

Scattering theory (Continued)

Next we consider full ~~rest~~ Hamiltonian with

- (a) spectrum bounded from below
- (b) lowest lying state not part of continuum
- (c) physical vacuum denoted by $|0\rangle_P$ satisfy

$$\begin{aligned} H|0\rangle_P &= 0 \\ P|0\rangle_P &= 0 \\ \langle 0|0\rangle_P &= 1 \end{aligned}$$

Introducing $\hat{H}_I \equiv \hat{H} - \hat{H}_0 \phi(x)$.

$$Z[S] \equiv \langle 0|S|0\rangle_P \rightarrow G^{(n)}(x_1 \dots x_n) = \frac{1}{i^n} \frac{\delta^n Z[S]}{\delta g(x_1) \dots \delta g(x_n)} \Big|_{g=0}$$

(i) We showed that

$G_i^{(n)}(x_1 \dots x_n) = G_F^n(x_1 \dots x_n) \rightarrow$ mindless sum of Feynman diagrams with free Hamiltonian

starting with

$$Z_F^n[S] = \frac{\langle 0|T\{\exp(-i\int_{-\infty}^{\infty} d^4x (H_I - g\phi_i))|0\rangle}{\langle 0|T\{\exp(-i\int_{-\infty}^{\infty} d^4x H_I)\}|0\rangle} \quad \text{the eigenstate}$$

(ii) LSZ

$$\langle l_1 \dots l_r | S^{-1} | k_1 \dots k_r \rangle = \prod_{a=1-r}^r \frac{l_a^2 - p_a^2}{i} \prod_{b=1-r}^r \frac{k_b^2 - m^2}{i} \mathcal{G}^{r/(r+s)}(-l_1, \dots, l_r, k_1, \dots, k_s)$$

also guaranteed for any Green's fn constructed with a field ϕ , satisfying

$$\langle 0|\phi'(x)|0\rangle = 0$$

$$\langle k|\phi'(x)|0\rangle = e^{ik \cdot x}$$

fields getting a new problematic $\rightarrow \phi'(x) = \frac{1}{\langle u|\phi(x)|0\rangle} \{ \phi(x) - \langle 0|\phi(x)|0\rangle \}$
also normalized

satisfied all

(iii) Toy model

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{\mu_0^2}{2} \phi^2 + \partial_\mu \psi^\dagger \partial^\mu \psi - m^2 \psi^\dagger \psi - g_0 \psi^\dagger \phi + \mathcal{L}_T$$

$$\mathcal{L}_T = \frac{1}{2} A \phi'^2 + \frac{B}{2} (\partial_\mu \phi)^2 - \frac{C}{2} \phi'^2 + D \partial_\mu \psi^\dagger \partial^\mu \psi - E \psi^\dagger \psi - F \psi^\dagger \phi' + \text{const.}$$

(1) we determined A perturbatively

(2) $\phi, B, C, D, E \rightarrow$ we derived the mass and normalization conditions from Green's fn
 $\rightarrow$ found spectral decomposition