

Recap

- moment generating fn^c : $\langle e^{\lambda x} \rangle$ with $\lambda \in \mathbb{R}$
- Characteristic fn^c : $\langle e^{ikx} \rangle$
- Cumulant gen fn^c : $\log \langle e^{\lambda x} \rangle$

Q. Is moment gen. fn^c unique and analytic?

IS exist for $\lambda \in [-\ell, \ell]$ for $\ell > 0$ then
it is unique and analytic.

[Example when not unique.

$$\text{take } p(x) = \frac{1}{x^{1+\alpha}}$$

$$\text{then } \langle e^{\lambda x} \rangle = \begin{cases} 1 & \text{for } \lambda = 0 \\ \infty & \text{for } |\lambda| > 0 \end{cases}$$

for all α .

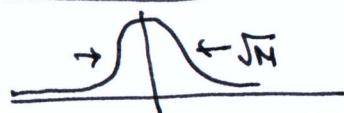
Central limit theorem.

$$M_N = \frac{\sum x_i}{N} \quad \text{where } x_i \text{ are iid with finite } \langle x \rangle \text{ and } \langle x^2 \rangle_c$$

Then,

$$\mathcal{P} \left(\frac{M_N - N \langle x \rangle}{\sqrt{N}} = z \right) \xrightarrow{N \rightarrow \infty} \frac{1}{\sqrt{2\pi \langle x^2 \rangle_c}} e^{-\frac{z^2}{2 \langle x^2 \rangle_c}}$$

$$\Rightarrow P(M_N) \approx \frac{1}{\sqrt{2\pi N \langle x^2 \rangle_c}} e^{-\frac{(M - N \langle x \rangle)^2}{2N \langle x^2 \rangle_c}}$$

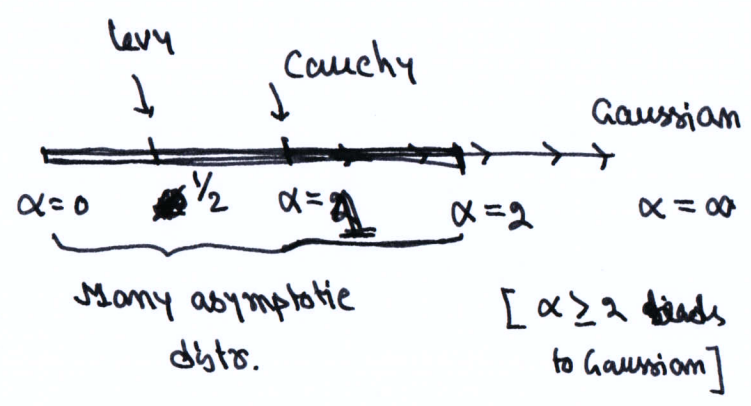


* What happens for $p(x)$ where variance diverge?

Example: $p(x) \sim \frac{1}{x^{3/2}}$ for $x \gg 1$. [first passage prob]

Is there any asymptotic distribution $M = \sum_{i=1}^N x_i$?

Answer: There are a large number of asymptotic distributions. (a two-parameter family)



Naive argument:

Let's redo the argument for Gaussian case.

Let $p(x)$ have finite variance. $\langle x^2 \rangle_c$

Then,

$$G_N(\lambda) = \int dx e^{i\lambda M} P_N(M) = \langle e^{i\lambda M} \rangle$$

← note the i.

$$= \langle e^{i\lambda x} \rangle^N$$

← characteristic function.

$M = \sum_{i=1}^N x_i$

$$\Rightarrow G_N(\lambda) = [g(\lambda)]^N$$

$g(\lambda) = \int dx e^{i\lambda x} p(x)$

$$= e^{iN\lambda \langle x \rangle} \cdot e^{-\frac{N\lambda^2}{2} \langle x^2 \rangle_c} \cdot e^{iN\lambda \frac{\lambda^3}{3} \langle x^3 \rangle_c + \dots}$$

(set zero for simplicity)

Put $\lambda = \frac{k}{\sqrt{N}}$

$$\Rightarrow G_N\left(\frac{k}{\sqrt{N}}\right) = e^{-\frac{k^2}{2} \langle x^2 \rangle_c} + o\left(\frac{1}{\sqrt{N}}\right)$$

$N \rightarrow \infty \rightarrow \hat{G}(k)$ ← characteristic for Gaussian

⇒ This means

$$P_N(M) = \frac{1}{2\pi} \int d\lambda e^{-i\lambda M} G_N(\lambda) \quad \text{Inverse Fourier.}$$

$$\approx \frac{1}{\sqrt{N}} \cdot \frac{1}{2\pi} \int dk e^{-ik \cdot \frac{M}{\sqrt{N}}} \hat{G}(k) \quad \text{large } N$$

$\lambda = \frac{k}{\sqrt{N}}$

$$= \frac{1}{\sqrt{N}} \cdot \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{M^2}{2N\sigma^2}} \quad [\sigma^2 = \langle x^2 \rangle]$$

⇒ $P\left(\frac{M}{\sqrt{N}} = z\right) = f(z)$ Central limit theorem.
↑ Gaussian

Remark: it does not matter if $p(x)$ is symm or not. OR how it looks!

Now, what if

$$g(\lambda) \approx e^{-c|\lambda|^\alpha + \dots}$$

↑ non-analytic at $\lambda \rightarrow 0$.

[this is what happens for $p(x) \sim \frac{1}{x^{1+\alpha}}$ with $\alpha \leq 2$]

Example: $p(x) \approx \frac{c}{x^2} \Rightarrow g(\lambda) = e^{-c|\lambda|} \text{ for } \lambda \rightarrow 0$

Repeat the exercise

$$G_N(\lambda) \approx e^{-cN|\lambda|^\alpha + \Theta(N \cdot \lambda^{1+\alpha})} \leftarrow (g(\lambda))^N$$

$\lambda = \frac{k}{\sqrt{N}}$

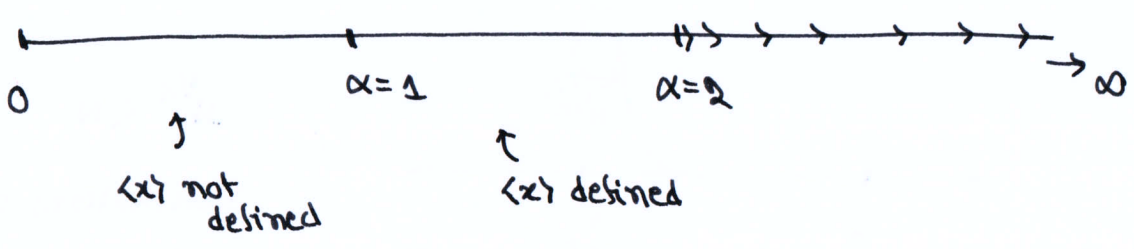
$$\Rightarrow G_N\left(\frac{k}{N^{1/\alpha}}\right) = e^{-c|k|^\alpha + \Theta\left(\frac{1}{N^{1/\alpha}}\right)}$$

⇒ $P\left(\frac{M}{N^{1/\alpha}} = z\right) = \mathcal{L}_{\alpha, \beta}(z)$ = $\frac{1}{2\pi} \int dk e^{-ikz} e^{-c|k|^\alpha}$
↑ very special case.

The general mathematical result:

see back of page 7.

what about $\alpha < 0$?



Need different scaling

flip page

$$z = \begin{cases} \frac{M_N - N \langle x \rangle}{N^{1/\alpha}} & \text{for } 1 < \alpha \leq 2, \\ \frac{M_N}{N^{1/\alpha}} & \text{for } 0 < \alpha \leq 1. \end{cases}$$

Then

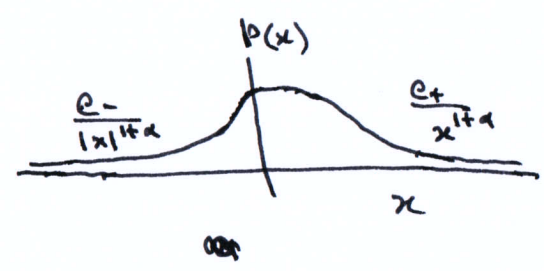
$$P(z) = \alpha_{\alpha, \beta}(z) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk e^{-ikz} \hat{G}_{\alpha, \beta}(k)$$

with $\hat{G}_{\alpha, \beta}(k) = e^{-|c_{\pm} k|^{\alpha}} \{ 1 + i \beta \text{sign}(k) \cdot \delta \} + ik \langle x \rangle$

Here α comes from $p(x) \sim \begin{cases} c_- / |x|^{1+\alpha} & \text{for } x \rightarrow -\infty \\ c_+ / |x|^{1+\alpha} & \text{for } x \rightarrow \infty \end{cases}$

$$\beta = \frac{c_+ - c_-}{c_+ + c_-}$$

$$\delta = \begin{cases} \tan \frac{\pi \alpha}{2} & \text{for } \alpha \neq 1 \\ -\frac{2}{\pi} \log |k| & \text{for } \alpha = 1 \end{cases}$$



* Only tails matter!

Remark: see, unlike in CLT, how the tail decides the asymptotic distribution. This is "black swan" phenomena, where rare events have dramatic consequences.

These are no explicit formula for the asymptotic distribution $P(z)$ for ~~an~~ arbitrary α, β . Can be written as an asymptotic series [see Kardar 1st book, eq 2.52]

~~Explicit results for~~

Explicit results for

** we choose $\langle x \rangle = 0$ in all cases **

① $\alpha = 2$: ~~($\beta = 0$)~~

$\Rightarrow \delta = 0 \Rightarrow \hat{G}(k) = e^{-|c||k|^2}$ (β does not matter)

$\Rightarrow P(z)$ is Gaussian with variance $\sigma^2 = 2c^2$

Imp Remark: note that for $p(x) \sim \frac{1}{x^3}$ variance does not exist!! still $\sigma^2 = 2e^2$.

② $\alpha = 1$ and $\beta = 0$ (only possible if $c_+ = c_- \Rightarrow$ symmetric $p(x)$)

$\langle z \rangle$ finite
but
 $\langle z^2 \rangle \rightarrow \infty$

$\hat{G}(k) = e^{-|c||k|}$

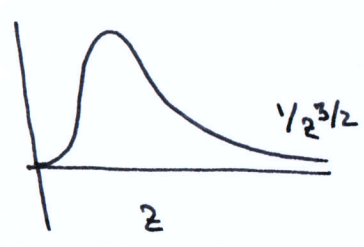
$\Rightarrow P(z) = \frac{1}{\pi} \frac{c}{z^2 + c^2}$ Cauchy distribution.

③ $\alpha = \frac{1}{2}$ and $\beta = 1$ (true if $c_- = 0 \Rightarrow x \in [0, \infty)$)

$\Rightarrow \delta = 1 \Rightarrow \hat{G}(k) = e^{-|c||k|^{\frac{1}{2}}} (1 + i \text{sign}(k))$

even $\langle z \rangle \rightarrow \infty$.

$\Rightarrow P(z)$ is Levy distribution (* comes in first passage time)



$$P(z) = \begin{cases} \sqrt{\frac{c}{2\pi}} \frac{e^{-c/2z}}{z^{3/2}} & \text{for } z \geq 0 \\ 0 & \text{for } z < 0 \end{cases}$$

Remark % One way to convince/test this asymptotic formula of $P(z)$ is by checking their stability. under operation $\sum_i x_i / b(N)$

These distributions are like stable fixed points, ~~are~~ in space of distributions.

They are also called "stable distributions".

Simple examples % (1) Gaussian is a stable distribution.

$$p(x) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{x^2}{2\sigma^2}} \rightarrow g(\lambda) = \int dx e^{i\lambda x} p(x) = \langle e^{i\lambda x} \rangle = e^{-\frac{\sigma^2}{2} \lambda^2}$$

$$\Rightarrow G(\lambda) = \langle e^{\lambda M} \rangle = [g(\lambda)]^N = e^{-\frac{N\sigma^2}{2} \lambda^2}$$

$$\Rightarrow P(M) = \frac{1}{\sqrt{2\pi N\sigma^2}} e^{-\frac{M^2}{2N\sigma^2}}$$

Only variance changed!

leave as exercise:

(2) ~~Cauchy~~ Cauchy distribution.

ux contour integration.

$$p(x) = \frac{1}{\pi} \frac{c}{(x-x_0)^2 + c^2} \rightarrow g(\lambda) = e^{i\lambda x_0 - c|\lambda|}$$

$$\Rightarrow G(\lambda) = e^{i\lambda(Nx_0) - (Nc)|\lambda|}$$

$$\rightarrow P(M) = \frac{1}{\pi} \frac{Nc}{(M-Nx_0)^2 + (Nc)^2}$$

③ Lévy distribution

$$p(x) = \sqrt{\frac{c}{2\pi}} \frac{e^{-c/2x}}{x^{3/2}} \quad \text{for } x \geq 0$$

$$\longrightarrow g(\lambda) = e^{-|c\lambda|^{1/2}} (1 - i \operatorname{sign}(\lambda))$$

$$\Rightarrow a(\lambda) = e^{-|N^2 c \cdot \lambda|^{1/2}} (1 - i \operatorname{sign}(\lambda))$$

$$\Rightarrow P(M) = \sqrt{\frac{N^2 c}{2\pi}} \frac{e^{-\frac{N^2 c}{2M}}}{M^{3/2}} \quad \text{for } M \geq 0.$$

④ For general

$$g_{\alpha, \beta}(\lambda) = e^{i\lambda \langle x \rangle - |c\lambda|^\alpha (1 - i\beta \operatorname{sign}(\lambda) \cdot \delta)}$$

$$\Rightarrow a_{\alpha, \beta}(\lambda) = e^{i(N\lambda) \langle x \rangle - |(N^{\frac{1}{\alpha}} c) \lambda|^\alpha (1 - i\beta \operatorname{sign}(\lambda) \delta)}$$

Term paper : describe CLT, stable distributions using R.G.

See (1) book by Sethna.

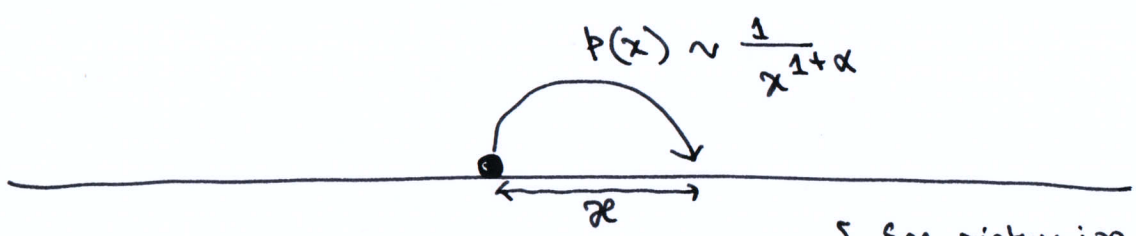
(2) An article by Jona-Lasinio -

Q. Where are these ideas of limit distributions important?

In normal and anomalous diffusion in nature.

Team paper: "Beyond Brownian Motion" Klafter et al, Phys. today 49, 33 (1996).

Simplest example: Random walk (discrete time, cont space)



[See picture in a paper by Ariel Amir

Position after N steps

$$M_N = \sum_{i=1}^N x_i$$

Case 1: normal diffusion. ($\alpha \geq 2$) including Gaussian.

$$\langle M_N^2 \rangle - \langle M_N \rangle^2 \sim N \Rightarrow P(M_N) \sim \frac{1}{\sqrt{2\pi N\sigma^2}} e^{-\frac{(M_N - \langle M_N \rangle)^2}{2N\sigma^2}}$$

and $P\left(\frac{M_N - \langle M_N \rangle}{\sqrt{N}} = z\right) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{z^2}{2\sigma^2}}$

typical displacement $\sim \sqrt{N}$

Case 2: anomalous diffusion ($\alpha < 2$.)

* Variance of M_N does not exist.

$$* P\left(\frac{M_N}{N^{1/\alpha}} = z\right) = \alpha_{\alpha, \beta}(z)$$

typical displacement

$$M_N \sim N^{1/\alpha} > N \text{ Super-diffusion.}$$

(for sub-diffusion see later)

Natural examples: (1) normal diffusion for many examples with finite correlation time. [see Sethna]

(1) Foraging of animals.



Flight path of an albatross fits $\alpha = 1$.
bumblebees & deer.

[Ref.] Viswanathan et al. Nature 381, 413 (1996).
Edwards et al. Nature 449, 1044 (2007)

(2) Stock prices.

degraded price of a stock. $M_N = \sum_i x_i$ → Price change at high temporal resolution.
(Fast trading high frequency)

$$p(x) = \frac{e^{-\lambda|x|}}{|x|^{1+\alpha}}$$

with $\alpha \approx 1.4$

Ref. Mantegna & Stanley, Nature 376, 46 (1995)

* A crossover happens. between Gaussian & Levy

Ref. Ismo Koponen, PRE 52, 1197 (1995)

(3) Living Polymers. $\alpha = 1.86$

↑ are chains of micelles that break apart & recombine to produce a stationary polymer length distribution.

One measures diffusion of a tracer micelles among other micelles.

Ref. Ott, Bouchaud, Langervin, Urbach 65, 2201, 1990.

Why super-diffusion? Any intuition?

④ Interstellar levý flights.

Radio waves that travel through interstellar medium (clouds of ionized gas) gets deflected randomly.

Simplest picture is that ray does a random walk due to random deflection.

But there are claims that these are levý flights.

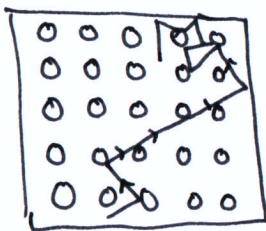
[Baldyrev & Gwinn, PRL 91, 131101 (2003)]

⑤ levý-flight in turbulence

[Phys Today article by Klatter

"Beyond Brownian motion"]

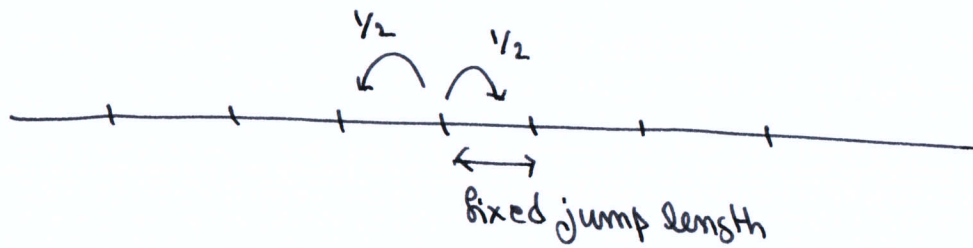
⑥ levý flight in 2-d Lorentz gas



See papers on moodle page.

Q. Any example for ~~super~~ ^{sub-} diffusion?

Continuous Time Random Walk (CTRW).



But waiting time chosen from $\beta(\tau) \sim \frac{1}{\tau^{1+\alpha}}$

In N -steps, displacement M_N

$$\Rightarrow \langle M_N^2 \rangle \sim N$$

Time in N -steps

$$t = \sum_{i=1}^N \tau_i$$

* note it is ~~not~~ also for $\alpha < 2$ (no variance!)

Case 1: Normal diffusion (for $1 < \alpha$)

$$t \approx N \langle \tau \rangle + N^{1/\alpha} \cdot \# \quad \text{for large } N.$$

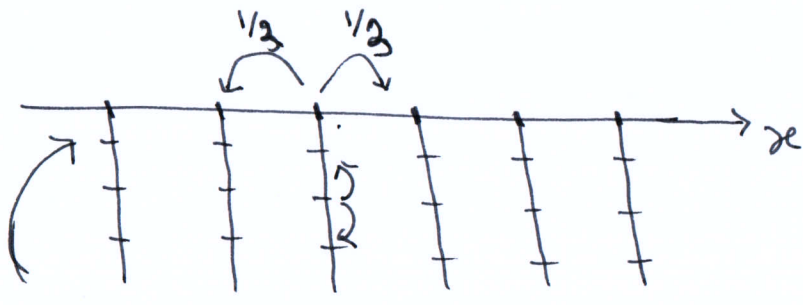
$$\Rightarrow \langle M_N^2 \rangle \approx N \sim \frac{t}{\langle \tau \rangle} \quad \leftarrow \text{normal diffusion.}$$

Case 2: ^{Sub} ~~super~~-diffusion (for $\alpha \leq 1$)

$$t \approx N^{1/\alpha} \quad \text{for large } N$$

$$\Rightarrow \langle M_N^2 \rangle \sim N \sim t^\alpha < t \quad \text{sub-diffusion.}$$

An example: Diffusion in Comb-lattice



time to return $\sim \frac{1}{z^{3/2}}$

[we will show this later. first return prob for random walk / diffusion]

\Rightarrow waiting time $P(z) \sim \frac{1}{z^{3/2}} \leftarrow \alpha = \frac{1}{2}$

gives $\langle \Delta z^2 \rangle \sim t^{1/2}$

Ref: diffusion in disordered media, [on fractals] OR

Hawlin & Ben-Avraham

Advances in Physics 36, 695, (1987)

A combination of both worlds

$P(z) \sim \frac{1}{z^{1+\alpha}}$

waiting time ~~$P(z)$~~ $P(z) \sim \frac{1}{z^{1+\alpha}}$

① normal diffusion for $\alpha_x \geq 2$ and $\alpha_t > 1$.

② Anomalous diffusion for else.

~~$x_t \sim t^{\alpha_x / (1 + \alpha_x)}$~~ $x_t \sim t^{\frac{\alpha_x}{2\alpha_x}}$

Remark: computation of the distribution is quite difficult!

How to get Stable distributions using RG?

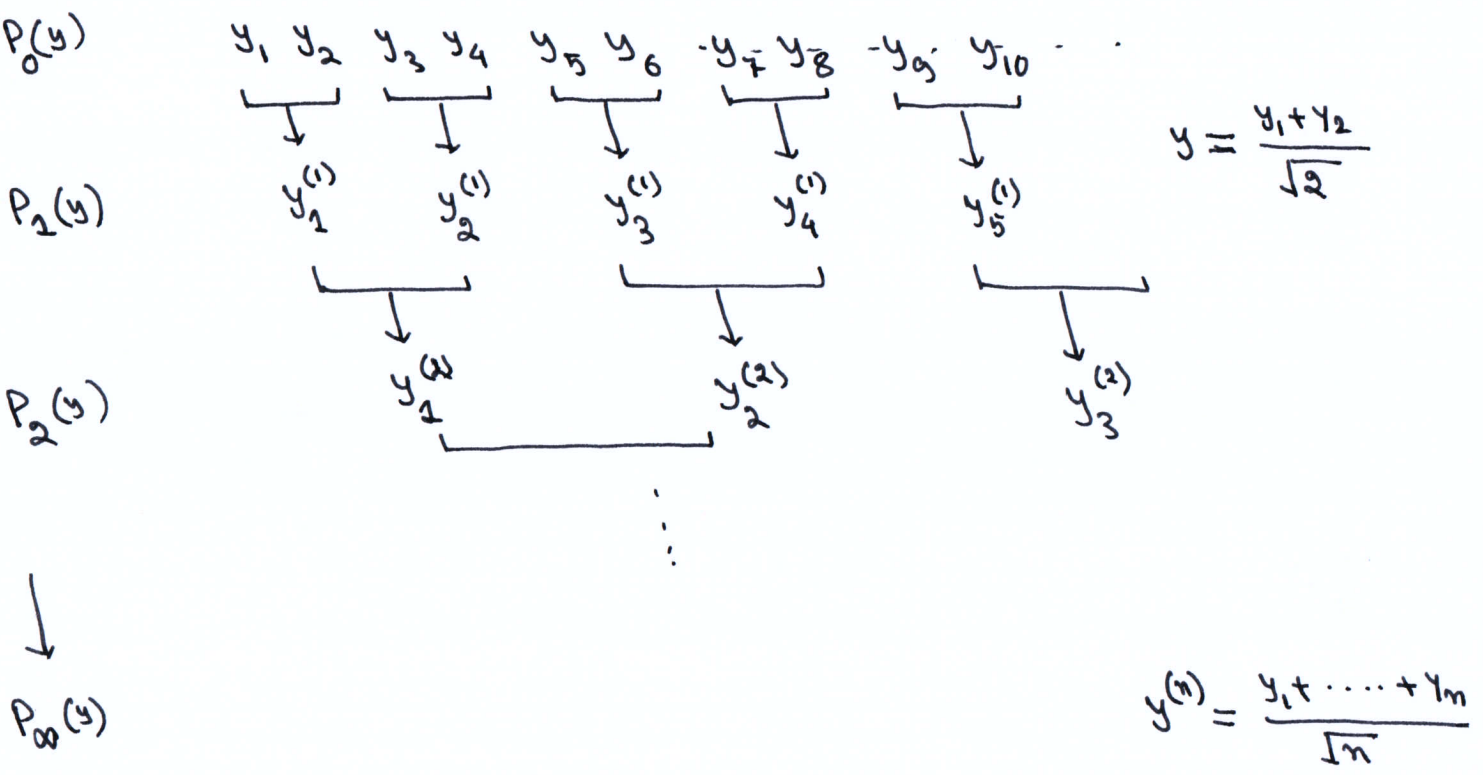
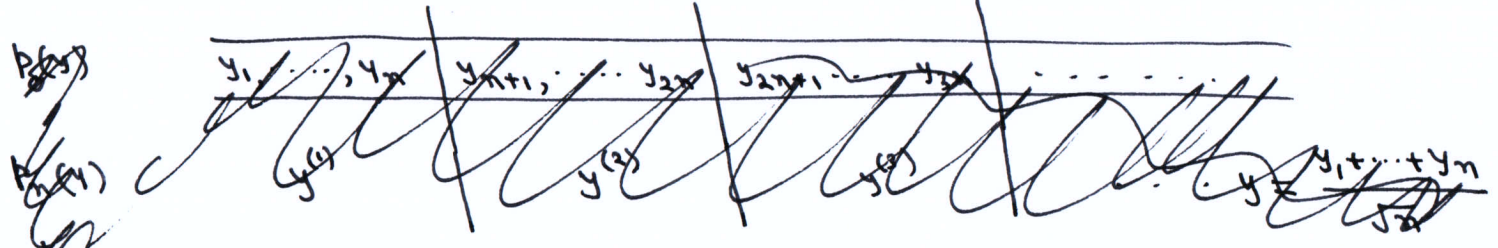
Ref 1. Exercise 12.11 in Book of Sethna.

Ref 2. "RG and probability theory", Jona-Lasinio, Phys. Rep. 352(2001) 439.

Ref 3. "An elementary renormalization-group approach to ...". Ariel Amir, J Stat mech (2020) 013214.

Term paper

Basic idea: For CLT case,



Show that

$$P_\infty(y) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{y^2}{2\sigma^2}} \quad (\text{take } \sigma^2 = 1 \text{ for simplicity})$$

(For CLT-case)

RA operation:

$$y = \frac{y_1 + y_2}{\sqrt{2}}$$

coarse-graining

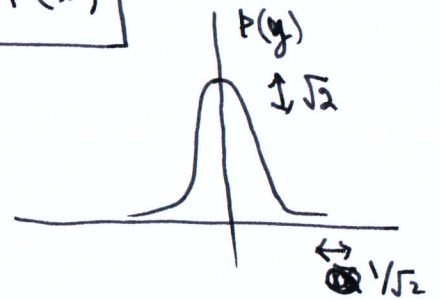
(2)

$$P_2(y) = \int dy_1 dy_2 P_0(y_1) P_0(y_2) \delta\left(\frac{y_1 + y_2}{\sqrt{2}} - y\right)$$

$$\begin{aligned} P_2(y) &= \int dy_1 dy_2 P_0(y_1) P_0(y_2) \delta\left(\frac{y_1 + y_2}{\sqrt{2}} - y\right) \\ &= \sqrt{2} \int dy_1 dy_2 P_0(y_1) P_0(y_2) \delta(y_1 + y_2 - \sqrt{2}y) \\ &= \sqrt{2} \int dy_2 P_0(\sqrt{2}y - y_2) P_0(y_2) \end{aligned}$$

$$\Rightarrow \boxed{R[P](y) = \sqrt{2} \int dx P(\sqrt{2}y - x) P(x)}$$

Rescaling



A fixed point:

$$R[P_n^*](y) = P_{n+1}^*(y) = P_n^*(y)$$

$$\boxed{P^*(y) = \sqrt{2} \int dx P^*(\sqrt{2}y - x) P^*(x)}$$

Easier to solve in Fourier-space.

$$\hat{g}(k) = \int dy e^{iky} P(y)$$

$$\Rightarrow \boxed{\hat{g}(k) = \left[\hat{g}\left(\frac{k}{\sqrt{2}}\right) \right]^2}$$

$$\Rightarrow \hat{g}(k) = e^{-\frac{k^2}{2}} \text{ is a fixed point}$$

\uparrow
 $P^*(x) = \text{Gaussian.}$

⊗ generalize:

$$y = \frac{y_1 + y_2}{b}$$

$$P^*(y) = b \int dx P^*(x) P^*(by-x)$$

b is the rescaling factor.

General: $b = 2^{1/\alpha}$

$$\hat{g}(k) = \left[\hat{g}\left(\frac{k}{2^{1/\alpha}}\right) \right]^2$$

$\Rightarrow \hat{g}(k) = e^{-|k|^\alpha}$ is a fixed point. \Leftrightarrow Lévy distr. Symmetric.

[choosing b is equivalent to ~~cho~~ keeping conserved quantities in R^d]

[Information of b comes from your parent $P_0(x)$, knowing what is your typical fluctuations]

If be chosen incorrectly, the limiting distribution may be trivial ($\delta(x)$) or it may not exist!

Is it an attractive fixed point? [Relevant, marginal, irrelevant perturbations.]

~~make~~ make a small perturbation around $P^*(y)$

$$\begin{aligned} R[P^* + \epsilon \delta P](y) &= b \int dx [P^* + \epsilon \delta P](by-x) (P^*(x) + \epsilon \delta P(x)) \\ &= P^*(y) + \epsilon [\mathcal{L} \cdot \delta P](y) \end{aligned}$$

$$[\mathcal{L} \cdot \delta P](y) = b \int dx [P^*(by-x) \cdot \delta P(x) + \delta P(by-x) \cdot P^*(x)]$$

Eigenbasis:

$$[\mathcal{L} \cdot \psi_\lambda](y) = \lambda \psi_\lambda(y)$$



Take perturbation along ^{any} one eigen basis. \circ $\delta p = \psi_\lambda$

(4)

Repeated \Rightarrow use \circ ~~$p^* + \epsilon \psi_\lambda$~~ $\rightarrow p^* + \epsilon \lambda \psi_\lambda$

~~Repeated operation~~ $\hat{R} [g^*(k) + \epsilon \hat{\psi}_\lambda(k)] = g^*(k) + \epsilon \cdot \lambda \cdot \hat{\psi}_\lambda(k)$

~~Re~~

\circ What are eigenvalues and eigenvectors. \circ Fourier basis.

$[\alpha \cdot \psi_\lambda](y) = \lambda \psi_\lambda(y)$
 $\Rightarrow [\hat{\alpha} \cdot \hat{\psi}_\lambda](k) = \lambda \hat{\psi}_\lambda(k)$

$\rightarrow 2 \cdot g^*(\frac{k}{b}) \hat{\psi}_\lambda(\frac{k}{b}) = \lambda \hat{\psi}_\lambda(k)$

Solution \circ

$\hat{\psi}_\lambda(k) = (ik)^{\alpha} g^*(k)$ with $\lambda = \frac{2}{b^\alpha}$

What it means \circ

- $\lambda_0 = 2 > 1 \leftrightarrow$ relevant \leftarrow but does not conserve probability.
- $\lambda_1 = \frac{2}{b} > 1 \leftrightarrow$ " for $b = \sqrt{2}$
- $\lambda_2 = \frac{2}{b^2} = 1 \leftrightarrow$ marginal
- $\lambda_3 = \frac{2}{b^3} < 1 \leftrightarrow$ irrelevant.
- \vdots

\rightarrow we don't perturb along this direction, because it breaks conservation that $\langle x \rangle, \langle x^2 \rangle$ are fixed!