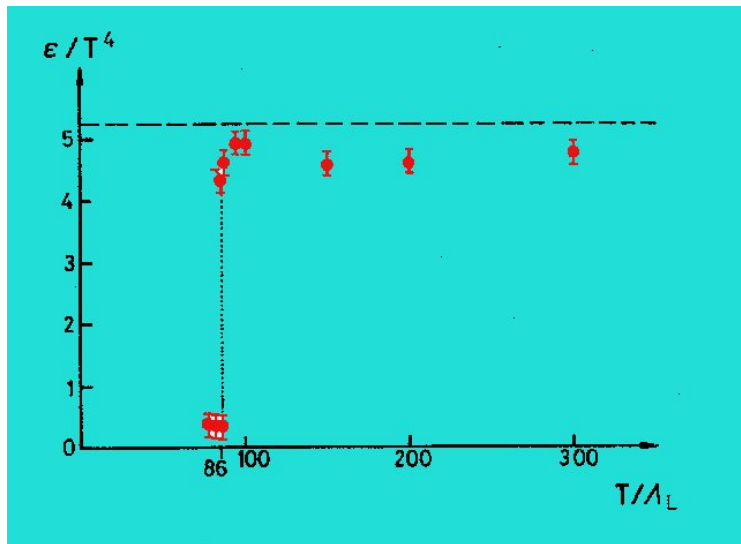


Nakamura and Sakai, PRL 94 (2005).

- Kubo's Linear Response Theory : Transport Coefficients in terms of equilibrium correlation functions.
- Obtain Energy-Momentum Correlation functions on Lattice (at discrete Matsubara frequencies).
- Continue them to get Retarded ones  $\rightsquigarrow$  Shear, Bulk Viscosities.
- Larger lattices and inclusion of dynamical quarks in future.

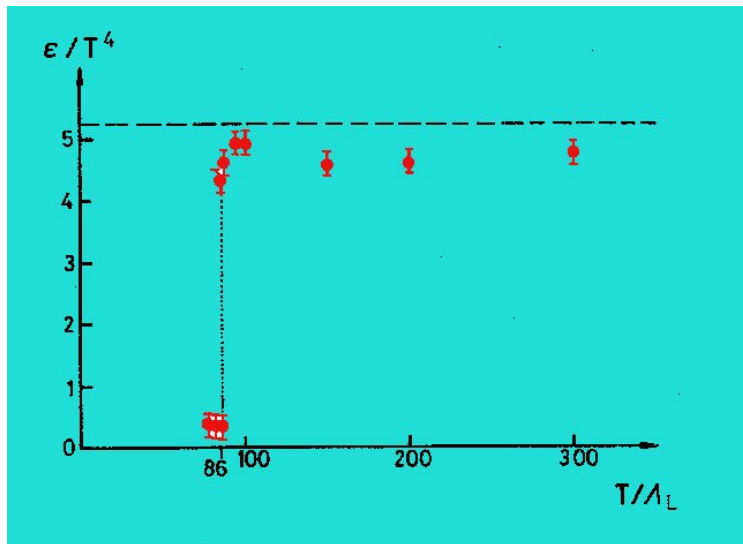
# EoS of QGP

- First results from Bielefeld :

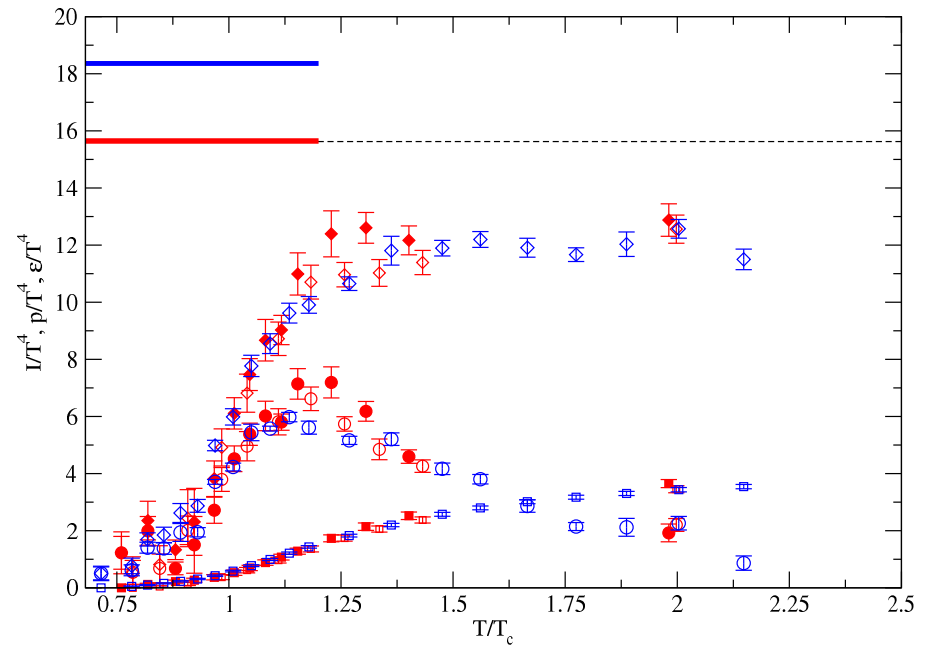


# EoS of QGP

- First results from Bielefeld :



Celik, Engels & Satz, PLB129, 323 1983

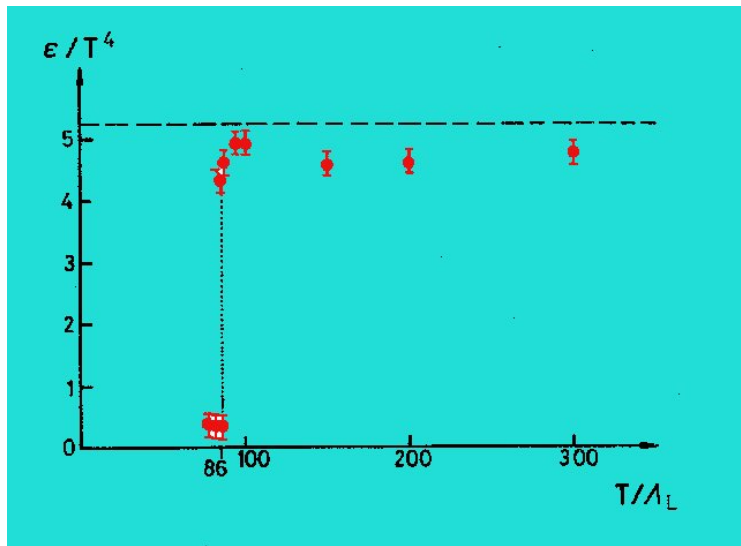


Bernard et al., MILC hep-lat/0509053

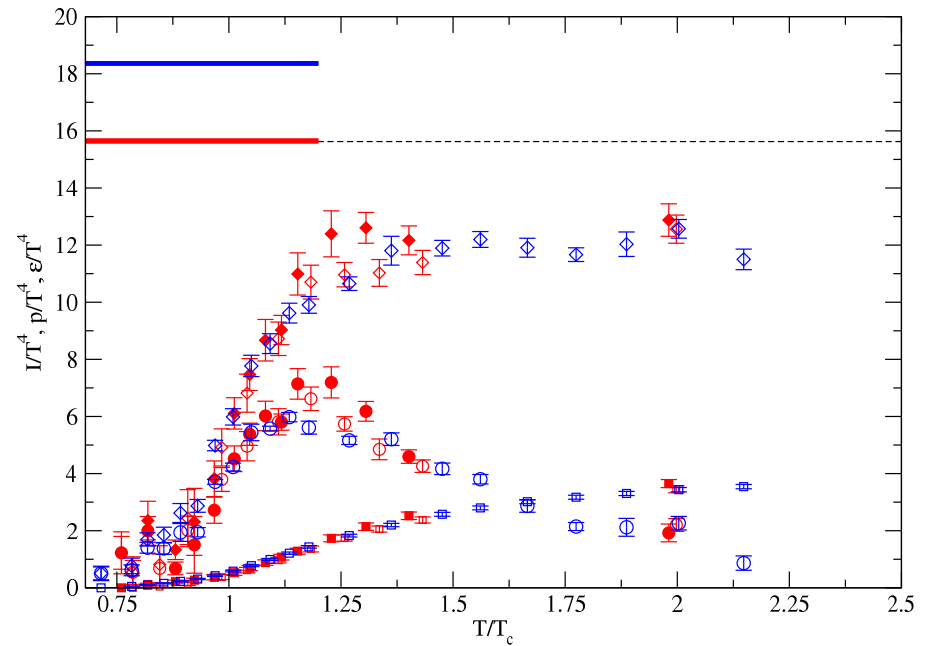
- Recent results for EoS :  $N_t=6$ , Smaller quark masses.

# EoS of QGP

- First results from Bielefeld :

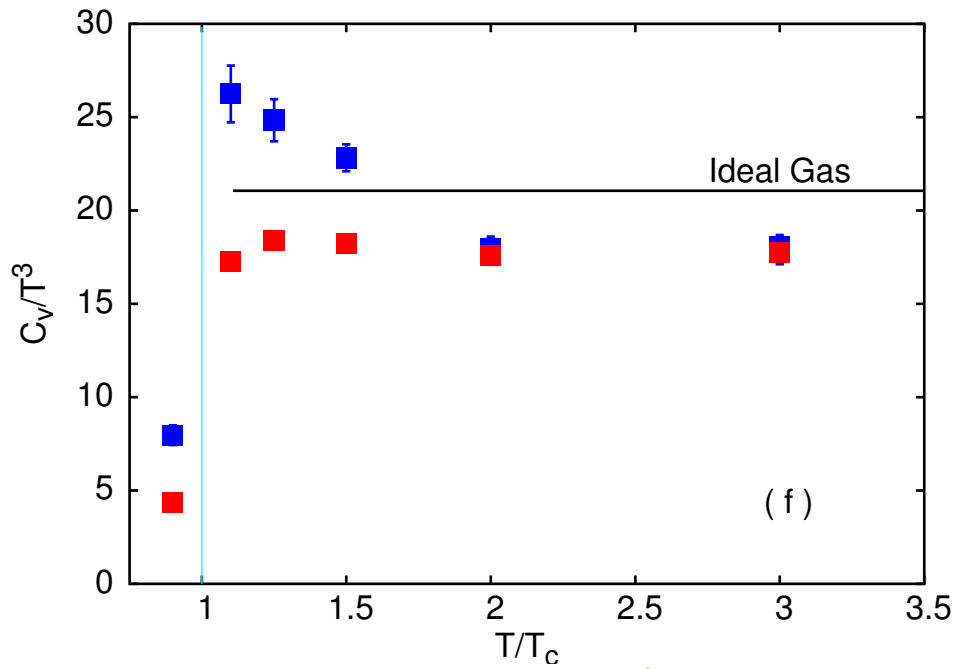


Celik, Engels & Satz, PLB129, 323 1983

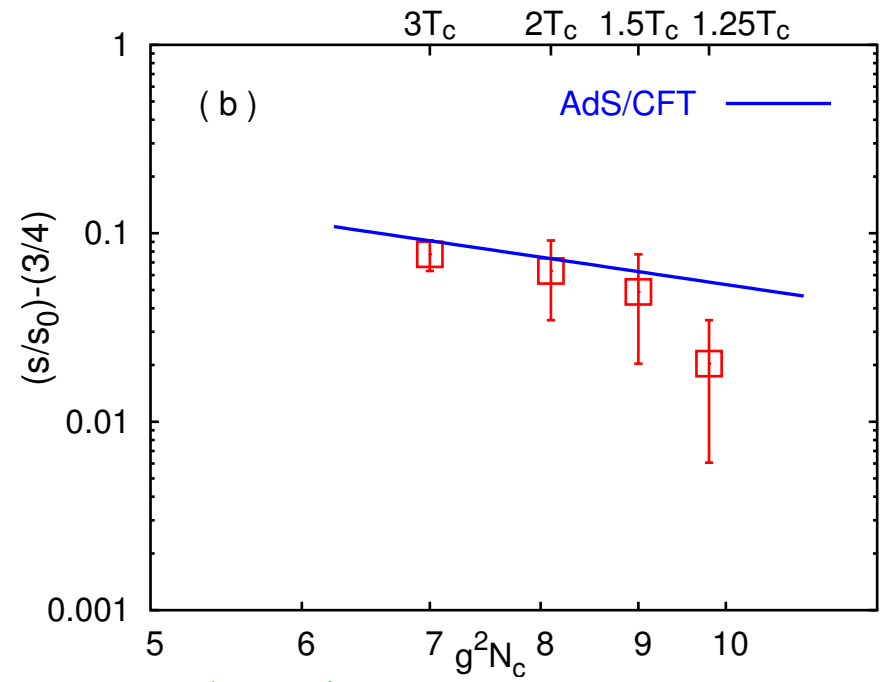


Bernard et al., MILC hep-lat/0509053

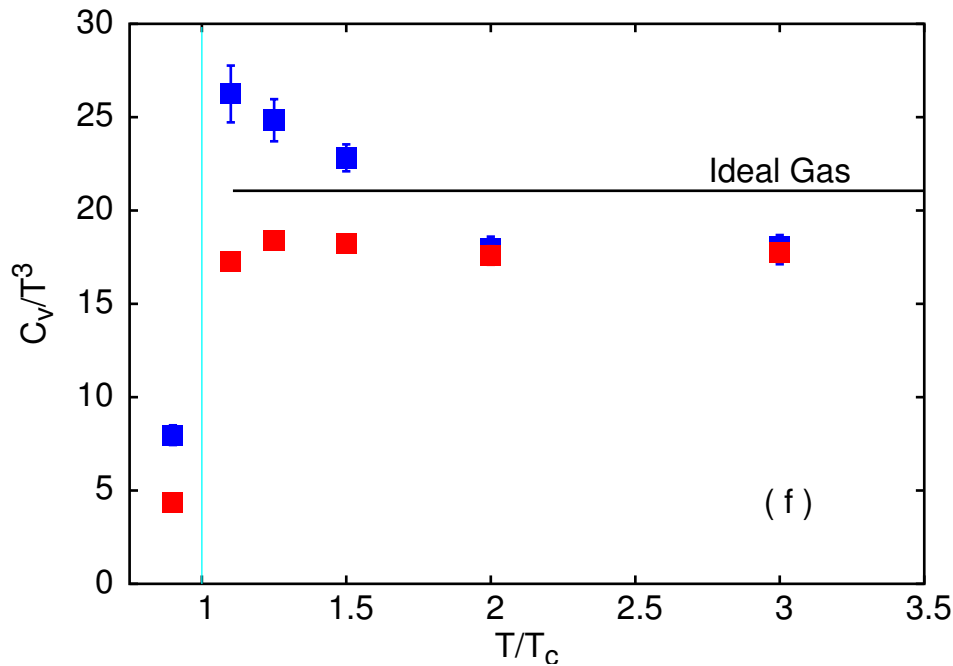
- Recent results for EoS :  $N_t=6$ , Smaller quark masses. Small differences for  $N_t = 4$  &  $6$ ;  $\epsilon(T_c) \sim 6T_c^4$  still. Too small volumes  $\longrightarrow$  Thermodynamic Limit ?



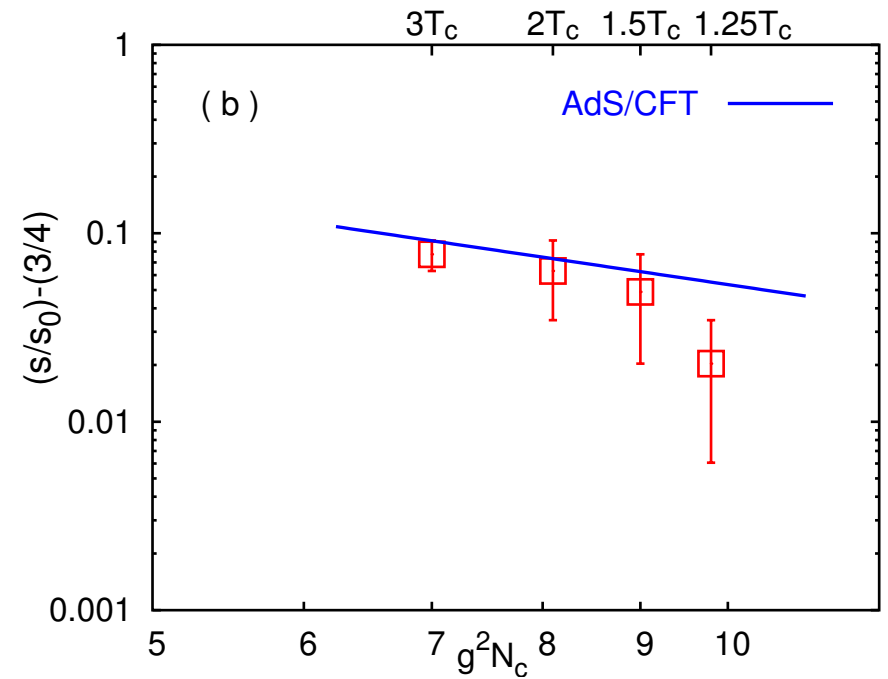
(RVG, S. Gupta and S. Mukherjee, hep-lat/0506015)



♠  $C_v \sim 4\epsilon$  for  $2T_c$  but No Ideal Gas limit.



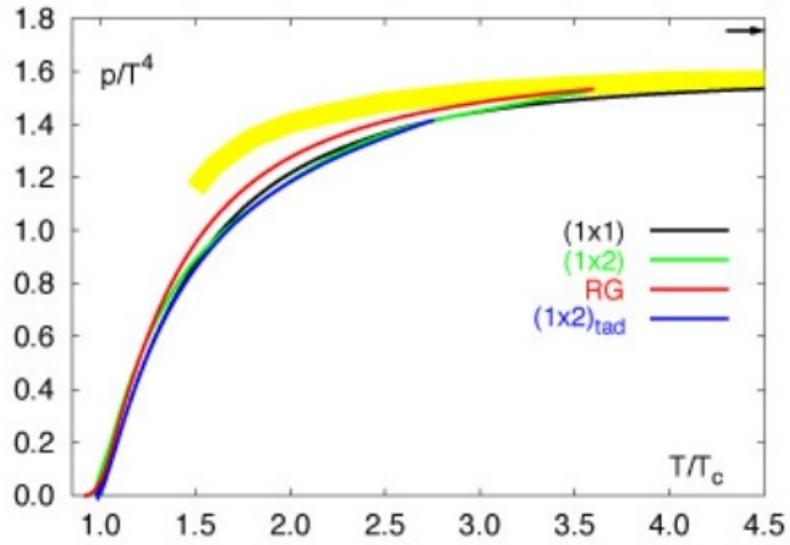
(RVG, S. Gupta and S. Mukherjee, hep-lat/0506015)



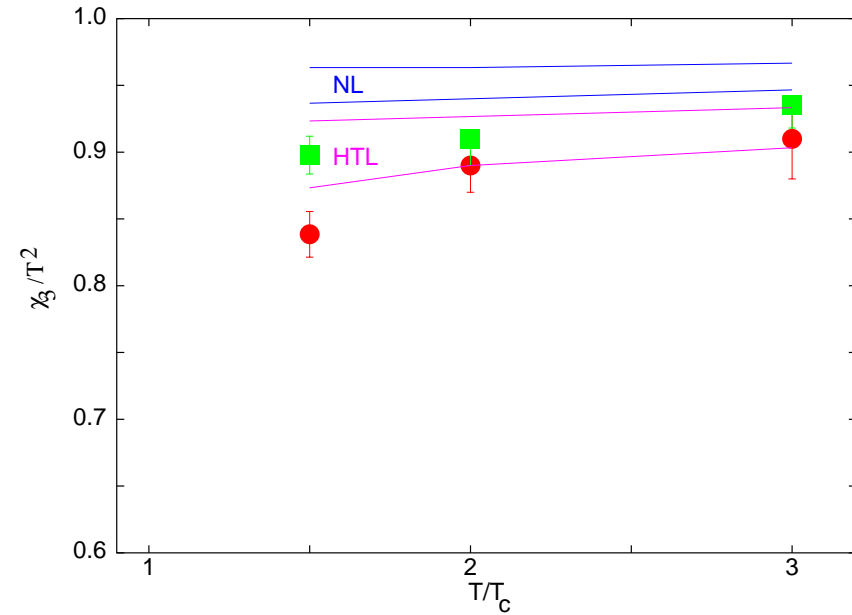
♠  $C_v \sim 4\epsilon$  for  $2T_c$  but No Ideal Gas limit.

♠ Entropy agrees with strong coupling SYM prediction (Gubser, Klebanov & Tseytlin, NPB '98, 202) for  $T = 1.5 - 3T_c$  but fails at lower  $T$ , as do various weak coupling schemes :  $\frac{s}{s_0} = f(g^2 N_c)$ , where  $f(x) = \frac{3}{4} + \frac{45}{32}\zeta(3)x^{-3/2} + \dots$  and  $s_0 = \frac{2}{3}\pi^2 N_c^2 T^3$ .

# Weak Coupling



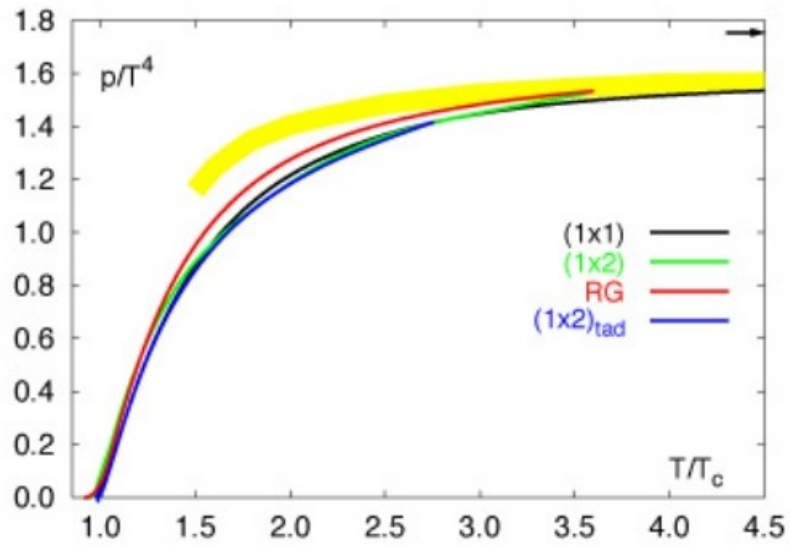
Blaizot, Iancu & Rebhan PRD 01, PLB 01



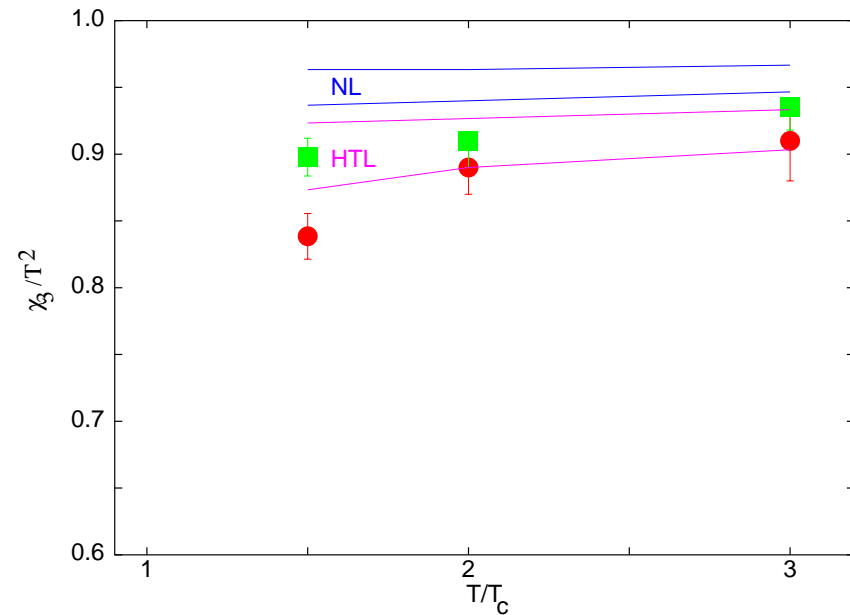
RVG & Gupta PRD 01, 03



# Weak Coupling



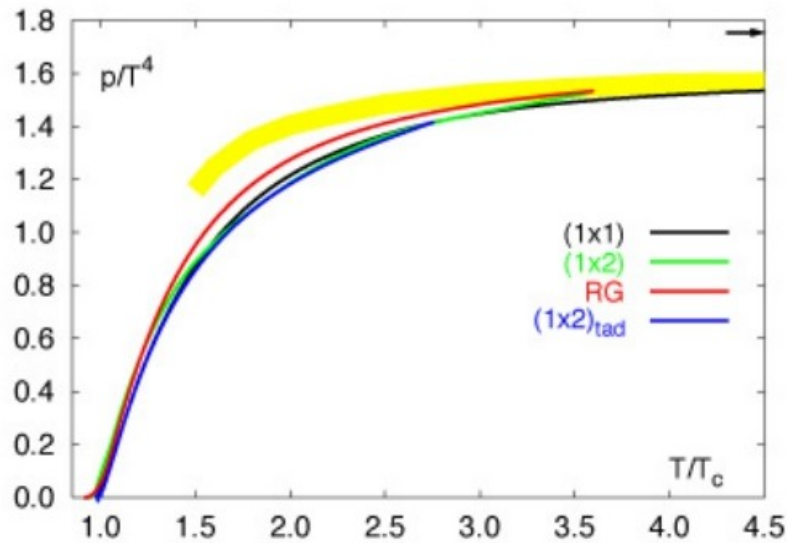
Blaizot, Iancu & Rebhan PRD 01, PLB 01



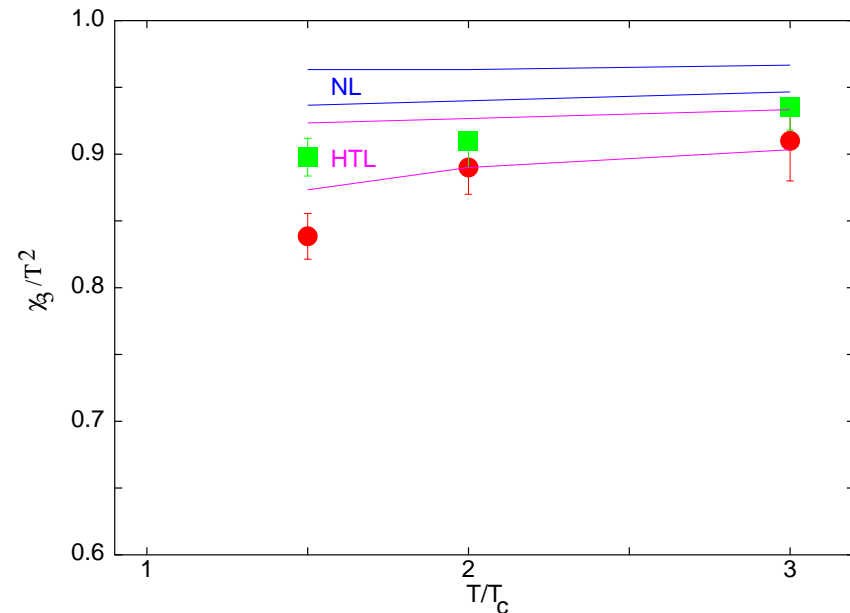
RVG & Gupta PRD 01, 03

♣ Re-summed weak coupling explains lattice results.

# Weak Coupling



Blaizot, Iancu & Rebhan PRD 01, PLB 01



RVG & Gupta PRD 01, 03

♣ Re-summed weak coupling explains lattice results.

♣ So does dimensional reduction (Kajantie et al, Vourinen)

♣ Quasiparticle, PNJL models (Kampfer et al., Wiese et al.).

# Baryon-Strangeness Correlation

- ♣ Correlation between quantum numbers  $K$  and  $L$  can be studied through the ratio  $C_{(KL)/L} = \frac{\langle KL \rangle - \langle K \rangle \langle L \rangle}{\langle L^2 \rangle - \langle L \rangle^2}$ .
- ♣ These are robust : theoretically & experimentally.

# Baryon-Strangeness Correlation

♣ Correlation between quantum numbers  $K$  and  $L$  can be studied through the ratio  $C_{(KL)/L} = \frac{\langle KL \rangle - \langle K \rangle \langle L \rangle}{\langle L^2 \rangle - \langle L \rangle^2}$ .

♣ These are robust : theoretically & experimentally.

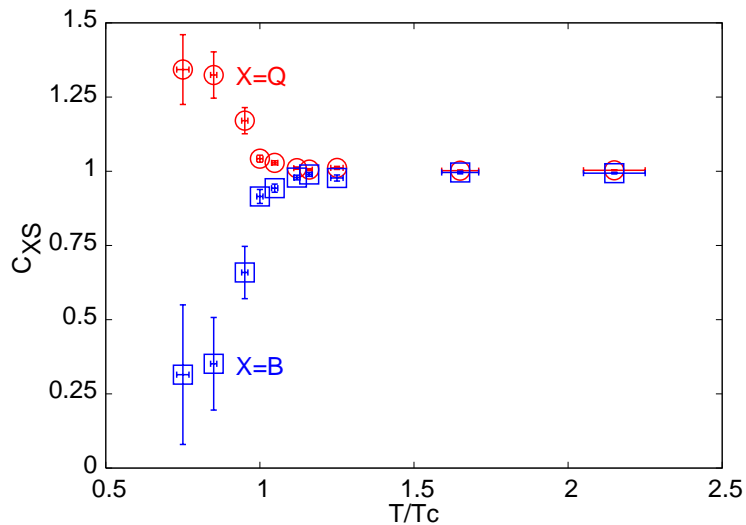
♣ Baryon Number(Charge)–Strangeness correlation :  $C_{(BS)/S}$  ( $C_{(QS)/S}$ ) (Koch, Majumdar and Randurp, PRL 95 (2005); RVG & Sourendu Gupta, PR D 2006; S. Mukherjee, hep-lat/0606018); *u-d* Correlation.

# Baryon-Strangeness Correlation

♣ Correlation between quantum numbers  $K$  and  $L$  can be studied through the ratio  $C_{(KL)/L} = \frac{\langle KL \rangle - \langle K \rangle \langle L \rangle}{\langle L^2 \rangle - \langle L \rangle^2}$ .

♣ These are robust : theoretically & experimentally.

♣ Baryon Number(Charge)–Strangeness correlation:  $C_{(BS)/S}$  ( $C_{(QS)/S}$ ) (Koch, Majumdar and Randurp, PRL 95 (2005); RVG & Sourendu Gupta, PR D 2006; S. Mukherjee, hep-lat/0606018);  $u$ - $d$  Correlation.

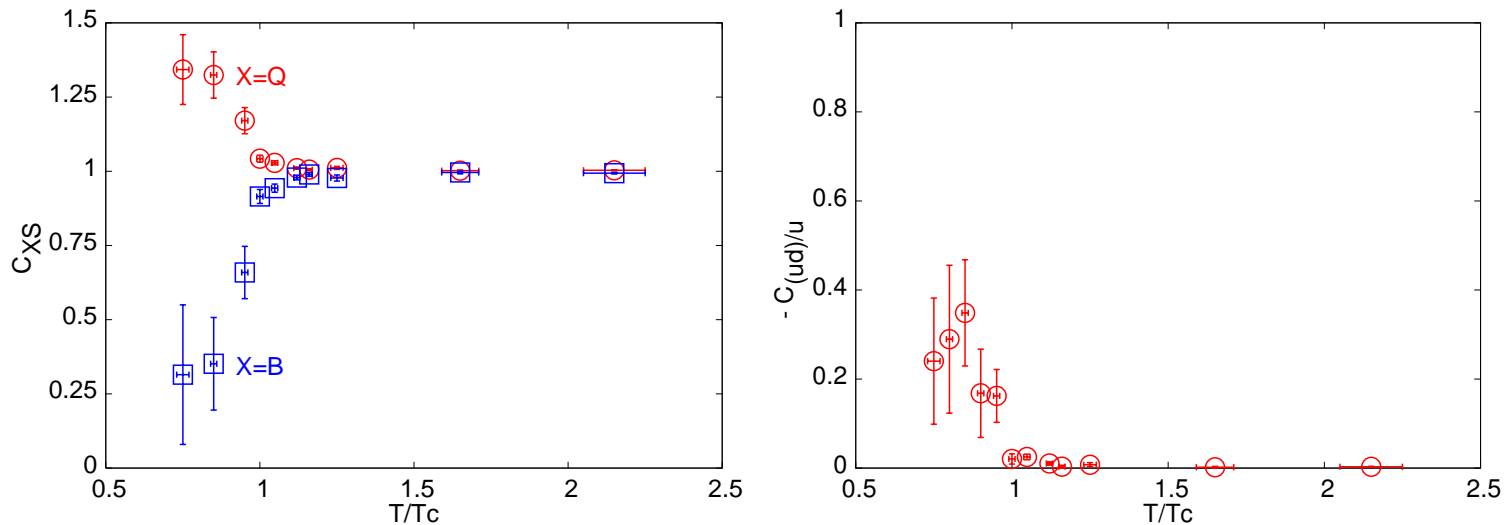


# Baryon-Strangeness Correlation

♣ Correlation between quantum numbers  $K$  and  $L$  can be studied through the ratio  $C_{(KL)/L} = \frac{\langle KL \rangle - \langle K \rangle \langle L \rangle}{\langle L^2 \rangle - \langle L \rangle^2}$ .

♣ These are robust : theoretically & experimentally.

♣ Baryon Number(Charge)–Strangeness correlation:  $C_{(BS)/S}$  ( $C_{(QS)/S}$ ) (Koch, Majumdar and Randurp, PRL 95 (2005); RVG & Sourendu Gupta, PR D 2006; S. Mukherjee, hep-lat/0606018);  $u$ - $d$  Correlation.



# Hadronic Screening Lengths

- DeTar & Kogut (PRD '87) advocated study of Hadronic Screening Lengths to explore the large scale composition of QGP : Long-range nonperturbative effects ?

# Hadronic Screening Lengths

- DeTar & Kogut (PRD '87) advocated study of Hadronic Screening Lengths to explore the large scale composition of QGP : Long-range nonperturbative effects ?
- Obtained from the long-distance behaviour of the correlator  $\langle C_{AB}(z) \rangle = \langle \bar{A}(z)\bar{B}(0) \rangle - \langle \bar{A}(0) \rangle \langle \bar{B}(0) \rangle \sim \exp(-\mu(T)z)$ , as  $z \rightarrow \infty$ .
- Here  $\bar{A}(z) = \sum_{x,y,t} A(x,y,z,t)/N_s^2 N_t$  and is typically taken as a local meson or baryon operator.  $\mu(T)^{-1}$  then is meson(baryon) screening length.



# Hadronic Screening Lengths

- DeTar & Kogut (PRD '87) advocated study of Hadronic Screening Lengths to explore the large scale composition of QGP : Long-range nonperturbative effects ?
- Obtained from the long-distance behaviour of the correlator  $\langle C_{AB}(z) \rangle = \langle \bar{A}(z)\bar{B}(0) \rangle - \langle \bar{A}(0) \rangle \langle \bar{B}(0) \rangle \sim \exp(-\mu(T)z)$ , as  $z \rightarrow \infty$ .
- Here  $\bar{A}(z) = \sum_{x,y,t} A(x,y,z,t)/N_s^2 N_t$  and is typically taken as a local meson or baryon operator.  $\mu(T)^{-1}$  then is meson(baryon) screening length.
- Their conclusion : Existence of hadronic modes in QGP, *unlike* expectations from naive pictures of deconfinement.

- $MT_c$ -collaboration (Born et al. PRL '89) pointed out that lowest Matsubara frequency for small  $N_t$  is much larger than in continuum  $\implies$  can explain  $\rho$  ( $N$ )-screening mass as that for free  $q\bar{q}$  ( $qqq$ )-pair. But  $\mu_\pi$  was still very different.

- $MT_c$ -collaboration (Born et al. PRL '89) pointed out that lowest Matsubara frequency for small  $N_t$  is much larger than in continuum  $\implies$  can explain  $\rho(N)$ -screening mass as that for free  $q\bar{q}$  ( $qqq$ )-pair. But  $\mu_\pi$  was still very different.
- Is  $\pi$  really different in QGP ? or are there “artifacts” of lattice formulation dominating it ?
- Similar results for  $N_f = 0$  (quenched), 2 and 4 flavours of dynamical quarks.

- $MT_c$ -collaboration (Born et al. PRL '89) pointed out that lowest Matsubara frequency for small  $N_t$  is much larger than in continuum  $\implies$  can explain  $\rho(N)$ -screening mass as that for free  $q\bar{q}$  ( $qqq$ )-pair. But  $\mu_\pi$  was still very different.
- Is  $\pi$  really different in QGP ? or are there “artifacts” of lattice formulation dominating it ?
- Similar results for  $N_f = 0$  (quenched), 2 and 4 flavours of dynamical quarks.
- Type of quarks ? Fermions on lattice have a well-known “No-Go” theorem due to Nielsen-Ninomiya :

- $MT_c$ -collaboration (Born et al. PRL '89) pointed out that lowest Matsubara frequency for small  $N_t$  is much larger than in continuum  $\implies$  can explain  $\rho(N)$ -screening mass as that for free  $q\bar{q}$  ( $qqq$ )-pair. But  $\mu_\pi$  was still very different.
- Is  $\pi$  really different in QGP ? or are there “artifacts” of lattice formulation dominating it ?
- Similar results for  $N_f = 0$  (quenched), 2 and 4 flavours of dynamical quarks.
- Type of quarks ? Fermions on lattice have a well-known “No-Go” theorem due to Nielsen-Ninomiya : Popular choices
  - Wilson Fermions – Break *all* chiral symmetries.
  - Kogut-Susskind Fermions – Break some chiral symmetries *but* break also flavour symmetry.
  - Overlap Fermions – *both* correct chiral and flavour symmetry on lattice.

# Overlap-Dirac Operator

♠ Neuberger (PLB 1998) proposed the overlap-Dirac operator :

$$aD = 1 + A(A^\dagger A)^{-1/2} \quad \text{with} \quad A = aD_w, \quad (2)$$

# Overlap-Dirac Operator

♠ Neuberger (PLB 1998) proposed the overlap-Dirac operator :

$$aD = 1 + A(A^\dagger A)^{-1/2} \quad \text{with} \quad A = aD_w, \quad (2)$$

♠ Here  $D_w$  is the Wilson-Dirac Operator given by,

$$aD_w = \frac{1}{2} \{ \gamma_\mu (\partial_\mu^* + \partial_\mu) - a \partial_\mu^* \partial_\mu \} + M, \quad (3)$$

with  $-2 < M < 0$  and  $\partial_\mu$  and  $\partial_\mu^*$  as forward and backward gauge-invariant difference operators.

# Overlap-Dirac Operator

♠ Neuberger ([PLB 1998](#)) proposed the overlap-Dirac operator :

$$aD = 1 + A(A^\dagger A)^{-1/2} \quad \text{with} \quad A = aD_w, \quad (2)$$

♠ Here  $D_w$  is the Wilson-Dirac Operator given by,

$$aD_w = \frac{1}{2} \{ \gamma_\mu (\partial_\mu^* + \partial_\mu) - a \partial_\mu^* \partial_\mu \} + M, \quad (3)$$

with  $-2 < M < 0$  and  $\partial_\mu$  and  $\partial_\mu^*$  as forward and backward gauge-invariant difference operators.

♠ Satisfies  $\{ \gamma_5, D \} = aD\gamma_5D \rightsquigarrow$  Exact Chiral Symmetry on lattice ([Lüscher, PLB 1999](#)).



# Overlap-Dirac Operator

♠ Neuberger ([PLB 1998](#)) proposed the overlap-Dirac operator :

$$aD = 1 + A(A^\dagger A)^{-1/2} \quad \text{with} \quad A = aD_w, \quad (2)$$

♠ Here  $D_w$  is the Wilson-Dirac Operator given by,

$$aD_w = \frac{1}{2} \{ \gamma_\mu (\partial_\mu^* + \partial_\mu) - a \partial_\mu^* \partial_\mu \} + M, \quad (3)$$

with  $-2 < M < 0$  and  $\partial_\mu$  and  $\partial_\mu^*$  as forward and backward gauge-invariant difference operators.

♠ Satisfies  $\{ \gamma_5, D \} = aD\gamma_5D \rightsquigarrow$  Exact Chiral Symmetry on lattice ([Lüscher, PLB 1999](#)).

♠ quark with a mass :  $D(ma) = ma + (1 - ma/2)D$ ; Use  $ma = 0.001 - 0.1$

# Computational Difficulties

- Quark Propagator,  $Y = D^{-1}X$ , needs inversion of  $D$ . Usually done iteratively (Conjugate Gradient).
- At each iteration for overlap, need  $M^{-1/2}X$  : Iterations within each iteration.

# Computational Difficulties

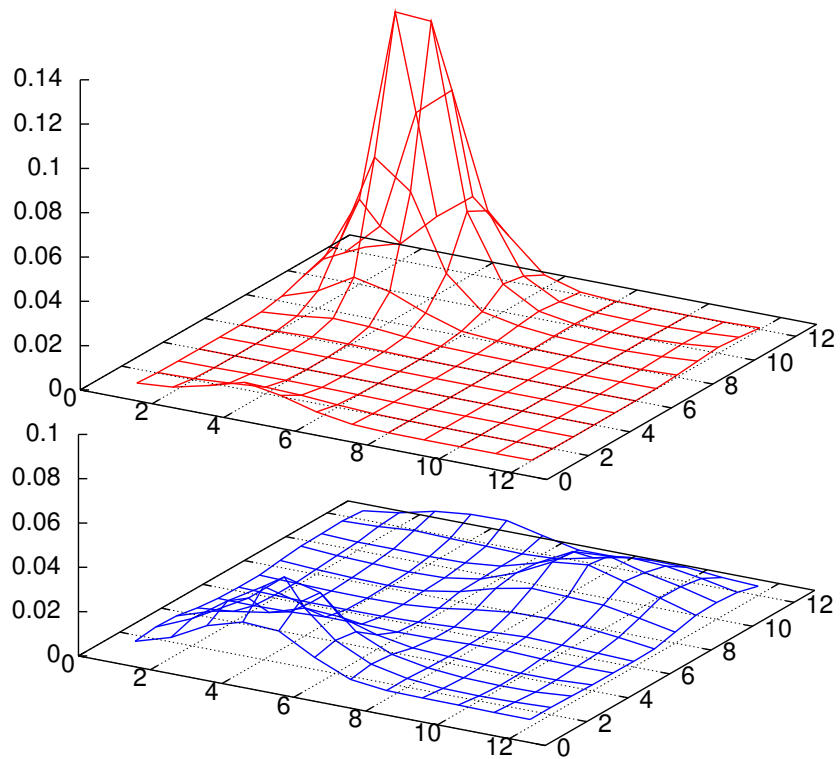
- Quark Propagator,  $Y = D^{-1}X$ , needs inversion of  $D$ . Usually done iteratively (Conjugate Gradient).
- At each iteration for overlap, need  $M^{-1/2}X$  : Iterations within each iteration.
- Quenched QCD with overlap quarks  $\equiv$  Full QCD with Wilson quarks in computational resources.  
Full QCD with overlap quarks  $\sim$  Square of that!

# Computational Difficulties

- Quark Propagator,  $Y = D^{-1}X$ , needs inversion of  $D$ . Usually done iteratively (Conjugate Gradient).
- At each iteration for overlap, need  $M^{-1/2}X$  : Iterations within each iteration.
- Quenched QCD with overlap quarks  $\equiv$  Full QCD with Wilson quarks in computational resources.  
Full QCD with overlap quarks  $\sim$  Square of that!
- Several methods for computing  $M^{-1/2}X$ , including one by us (PRD 2002, CPC 2003).
- We use two algorithms : Conjugate Gradient based CGA, and Zolotarev Approximation.

# Our Results

Gvai, Gupta, Lacaze PRD 2002

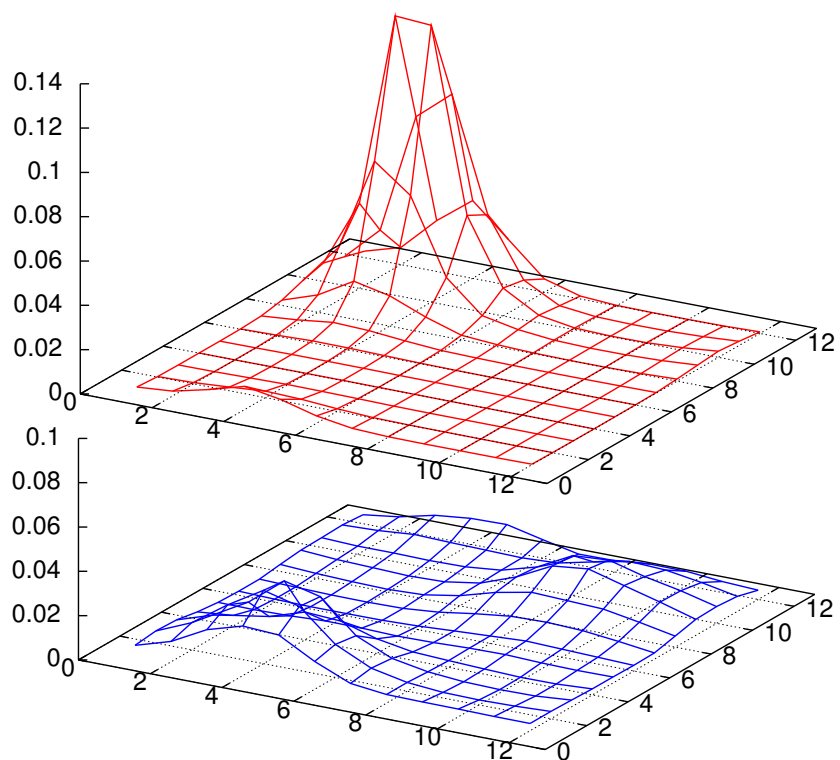


# Our Results

Eigenvalues of  $D$  come in pairs of opposite chiralities except zero.

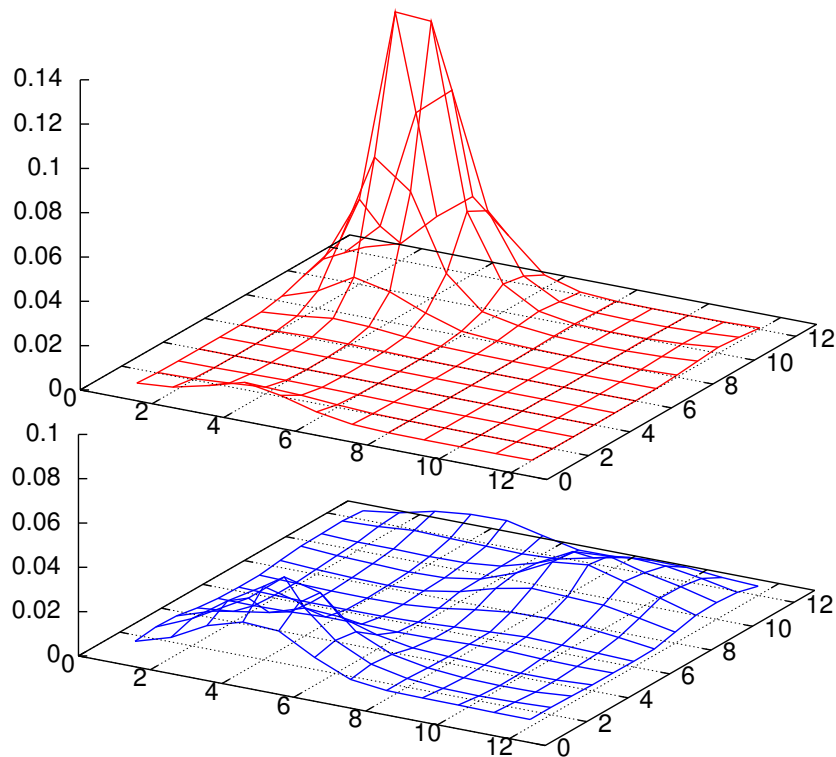
Index theorem for Overlap quarks (P. Hasenfratz et al. PLB 1998)  $\rightsquigarrow$  Instanton - zero modes linkage

Gvai, Gupta, Lacaze PRD 2002



# Our Results

Gvai, Gupta, Lacaze PRD 2002



Eigenvalues of  $D$  come in pairs of opposite chiralities except zero.

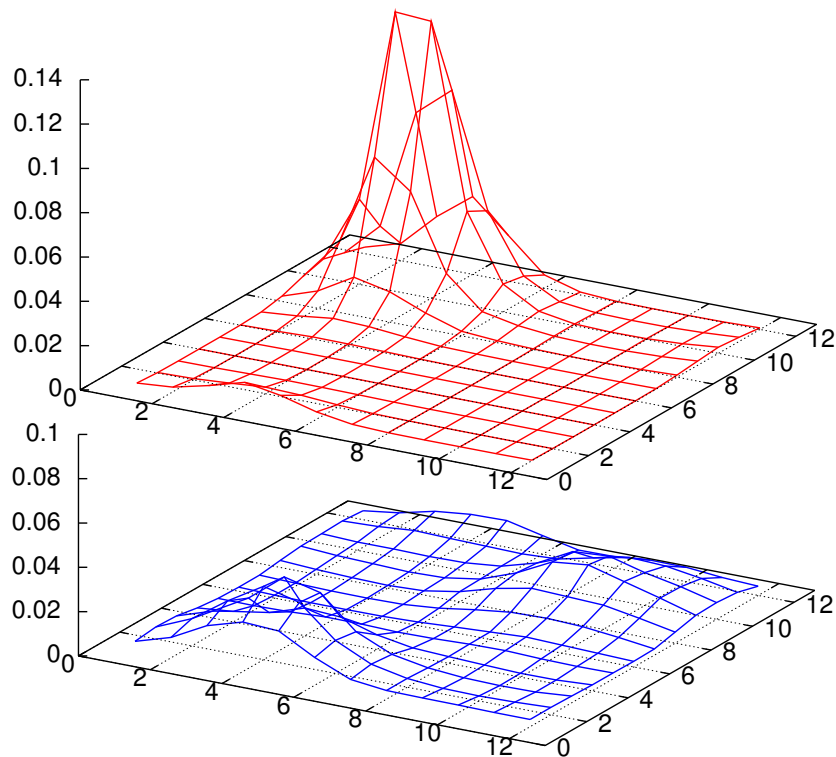
Index theorem for Overlap quarks (P. Hasenfratz et al. PLB 1998)  $\rightsquigarrow$  Instanton - zero modes linkage

Found zero modes all the way up to  $2T_c$ , but in decreasing numbers.

$\langle (n_+ - n_-)^2 \rangle / V$  falls as power of  $T/T_c$  :  $U_A(1)$  continues to be broken up to  $2T_c$ .

# Our Results

Gvai, Gupta, Lacaze PRD 2002



Eigenvalues of  $D$  come in pairs of opposite chiralities except zero.

Index theorem for Overlap quarks (P. Hasenfratz et al. PLB 1998)  $\rightsquigarrow$  Instanton - zero modes linkage

Found zero modes all the way up to  $2T_c$ , but in decreasing numbers.

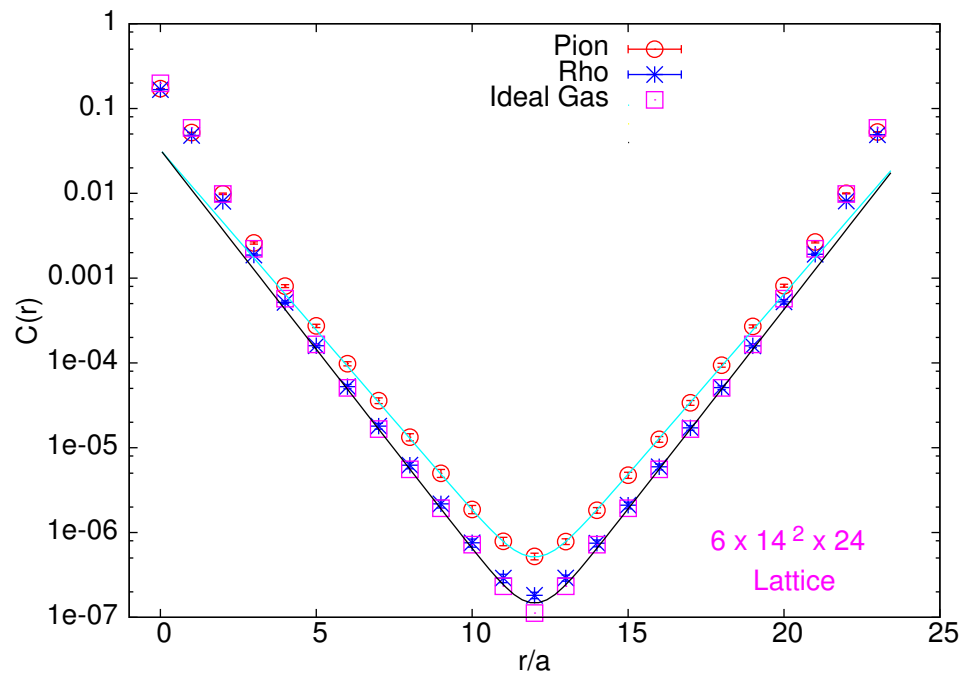
$\langle (n_+ - n_-)^2 \rangle / V$  falls as power of  $T/T_c$  :  $U_A(1)$  continues to be broken up to  $2T_c$ .

$C_V = C_A$  &  $C_{PS} = -C_S$  after subtraction of zero modes.

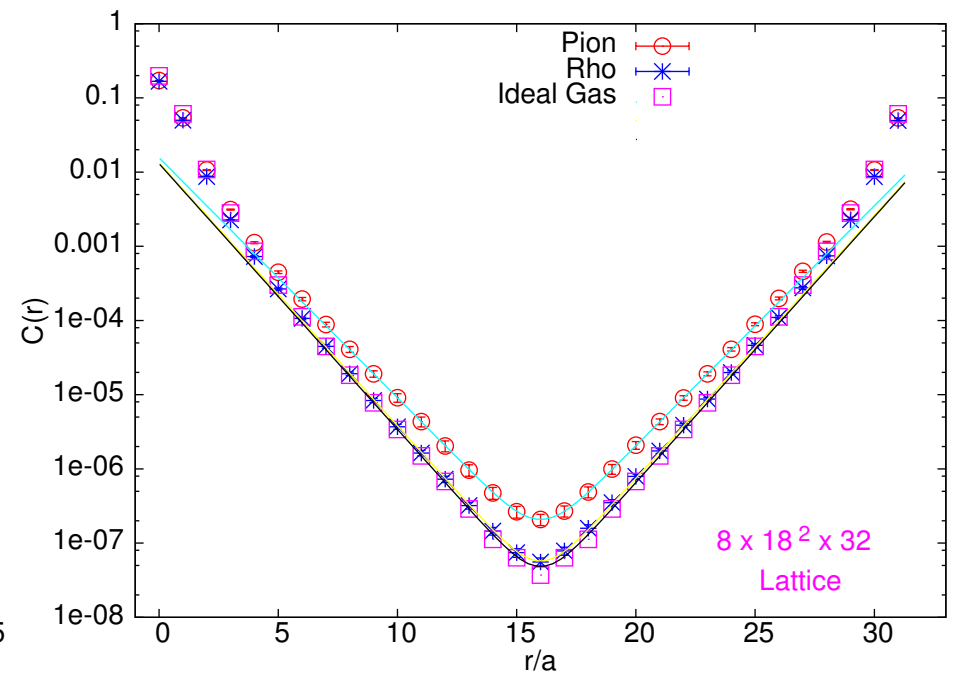
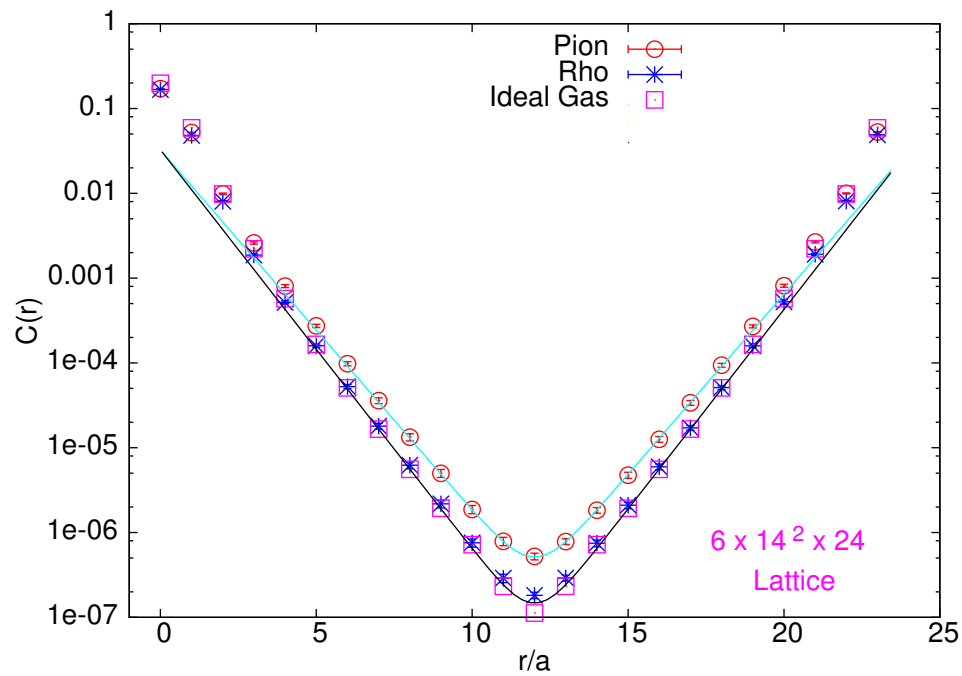


♡ Screening lengths ( $\mu/T$ ) essentially  $T$ -independent for  $1.25 \leq T/T_c \leq 2$  for  $N_t = 4$ . (GGL, PRD 2002). Investigating now continuum limit at  $2T_c$  :

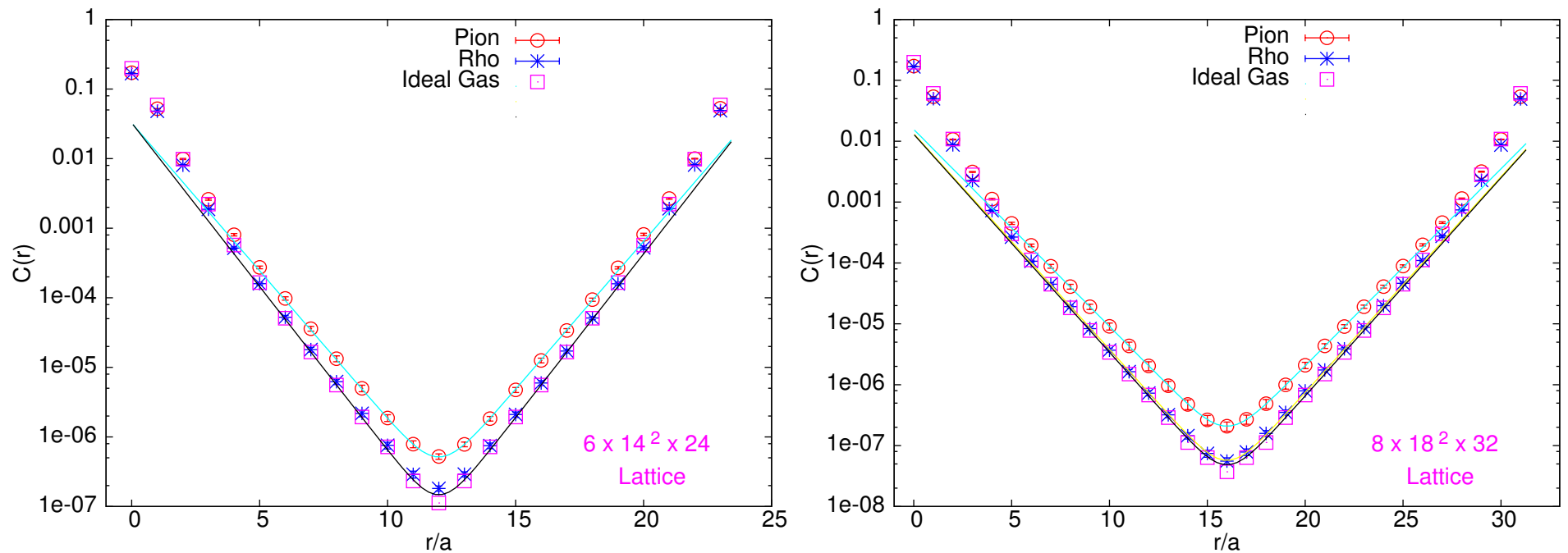
♡ Screening lengths ( $\mu/T$ ) essentially  $T$ -independent for  $1.25 \leq T/T_c \leq 2$  for  $N_t = 4$ . (GGL, PRD 2002). Investigating now continuum limit at  $2T_c$  :



♡ Screening lengths ( $\mu/T$ ) essentially  $T$ -independent for  $1.25 \leq T/T_c \leq 2$  for  $N_t = 4$ . (GGL, PRD 2002). Investigating now continuum limit at  $2T_c$  :



♡ Screening lengths ( $\mu/T$ ) essentially  $T$ -independent for  $1.25 \leq T/T_c \leq 2$  for  $N_t = 4$ . (GGL, PRD 2002). Investigating now continuum limit at  $2T_c$  :

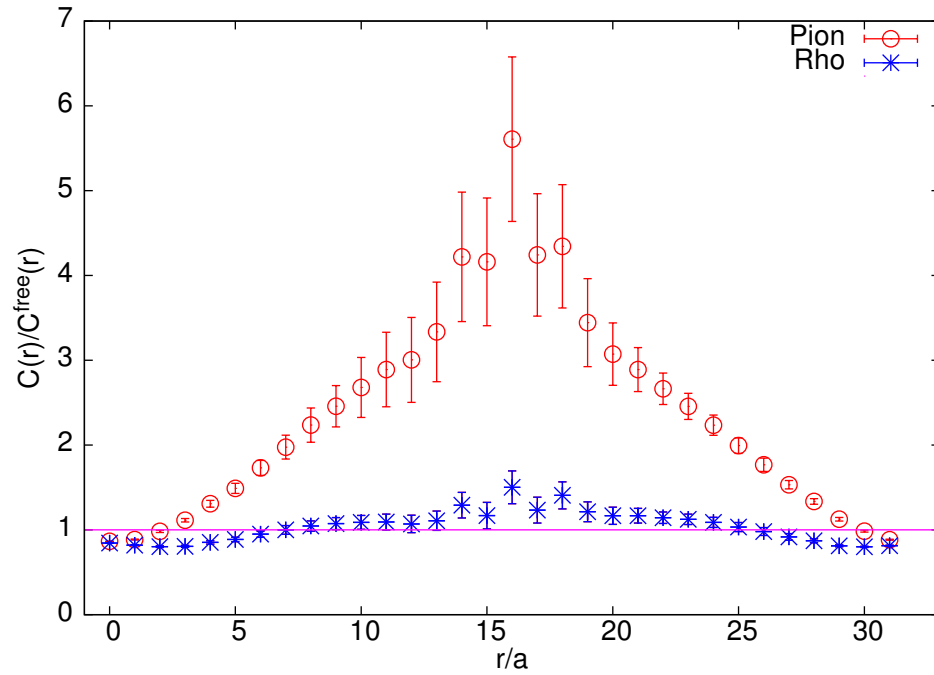


♣ On both  $N_t = 6$  and  $8$ , cosh-like behaviour is seen.

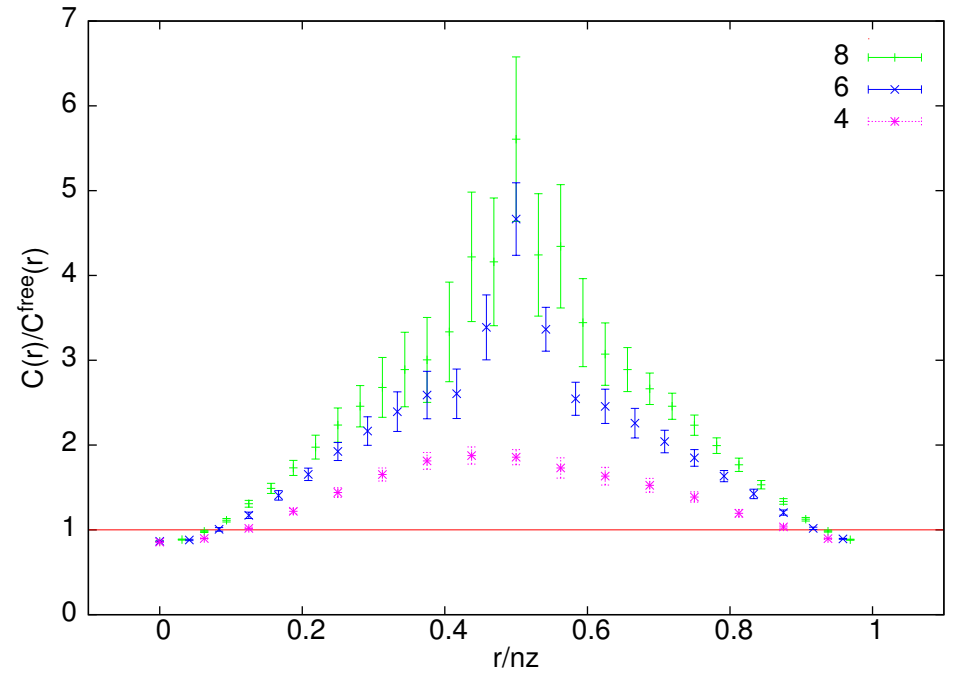
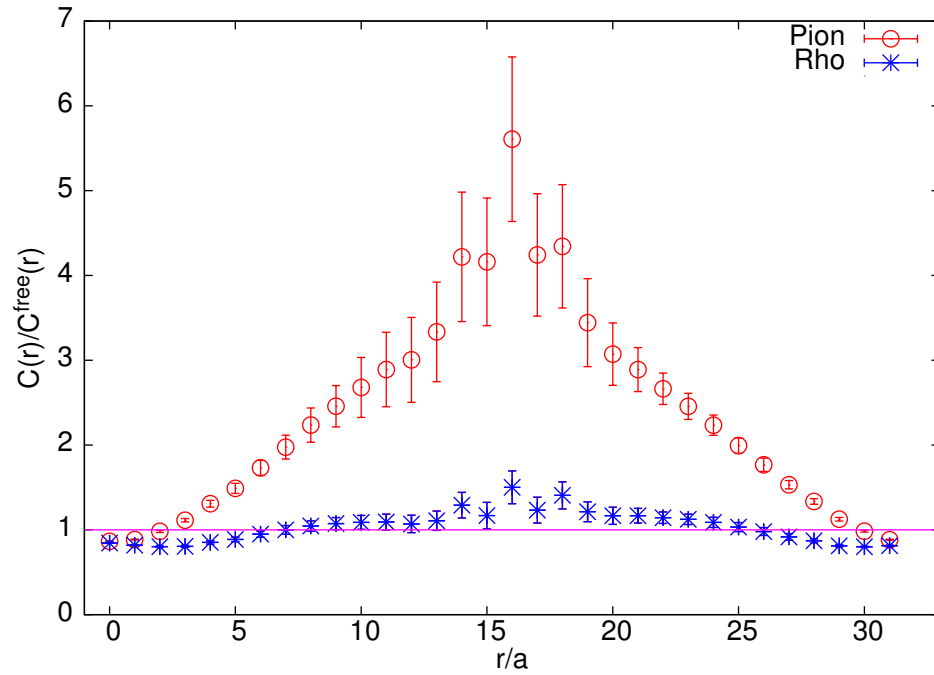
♣ Ideal gas correlator very close in each case.

♣ Pion seems to deviate from FFT much more than rho.

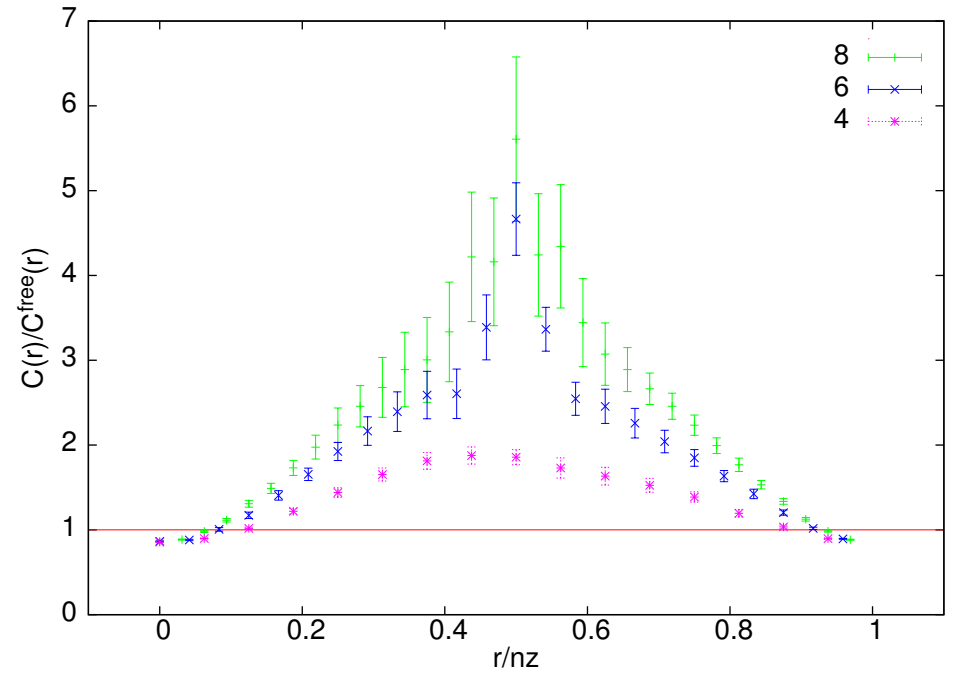
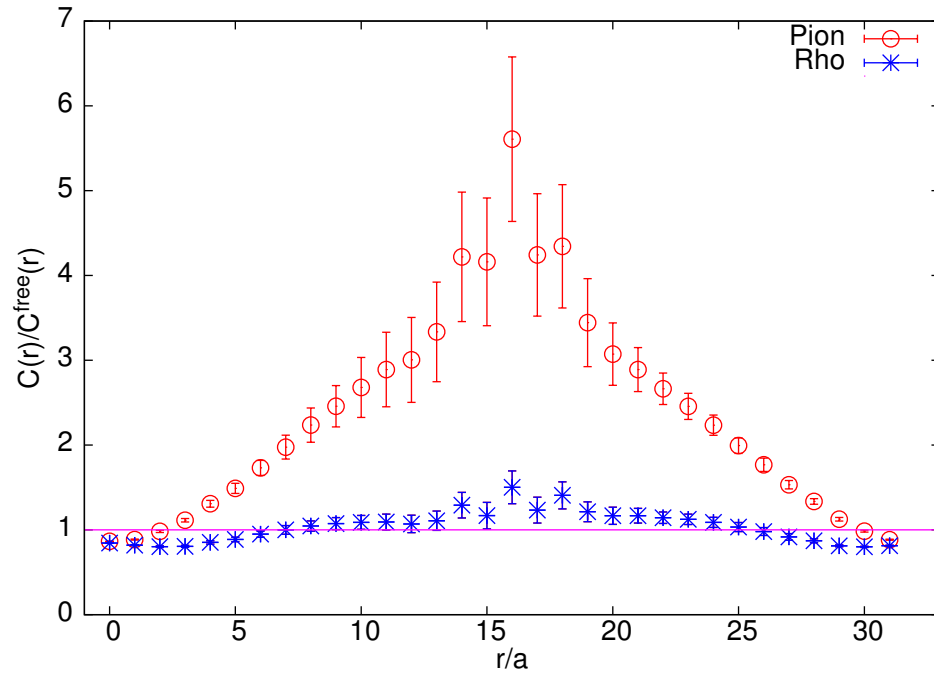
♣ Pion seems to deviate from FFT much more than rho.



♣ Pion seems to deviate from FFT much more than rho.



♣ Pion seems to deviate from FFT much more than rho.

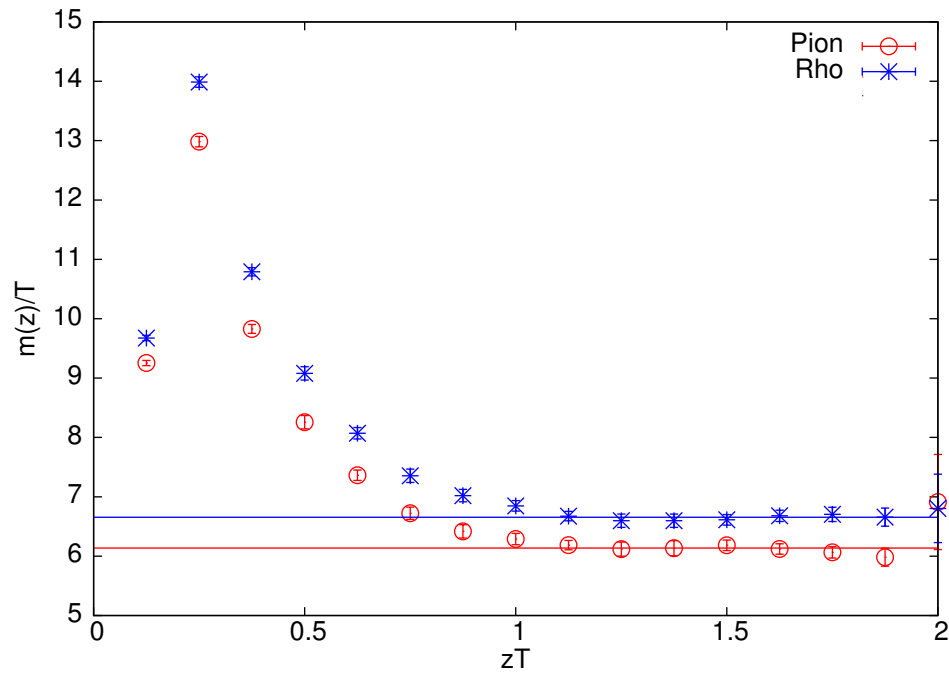


♣ As  $a$  gets smaller, the pion deviations increase.

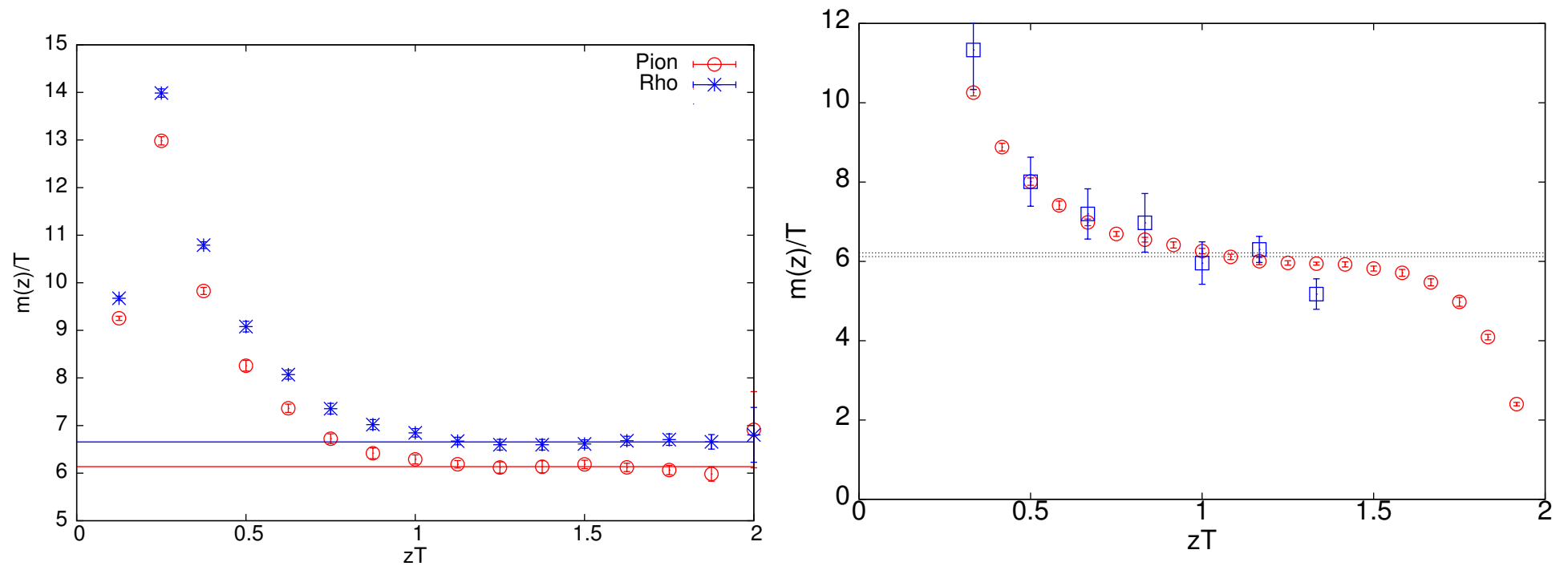


♣ Local masses [ $\sim \ln(C(r)/C(r+1))$ ] show nice plateau behaviour for pi & rho.

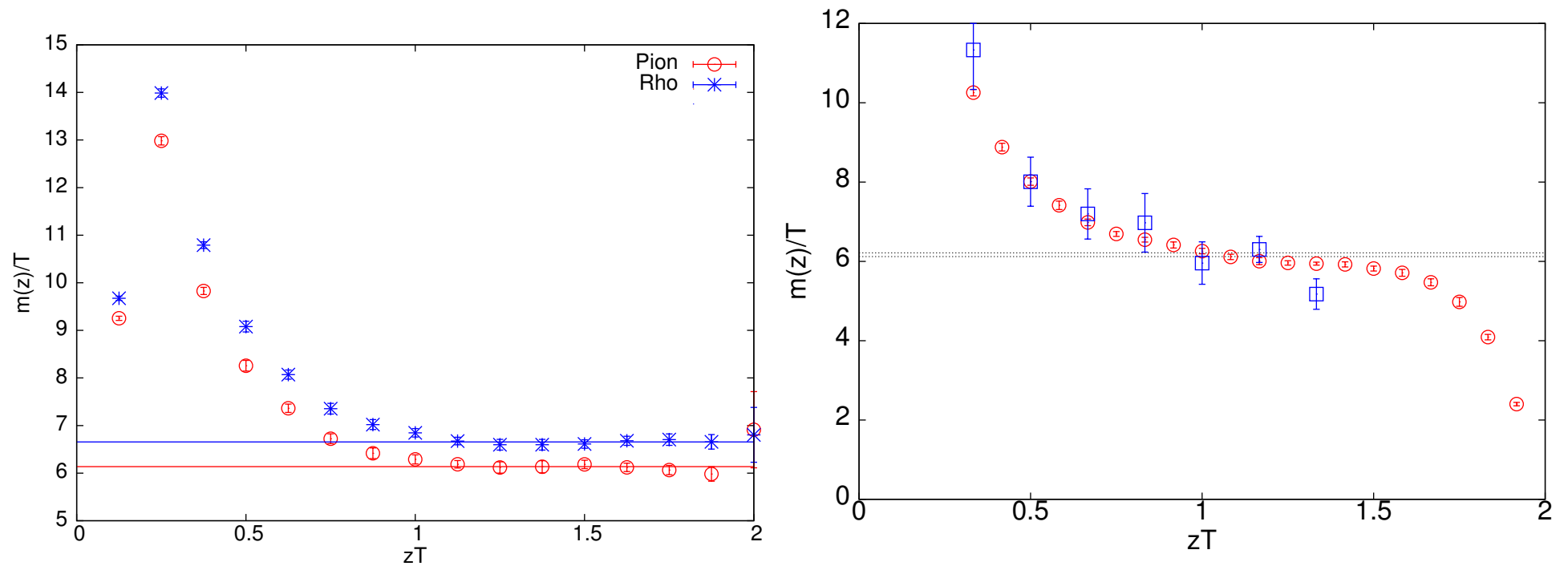
♣ Local masses [ $\sim \ln(C(r)/C(r+1))$ ] show nice plateau behaviour for pi & rho.



♣ Local masses [ $\sim \ln(C(r)/C(r + 1))$ ] show nice plateau behaviour for pi & rho.



♣ Local masses [ $\sim \ln(C(r)/C(r+1))$ ] show nice plateau behaviour for pi & rho.



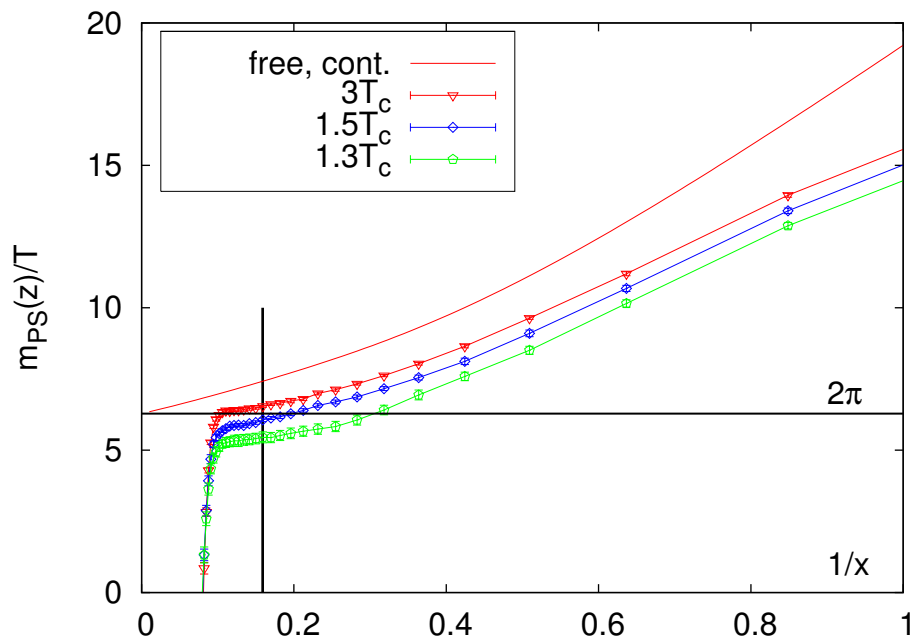
♣ Contrast this with the staggered effective mass (Gavai & Gupta PRD 2002).

# Comparison with Wilson Fermions

♣ Wilson Fermions (Figure from PoS Lattice 2005, 164. (Bielefeld Group))

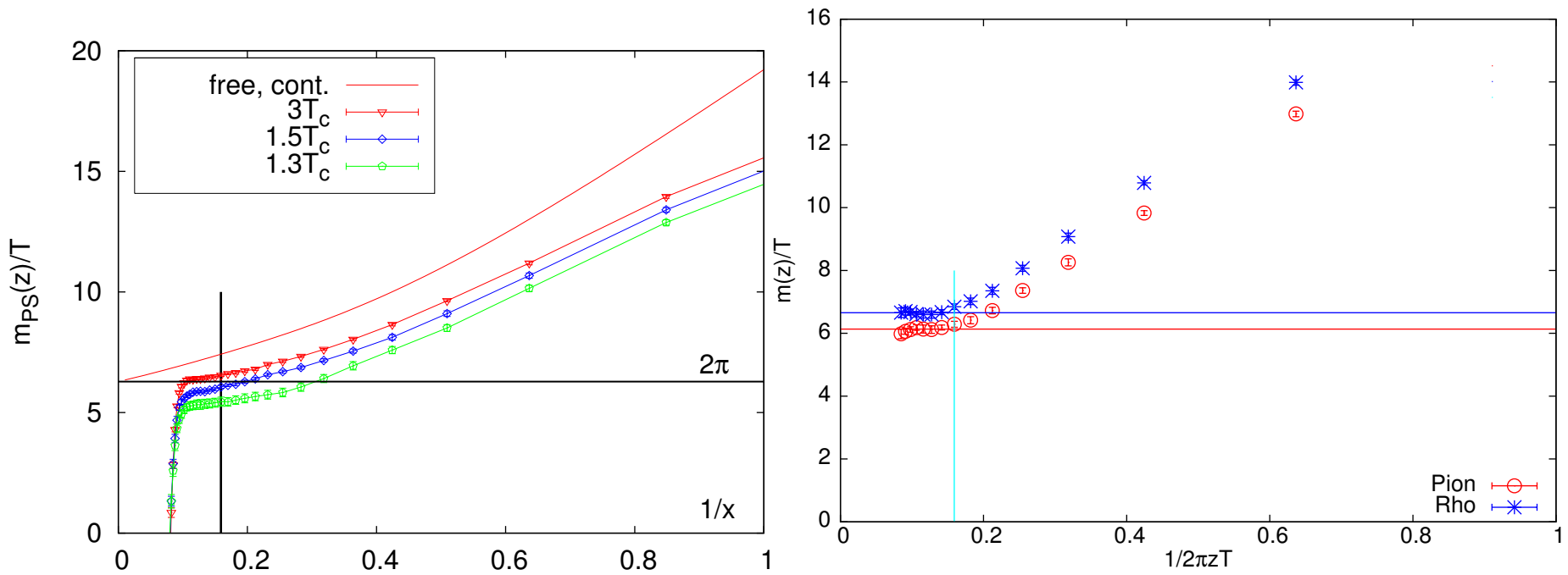
# Comparison with Wilson Fermions

♣ Wilson Fermions (Figure from PoS Lattice 2005, 164. (Bielefeld Group))



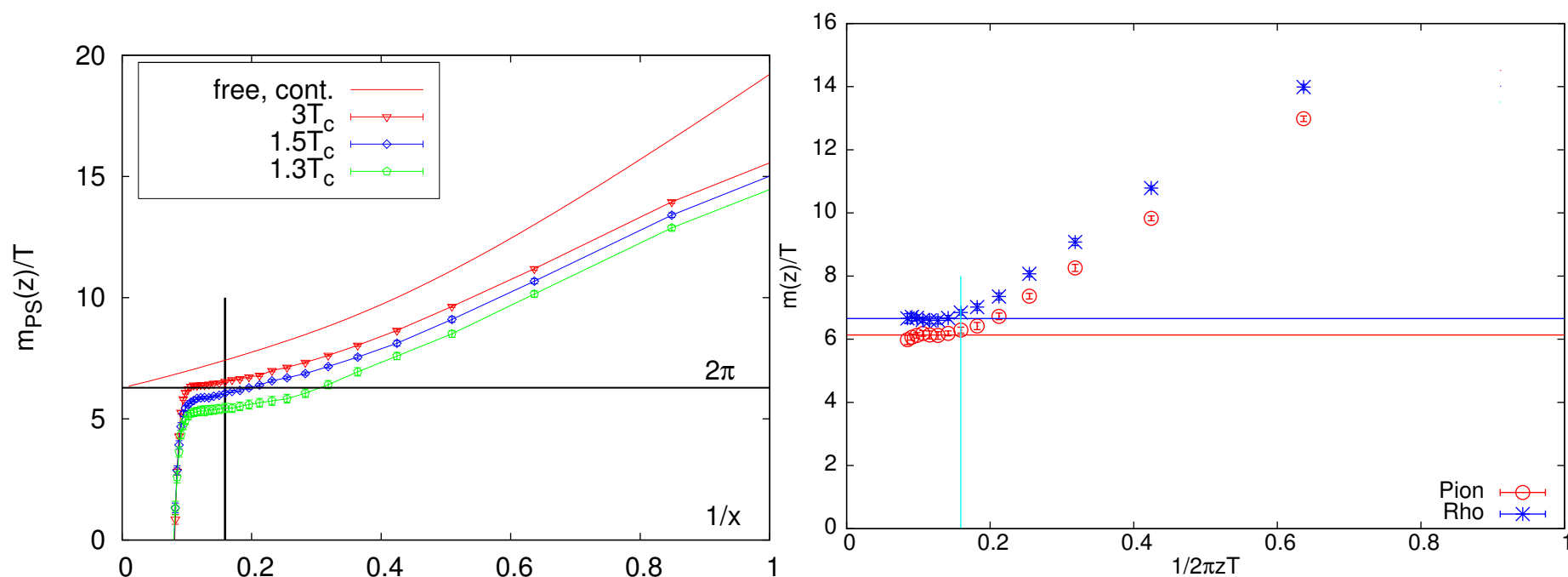
# Comparison with Wilson Fermions

♣ Wilson Fermions (Figure from PoS Lattice 2005, 164. (Bielefeld Group))



# Comparison with Wilson Fermions

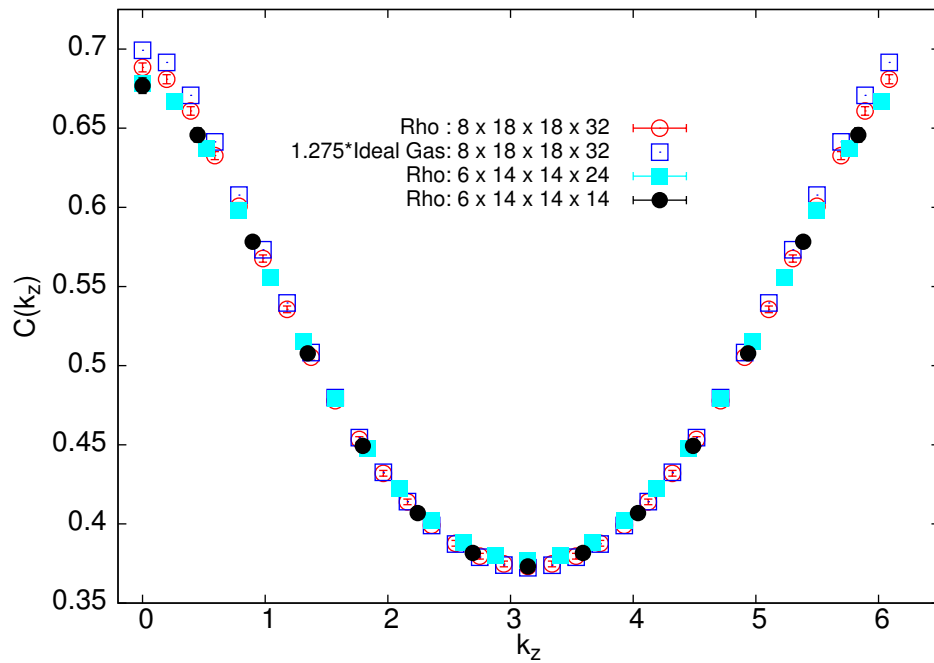
♣ Wilson Fermions (Figure from PoS Lattice 2005, 164. (Bielefeld Group))



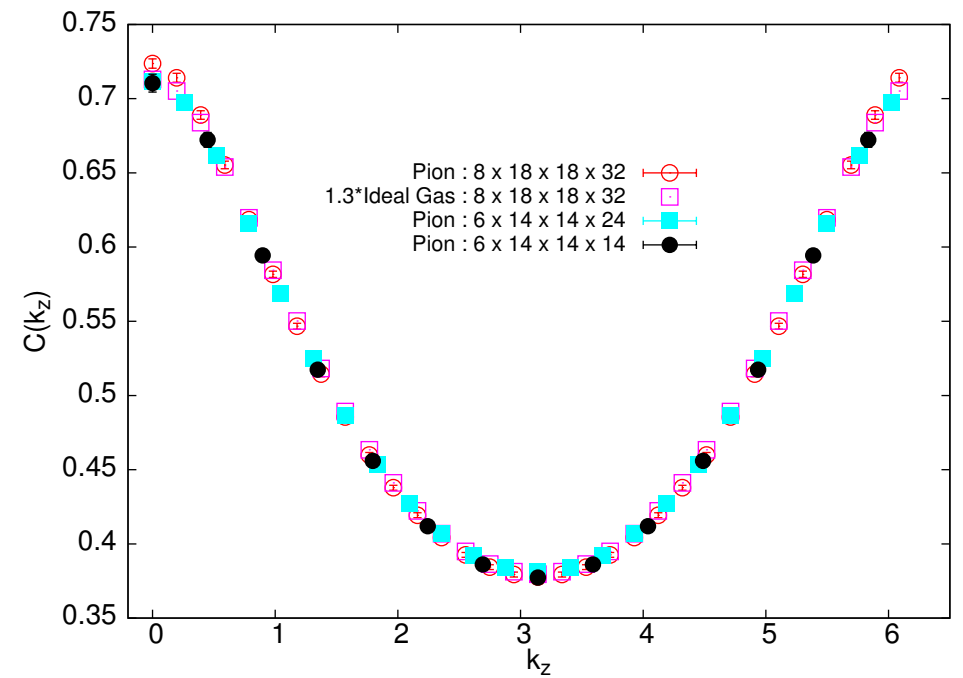
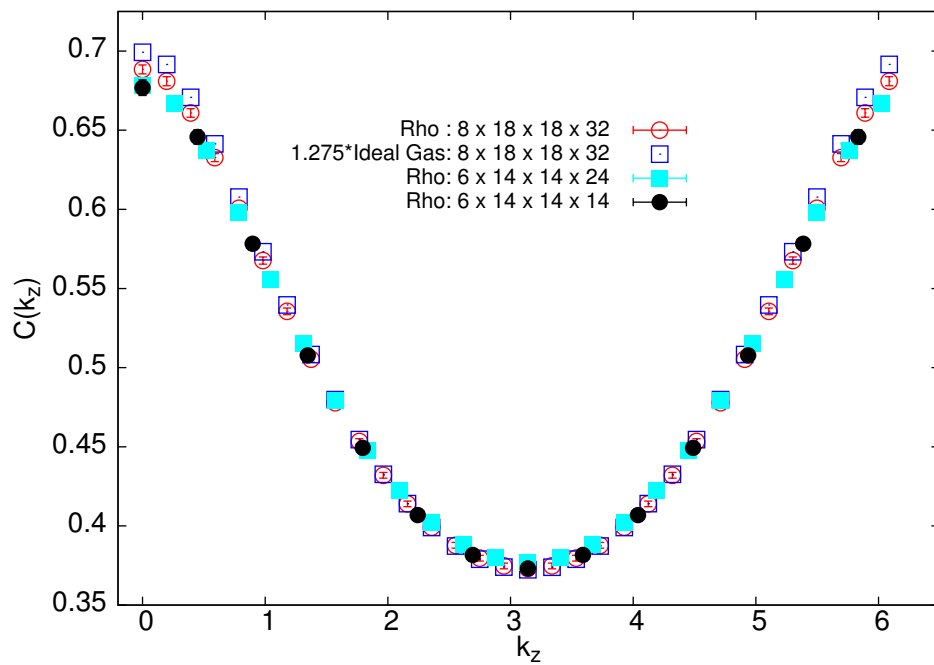
♣ Nice plateau behaviour for Overlap fermions.



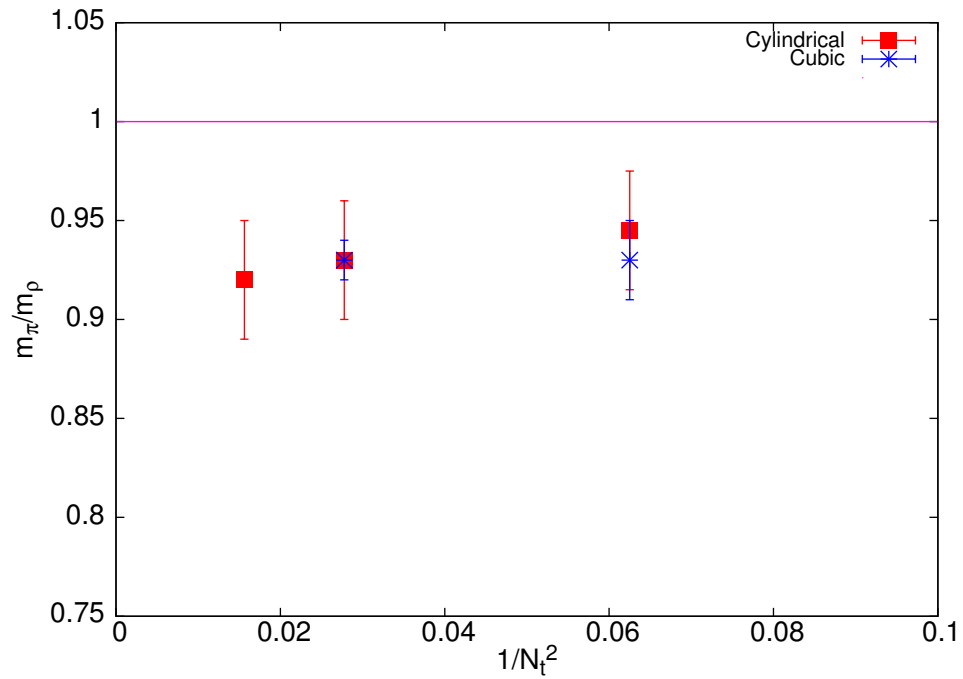
# Momentum Space Correlators



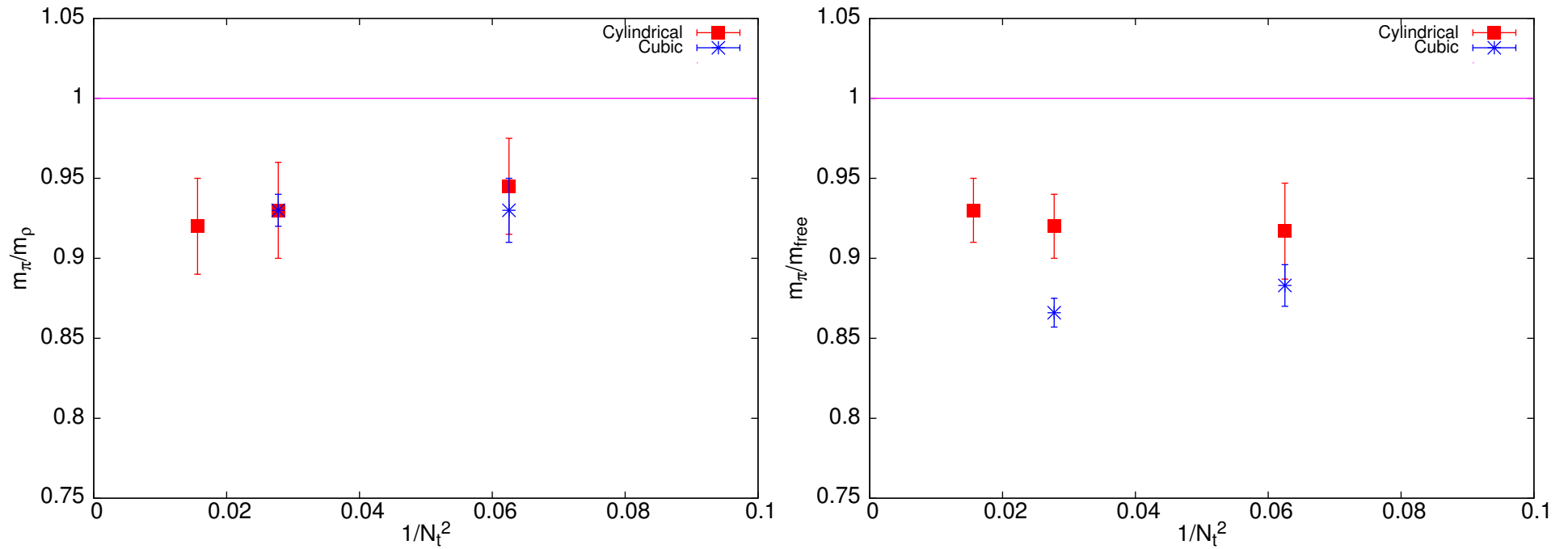
# Momentum Space Correlators



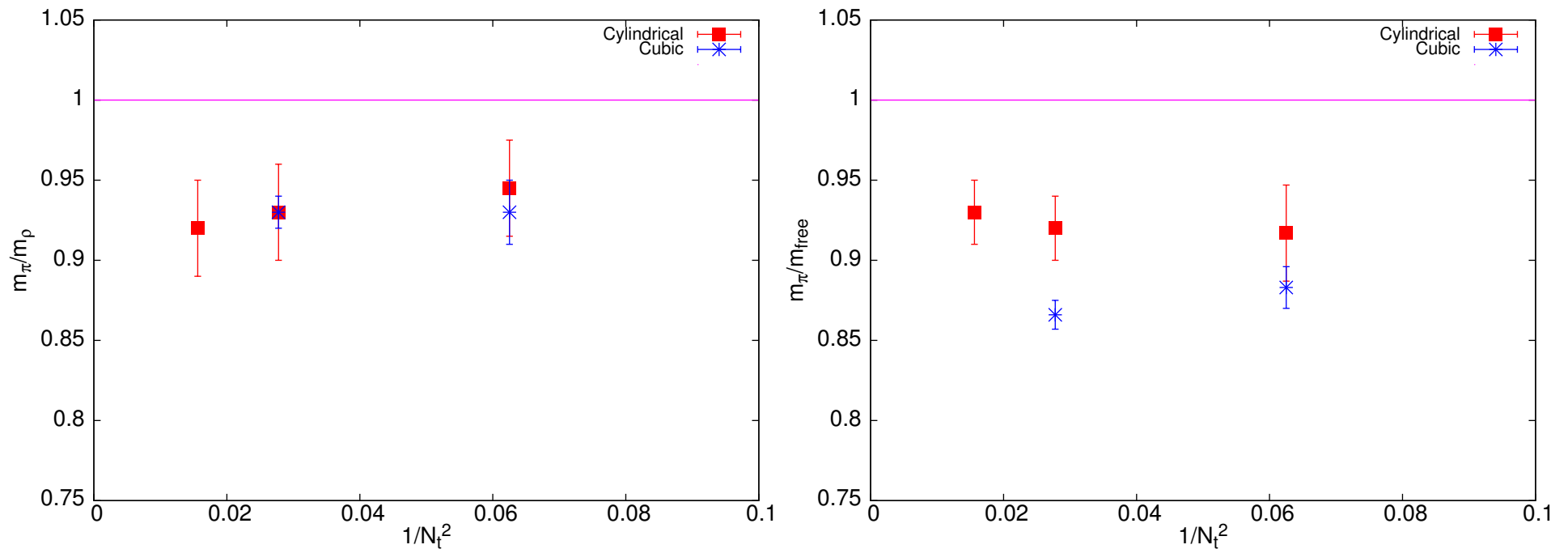
# Screening masses vs. $a$



# Screening masses vs. $a$



# Screening masses vs. $a$



♣ Very small  $a$  dependence.

♣  $m_\rho$  consistent with Ideal Gas but  $m_\pi$  smaller by about 10 %.

# Summary

- Single *cosh* behaviour, leading to nice plateau in local masses, seen on *ALL*  $N_t = 4, 6$  and  $8$ .
- Rho correlator in very good agreement with ideal gas one, but pion differs on all  $N_t$ ; Deviations increase in continuum limit.

# Summary

- Single *cosh* behaviour, leading to nice plateau in local masses, seen on *ALL*  $N_t = 4, 6$  and  $8$ .
- Rho correlator in very good agreement with ideal gas one, but pion differs on all  $N_t$ ; Deviations increase in continuum limit.
- Pion screening mass remained different from the ideal gas at  $\sim 10\%$  or  $3\sigma$  level, while rho mass was in agreement.
- Very little, if any,  $a$  dependence  $\implies$  difference to persist on very large  $N_t$ .

# Summary II

- Lattice QCD **predicts** transition to Quark-Gluon Plasma and several of its properties,  $T_c$ , EoS,  $\mu_\pi$ ,  $\lambda_s$ ,  $\eta$  ...
- $\pi$ -screening length appears *nontrivial* even in continuum limit.



## Summary II

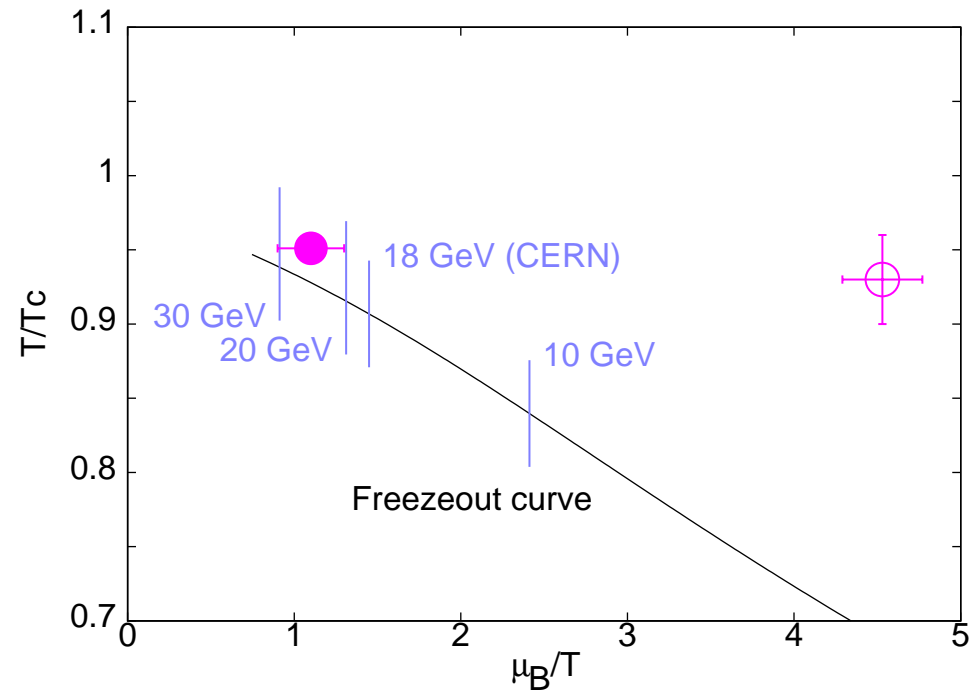
- Lattice QCD **predicts** transition to Quark-Gluon Plasma and several of its properties,  $T_c$ , EoS,  $\mu_\pi$ ,  $\lambda_S$ ,  $\eta$  ...
- $\pi$ -screening length appears *nontrivial* even in continuum limit.
- Our results on correlations of quantum numbers suggest quark-like excitations in QGP.

# Summary II

- Lattice QCD **predicts** transition to Quark-Gluon Plasma and several of its properties,  $T_c$ , EoS,  $\mu_\pi$ ,  $\lambda_S$ ,  $\eta$  ...
- $\pi$ -screening length appears *nontrivial* even in continuum limit.
- Our results on correlations of quantum numbers suggest quark-like excitations in QGP.
- Phase diagram in  $T - \mu_B$  plane has begun to emerge: Our estimate for the critical point is  $\mu_B/T \sim 1 - 2$ .

# Summary II

- Lattice QCD **predicts** transition to Quark-Gluon Plasma and several of its properties,  $T_c$ , EoS,  $\mu_\pi$ ,  $\lambda_S$ ,  $\eta$  ...
- $\pi$ -screening length appears *nontrivial* even in continuum limit.
- Our results on correlations of quantum numbers suggest quark-like excitations in QGP.
- Phase diagram in  $T - \mu_B$  plane has begun to emerge: Our estimate for the critical point is  $\mu_B/T \sim 1 - 2$ .



# Wróblewski Parameter

- Measure of Strangeness produced. Quark number susceptibilities,  $\chi_{ij} \sim \partial \ln Z / \partial \mu_i \partial \mu_j$ , can provide a handle; QNS also useful theoretical check on models.

# Wróblewski Parameter

- Measure of Strangeness produced. Quark number susceptibilities,  $\chi_{ij} \sim \partial \ln Z / \partial \mu_i \partial \mu_j$ , can provide a handle; QNS also useful theoretical check on models.
- Fluctuation-Dissipation Theorem  $\longrightarrow$  Production of Strange quark-antiquark pair  $\sim$  imaginary part of generalized strange quark susceptibility.
- Kramers - Krönig relation can be used to relate it to the real part of the susceptibility, which we obtain from lattice QCD simulations.

# Wróblewski Parameter

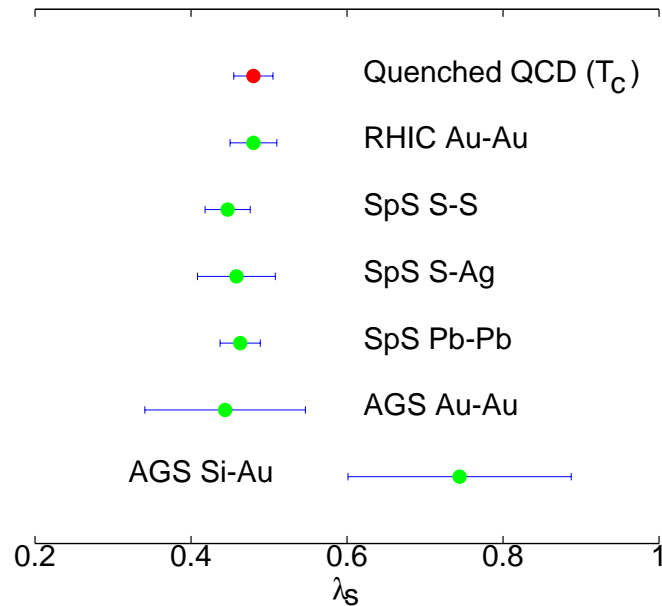
- Measure of Strangeness produced. Quark number susceptibilities,  $\chi_{ij} \sim \partial \ln Z / \partial \mu_i \partial \mu_j$ , can provide a handle; QNS also useful theoretical check on models.
- Fluctuation-Dissipation Theorem  $\longrightarrow$  Production of Strange quark-antiquark pair  $\sim$  imaginary part of generalized strange quark susceptibility.
- Kramers - Krönig relation can be used to relate it to the real part of the susceptibility, which we obtain from lattice QCD simulations.
- Finally, make a relaxation time approximation ( $\omega\tau \gg 1$ )  $\rightsquigarrow$  ratio of real parts is the same as the ratio of imaginary parts.

We use  $m/T_c = 0.03$  for  $u, d$  and  $m/T_c = 1$  for  $s$  quark;  
At each  $T$ , ratio of  $\chi_s$  and  $\chi_{ud} \rightarrow \lambda_s(T)$ .

Extrapolate it to  $T_c$ . (RVG & Sourendu Gupta, PRD 2002, PRD 2003 and PRD 2006)

We use  $m/T_c = 0.03$  for  $u, d$  and  $m/T_c = 1$  for  $s$  quark;  
 At each  $T$ , ratio of  $\chi_s$  and  $\chi_{ud} \rightarrow \lambda_s(T)$ .

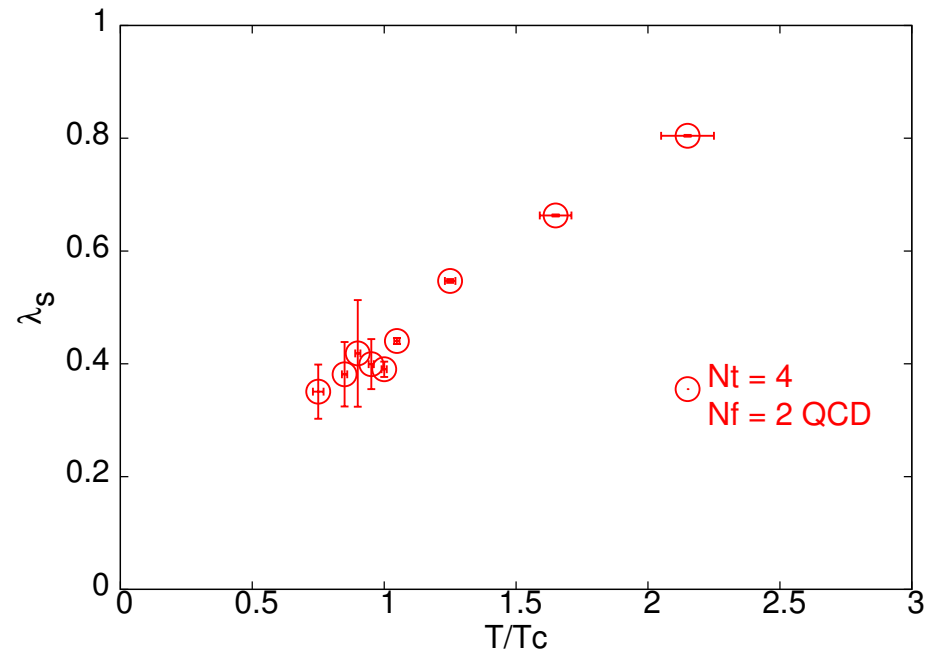
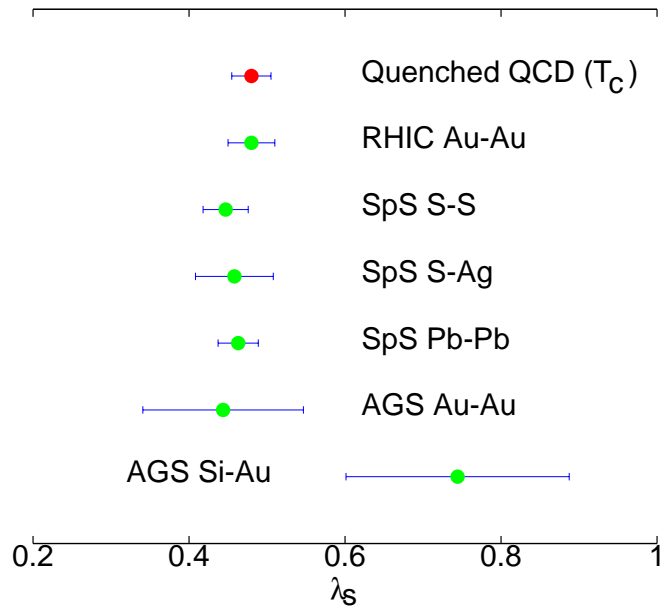
Extrapolate it to  $T_c$ . (RVG & Sourendu Gupta, PRD 2002, PRD 2003 and PRD 2006)

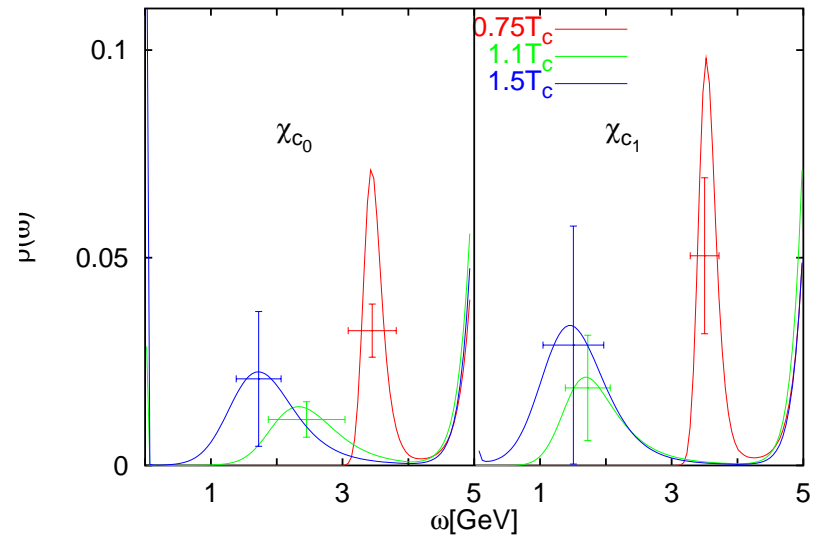
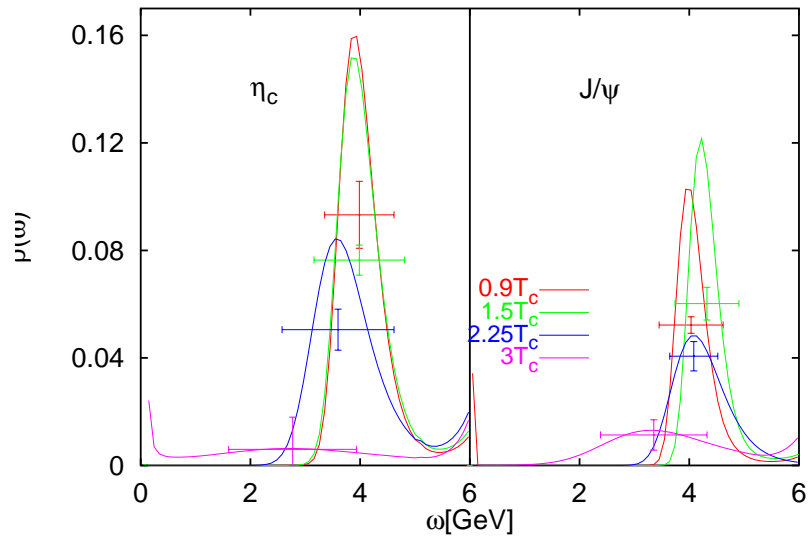




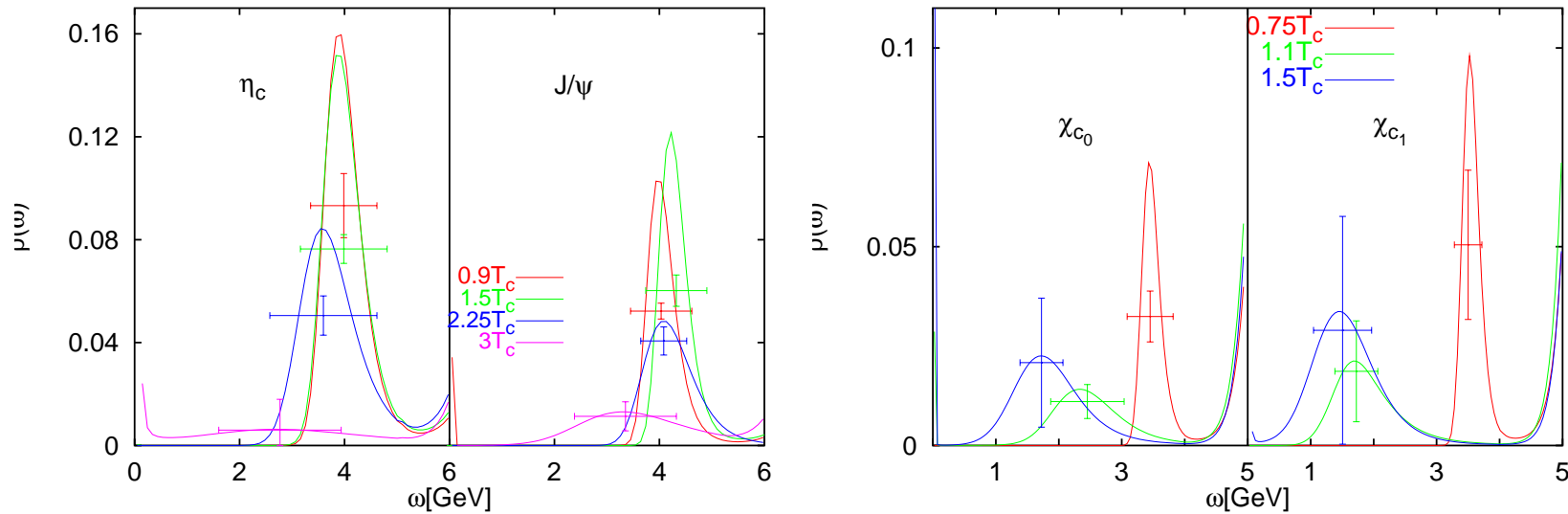
We use  $m/T_c = 0.03$  for  $u, d$  and  $m/T_c = 1$  for  $s$  quark;  
 At each  $T$ , ratio of  $\chi_s$  and  $\chi_{ud} \rightarrow \lambda_s(T)$ .

Extrapolate it to  $T_c$ . (RVG & Sourendu Gupta, PRD 2002, PRD 2003 and PRD 2006)



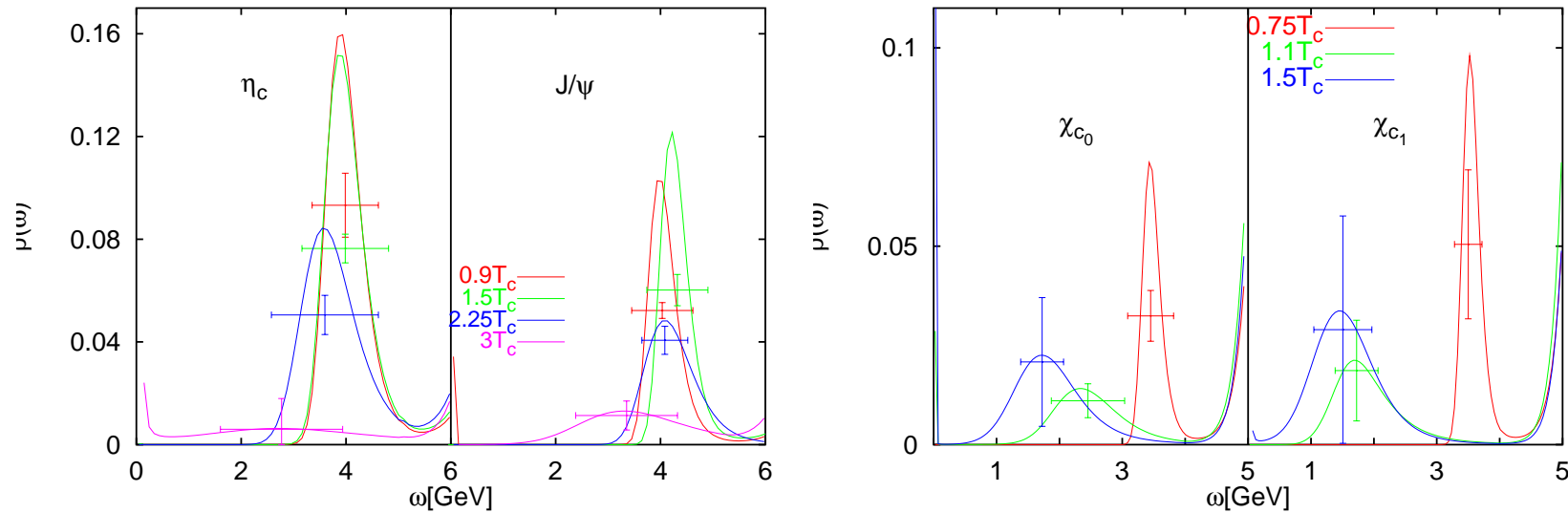


(S. Datta et al., Phys. Rev. D 69, 094507 (2004).)



(S. Datta et al., Phys. Rev. D 69, 094507 (2004).)

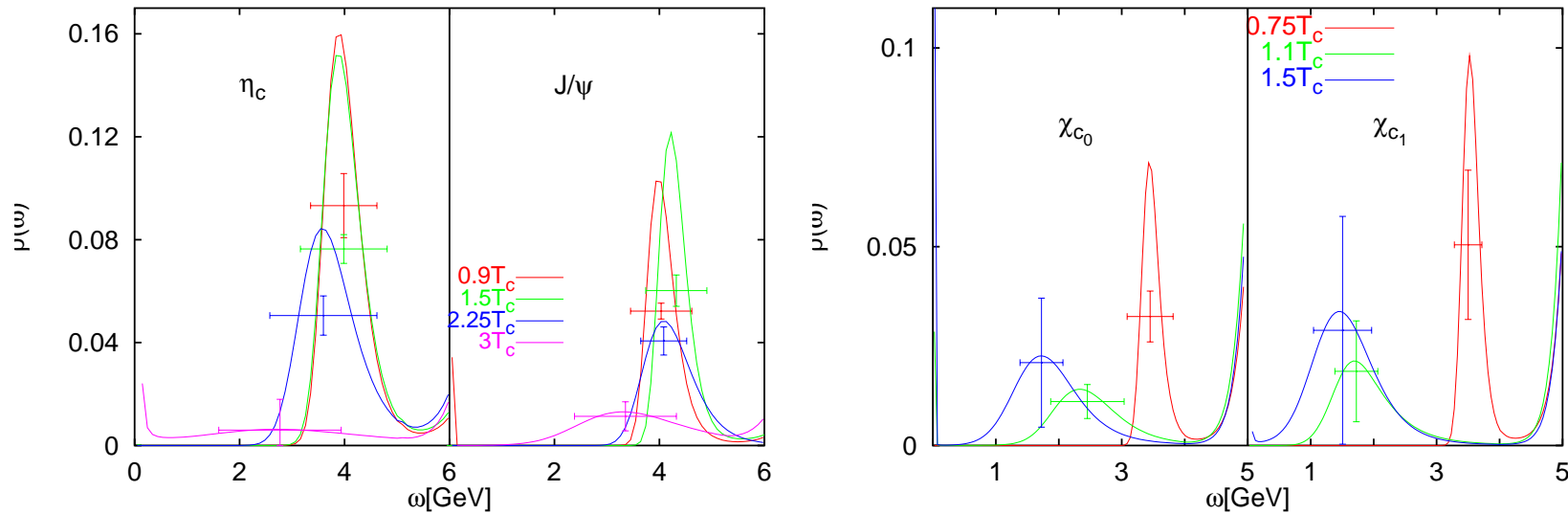
♠  $\chi_c$  seems to indeed dissolve by  $1.1T_c$ , however,  $J/\psi$  and  $\eta_c$  persist up to  $2.25T_c$  and are gone at  $3T_c$ ; Similar results by Asakawa-Hatsuda and Matsufuru.



(S. Datta et al., Phys. Rev. D 69, 094507 (2004).)

♠  $\chi_c$  seems to indeed dissolve by  $1.1T_c$ , however,  $J/\psi$  and  $\eta_c$  persist up to  $2.25T_c$  and are gone at  $3T_c$ ; Similar results by Asakawa-Hatsuda and Matsufuru.

♠ Since about 30-40 % observed  $J/\psi$  come through  $\chi$  and  $\psi'$  decays, expect changes of suppression patterns as a function of  $T$  or  $\sqrt{s}$ .



(S. Datta et al., Phys. Rev. D 69, 094507 (2004).)

♠  $\chi_c$  seems to indeed dissolve by  $1.1T_c$ , however,  $J/\psi$  and  $\eta_c$  persist up to  $2.25T_c$  and are gone at  $3T_c$ ; Similar results by Asakawa-Hatsuda and Matsufuru.

♠ Since about 30-40 % observed  $J/\psi$  come through  $\chi$  and  $\psi'$  decays, expect changes of suppression patterns as a function of  $T$  or  $\sqrt{s}$ .

♠ No Significant Effect of inclusion of dynamical fermions ?

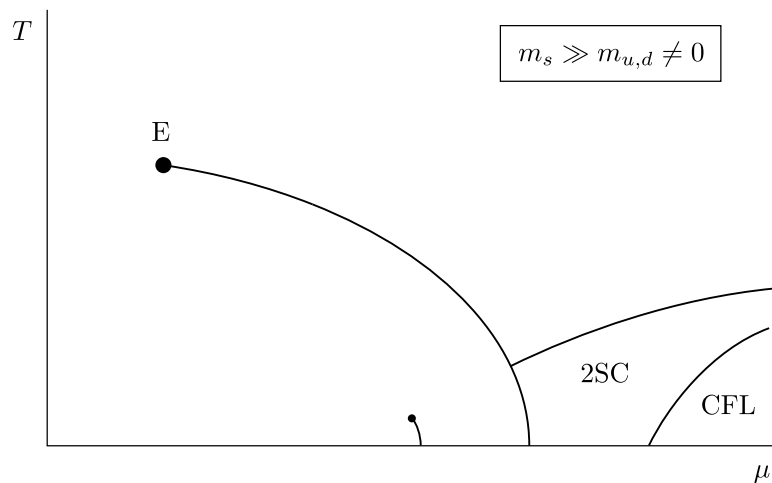
# QCD Phase diagram

♠ Another fundamental aspect – Critical Point in  $T-\mu_B$  plane; based on symmetries and models.

# QCD Phase diagram

♠ Another fundamental aspect – Critical Point in  $T-\mu_B$  plane; based on symmetries and models.

## Expected QCD Phase Diagram

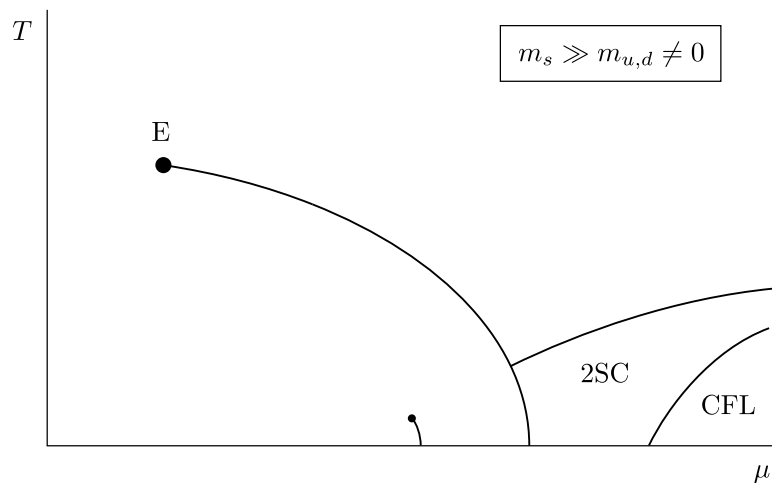


# QCD Phase diagram

♠ Another fundamental aspect – Critical Point in  $T$ - $\mu_B$  plane; based on symmetries and models.

Expected QCD Phase Diagram

... but could, however, be ...



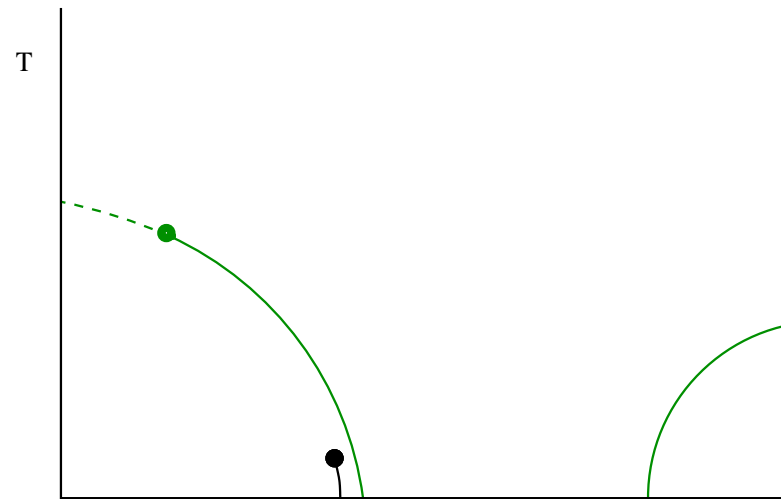
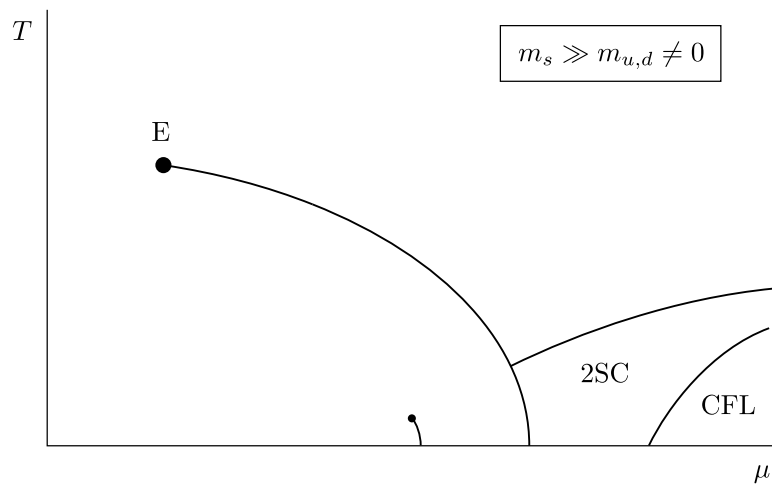


# QCD Phase diagram

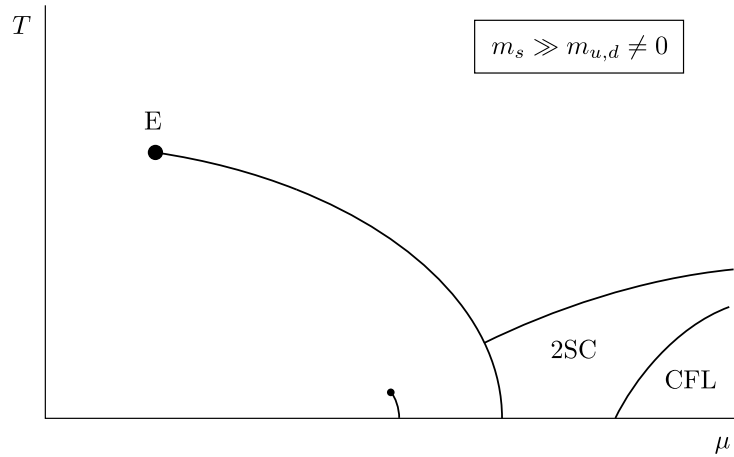
♠ Another fundamental aspect – Critical Point in  $T-\mu_B$  plane; based on symmetries and models.

Expected QCD Phase Diagram

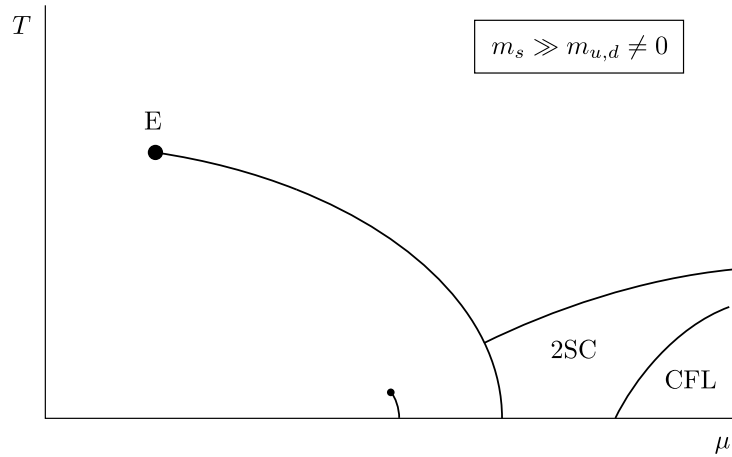
... but could, however, be ...



## Expected QCD Phase Diagram and Lattice Approaches to unravel it.

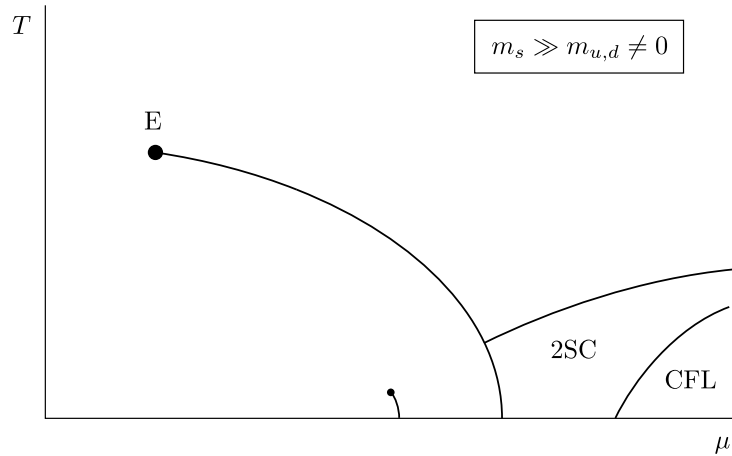


## Expected QCD Phase Diagram and Lattice Approaches to unravel it.



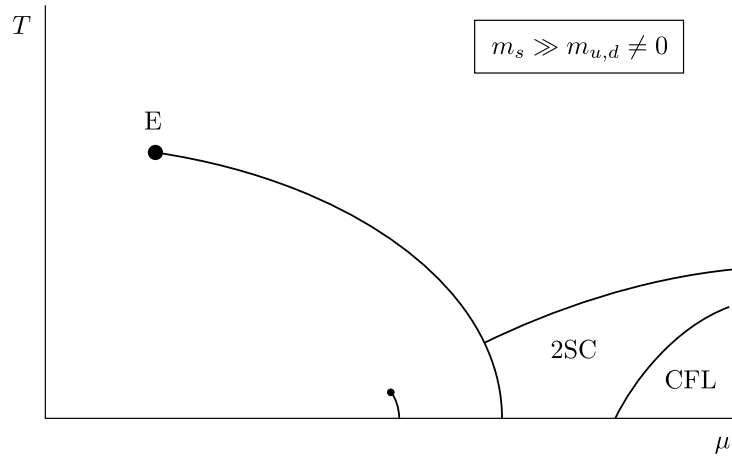
- Lee-Yang Zeroes and Two parameter Re-weighting (Z. Fodor & S. Katz, JHEP 0203 (2002) 014 ).

## Expected QCD Phase Diagram and Lattice Approaches to unravel it.



- Lee-Yang Zeroes and Two parameter Re-weighting (Z. Fodor & S. Katz, JHEP 0203 (2002) 014 ).
- Imaginary Chemical Potential (Ph. de Forcrand & O. Philipsen, NP B642 (2002) 290; M.-P. Lombardo & M. D'Elia PR D67 (2003) 014505 ).

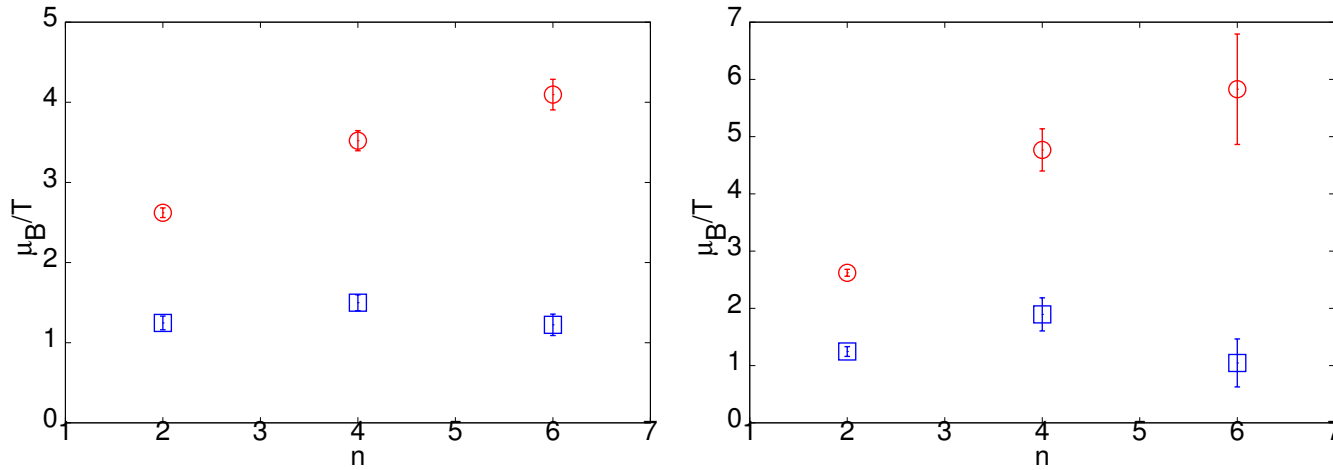
## Expected QCD Phase Diagram and Lattice Approaches to unravel it.



- Lee-Yang Zeroes and Two parameter Re-weighting (Z. Fodor & S. Katz, JHEP 0203 (2002) 014 ).
- Imaginary Chemical Potential (Ph. de Forcrand & O. Philipsen, NP B642 (2002) 290; M.-P. Lombardo & M. D'Elia PR D67 (2003) 014505 ).
- Taylor Expansion (C. Allton et al., PR D66 (2002) 074507 & D68 (2003) 014507; R.V. Gavai and S. Gupta, PR D68 (2003) 034506 ).

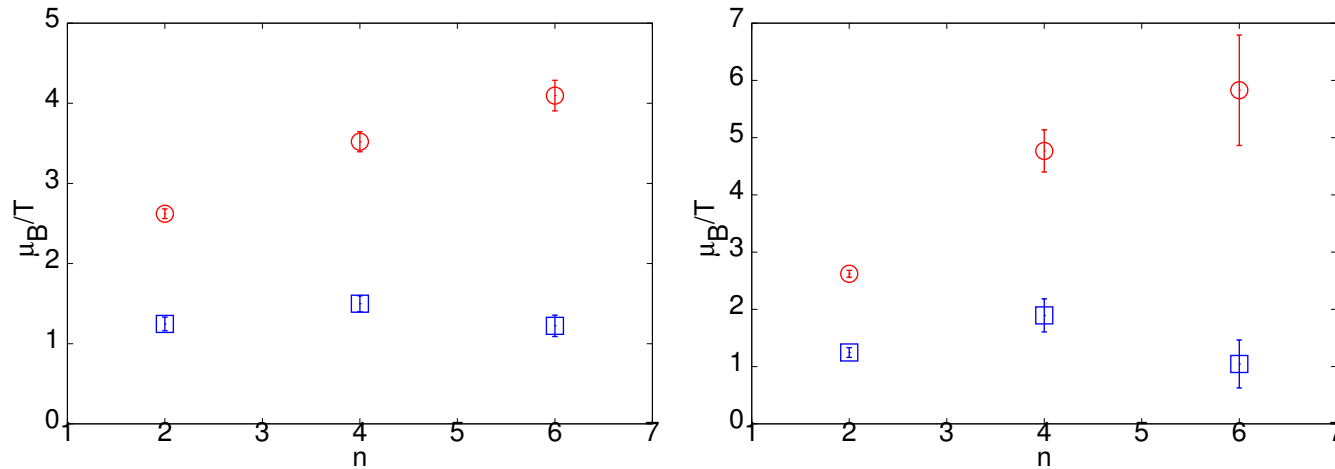
# Critical Point Estimate

RVG & S. Gupta, PR D 71 2005.



# Critical Point Estimate

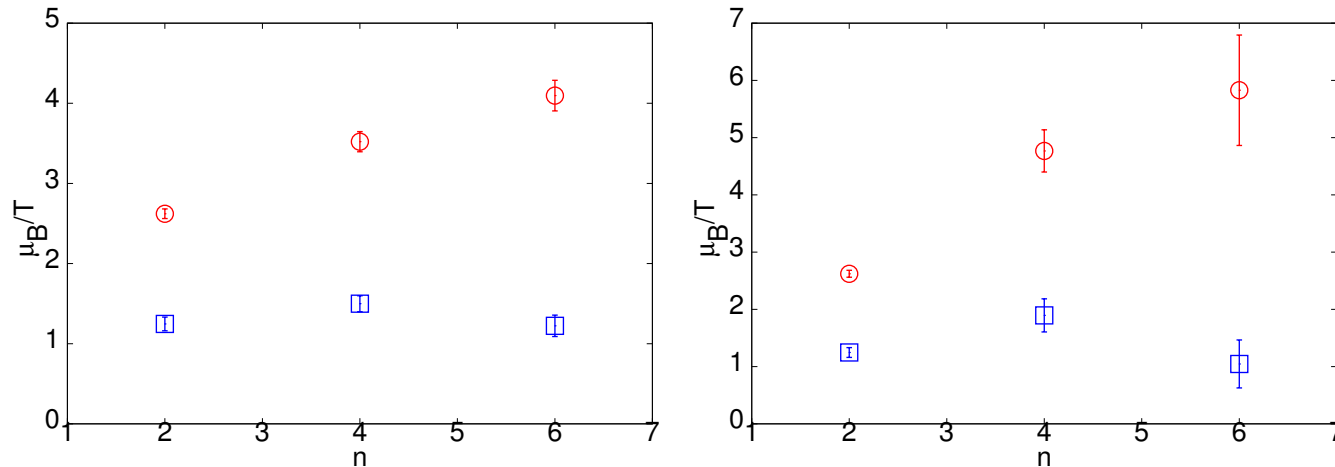
RVG & S. Gupta, PR D 71 2005.



♠ Radii of convergence as a function of the order of expansion at  $T = 0.95T_c$  on  $N_s = 8$  (circles) and 24 (boxes). Left panel for  $\rho_n$  and right one for  $r_n$ .

# Critical Point Estimate

RVG & S. Gupta, PR D 71 2005.



♠ Radii of convergence as a function of the order of expansion at  $T = 0.95T_c$  on  $N_s = 8$  (circles) and 24 (boxes). Left panel for  $\rho_n$  and right one for  $r_n$ .

♠ Extrapolation in  $n \rightsquigarrow \mu^E/T^E = 1.1 \pm 0.2$  at  $T^E = 0.95T_c$ . Finite volume shift consistent with Ising Universality class.



$m_\rho/T_c$	$m_\pi/m_\rho$	$m_N/m_\rho$	$N_s m_\pi$	flavours	$T^E/T_c$	$\mu_B^E/T^E$
5.372 (5)	0.185 (2)	—	1.9–3.0	2+1	0.99 (2)	2.2 (2)
5.12 (8)	0.307 (6)	—	3.1–3.9	2+1	0.93 (3)	4.5 (2)
5.4 (2)	0.31 (1)	1.8 (2)	3.3–10.0	2	0.95 (2)	1.1 (2)
5.4 (2)	0.31 (1)	1.8 (2)	3.3	2	—	—
5.5 (1)	0.70 (1)	—	15.4	2	—	—

Table 1: Summary of critical end point estimates—the lattice spacing is  $a = 1/4T$ .  $N_s$  is the spatial size of the lattice and  $N_s m_\pi$  is the size in units of the pion Compton wavelength, evaluated for  $T = \mu = 0$ . The ratio  $m_\pi/m_K$  sets the scale of the strange quark mass.

Results are sequentially from Fodor-Katz '04, Fodor-Katz '02, Gavai-Gupta, de Forcrand- Philippsen and Bielefeld-Swansea.