

Lecture 20

Summary

spin 1 system:

gauge transformation:

$$A_\mu \rightarrow U A_\mu U^{-1} - (\partial_\mu U) U^{-1} \quad \text{for } U = e^{iX(x)}$$

given a scalar interacting with A , it must have $\phi \rightarrow U \phi$
(ϕ) under gauge transformation.

In a theory with n massless spin 1 fields interacting among themselves $\Rightarrow$ must furnish a ' n ' dimensional real representation of the Lie group.

Ex, for $SU(N)$ Lie group ($\frac{N(N-1)}{2}$) $n = N^2 - 1 \rightarrow$ adjoint representation.

$$U = e^{i\alpha_a t_a}$$

$\Phi \rightarrow$ say in the fundamental rep. of $SU(N)$

$$\rightarrow U \Phi \rightarrow$$

$$A^a \rightarrow U A^a U^{-1} - (\partial_\mu U) U^{-1} \quad \text{extra piece for massless gauge fields}$$

$$F_{\mu\nu}^a \rightarrow U F_{\mu\nu}^a U^{-1} \Rightarrow \text{no homogeneous piece.}$$

$$\text{where } F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + f_{bc}^a A_\mu^b A_\nu^c$$

$$A \equiv A^a t_a$$

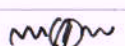
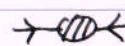
$$\Rightarrow \mathcal{L} \supset -\frac{1}{4} (F_{\mu\nu})^2 \rightarrow \text{gauge invariant.}$$

We constructed the Lagrangian of fermions and charged under a $U(1)$ gauge transformation $\rightarrow$

$\rightarrow$ introduced Feynman rules in 2-component notation.

Calculated

Green's function:

(a) 2 pt.  $\rightarrow$ 

(b) 3 point.

